

ANALYSIS & PDE

Volume 1

No. 1

2008

AN INVERSE SOURCE PROBLEM IN OPTICAL MOLECULAR IMAGING

PLAMEN STEFANOV AND GUNTHER UHLMANN



mathematical sciences publishers

AN INVERSE SOURCE PROBLEM IN OPTICAL MOLECULAR IMAGING

PLAMEN STEFANOV AND GUNTHER UHLMANN

We study the direct and an inverse source problem for the radiative transfer equation arising in optical molecular imaging. We show that for generic absorption and scattering coefficients, the direct problem is well-posed and the inverse one is uniquely solvable, with a stability estimate.

1. Introduction

We consider an inverse source problem arising in optical molecular imaging (OMI) which is currently undergoing a rapid expansion. The design of new biochemical markers that can detect faulty genes and other molecular processes allows us to detect diseases before macroscopic symptoms appear. This has been studied extensively in the bioengineering literature. See for instance [Chang et al. 1997; Contag et al. 1998; Jang et al. 2000]. Unlike higher-energetic markers used in classical nuclear imaging techniques such as single photon emission computed tomography (SPECT), positron emission tomography (PET), as well as magnetic resonance imaging (MRI), optical markers emit relatively low-frequency photons. The objective of OMI is to reconstruct the concentration of such markers from their radiations measured at the boundary of the domain. The radiations in OMI are governed by the equations of radiative transfer and the inverse problem in OMI is thus an inverse transport source problem, at least once the optical properties of the underlying medium are known. We now describe more precisely the mathematical problem.

We assume that Ω is a bounded domain of \mathbb{R}^n with smooth boundary. We will assume also that Ω is strictly convex. This is not an essential assumption since for the problem that we study, one can always push the boundary away and make it strictly convex, without losing generality. In our main result [Theorem 3.1](#), we assume that the data is given on the boundary of a larger $\Omega_1 \ni \Omega$. This is not essential for the uniqueness result but it is essential for the stability estimate [\(9\)](#).

The radiative transport equation is given by

$$\theta \cdot \nabla_x u(x, \theta) + \sigma(x, \theta)u(x, \theta) - \int_{S^{n-1}} k(x, \theta, \theta')u(x, \theta') d\theta' = f(x), \quad u|_{\partial_- S\Omega} = 0, \quad (1)$$

where the *absorption* σ and the *collision kernel* k are functions with a regularity that will be specified below. The source term f is assumed to depend on x only.

In [Section 2](#) we study the direct problem. We show that for an open and dense set of absorption and scattering coefficients the direct problem [\(1\)](#) is well-posed. See [Theorem 2.1](#) for details.

MSC2000: primary 35R30; secondary 35Q60, 35S05.

Keywords: transport equation, optical molecular imaging, inverse problem, tomography.

First author partly supported by an NSF FRG Grant No. 0554065. Second author partly supported by an NSF FRG grant No. 0554571 and a Walker Family Endowed Professorship.

The boundary measurements are modeled by

$$Xf(x, \theta) = u|_{\partial_+ S\Omega}, \quad (x, \theta) \in \partial_+ S\Omega,$$

where $u(x, \theta)$ is a solution of (1), and $\partial_+ S\Omega$ denotes the points $x \in \partial\Omega$ with direction θ pointing outwards.

In Section 3 we consider the inverse source problem, that consists in determining the source term f from measuring Xf . Notice that in the case $\sigma = k = 0$ the linear operator X is the standard X-ray transform and when $k = 0$, X is a weighted X-ray transform (see Section 2).

This inverse problem has been considered in several papers in the mathematical and engineering community [Bal and Tamasan 2007; Larsen 1975; Panchenko 1993; Sharafutdinov 1997; Siewert 1993; Yi et al. 1992]. In particular in [Bal and Tamasan 2007] it is shown that one can prove uniqueness when $k = k(x, \theta \cdot \theta')$, and k is small enough in a suitable norm. We show that for the absorption and scattering in an dense and open subset we can uniquely determine the source f from the boundary measurements. We also prove a stability estimate. See Theorem 3.1 for details.

2. The direct problem

Set

$$T_0 = \theta \cdot \nabla_x, \quad T_1 = T_0 + \sigma, \quad T = T_0 + \sigma - K,$$

where σ is viewed as the operator of multiplication by $\sigma(x, \theta)$, and K is the integral operator in (1).

Let u solve

$$Tu = f, \quad u|_{\partial_- S\Omega} = 0. \tag{2}$$

As mentioned in the introduction the operator X is the X-ray transform, if $\sigma = k = 0$,

$$Xf(x, \theta) = If(x, \theta) := \int_{\tau_-(x, \theta)}^0 f(x + t\theta) dt, \quad (x, \theta) \in \partial_+ S\Omega,$$

where $\pm\tau_{\pm}(x, \theta) \geq 0$ are defined by $(x, x + \tau_{\pm}(x, \theta)) \in \partial_{\pm} S\Omega$. We will always extend f as 0 outside Ω so we can assume that we integrate above over \mathbb{R} . If $k = 0$, then X reduces to the following weighted X-ray transform

$$Xf(x, \theta) = I_{\sigma} f(x, \theta) := \int E(x + t\theta, \theta) f(x + t\theta) dt, \quad (x, \theta) \in \partial_+ S\Omega, \tag{3}$$

where

$$E(x, \theta) = \exp\left(-\int_0^{\infty} \sigma(x + s\theta, \theta) ds\right).$$

If $\sigma > 0$ depends on x only, this is known as the attenuated X-ray transform, that is injective, and there is an explicit inversion formula (see [Novikov 2002; Arbuзов et al. 1997]).

We define the adjoint X^* of X with respect to the measure $d\Sigma$ defined above. We will view X as a perturbation of I_{σ} , and our goal is to show that X^*X is a relatively compact perturbation of $I_{\sigma}^*I_{\sigma}$.

First we will analyze the direct problem. Some conditions are needed for its well-posedness, that usually involve smallness of k with respect to σ ; see, for example, [Dautray and Lions 1993; Reed and Simon 1979] and [Sharafutdinov 1997] for the Riemannian case. In the next theorem, f is allowed to depend on θ as well and we show that the direct problem is generically solvable.

Theorem 2.1. *There exists an open and dense set of pairs $(\sigma, k) \in C^2(\bar{\Omega} \times S^{n-1}) \times C^2(\bar{\Omega} \times S^{n-1} \times S^{n-1})$, including a neighborhood of $(0, 0)$, so that for each (σ, k) in that set,*

- (a) *the direct problem (2) has a unique solution $u \in L^2(\Omega \times S^{n-1})$ for any $f \in L^2(\Omega \times S^{n-1})$ depending both on x and θ ;*
- (b) *X extends to a bounded operator*

$$X : L^2(\Omega \times S^{n-1}) \rightarrow L^2(\partial_+ S\Omega, d\Sigma).$$

Proof. We start with the analysis of the direct problem (2). In what follows, let T_0 , T_1 and T denote the operators given by (1) in $L^2(\Omega \times S^{n-1})$ with domain

$$D(T_0) = D(T_1) = D(T) = \left\{ f \in L^2(\Omega \times S^{n-1}); \theta \cdot \nabla_x u \in L^2(\Omega \times S^{n-1}), u|_{\partial_- S\Omega} = 0 \right\}.$$

We will assume here that f depends both on x and θ . Note first that the solution to the problem (2) with $k = 0$ is given by $u = T_1^{-1}f$, where

$$[T_1^{-1}f](x, \theta) = \int_{-\infty}^0 \exp\left(-\int_s^0 \sigma(x + \tau\theta, \theta) d\tau\right) f(x + s\theta, \theta) ds.$$

This follows easily from the fact that E is an integrating factor, that is, $T_0 = E^{-1}T_1E$.

Apply T_1^{-1} to both sides of (2) to get

$$u = T_1^{-1}(Ku + f).$$

We therefore see that (2) is equivalent to the integral equation

$$(\text{Id} - T_1^{-1}K)u = T_1^{-1}f.$$

Therefore, if $(\text{Id} - T_1^{-1}K)$ is invertible, (2) is uniquely solvable, and the solution is given by

$$u = T^{-1}f = (\text{Id} - T_1^{-1}K)^{-1}T_1^{-1}f. \quad (4)$$

When f depends on x only, set

$$[Jf](x, \theta) := f(x).$$

Then

$$u = T^{-1}Jf = (\text{Id} - T_1^{-1}K)^{-1}T_1^{-1}Jf.$$

Lemma 2.2. *The operator $KT_1^{-1}J : L^2(\Omega) \rightarrow L^2(\Omega \times S^{n-1})$ is compact.*

Proof. Let first f depend both on x and θ . Then

$$\begin{aligned} [KT_1^{-1}f](x, \theta) &= \int_{S^{n-1}} k(x, \theta, \theta') \int_{-\infty}^0 \exp\left(-\int_s^0 \sigma(x + \tau\theta', \theta') d\tau\right) f(x + s\theta', \theta') ds d\theta' \\ &= \int \frac{\Sigma(x, |x-y|, \frac{x-y}{|x-y|}) k(x, \theta, \frac{x-y}{|x-y|})}{|x-y|^{n-1}} f\left(y, \frac{x-y}{|x-y|}\right) dy, \end{aligned} \quad (5)$$

where

$$\Sigma(x, s, \theta') = \exp\left(-\int_{-s}^0 \sigma(x + \tau\theta', \theta') d\tau\right)$$

(we replaced s by $-s$ and then made the change $x - s\theta' = y$).

Assume now that f depends on x only, that is, we have Jf above with such an f . Then

$$[KT_1^{-1}J]f(x, \theta) = \int_{\Omega} \frac{\Sigma(x, |x-y|, \frac{x-y}{|x-y|})k(x, \theta, \frac{x-y}{|x-y|})}{|x-y|^{n-1}} f(y) dy. \quad (6)$$

The integral above is a typical singular operator with a weakly singular kernel, and an additional parameter θ ; see [Michlin and Prössdorf 1980; Stein 1970]. Under the smoothness assumptions on σ and k , it is easy to see that $\partial_{\theta}KT_1^{-1}$ and $\partial_xKT_1^{-1}$ are bounded operators; see Proposition 3.4 below. This completes the proof of the lemma. \square

Remark 2.3. The arguments above do not prove that KT_1^{-1} is compact in $L^2(\Omega \times S^{n-1})$ because there are not enough integrations in this case to apply the same arguments. Its square however is compact, as the next lemma shows. On the other hand, under appropriate smoothness assumptions on k , similar to those in Theorem 3.1 (see (9)), the operator KT_1^{-1} is compact, indeed. This is a consequence of the velocity averaging lemma that states that if $k = k(\theta')$ with k of appropriate regularity, then KT^{-1} is compact in $L^2(\Omega \times S^{n-1})$. The gained regularity then is $\frac{1}{2}$ only, not 1. Now, for $k = k(x, \theta', \theta)$ smooth enough, one can approximate K uniformly with finite sums of operators with kernels $\kappa(x)\Theta'(\theta')\Theta(\theta)$, each one of which is compact. For more details, we refer to [Mokhtar-Kharroubi 1997] and the references there.

Lemma 2.4. *The operator $KT_1^{-1}K : L^2(\Omega \times S^{n-1}) \rightarrow L^2(\Omega \times S^{n-1})$ is compact.*

Proof. Replace $f(y, \frac{x-y}{|x-y|})$ in (5) by

$$[Kf]\left(y, \frac{x-y}{|x-y|}\right) = \int_{S^{n-1}} k\left(y, \frac{x-y}{|x-y|}, \theta'\right) f(y, \theta') d\theta'.$$

Then the compactness follows from the same arguments as in Lemma 2.2. Indeed, we have

$$[KT_1^{-1}Kf](x, \theta) = \iint_{\Omega \times S^{n-1}} \frac{\alpha(x, y, |x-y|, \frac{x-y}{|x-y|}, \theta, \theta')}{|x-y|^{n-1}} f(y, \theta') dy d\theta',$$

with an obvious definition of α . In particular, all second order derivatives of α are bounded. Let $g(x, \theta, \theta')$ be the y -integral above, that is, the right hand side above becomes $\int g(x, \theta, \theta') d\theta'$. Then by Proposition 3.4 below,

$$\int_{\Omega} |\partial_x g(x, \theta, \theta')|^2 dx \leq C \int_{\Omega} |f(x, \theta')|^2 dx$$

for any θ, θ' . In particular,

$$\iint_{\Omega \times S^{n-1}} |\partial_x g(x, \theta, \theta')|^2 dx d\theta' \leq C \|f\|_{L^2}^2.$$

Then

$$\begin{aligned} \|\partial_x KT_1^{-1}Kf\|^2 &= \iint_{\Omega \times S^{n-1}} \left| \int_{S^{n-1}} \partial_x g(x, \theta, \theta') d\theta' \right|^2 dx d\theta \\ &\leq C \iint_{\Omega \times S^{n-1}} \int_{S^{n-1}} |\partial_x g(x, \theta, \theta')|^2 d\theta' dx d\theta \\ &\leq C' \|f\|_{L^2}^2. \end{aligned}$$

It is easy to see that $\partial_\theta K T_1^{-1} K f \in L^2$ as well. This, and the estimate above, imply the compactness of $K T_1^{-1} K$. \square

We proceed with the proof of part (a) of the theorem. We are looking for k so that T^{-1} exists. Consider

$$A(\lambda) = (\text{Id} - (\lambda K T_1^{-1})^2)^{-1}$$

in $L^2(\Omega \times S^{n-1})$. The operator $(K T_1^{-1})^2$ is compact, and for $\lambda = 0$, the resolvent above exists. By the analytic Fredholm theorem [Reed and Simon 1980], $A(\lambda)$ is a meromorphic family of bounded operators. In particular, it exists for all but a discrete set of λ 's. Thus for the those λ 's, the resolvent $(\text{Id} - \lambda K T_1^{-1})^{-1}$ exists and is given by

$$(\text{Id} - \lambda K T_1^{-1})^{-1} = (\text{Id} + \lambda K T_1^{-1}) A(\lambda). \tag{7}$$

Indeed, it is obvious that the operator on the right hand side above is a right inverse to $\text{Id} - \lambda K T_1^{-1}$. For $|\lambda| \ll 1$, one can use Neumann series to show that it is left inverse as well. One can use analytic continuation around the poles to show that this remains true for all λ that are not poles.

By (4), then T^{-1} exists for such λ 's and k replaced by λk . In particular, this shows that the set of such (k, σ) is dense. Standard perturbation arguments show that the set of k 's for which $\text{Id} - \lambda K T_1^{-1}$ is invertible, is open in C^0 for a fixed σ and the set of pairs $(\sigma, k) \in C^0 \times C^0$ with the same property is open, too. Since we just showed that it is dense as well in $C^0 \times C^0$, this completes the proof of (a).

We proceed with the proof of (b). For X we get (see (4)),

$$Xf = R_+ T^{-1} f = R_+ (\text{Id} - T_1^{-1} K)^{-1} T_1^{-1} f,$$

where $R_+ h = h|_{\partial_+ S\Omega}$. If f depends on x only, then

$$Xf = R_+ T^{-1} Jf = R_+ (\text{Id} - T_1^{-1} K)^{-1} T_1^{-1} Jf. \tag{8}$$

Notice first that

$$(\text{Id} - T_1^{-1} K)^{-1} T_1^{-1} = T_1^{-1} (\text{Id} - K T_1^{-1})^{-1},$$

and in particular, the resolvent on the left exists if and only if the resolvent in the right hand side does. We therefore have

$$Xf = R_+ T_1^{-1} (\text{Id} - K T_1^{-1})^{-1} Jf.$$

To prove (b), it is enough to show that

$$R_+ T_1^{-1} : L^2(\Omega \times S^{n-1}) \rightarrow L^2(\partial_+ S\Omega, d\Sigma)$$

is bounded. A straightforward computation (see also [Choulli and Stefanov 1999]) shows that

$$\int_{\partial_+ S\Omega} \int_{\tau_-(x,\theta)}^0 f(x - t\theta, \theta) dt d\Sigma = \int_{\Omega \times S^{n-1}} f(x, \theta) dx d\theta$$

for any $f \in L^1(\Omega \times S^{n-1})$. Therefore,

$$\begin{aligned} \|R_+ T_1^{-1} f\|_{L^2(\partial_+ S\Omega, d\Sigma)}^2 &= \int_{\partial_+ S\Omega} |R_+ T_1^{-1} f(x, \theta)|^2 d\Sigma \leq \int_{\partial_+ S\Omega} \left| \int_{\tau_-(x, \theta)}^0 f(x + t\theta, \theta) dt \right|^2 d\Sigma \\ &\leq \int_{\partial_+ S\Omega} \left(|\tau_-(x, \theta)| \int_{\tau_-(x, \theta)}^0 |f(x + t\theta, \theta)|^2 dt \right) d\Sigma \\ &\leq \text{diam}(\Omega) \|f\|_{L^2(\Omega \times S^{n-1})}^2. \end{aligned} \quad \square$$

3. The inverse source problem

In this section we consider the inverse source problem. The next theorem shows that for generic (σ, k) there is uniqueness and stability. As mentioned in the introduction a similar result has been proven in [Bal and Tamasan 2007] in the case where $k = k(x, \theta \cdot \theta')$, and k is small enough in a suitable norm.

Fix another strictly convex bounded domain Ω_1 so that $\Omega_1 \ni \Omega$. Extend (σ, k) with regularity as below to functions in Ω_1 with the same regularity. We chose and fix that extension as a continuous operator in those spaces. Define the operator $X_1 : L^2(\Omega_1) \rightarrow L^2(\partial_+ S\Omega_1)$ in the same way as X . We will be interested in the restriction of X_1 to functions f supported in $\bar{\Omega}$. We always extend such f as zero to $\Omega_1 \setminus \Omega$. This corresponds to taking measurements on $\partial\Omega_1$ instead of $\partial\Omega$.

Theorem 3.1. *There exists an open and dense set of pairs*

$$(\sigma, k) \in C^2(\bar{\Omega} \times S^{n-1}) \times C^2(\bar{\Omega} \times S_{\theta'}^{n-1}; C^{n+1}(S_\theta^{n-1})), \quad (9)$$

including a neighborhood of $(0, 0)$, so that for each (σ, k) in that set, the conclusions of [Theorem 2.1](#) hold in Ω_1 , and

- (a) the map X_1 is injective on $L^2(\Omega)$,
- (b) the following stability estimate holds:

$$\|f\|_{L^2(\Omega)} \leq C \|X_1^* X_1 f\|_{H^1(\Omega_1)}, \quad (10)$$

for all $f \in L^2(\Omega)$, with a constant $C > 0$ locally uniform in (σ, k) .

Remark 3.2. The smoothness requirement on k can be reduced to $k \in C^2$ if k is of a special form, like $k = \Theta(\theta)\kappa(x, \theta')$ or a finite sum of such; see (15), (16).

From now on, we will drop the subscript 1, and all operators below are as defined before but in the domain Ω_1 . We assume that (σ, k) are such that T^{-1} exists. We assume now that X is applied to f that depends on x only. For now, it is not important that f is supported in $\bar{\Omega}$; that will be needed in (20) and after that; so we apply X to functions in $L^2(\Omega_1)$. By (8),

$$X = I_\sigma + L, \quad L := R_+(-\text{Id} + (\text{Id} - T_1^{-1}K)^{-1})T_1^{-1}J \quad (11)$$

(see also (3)). Then

$$X^* X = I_\sigma^* I_\sigma + \mathcal{L}, \quad \mathcal{L} := I_\sigma^* L + L^* I_\sigma + L^* L. \quad (12)$$

In our analysis, we will apply a parametrix of $I_\sigma^* I_\sigma$ to $X^* X$. That parametrix is a first order operator. For this reason, we study $\partial_x I_\sigma^* L$.

Lemma 3.3. *The operators*

$$\partial_x I_\sigma^* L, \quad \partial_x L^* I_\sigma, \quad \partial_x L^* L$$

are compact as operators mapping $L^2(\Omega_1)$ into $L^2(\Omega_1)$.

Proof. To analyze $I_\sigma^* L$, note that L also admits the following representation

$$L = R_+ T_1^{-1} K T_1^{-1} (\text{Id} - K T_1^{-1})^{-1} J. \quad (13)$$

We need to study $I_\sigma^* R_+ T_1^{-1} K T_1^{-1} h$, where $h = h(x, \theta)$. Notice first that

$$[I_\sigma^* h](x) = \int_{S^{n-1}} \bar{E}(x, \theta) h^\sharp(x, \theta) d\theta,$$

where \bar{E} denotes complex conjugate, and h^\sharp is the extension of $h \in C(\partial_+ S\Omega_1)$ as a constant along the lines originating from x in the direction $-\theta$; see, for example, [Frigyik et al. 2008, Section 4]. In other words,

$$h^\sharp(x, \theta) = h(x + \tau_+(x, \theta), \theta).$$

Next, $R_+ T_1^{-1} h$ looks just like I_σ (see (3)) but with f there depending on θ as well. Therefore,

$$[I_\sigma^* R_+ T_1^{-1} g](x) = \int_{S^{n-1}} \bar{E}(x, \theta) \left[\int_{-\infty}^0 E(x + t\theta, \theta) g(x + t\theta, \theta) dt \right]^\sharp d\theta.$$

This yields (see [Frigyik et al. 2008] again):

$$\begin{aligned} [I_\sigma^* R_+ T_1^{-1} g](x) &= \iint_{S^{n-1}} \bar{E}(x, \theta) (Eg)(x + t\theta, \theta) d\theta dt \\ &= 2 \int_{\Omega_1} \frac{[\bar{E}(x, \frac{y-x}{|y-x|}) (Eg)(y, \frac{y-x}{|y-x|})]_{\text{even}}}{|y-x|^{n-1}} dy, \end{aligned} \quad (14)$$

where $F_{\text{even}}(x, \theta)$ is the even part of F as a function of θ . If we set $g = K T_1^{-1} h$, that would give us $I_\sigma^* R_+ T_1^{-1} K T_1^{-1} h$.

Instead of assuming (9), we will make the following weaker assumption at this point: k can be written as the infinite sum

$$k(x, \theta, \theta') = \sum_{j=1}^{\infty} \Theta_j(\theta) \kappa_j(x, \theta') \quad (15)$$

with some functions Θ_j and κ_j so that

$$\sum_{j=1}^{\infty} \|\Theta_j\|_{H^1(S^{n-1})} \|\kappa_j\|_{L^\infty(\Omega_1 \times S^{n-1})} < \infty. \quad (16)$$

One such way to do this is to choose Θ_j to be the spherical harmonics Y_j ; then κ_j are the corresponding Fourier coefficients. Then $\|Y_j\|_{H^1(S^{n-1})} \leq C(1 + \lambda_j)$, where λ_j^2 are the eigenvalues of the positive Laplacian on S^{n-1} . Since $\lambda_j = O(j^{1/(n-1)})$, for the uniform convergence of (15) it is enough to have $\|\kappa_j\|_{L^\infty} \leq C(1 + \lambda_j)^{-n-\varepsilon}$ with $\varepsilon > 0$. This would be guaranteed if $k \in L^\infty(\Omega_1 \times S_{\theta'}^{n-1}; C_\theta^{n+1}(S^{n-1}))$ by standard integration by parts arguments. Therefore, the hypothesis (9) of the theorem implies (15) and (16).

Under this assumption, for $K_j T_1^{-1} h$, where K_j has kernel $\Theta_j \kappa_j$, we have (see (5)):

$$\begin{aligned} [K_j T_1^{-1} h](x, \theta) &= \Theta_j(\theta) [B_j h](x), \\ B_j h(x) &:= \int_{\Omega_1} \frac{\Sigma(x, |x-y|, \frac{x-y}{|x-y|}) \kappa_j(x, \frac{x-y}{|x-y|})}{|x-y|^{n-1}} h\left(y, \frac{x-y}{|x-y|}\right) dy. \end{aligned} \quad (17)$$

We claim now that $B_j(\text{Id} - K T_1^{-1})^{-1} J : L^2(\Omega_1) \rightarrow L^2(\Omega_1)$ is compact. We have

$$(\text{Id} - K T_1^{-1})^{-1} J = J + (\text{Id} - K T_1^{-1})^{-1} K T_1^{-1} J.$$

By Lemma 2.2, the second term on the right is compact. Therefore, it remains to show that $B_j J$ is compact. Observe that $B_j J h$ is given by (17) with $h = h(x)$. The compactness then follows from Proposition 3.4, assuming that $\kappa_j \in C^2$. On the other hand, $B_j J$ is compact under the assumption that $\kappa_j \in L^\infty$ only, by [Michlin and Prössdorf 1980, Theorem VII.3.3]. Moreover, its norm is bounded by $C \|\kappa_j\|_{L^\infty}$.

We can now write

$$\begin{aligned} \partial_x I_\sigma^* L &= \partial_x I_\sigma^* R_+ T_1^{-1} K T_1^{-1} (\text{Id} - K T_1^{-1})^{-1} J \\ &= \sum_{j=1}^{\infty} (\partial_x I_\sigma^* R_+ T_1^{-1} \Theta_j J) (B_j (\text{Id} - K T_1^{-1})^{-1} J). \end{aligned} \quad (18)$$

We notice first that $\partial_x I_\sigma^* R_+ T_1^{-1} \Theta_j J : L^2(\Omega_1) \rightarrow L^2(\Omega_1)$ is bounded by Proposition 3.4 (b), compare to (14), with a norm bounded by $C \|\sigma\|_{C^2} \|\Theta_j\|_{H^1}$. The operator $B_j (\text{Id} - K T_1^{-1})^{-1} J$ on the right is compact, as we have just seen. Therefore, each summand in the right hand side of (18) is a compact operator with a norm not exceeding $C \|\Theta_j\|_{H^1} \|\kappa_j\|_{L^\infty}$, where C depends on σ as well. Then the series in (18) converges uniformly by (16). Under this condition, $\partial_x I_\sigma^* L$ is compact.

To analyze $\partial_x L^* L$, we will follow the proof above. It is enough to show that $\partial_x L^* R_+ T_1^{-1} \Theta_j J : L^2(\Omega) \rightarrow L^2(\Omega_1)$ is bounded. We have (see (13)):

$$\begin{aligned} \partial_x L^* R_+ T_1^{-1} \Theta_j J &= \partial_x (R_+ T_1^{-1} (\text{Id} - K T_1^{-1})^{-1} K T_1^{-1} J)^* R_+ T_1^{-1} \Theta_j J \\ &= \partial_x (K T_1^{-1} J)^* (R_+ T_1^{-1} (\text{Id} - K T_1^{-1})^{-1})^* R_+ T_1^{-1} \Theta_j J. \end{aligned} \quad (19)$$

Since $R_+ T_1^{-1}$ is bounded, it remains to show that the operator $\partial_x (K T_1^{-1} J)^* : L^2(\Omega_1 \times S^{n-1}) \rightarrow L^2(\Omega)$ is bounded, as well. The kernel of the latter is (see (6))

$$(x, (y, \theta)) \mapsto \partial_x \frac{\Sigma(y, |y-x|, \frac{y-x}{|y-x|}) k(y, \theta, \frac{y-x}{|y-x|})}{|y-x|^{n-1}}.$$

Then the boundedness of $\partial_x (K T_1^{-1} J)^*$ then follows as in Lemma 2.4.

Finally, the fact that $\partial_x L^* I_\sigma$ is bounded follows from the proof for $\partial_x L^* L$. Indeed,

$$\partial_x L^* I_\sigma = \partial_x L^* R_+ T_1^{-1} J,$$

compare with (19), where we can set $\Theta_j = 1$.

This completes the proof of Lemma 3.3. □

Proof of Theorem 3.1. We return to the analysis of the operator X^*X ; see (12). We showed in Lemma 3.3 that, up to a relative compact operator, X^*X coincides with $I_\sigma^*I_\sigma$. Assume that σ and k are C^∞ . Let Q be a parametrix (of order 1) to the elliptic Ψ DO $I_\sigma^*I_\sigma$ in Ω_1 . We restrict the image of Q to $L^2(\Omega)$, that is, we view Q as an operator $Q : H^1(\Omega_1) \rightarrow L^2(\Omega)$. Then for any f supported in $\bar{\Omega}$, we have

$$QI_\sigma^*I_\sigma f = f + K_1 f, \quad (20)$$

where K_1 is of order -1 near Ω . Apply Q to X^*X to get

$$QX^*Xf = f + K_2 f, \quad K_2 := K_1 + Q\mathcal{L}. \quad (21)$$

Then $K_2 : L^2(\Omega) \rightarrow L^2(\Omega)$ is compact. We get that the problem of inverting X^*X is reduced to a Fredholm equation. We will show that it is generically solvable, as in the theorem.

We show first that the set of pairs for which X is injective is dense.

By the results of [Frigyik et al. 2008, Theorems 1 and 2], if σ is real analytic in a $\bar{\Omega}_1$, then I_σ is injective, and therefore $I_\sigma^*I_\sigma$, is injective as well. Moreover, for a small $C^2(\bar{\Omega})$, perturbation preserves that property. Actually, the remark after [Frigyik et al. 2008, Theorem 2] shows that this is true even for small enough C^1 perturbations. Fix σ real analytic in $\bar{\Omega}_1$. Fix k as well so that (σ, k) belongs to the generic set in Theorem 2.1, related to Ω_1 , and the regularity assumption (9) is satisfied. That can be done for an open dense set of k 's by the proof of Theorem 2.1. Consider X related to $(\sigma, \lambda k)$ with λ belonging to some complex neighborhood \mathcal{C} of $[0, 1]$. The operator K_2 in (21) depends meromorphically on $\lambda \in \mathcal{C}$. Indeed, K_1 is related to $(\sigma, 0)$ (that is, to $\lambda = 0$), and is therefore independent of λ . The parametrix Q is also independent of λ . The analysis above shows that \mathcal{L} is a meromorphic function of λ because L has that property; see (7) and (11). For $\lambda = 0$, we have $\mathcal{L} = 0$, and then $K_2 = K_1$. By adding a finite rank operator to Q , we can arrange that $\text{Id} + K_1$ (see (20)) is injective; see also the proof of [Stefanov and Uhlmann 2005, Proposition 4]. We can then apply the analytic Fredholm theorem again in \mathcal{C} with the poles of $(\text{Id} - \lambda K)^{-1}T_1^{-1}$ removed. The latter is a connected set, containing $\lambda = 0$ and $\lambda = 1$. The analytic Fredholm theorem then implies that QX^*X is invertible for all λ in that set with the possible exception of a discrete set. In particular, there are λ 's as close to $\lambda = 1$ as needed with that property. For those λ 's, X^*X and X are injective as well. This shows that there is a dense set of pairs (σ, k) in the space (9) so that X is injective. Let us call that set \mathcal{U} .

We show next that for (σ, k) in some neighborhood of \mathcal{U} , X is still injective.

Let $(k, \sigma) \in \mathcal{U}$. Then $X : L^2(\Omega) \rightarrow L^2(\partial\Omega_1, d\Sigma)$ is injective. Then $X^*X : L^2(\Omega) \rightarrow H^1(\Omega_1)$ is injective as well, as an integration by parts shows. By adding a finite rank operator to Q , we can arrange that $\text{Id} + K_1$ (see (20)) is injective, as above. Then $\text{Id} + K_1$ is invertible on $L^2(\Omega)$, and we deduce that (10) holds.

The analysis above implies that the norm $\|X^*X\|_{L^2(\Omega) \rightarrow H^1(\Omega_1)}$ depends continuously on (σ, k) as in (9). Therefore, we can perturb (σ, k) , and (10) would remain true because the perturbation of the right hand side will be absorbed by the left hand side. On the other hand, injectivity of X^*X implies injectivity of X .

This proves that the set of pairs (σ, k) , for which X is injective, is open subset of the (generic) set of pairs, for which the direct problem is guaranteed to be uniquely solvable by Theorem 2.1. Moreover, (10) holds with C locally uniform.

This completes the proof of Theorem 3.1. □

In the proof of the theorem, we used the following proposition about singular operators.

Proposition 3.4. *Let A be the operator*

$$[Af](x) = \int \frac{\alpha(x, y, |x-y|, \frac{x-y}{|x-y|})}{|x-y|^{n-1}} f(y) dy$$

with $\alpha(x, y, r, \theta)$ compactly supported in x, y . Then

- (a) if $\alpha \in C^2$, then $A : L^2 \rightarrow H^1$ is continuous with a norm not exceeding $C\|\alpha\|_{C^2}$;
- (b) let $\alpha(x, y, r, \theta) = \alpha'(x, y, r, \theta)\phi(\theta)$ then

$$\|A\|_{L^2 \rightarrow H^1} \leq C\|\alpha'\|_{C^2}\|\phi\|_{H^1(S^{n-1})}.$$

Proof. We recall some facts about the Calderón–Zygmund theory of singular operators; see [Michlin and Prössdorf 1980]. First, if K is an integral operator with singular kernel $k(x, y) = \phi(x, \theta)r^{-n}$, where $\theta = \frac{x-y}{|x-y|}$, $r = |x-y|$, and if the “characteristic” ϕ has a mean value 0 as a function of θ , for any x , then K is a well-defined operator on test functions, where the integral has to be understood in the principle value sense. Moreover, K extends to a bounded operator to L^2 with a norm not exceeding $C \sup_x \|\phi(x, \cdot)\|_{L^2(S^{n-1})}$; see [Michlin and Prössdorf 1980, Theorem XI.3.1]. The characteristic ϕ does not need to have zero mean value in θ but then the integral has to be considered as a convolution in distribution sense. The latter is well defined because the Fourier transform of the kernel with respect to the variable $z = r\theta$ is homogeneous of order 0, thus bounded.

Also, see [Michlin and Prössdorf 1980, Theorem XI.11.1]; if B is an operator with a weakly singular kernel $\psi(x, \theta)r^{-n+1}$, then $\partial_x B$ is an integral operator with singular kernel $\partial_x[\psi(x, \theta)r^{-n+1}]$. The latter, up to a weakly singular operator, has a singular kernel of the type ϕr^{-n} , and the integration is again understood in the principle value sense; see the next paragraph. In particular, the zero mean value condition is automatically satisfied.

In our case, $\beta = \alpha$ depends on y and r as well. Assume first that it does not, that is, B is as above. Extend β as a homogeneous function of θ of order 0 near S^{n-1} . Then

$$\begin{aligned} \partial_{x_i} \frac{\beta(x, \theta)}{r^{n-1}} &= (1-n) \frac{\theta_i}{r^n} \beta + \sum_j \frac{\frac{\partial \beta}{\partial \theta_j}}{r^{n-1}} \frac{\partial \theta_j}{\partial x_i} + \frac{\beta_{x_i}(x, \theta)}{r^{n-1}} \\ &= (1-n) \frac{\theta_i}{r^n} \beta + \sum_j \frac{\frac{\partial \beta}{\partial \theta_j}}{r^n} (\delta_{ij} - \theta_i \theta_j) + \frac{\beta_{x_i}(x, \theta)}{r^{n-1}} \\ &= \frac{(1-n)\theta_i \beta + \frac{\partial \beta}{\partial \theta_i}}{r^n} + \frac{\beta_{x_i}(x, \theta)}{r^{n-1}}. \end{aligned} \tag{22}$$

We used the fact that $\sum_j \theta_j \frac{\partial \beta}{\partial \theta_j} = 0$ because β is homogeneous of order 0 in θ . It is not hard to show that the “characteristic”

$$\phi(x, \theta) = (1-n)\theta_i \beta + \frac{\partial \beta}{\partial \theta_i}$$

has zero mean over S^{n-1} ; see [Michlin and Prössdorf 1980, p. 243]. In this particular case where $\alpha(x, y, \theta) = \beta(x, \theta)$, independent of y and r , statement (a) can be proven as follows. Choose a finite

atlas of charts for S^{n-1} so that for each chart, $n-1$ of the θ coordinates (that we keep fixed in \mathbb{R}^n) can be chosen as local coordinates. By rearranging the x , and respectively, the θ coordinates, in each fixed chart, we can assume that they are $\theta' = (\theta_1, \dots, \theta_{n-1})$. Then $\frac{\partial \beta}{\partial \theta_n} = -\sum_{i=1}^{n-1} \frac{\partial \beta}{\partial \theta_i}$. Then in (22), we have derivatives of β with respect to θ' (and x) with smooth coefficients. The contribution of the first term then can be estimated by the Calderón–Zygmund theorem. The second term is a kernel of a weakly singular operator. The following criterion can be applied to it: If K has an integral kernel $k(x, y)$ with the property

$$\sup_x \int |k(x, y)| dx \leq M, \quad \sup_y \int |k(x, y)| dy \leq M, \quad (23)$$

then K is bounded in L^2 with a norm not exceeding M [Taylor 1996, Proposition A.5.1].

This proves (a) for $\alpha = \beta$.

To replace $\beta(x, \theta)$ above by $\alpha(x, y, \theta)$, write $\alpha(x, y, r, \theta) = \alpha(x, x, 0, \theta) + r\gamma(x, y, r, \theta)$.

To prove (b), write first as above,

$$\alpha(x, y, r, \theta) = \beta'(x, \theta)\phi(\theta) + r\gamma(x, y, r, \theta)\phi(\theta), \quad \beta'(x, \theta) := \alpha_1(x, x, 0, \theta),$$

where $\gamma \in C^1$. Notice then that in (22), with $\beta = \beta'\phi$, we have

$$(1-n)\theta_i\beta + \frac{\partial \beta}{\partial \theta_i} = (1-n)\theta_i\beta'\phi + \phi \frac{\partial \beta'}{\partial \theta_i} + \beta' \frac{\partial \phi}{\partial \theta_i}.$$

Choosing local coordinates as above, and applying the Calderón–Zygmund theorem again, we get that the first term above contributes a singular operator with a norm not exceeding $\|\alpha_1\|_{C^1}\|\phi\|_{H^1}$. The second term $r\gamma$ generates an operator with a kernel $\gamma(x, y, r, \theta)\phi(\theta)r^{-n+2}$. Differentiate with respect to x , and we still get a weakly singular operator whose norm can be estimated as in (23) to give a norm not exceeding $\|\gamma_1\|_{C^1}\|\phi\|_{H^1}$. \square

Remark 3.5. The only second order derivatives of α that were needed in the proof of (a) were $\partial_{(x,\theta)}\partial_{(t,r)}\alpha$.

References

- [Arbuzov et al. 1997] È. V. Arbuzov, A. L. Bukhgeim, and S. G. Kazantsev, “Two-dimensional tomography problems and the theory of A -analytic functions”, pp. 6–20 in *Алгебра, геометрия, анализ и математическая физика: 10-я сиб. шк.* (Novosibirsk, 1996), Math. Inst. Russian Acad. of Sciences, Novosibirsk, 1997. In Russian; translated in *Siberian Adv. Math.* **8**:4 (1998), 1–20. [MR 2000a:44002](#) [Zbl 0912.35154](#)
- [Bal and Tamasan 2007] G. Bal and A. Tamasan, “Inverse source problems in transport equations”, *SIAM J. Math. Anal.* **39**:1 (2007), 57–76. [MR 2318375](#) [Zbl 1135.35094](#)
- [Chang et al. 1997] J. Chang, R. L. Barbour, and H. Graber, “Imaging of fluorescence in highly scattering media”, *IEEE Trans. Biomed. Eng.* **44**:9 (1997), 810–822.
- [Choulli and Stefanov 1999] M. Choulli and P. Stefanov, “An inverse boundary value problem for the stationary transport equation”, *Osaka J. Math.* **36**:1 (1999), 87–104. [MR 2000g:35228](#) [Zbl 0998.35064](#)
- [Contag et al. 1998] P. R. Contag, I. N. Olomu, D. K. Stevenson, and C. H. Contag, “Bioluminescent indicators in living mammals”, *Nature Medicine* **4**:2 (1998), 245–247.
- [Dautray and Lions 1993] R. Dautray and J.-L. Lions, *Mathematical analysis and numerical methods for science and technology*, vol. 6, Springer, Berlin, 1993. [MR 95e:00001](#) [Zbl 0802.35001](#)
- [Frigyik et al. 2008] B. Frigyik, P. Stefanov, and G. Uhlmann, “The X-ray transform for a generic family of curves and weights”, *J. Geom. Anal.* **18**:1 (2008), 89–108. [MR 2008j:53128](#)

- [Jang et al. 2000] H. B. Jang, S. Ramesh, and M. Barlett, “Combined optical and fluorescence imaging for breast cancer detection and diagnosis”, *Crit. Rev. Biom. Eng.* **28**:3-4 (2000), 371–375.
- [Larsen 1975] E. W. Larsen, “The inverse source problem in radiative transfer”, *J. Quant. Spectrosc. Radiat. Transfer* **15**:1 (1975), 1–5. [MR 54 #2128](#)
- [Michlin and Prössdorf 1980] S. G. Michlin and S. Prössdorf, *Singuläre Integraloperatoren*, Mathematische Monographien **52**, Akademie-Verlag, Berlin, 1980. [MR 82c:45002](#) [Zbl 0442.47027](#)
- [Mokhtar-Kharroubi 1997] M. Mokhtar-Kharroubi, *Mathematical topics in neutron transport theory*, Series on Advances in Mathematics for Applied Sciences **46**, World Scientific, River Edge, NJ, 1997. [MR 99c:82078](#) [Zbl 0997.82047](#)
- [Novikov 2002] R. G. Novikov, “An inversion formula for the attenuated X-ray transformation”, *Ark. Mat.* **40**:1 (2002), 145–167. [MR 2003k:44004](#) [Zbl 1036.53056](#)
- [Panchenko 1993] A. N. Panchenko, “Inverse source problem of radiative transfer: a special case of the attenuated Radon transform”, *Inverse Problems* **9**:2 (1993), 321–337. [MR 93k:35271](#) [Zbl 0829.45010](#)
- [Reed and Simon 1979] M. Reed and B. Simon, *Methods of modern mathematical physics, III: Scattering theory*, Academic Press, New York, 1979. [MR 80m:81085](#) [Zbl 0405.47007](#)
- [Reed and Simon 1980] M. Reed and B. Simon, *Methods of modern mathematical physics, I: Functional analysis*, 2nd ed., Academic Press, New York, 1980. [MR 85e:46002](#) [Zbl 0459.46001](#)
- [Sharafutdinov 1997] V. A. Sharafutdinov, “Inverse problem of determining a source in the stationary transport equation on a Riemannian manifold”, *Zap. Nauchn. Sem. S.-Peterburg. Otdel. Mat. Inst. Steklov. (POMI)* **239** (1997), 236–242. In Russian; translated in *J. Math. Sci. New York* **96**:4 (1999), 3430–3433. [MR 2000g:58056](#) [Zbl 0935.58013](#)
- [Siewert 1993] C. E. Siewert, “An inverse source problem in radiative transfer”, *J. Quant. Spect. Radiat. Transfer* **50** (1993), 603–609.
- [Stefanov and Uhlmann 2005] P. Stefanov and G. Uhlmann, “Boundary rigidity and stability for generic simple metrics”, *J. Amer. Math. Soc.* **18**:4 (2005), 975–1003. [MR 2006h:53031](#) [Zbl 1079.53061](#)
- [Stein 1970] E. M. Stein, *Singular integrals and differentiability properties of functions*, Princeton Mathematical Series **30**, Princeton University Press, Princeton, N.J., 1970. [MR 44 #7280](#) [Zbl 0207.13501](#)
- [Taylor 1996] M. E. Taylor, *Partial differential equations, I*, Applied Mathematical Sciences **115**, Springer, New York, 1996. [MR 98b:35002b](#) [Zbl 0869.35002](#)
- [Yi et al. 1992] H. C. Yi, R. Sanchez, and N. J. McCormick, “Bioluminescence estimation from ocean in situ irradiances”, *Applied Optics* **31** (1992), 822–830.

Received 28 Mar 2008. Revised 27 Jun 2008. Accepted 12 Aug 2008.

PLAMEN STEFANOV: stefanov@math.purdue.edu

Department of Mathematics, Purdue University, 150 N. University Street, West Lafayette, IN 47907-2067, United States
<http://www.math.purdue.edu/~stefanov/>

GUNTHER UHLMANN: gunther@math.washington.edu

Department of Mathematics, University of Washington, C-449 Padelford Hall, Box 354350, Seattle, WA 98195-4350, United States
<http://www.math.washington.edu/~gunther/>