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We improve on several weighted inequalities of recent interest by replacing a part of the A_p bounds by weaker A_∞ estimates involving Wilson’s A_∞ constant

$$[w]_{A_\infty}' := \sup_Q \frac{1}{w(Q)} \int_Q M(w\chi_Q).$$

In particular, we show the following improvement of the first author’s A_2 theorem for Calderón–Zygmund operators T :

$$\|T\|_{\mathcal{B}(L^2(w))} \leq c_T [w]_{A_2}'^{1/2} ([w]_{A_\infty}' + [w^{-1}]_{A_\infty}')^{1/2}.$$

Corresponding A_p type results are obtained from a new extrapolation theorem with appropriate mixed A_p - A_∞ bounds. This uses new two-weight estimates for the maximal function, which improve on Buckley’s classical bound.

We also derive mixed A_1 - A_∞ type results of Lerner, Ombrosi and Pérez (2009) of the form

$$\begin{aligned} \|T\|_{\mathcal{B}(L^p(w))} &\leq c p p' [w]_{A_1}'^{1/p} ([w]_{A_\infty}')^{1/p'}, \quad 1 < p < \infty, \\ \|Tf\|_{L^{1,\infty}(w)} &\leq c [w]_{A_1} \log(e + [w]_{A_\infty}') \|f\|_{L^1(w)}. \end{aligned}$$

An estimate dual to the last one is also found, as well as new bounds for commutators of singular integrals.

1. Introduction and statements of the main results

The weights w for which the usual operators T of classical analysis (like the Hardy–Littlewood maximal operator, the Hilbert transform, and general classes of Calderón–Zygmund operators) act boundedly on $L^p(w)$ were identified in works of Muckenhoupt [1972], Hunt, Muckenhoupt and Wheeden [Hunt et al. 1973], and Coifman and Fefferman [1974]. This class consists of the Muckenhoupt A_p weights, defined by the condition that (see [García-Cuerva and Rubio de Francia 1985])

$$[w]_{A_p} := \sup_Q \left(\int_Q w \right) \left(\int_Q w^{-1/(p-1)} \right)^{p-1} < \infty, \quad p \in (1, \infty),$$

where the supremum is over all cubes in \mathbb{R}^d . Hence it is shown for any of these important operators T , whether it is linear or not, that

$$\|T\|_{\mathcal{B}(L^p(w))} := \sup_{f \neq 0} \frac{\|Tf\|_{L^p(w)}}{\|f\|_{L^p(w)}}$$

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is finite if and only if $[w]_{A_p} < \infty$.

It is a natural question to look for optimal quantitative bounds of $\|T\|_{\mathcal{B}(L^p(w))}$ in terms of $[w]_{A_p}$. The first author who studied that question was S. Buckley [1993], who proved

$$\|M\|_{\mathcal{B}(L^p(w))} \leq c_{p,d} [w]_{A_p}^{1/(p-1)}, \quad 1 < p < \infty, \quad (1.1)$$

where M is the usual Hardy–Littlewood maximal function on \mathbb{R}^d . However, there has been a great impetus toward finding such precise dependence for more singular operators after the work of Astala, Iwaniec and Saksman [Astala et al. 2001], due to the connections with sharp regularity results for solutions to the Beltrami equation. The key fact was to prove that the operator norm of the Beurling–Ahlfors transform on $L^2(w)$ grows linearly in terms of the A_2 constant of w . This was proved by S. Petermichl and A. Volberg [2002] and by Petermichl [2007; 2008] for the Hilbert transform and the Riesz transforms. To be precise, in these papers it has been shown that if T is any of these operators, then

$$\|T\|_{\mathcal{B}(L^p(w))} \leq c_{p,T} [w]_{A_p}^{\max\{1, 1/(p-1)\}}. \quad (1.2)$$

The exponents are optimal in the sense that the exponent cannot be replaced by any smaller quantity. It was conjectured then that the same estimate holds for any Calderón–Zygmund operator T . This was proven first for special classes of integral transforms in [Cruz-Uribe et al. 2010; Lacey et al. 2010b], and eventually for general Calderón–Zygmund operators by the first author in [Hytönen 2012], using the main result from [Pérez et al. 2010], where it is shown that a weak type estimate is enough to prove the strong type. A direct proof of this result can be found in [Hytönen et al. 2010]. Other related works are [Cruz-Uribe et al. 2012; Hytönen et al. 2011; Lacey et al. 2010a; Lerner 2011; Vagharshakyan 2010].

The main purpose of this paper is to show that these results can be further improved. To do this, we recall the following definitions of the A_∞ constant of a weight w . First, there is the notion introduced by Hruščev [1984] (see also [García-Cuerva and Rubio de Francia 1985]),

$$[w]_{A_\infty} := \sup_Q \left(\int_Q w \right) \exp \left(\int_Q \log w^{-1} \right);$$

and second, there is the (as it turns out) smaller quantity that appeared with a different notation in the work of Wilson [1987; 1989; 2008] and was recently termed the “ A_∞ constant” by Lerner [2011, Section 5.5]:

$$[w]'_{A_\infty} := \sup_Q \frac{1}{w(Q)} \int_Q M(w \chi_Q).$$

Observe that

$$c_d [w]'_{A_\infty} \leq [w]_{A_\infty} \leq [w]_{A_p} \quad \text{for all } p \in [1, \infty),$$

where the second estimate is elementary, and the first will be checked in Proposition 2.2. While the constant $[w]_{A_\infty}$ is more widely used in the literature, and is also more flexible for our purposes, it is of interest to observe situations where the smaller constant $[w]'_{A_\infty}$ is sufficient for our estimates, thereby giving a sharper bound.

Now, if $\sigma = w^{-1/(p-1)}$ is the dual weight of w , we also have $[\sigma]_{A_\infty}^{p-1} \leq [\sigma]_{A_{p'}}^{p-1} = [w]_{A_p}$. The point here is that these quantities can be much smaller for some classes of weights. Our results will be of the form

$$\|T\|_{\mathfrak{B}(L^p(w))} \leq c_{p,T} \sum [w]_{A_p}^{\alpha(p)} [w]_{A_\infty}^{\beta(p)} [\sigma]_{A_\infty}^{(p-1)\gamma(p)},$$

sometimes even with the smaller $[\]'_{A_\infty}$ constant instead of $[\]_{A_\infty}$, where the sum is over at most two triplets (α, β, γ) , and the exponents satisfy $\alpha(p) + \beta(p) + \gamma(p) = \tau(p)$, where $\tau(p)$ is the exponent from the earlier sharp results. However, we will have $\alpha(p) < \tau(p)$, which shows that part of the necessary A_p control may in fact be replaced by weaker A_∞ control.

We now turn to a more detailed discussion of our results.

1A. The A_2 theory for Calderón–Zygmund operators. Our main result for Calderón–Zygmund operators is the following:

Theorem 1.3. *Let T be a Calderón–Zygmund operator and let $w \in A_2$ and $\sigma = w^{-1}$. Then there is a constant $c = c_{d,T}$ such that*

$$\|T\|_{\mathfrak{B}(L^2(w))} \leq c [w]_{A_2}^{1/2} ([w]'_{A_\infty} + [w^{-1}]'_{A_\infty})^{1/2} \leq c [w]_{A_2}^{1/2} ([w]_{A_\infty} + [w^{-1}]_{A_\infty})^{1/2}. \tag{1.4}$$

We will prove this by following the approach from [Hytönen 2012; Hytönen et al. 2010] to the A_2 theorem $\|T\|_{\mathfrak{B}(L^2(w))} \leq c_T [w]_{A_2}$, and modifying the proof at some critical points. Indeed, the original argument uses the A_2 property basically *twice*, each time producing the factor $[w]_{A_2}^{1/2}$, and it suffices to observe that only the A_∞ property is actually needed in one of these estimates.

An interesting consequence of this theorem is the following: for any fixed Calderón–Zygmund operator T , we have

$$\inf_{w \in A_2} \frac{\|T\|_{\mathfrak{B}(L^2(w))}}{[w]_{A_2}} = 0. \tag{1.5}$$

This follows once we describe, in Section 8, a family of weights $w \in A_2$ for which both $[w]'_{A_\infty}$ and $[\sigma]'_{A_\infty}$ (and even $[w]_{A_\infty}$ and $[\sigma]_{A_\infty}$) grow slower than $[w]_{A_2}$. In particular, the “reverse A_2 conjecture” $[w]_{A_2} \leq c_T \|T\|_{\mathfrak{B}(L^2(w))}$ is false.

1B. The maximal function. We next discuss the sharp weighted bounds for the Hardy–Littlewood maximal function, which we first do in a two-weight setting. We need a new two-weight constant $B_p[w, \sigma]$ defined by the functional

$$B_p[w, \sigma] := \sup_Q \left(\int_Q w \right) \left(\int_Q \sigma \right)^p \exp \left(\int_Q \log \sigma^{-1} \right), \tag{1.6}$$

which clearly satisfies

$$[w]_{A_p} \leq B_p[w, \sigma] \leq [w]_{A_p} [\sigma]_{A_\infty}.$$

Theorem 1.7. *Let M be the Hardy–Littlewood maximal operator and let $p \in (1, \infty)$. Then we have the estimates*

$$\|M(f\sigma)\|_{L^p(w)} \leq C_d \cdot p' \cdot (B_p[w, \sigma])^{1/p} \|f\|_{L^p(\sigma)} \tag{1.8}$$

and

$$\|M(f\sigma)\|_{L^p(w)} \leq C_d \cdot p' \cdot ([w]_{A_p} [\sigma]_{A_\infty}')^{1/p} \|f\|_{L^p(\sigma)}. \tag{1.9}$$

We refer to [Section 4](#) for the proof and for more information and background about this two-weight estimate for M . By a well-known change-of-weight argument, [\(1.9\)](#) implies:

Corollary 1.10. *For M and p as above, and $\sigma = w^{-1/(p-1)}$, we have*

$$\|M\|_{\mathfrak{B}(L^p(w))} \leq C_d \cdot p' \cdot ([w]_{A_p} [\sigma]_{A_\infty}')^{1/p}. \tag{1.11}$$

This improves on Buckley’s theorem $\|M\|_{\mathfrak{B}(L^p(w))} \leq C_d \cdot p' \cdot [w]_{A_p}^{1/(p-1)}$. [Corollary 1.10](#), at least for $p = 2$, was also independently discovered by A. Lerner and S. Ombrosi [\[2008\]](#).

1C. The A_1 theory for Calderón–Zygmund operators. It is an interesting fact that if we assume that the weight satisfies the stronger condition $w \in A_1$, then the estimate [\(1.2\)](#) can be considerably improved. Indeed, if T is any Calderón–Zygmund operator, then T is of course bounded on $L^p(w)$, because $A_1 \subset A_p$, but with a much better bound, namely

$$\|T\|_{\mathfrak{B}(L^p(w))} \leq c p p' [w]_{A_1}, \quad 1 < p < \infty. \tag{1.12}$$

Observe that the dependence on the A_1 constant is linear for any p , while in the A_p case it is highly nonlinear for $1 < p < 2$; see [\(1.2\)](#). The result is sharp both in terms of the dependence on $[w]_{A_1}$, and in terms of the dependence on p when taking $w = 1$ by the classical theory. This fact was used to get the following endpoint result:

$$\|Tf\|_{L^{1,\infty}(w)} \leq c [w]_{A_1} \log(e + [w]_{A_1}) \|f\|_{L^1(w)}. \tag{1.13}$$

See [\[Lerner et al. 2009a\]](#) and also [\[Lerner et al. 2008\]](#) for these results and for more information about the problem. It was conjectured in the first of these works that the growth of this bound would be linear; however, it was shown in [\[Nazarov et al. 2010\]](#) that the growth of the bound is worse than linear. It seems that most probably the $L \log L$ result [\(1.13\)](#) is the best possible.

On the other hand, in [\[Lerner et al. 2009b\]](#), a sort of “dual” estimate to the last bound was found, which is also of interest for related matters:

$$\left\| \frac{Tf}{w} \right\|_{L^{1,\infty}(w)} \leq c [w]_{A_1} \log(e + [w]_{A_1}) \int_{\mathbb{R}^d} |f| \, dx.$$

In this paper, we improve these results following our new quantitative estimates, this time involving A_1 and A_∞ control. To be precise, we will prove the following new results:

Theorem 1.14. *Let T be a Calderón–Zygmund operator and let $1 < p < \infty$. Then*

$$\|T\|_{\mathfrak{B}(L^p(w))} \leq c p p' [w]_{A_1}^{1/p} ([w]_{A_\infty}')^{1/p'},$$

where $c = c(d, T)$.

We will prove this by following the approach from [\[Lerner et al. 2008; 2009a\]](#) to [\(1.12\)](#), modifying the proof at several points. In analogy to [\(1.5\)](#), [Theorem 1.14](#) disproves the “reverse A_1 conjecture”

$[w]_{A_1} \leq c_T \|T\|_{\mathcal{B}(L^p(w))}$ for all $p \in (1, \infty)$: considering a family of weights $w \in A_1$ for which $[w]_{A_\infty}$ grows slower than $[w]_{A_1}$, for any fixed Calderón–Zygmund operator T , we have (see [Section 8](#) for details)

$$\inf_{w \in A_1} \frac{\|T\|_{\mathcal{B}(L^p(w))}}{[w]_{A_1}} = 0, \quad 1 < p < \infty.$$

Finally, we will also use the approach from [[Lerner et al. 2009a](#); [2009b](#)] to prove the following theorems, respectively.

Theorem 1.15. *Let T be a Calderón–Zygmund operator. Then*

$$\|Tf\|_{L^{1,\infty}(w)} \leq c_{d,T} [w]_{A_1} \log(e + [w]_{A_\infty}') \|f\|_{L^1(w)}.$$

Theorem 1.16. *Let T be a Calderón–Zygmund operator. Then*

$$\left\| \frac{Tf}{w} \right\|_{L^{1,\infty}(w)} \leq c_{d,T} [w]_{A_\infty}' \log(e + [w]_{A_1}) \|f\|_{L^1(\mathbb{R}^d)}.$$

1D. Commutators with BMO functions. We further pursue the A_∞ point of view by proving a result in the spirit of [Theorem 1.3](#) for commutators of linear operators T with BMO functions. These operators are defined formally by the expression

$$[b, T]f = bT(f) - T(bf).$$

More generally, we can consider the k -th order commutator defined by

$$T_b^k := [b, T_b^{k-1}].$$

When T is a singular integral operator, these operators were considered by Coifman, Rochberg and Weiss [[Coifman et al. 1976](#)], and since then many results have been obtained. We refer to [[Chung et al. 2012](#)] for more information about these operators; it is shown there that if T is a linear operator bounded on $L^2(w)$ for any $w \in A_2$ with bound

$$\|T\|_{\mathcal{B}(L^2(w))} \leq \varphi([w]_{A_2}),$$

where φ is an increasing function $\varphi : [1, \infty) \rightarrow [0, \infty)$, then there is a dimensional constant c such that

$$\|[b, T]\|_{\mathcal{B}(L^2(w))} \leq c\varphi(c[w]_{A_2}) [w]_{A_2} \|b\|_{\text{BMO}}.$$

In particular, if T is any Calderón–Zygmund operator, we can use the linear A_2 theorem for T to deduce

$$\|[b, T]\|_{\mathcal{B}(L^2(w))} \leq c[w]_{A_2}^2 \|b\|_{\text{BMO}},$$

and the quadratic exponent cannot be improved.

An analogous result adapted to the A_∞ control reads as follows:

Theorem 1.17. *Let T be a linear operator bounded on $L^2(w)$ for any $w \in A_2$ and let $b \in \text{BMO}$. Suppose further that there is a function $\varphi : [1, \infty)^3 \rightarrow [0, \infty)$, increasing with respect to each component, such that*

$$\|T\|_{\mathcal{B}(L^2(w))} \leq \varphi([w]_{A_2}, [w]_{A_\infty}', [\sigma]_{A_\infty}').$$

Then there is a dimensional constant c such that

$$\| [b, T] \|_{\mathfrak{B}(L^2(w))} \leq c\varphi(c[w]_{A_2}, c[w]_{A_\infty}', c[\sigma]_{A_\infty}') ([w]_{A_\infty}' + [\sigma]_{A_\infty}') \|b\|_{\text{BMO}},$$

or more generally,

$$\| T_b^k \|_{\mathfrak{B}(L^2(w))} \leq c\varphi(c[w]_{A_2}, c[w]_{A_\infty}', c[\sigma]_{A_\infty}') ([w]_{A_\infty}' + [\sigma]_{A_\infty}')^k \|b\|_{\text{BMO}}^k.$$

We can now apply [Theorem 1.3](#).

Corollary 1.18. *Let T be any Calderón–Zygmund operator, and let $b \in \text{BMO}$. Then*

$$\| [b, T] \|_{\mathfrak{B}(L^2(w))} \leq c[w]_{A_2}^{1/2} ([w]_{A_\infty}' + [w^{-1}]_{A_\infty}')^{3/2} \|b\|_{\text{BMO}},$$

or more generally,

$$\| T_b^k \|_{\mathfrak{B}(L^2(w))} \leq c[w]_{A_2}^{1/2} ([w]_{A_\infty}' + [\sigma]_{A_\infty}')^{k+1/2} \|b\|_{\text{BMO}}^k.$$

1E. An end-point estimate when $p = \infty$. We next discuss the limiting form of the estimate (1.2) as $p \rightarrow \infty$, that is, the sharp bounds for the norm of Calderón–Zygmund operators

$$T : L^\infty(w) \rightarrow \text{BMO}(w), \quad w \in A_\infty.$$

Qualitatively, this situation seems slightly uninteresting, as these end-point spaces simply reduce to their unweighted analogues: that $L^\infty(w) = L^\infty$ with equal norms is immediate from the fact that w and the Lebesgue measure share the same zero sets for $w \in A_\infty$. That the weighted norm

$$\|f\|_{\text{BMO}(w)} := \sup_Q \inf_c \frac{1}{w(Q)} \int_Q |f - c| w < \infty$$

is equivalent to the usual $\|f\|_{\text{BMO}}$ for $w \in A_\infty$ was proven by Muckenhoupt and Wheeden [1975, Theorem 5]. However, one may still investigate the quantitative bound of operators $T : L^\infty \rightarrow \text{BMO} = \text{BMO}(w)$, when the latter space is equipped with the norm $\| \cdot \|_{\text{BMO}(w)}$. We start with:

Theorem 1.19. *For $w \in A_\infty$, we have a bounded embedding $I : \text{BMO} \hookrightarrow \text{BMO}(w)$ of norm at most $c[w]_{A_\infty}'$, where c is dimensional. This estimate is sharp in the following sense: if the norm of the embedding is bounded by $\phi([w]_{A_\infty}')$, or just by $\phi([w]_{A_\infty})$, for all $w \in A_\infty$, then $\phi(t) \geq ct$.*

The following corollary for Calderón–Zygmund operators can be seen as an easy endpoint estimate of the bound $\|T\|_{\mathfrak{B}(L^p(w))} \leq c_{p,T} [w]_{A_p}$ for $p \in [2, \infty)$.

Corollary 1.20. *Let T be any Calderón–Zygmund operator and let $w \in A_\infty$; then $T : L^\infty \rightarrow \text{BMO}(w)$ with norm at most $c_T [w]_{A_\infty}'$. Furthermore, this estimate is sharp in terms of the dependence on $[w]_{A_\infty}'$ in the same way as [Theorem 1.19](#).*

A related observation quantifying the known relation of A_∞ and BMO is as follows:

Proposition 1.21. *If $w \in A_\infty$, then $\log w \in \text{BMO}$ with*

$$\| \log w \|_{\text{BMO}} \leq \log(2e[w]_{A_\infty}).$$

1F. Extrapolation with A_∞ control. We recall the following quantitative version of Rubio de Francia's classical extrapolation theorem due to Dragičević, Grafakos, Pereyra, and Petermichl [Dragičević et al. 2005]: if an operator T satisfies

$$\|T\|_{\mathcal{B}(L^r(w))} \leq \varphi([w]_{A_r})$$

for a fixed increasing function φ and for all $w \in A_r$, then it satisfies a similar estimate for all $p \in (1, \infty)$,

$$\|T\|_{\mathcal{B}(L^p(w))} \leq 2\varphi(c_{p,r,d}[w]_{A_p}^{\max\{1,(r-1)/(p-1)\}});$$

in particular, $\|T\|_{\mathcal{B}(L^r(w))} \lesssim [w]_{A_r}^{\tau(r)}$ implies that

$$\|T\|_{\mathcal{B}(L^p(w))} \lesssim [w]_{A_p}^{\tau(r) \max\{1,(r-1)/(p-1)\}}.$$

With our new quantitative estimates involving both A_2 and A_∞ control, it seems of interest to extrapolate such bounds as well. Hence we consider weighted estimates of the form

$$\|Tf\|_{L^r(w)} \leq \varphi([w]_{A_r}, [w]_{A_\infty}, [w^{-1/(r-1)}]_{A_\infty}^{(r-1)}) \|f\|_{L^r(w)}, \quad (1.22)$$

where $\varphi : [1, \infty)^3 \rightarrow [0, \infty)$ is an increasing function with respect to each of the variables. An example is our bound for singular integrals 1.3, where

$$\varphi(x, y, z) = Cx^{1/2}(y+z)^{1/2}. \quad (1.23)$$

We now aim to extrapolate bounds like (1.22) from the given $r \in (1, \infty)$ to other exponents $p \in (1, \infty)$.

Theorem 1.24 (lower extrapolation). *Suppose that for some r and every $w \in A_r$, an operator T satisfies (1.22). Then for every $p \in (1, r)$, it satisfies*

$$\begin{aligned} \|Tf\|_{L^p(w)} &\leq 2\varphi\left((2\|M\|_{\mathcal{B}(L^p(w))})^{r-p}([w]_{A_r}, [w]_{A_\infty}, [w^{-1/(p-1)}]_{A_\infty}^{(p-1)})\right) \|f\|_{L^p(w)} \\ &\leq 2\varphi\left((c_d([w]_{A_p}[w^{-1/(p-1)}]_{A_\infty}')^{1/p})^{r-p}([w]_{A_p}, [w]_{A_\infty}, [w^{-1/(p-1)}]_{A_\infty}^{(p-1)})\right) \|f\|_{L^p(w)}. \end{aligned}$$

In typical applications, like (1.23), the function φ will have a homogeneity of the form $\varphi(\lambda x, \lambda y, \lambda z) = \lambda^s \varphi(x, y, z)$, and hence the common factor

$$(2\|M\|_{\mathcal{B}(L^p(w))})^{r-p} \leq (c_d([w]_{A_p}[w^{-1/(p-1)}]_{A_\infty}')^{1/p})^{r-p}$$

may be extracted out of φ .

Observe that the condition (1.22) is of course implied by the stronger inequality

$$\|Tf\|_{L^r(w)} \leq \varphi([w]_{A_r}, c_d^{-1}[w]_{A_\infty}', (c_d^{-1}[w^{-1/(r-1)}]_{A_\infty}')^{(r-1)}) \|f\|_{L^r(w)};$$

however, even if we have this stronger inequality to start with (as is the case with the A_2 theorem for Calderón–Zygmund operators), we do not know how to exploit it to get a stronger conclusion than what we can derive from (1.22). A related difficulty will be pointed out in the proof. This is why we restrict to the assumption (1.22) only.

Theorem 1.25 (upper extrapolation). *Suppose that for some r and every $w \in A_r$, an operator T satisfies (1.22). Then for every $p \in (r, \infty)$, it satisfies*

$$\begin{aligned} \|Tf\|_{L^p(w)} &\leq 2\varphi\left(\left(2\|M\|_{\mathfrak{B}(L^{p'}(w^{1-p'}))}\right)^{(p-r)/(p-1)}\right. \\ &\quad \left.\times ([w]_{A_p}^{(r-1)/(p-1)}, [w]_{A_\infty}^{(r-1)/(p-1)}, [w^{-1/(p-1)}]_{A_\infty}^{(r-1)})\right)\|f\|_{L^p(w)} \\ &\leq 2\varphi\left(\left(c_d[w]_{A_p}^{1/p}([w]_{A_\infty}')^{1/p'}\right)^{(p-r)/(p-1)}\right. \\ &\quad \left.\times ([w]_{A_p}^{(r-1)/(p-1)}, [w]_{A_\infty}^{(r-1)/(p-1)}, [w^{-1/(p-1)}]_{A_\infty}^{(r-1)})\right)\|f\|_{L^p(w)}. \end{aligned}$$

1G. The A_p theory for Calderón–Zygmund operators. As an application of the extrapolation theorems, we can deduce weighted L^p estimates for Calderón–Zygmund operators with mixed A_p and A_∞ control, akin to the A_2 bounds of Theorem 1.3. The same strategy has been earlier employed to prove the original A_p theorem (1.2) as a corollary of its A_2 version. However, in contrast to the “pure” A_p estimates, where the extrapolated result still exhibits the sharp dependence on the weight, it seems that the extrapolation of the mixed bounds is not equally efficient: the extrapolated bounds given below can be improved by methods directly adapted to L^p . Since the first public distribution of our present results, such further developments have been carried out in [Hytönen et al. 2011, Section 12; Lacey 2012; Hytönen and Lacey 2011]. Nevertheless, it seems worth recording the form of the A_p estimates directly delivered by the extrapolation method:

Corollary 1.26. *Let T be a Calderón–Zygmund operator and let $p \in (1, \infty)$. Then if $w \in A_p$ and $\sigma = w^{-1/(p-1)}$, we have*

$$\begin{aligned} \|T\|_{\mathfrak{B}(L^p(w))} &\lesssim [w]_{A_p}^{2/p-1/2}([w]_{A_\infty}^{1/2} + [\sigma]_{A_\infty}^{(p-1)/2})([\sigma]_{A_\infty}')^{2/p-1} \\ &\lesssim [w]_{A_p}^{2/p}([\sigma]_{A_\infty}')^{2/p-1} \quad \text{for } p \in (1, 2], \end{aligned} \tag{1.27}$$

and

$$\begin{aligned} \|T\|_{\mathfrak{B}(L^p(w))} &\lesssim [w]_{A_p}^{2/p-1/[2(p-1)]}([w]_{A_\infty}^{1/[2(p-1)]} + [\sigma]_{A_\infty}^{1/2})([w]_{A_\infty}')^{1-2/p} \\ &\lesssim [w]_{A_p}^{2/p}([\sigma]_{A_\infty}')^{1-2/p} \quad \text{for } p \in [2, \infty). \end{aligned} \tag{1.28}$$

Here the simpler forms of the estimates in (1.27) and (1.28) are almost as good as the more complicated ones, since for many common weights, like power weights, we have $[w]_{A_\infty} + [\sigma]_{A_\infty}^{p-1} \approx [w]_{A_p}$; see Section 8.

It is immediate to check that Theorems 1.24 and 1.25, in combination with Theorem 1.3, give Corollary 1.26. Actually, the two statements (1.27) and (1.28) are equivalent to each other by using

$$\|T\|_{\mathfrak{B}(L^p(w))} = \|T^*\|_{\mathfrak{B}(L^{p'}(\sigma))}$$

and the fact that T^* is also a Calderón–Zygmund operator. Thanks to this equivalence, we would only need one of Theorems 1.24 and 1.25 to deduce this corollary. But for other classes of operators without a self-dual structure, it is useful to have both upper and lower extrapolation results available.

2. The two different A_∞ constants

Before pursuing further our analysis of inequalities with A_∞ control, we include this short section to compare the two A_∞ constants

$$[w]_{A_\infty} := \sup_Q \left(\int_Q w \right) \exp \left(\int_Q \log w^{-1} \right), \quad [w]_{A_\infty}' := \sup_Q \frac{1}{w(Q)} \int_Q M(w\chi_Q).$$

We need the following auxiliary estimate, which is also used later in the paper:

Lemma 2.1. *The logarithmic maximal function*

$$M_0 f := \sup_Q \exp \left(\int_Q \log |f| \right) \chi_Q$$

satisfies

$$\|M_0 f\|_{L^p} \leq c_d^{1/p} \|f\|_{L^p}$$

for all $p \in (0, \infty)$. For the dyadic version, we can take $c_d = e$, independent of dimension d .

Proof. By Jensen’s inequality and the basic properties of the logarithm, we have

$$M_0 f \leq Mf, \quad M_0 f = (M_0 |f|^{1/q})^q \leq (M |f|^{1/q})^q, \quad q \in (0, \infty),$$

where M is the Hardy–Littlewood maximal operator, or the dyadic maximal operator in the case of dyadic M_0 . By the L^q boundedness of the usual maximal function for $q > 1$, we have

$$\int [M_0 f]^p \leq \int [M |f|^{p/q}]^q \leq (C_d \cdot q')^q \int (|f|^{p/q})^q = (C_d \cdot q')^q \int |f|^p.$$

In the nondyadic case, we simply take, say, $q = 2$, giving the claim with $c_d = (2C_d)^2$. In the dyadic case, we have $C_d = 1$, and we can take the limit $q \rightarrow \infty$, which gives

$$(q')^q = \left(\frac{q}{q-1} \right)^q = \left(1 + \frac{1}{q-1} \right)^q \rightarrow e,$$

and hence $\|M_0 f\|_{L^p}^p \leq e \|f\|_{L^p}^p$. □

Proposition 2.2. *We have $[w]_{A_\infty}' \leq c_d [w]_{A_\infty}$, where c_d is as in [Lemma 2.1](#).*

Proof. For $x \in Q$, it is not difficult to see that for the computation of $M(w\chi_Q)(x)$, it suffices to take the supremum over cubes $R \ni x$ with $R \subseteq Q$:

$$M(w\chi_Q)(x) = \sup_{\substack{R \ni x \\ R \subseteq Q}} \int_R w \quad \text{for all } x \in Q.$$

By the definition of $[w]_{A_\infty}$, we have

$$\int_R w \leq [w]_{A_\infty} \exp \left(\int_R \log w \right),$$

and hence, taking the supremum over R ,

$$M(w\chi_Q)(x) \leq [w]_{A_\infty} M_0(w\chi_Q)(x) \quad \text{for all } x \in Q.$$

Integration over Q and application of [Lemma 2.1](#) now give

$$\int_Q M(w\chi_Q) \leq [w]_{A_\infty} \int M_0(w\chi_Q) \leq [w]_{A_\infty} c_d \int w\chi_Q = c_d [w]_{A_\infty} w(Q);$$

thus $[w]'_{A_\infty} \leq c_d [w]_{A_\infty}$. □

It is a well-known fact that any A_∞ weight satisfies a reverse Hölder inequality playing a central role in the area. In this paper, a sharp version of this property will also play a fundamental role. To be precise, if $w \in A_\infty$, we define

$$r(w) := 1 + \frac{1}{\tau_d [w]'_{A_\infty}},$$

where τ_d is a dimensional constant that we may take to be $\tau_d = 2^{11+d}$. Note that $r(w)' \approx [w]'_{A_\infty}$. The result we need is the following.

Theorem 2.3 (a new sharp reverse Hölder inequality). (a) *If $w \in A_\infty$, then*

$$\left(\int_Q w^{r(w)} \right)^{1/r(w)} \leq 2 \int_Q w.$$

(b) *Furthermore, the result is optimal up to a dimensional factor: If a weight w satisfies the reverse Hölder inequality*

$$\left(\int_Q w^r \right)^{1/r} \leq K \int_Q w,$$

then $[w]'_{A_\infty} \leq c_d \cdot K \cdot r'$.

This result is new in the literature and has its own interest. In the classical situation, most of the available proofs do not give such explicit constants, which are important for us. Only under the stronger condition of A_1 was such a result found and used in a crucial way in [\[Lerner et al. 2009a\]](#). Recently a very nice proof by A. de la Torre for the case $[w]_{A_\infty}$ was sent to us (personal communication, 2010). Another less precise proof, for the A_p case, $1 < p < \infty$, can be found in [\[Pérez 2013\]](#).

Part (b) follows from the boundedness of the maximal function in L^r with constant $c_d r'$:

$$\int_Q M(\chi_Q w) \leq \left(\int_Q M(\chi_Q w)^r \right)^{1/r} \leq c_d \cdot r' \left(\int_Q w^r \right)^{1/r} \leq c_d \cdot r' \cdot K \int_Q w.$$

Remark 2.4. Results analogous to [Proposition 2.2](#) and [Theorem 2.3](#) have been independently obtained by O. Beznosova and A. Reznikov [\[2011\]](#). Their formulation is slightly different, and involves yet another weight constant closely related to $[w]'_{A_\infty}$.

3. The A_2 theorem for Calderón–Zygmund operators

The purpose of this section is to prove [Theorem 1.3](#), namely, the estimate

$$\|T\|_{\mathcal{B}(L^2(w))} \leq c [w]_{A_2}^{1/2} ([w]'_{A_\infty} + [\sigma]'_{A_\infty})^{1/2},$$

where $c = c_{d,T}$ is a constant depending on the dimension and the operator T .

Here and throughout this section, $\sigma = w^{-1}$. This improves on the A_2 theorem [Hytönen 2012]:

$$\|T\|_{\mathfrak{B}(L^2(w))} \leq c[w]_{A_2},$$

and its proof follows the same outline, with the implementation of the A_∞ philosophy at key points.

3A. Reduction to a dyadic version. Fundamental to this proof strategy is the notion of *dyadic shifts*, which we recall. We work with a *general dyadic system* \mathfrak{D} , this being a collection of axis-parallel cubes Q , whose sidelengths $\ell(Q)$ are of the form 2^k , $k \in \mathbb{Z}$, where moreover $Q \cap R \in \{Q, R, \emptyset\}$ for any two $Q, R \in \mathfrak{D}$, and the cubes of a fixed sidelength 2^k form a partition of \mathbb{R}^d . Given such a dyadic system, a dyadic shift with parameters (m, n) is an operator of the form

$$\mathbb{H}f = \sum_{K \in \mathfrak{D}} A_K f, \quad A_K f = \frac{1}{|K|} \sum_{\substack{I, J \in \mathfrak{D}; I, J \subseteq K \\ \ell(I) = 2^{-m} \ell(K) \\ \ell(J) = 2^{-n} \ell(K)}} \langle h_I^J, f \rangle k_J^I,$$

where h_I^J is a generalized Haar function on I (supported on I , constant on its dyadic subcubes, and normalized by $\|h_I^J\|_\infty \leq 1$), and k_J^I on J . This implies that $|A_K f| \leq \chi_K \cdot |K|^{-1} \cdot \int_K |f|$. For any subcollection $\mathfrak{Q} \subset \mathfrak{D}$, we write

$$\mathbb{H}_{\mathfrak{Q}} f := \sum_{K \in \mathfrak{Q}} A_K f, \tag{3.1}$$

and we require that $\|\mathbb{H}_{\mathfrak{Q}} f\|_{L^2} \leq \|f\|_{L^2}$ for all $\mathfrak{Q} \subset \mathfrak{D}$. This is automatic from straightforward orthogonality considerations in case we only have cancellative Haar functions with $\int h_I^J = \int k_J^I = 0$.

Dyadic shifts with parameters $(0, 0)$ are well known in dyadic harmonic analysis under different names. Auscher et al. [2002] study such operators under the name *perfect dyadic operators*, which they decompose into a sum of a *Haar multiplier* (or *martingale transform*), a *paraproduct*, and a *dual paraproduct*. These three types of operators have of course been well known for a long time. The first dyadic shift (and this name) with parameters $(0, 1)$ was introduced in [Petermichl 2000], and the definition in the above generality was given by Lacey, Petermichl and Reguera [Lacey et al. 2010b].

The importance of these dyadic shifts for the analysis of Calderón–Zygmund operators comes from the following:

Theorem 3.2 (dyadic representation theorem [Hytönen 2012, Theorem 4.2; Hytönen et al. 2010, Theorem 4.1]). *Let $T \in \mathfrak{B}(L^2(\mathbb{R}^d))$ be a Calderón–Zygmund operator satisfying the standard estimates with the Hölder continuity exponent $\alpha \in (0, 1]$. Then T has the representation*

$$\langle g, T f \rangle = c_{T,d} \mathbb{E}_{\mathfrak{D}} \sum_{m,n=0}^{\infty} 2^{-(m+n)\alpha/2} \langle g, \mathbb{H}_{\mathfrak{D}}^{mn} f \rangle,$$

valid for all bounded and compactly supported functions f and g , where $\mathbb{H}_{\mathfrak{D}}^{mn}$ is a dyadic shift with parameters (m, n) related to the dyadic system \mathfrak{D} , and $\mathbb{E}_{\mathfrak{D}}$ is the expectation with respect to a probability measure on the space of all generalized dyadic systems; see [Hytönen 2012] for the details of the construction of this probability space.

This result was preceded by several versions restricted to special operators T : the Beurling–Ahlfors transform by Dragičević and Volberg [2003], the Hilbert transform by Petermichl [2000], the Riesz transforms by Petermichl, Treil and Volberg [Petermichl et al. 2002], and all one-dimensional convolution operators with an odd, smooth kernel by Vagharshakyan [2010]. An immediate consequence of the dyadic representation theorem is that Theorem 1.3 will be a consequence of the following dyadic version. (Similarly, the special cases of the representation theorem all played a role in proving the A_2 theorem for the mentioned particular operators.)

Theorem 3.3. *Let \mathbb{H} be a dyadic shift with parameters (m, n) , and $r = \max\{m, n\}$. For $w \in A_2$ and $\sigma = w^{-1}$, we have*

$$\|\mathbb{H}f\|_{L^2(w)} \leq C(r + 1)^2 [w]_{A_2}^{1/2} ([w]_{A_\infty}' + [\sigma]_{A_\infty}')^{1/2} \|f\|_{L^2(w)}.$$

The weighted norm of the shifts, in turn, is most conveniently deduced with the help of the following characterization of their boundedness in a two-weight situation:

Theorem 3.4 [Hytönen et al. 2010, Theorem 3.4]. *Let \mathbb{H} be a dyadic shift with parameters (m, n) , and let $r = \max\{m, n\}$. If for all $Q \in \mathcal{D}$ and some B there holds*

$$\left(\int_Q |\mathbb{H}(\chi_Q \sigma)|^2 w \right)^{1/2} \leq B \sigma(Q)^{1/2}, \quad \left(\int_Q |\mathbb{H}^*(\chi_Q w)|^2 \sigma \right)^{1/2} \leq B w(Q)^{1/2},$$

then for a dimensional constant c , we have

$$\|\mathbb{H}(f\sigma)\|_{L^2(w)} \leq c((r + 1)B + (r + 1)^2(A_2[w, \sigma])^{1/2}) \|f\|_{L^2(\sigma)},$$

where $A_2[w, \sigma]$ is defined by the functional

$$A_2[w, \sigma] := \sup_Q \left(\int_Q w \right) \left(\int_Q \sigma \right).$$

Since the last bound is equivalent to

$$\|\mathbb{H}f\|_{L^2(w)} \leq c((r + 1)B + (r + 1)^2[w]_{A_2}^{1/2}) \|f\|_{L^2(w)}$$

if $\sigma = w^{-1}$, and since $[w]_{A_\infty}, [\sigma]_{A_\infty} \geq 1$, we are reduced to estimating the quantity B for $\sigma = w^{-1}$. Since \mathbb{H} and \mathbb{H}^* are operators of the same form, and by the symmetry of w and σ , Theorem 3.4 shows that proving Theorem 3.3 amounts to showing that

$$\left(\int_Q |\mathbb{H}(w\chi_Q)|^2 \sigma \right)^{1/2} \leq c(r + 1) ([w]_{A_2} [w]_{A_\infty}' w(Q))^{1/2}.$$

We observe that

$$\mathbb{H}(w\chi_Q) = \sum_{K \subseteq Q} A_K(w\chi_Q) + \sum_{K \supset Q} A_K(w\chi_Q),$$

and it suffices to consider the two parts separately. The big cubes are immediately handled by the maximal

function estimate (see [Corollary 1.10](#)):

$$\begin{aligned} \int_Q \left| \sum_{K \supset Q} A_K(w\chi_Q) \right|^2 \sigma &\leq \int_Q \left(\sum_{K \supset Q} \frac{w(Q)}{|K|} \chi_K \right)^2 \sigma \lesssim \int_Q M_d(w\chi_Q)^2 \sigma \\ &\leq [\sigma]_{A_2} [w]_{A_\infty}' w(Q) = [w]_{A_2} [w]_{A_\infty}' w(Q). \end{aligned} \tag{3.5}$$

Hence, to prove [Theorem 3.3](#), we are reduced to showing that

$$\left(\int_Q \left| \sum_{K \subseteq Q} A_K(w\chi_Q) \right|^2 \sigma \right)^{1/2} \leq c(r+1) ([w]_{A_2} [w]_{A_\infty}' w(Q))^{1/2}. \tag{3.6}$$

This is the goal for the rest of this section.

3B. Proof of the key estimate (3.6). We follow the key steps from [[Hytönen 2012](#); [Hytönen et al. 2010](#); [Lacey et al. 2010b](#)]. The collection $\{K \in \mathcal{D} : K \subseteq Q\}$ is first split into $(r+1)$ subcollections according to the value of $\log_2 \ell(K) \pmod{(r+1)}$; we henceforth work with one of these subcollections, which we denote by \mathcal{H} . This is the step which introduces the factor $(r+1)$, and we will estimate $\mathbb{I}\mathbb{I}\mathbb{I}_{\mathcal{H}}(w\chi_Q)$ with a bound independent of r .

The collection \mathcal{H} is further divided into the sets \mathcal{H}^a of those cubes with

$$2^a < \frac{w(Q)}{|Q|} \frac{\sigma(Q)}{|Q|} \leq 2^{a+1}, \tag{3.7}$$

where $a \leq \log_2 [w]_{A_2}$.

Among the cubes $K \in \mathcal{H}^a$, we choose the principal cubes $\mathcal{S}^a = \bigcup_{k=0}^\infty \mathcal{S}_k^a$ so that \mathcal{S}_0^a consists of the maximal cubes in \mathcal{H}^a , and \mathcal{S}_k^a the maximal cubes $S \in \mathcal{H}^a$ contained in some $S' \in \mathcal{S}_{k-1}^a$ with $\sigma(S)/|S| > 2\sigma(S')/|S'|$. Then

$$\mathcal{H}^a = \bigcup_{S \in \mathcal{S}^a} \mathcal{H}^a(S), \quad \mathcal{H}^a(S) := \{K \in \mathcal{H}^a \mid K \subseteq S, \text{ there exists no } S' : K \subseteq S' \subset S\}.$$

It follows that, in the notation from [\(3.1\)](#),

$$\mathbb{I}\mathbb{I}\mathbb{I}_{\mathcal{H}}(w\chi_Q) = \sum_{a \leq \log_2 [w]_{A_2}} \sum_{S \in \mathcal{S}^a} \mathbb{I}\mathbb{I}\mathbb{I}_{\mathcal{H}^a(S)}(w\chi_Q). \tag{3.8}$$

To proceed, we recall the following distributional estimate:

Lemma 3.9 [[Hytönen et al. 2010](#), (5.26)]. *With notation as above, we have*

$$\sigma\left(\left|\mathbb{I}\mathbb{I}\mathbb{I}_{\mathcal{H}^a(S)}(w\chi_Q)\right| > t \langle w \rangle_S\right) \leq C e^{-ct} \sigma(S) \quad \text{for all } S \in \mathcal{S}^a, \tag{3.10}$$

where the constants C and c are at worst dimensional.

This is a powerful estimate which readily leads to norm bounds for [\(3.8\)](#). The following computation, simplifying the corresponding ones from [[Hytönen 2012](#); [Hytönen et al. 2010](#); [Lacey et al. 2010b](#)], is borrowed from [[Hytönen et al. 2011](#)]: writing

$$E_j(S) := \{j \leq \left|\mathbb{I}\mathbb{I}\mathbb{I}_{\mathcal{H}^a(S)}(w\chi_Q)\right| / \langle w \rangle_S < j+1\} \subseteq S,$$

we have

$$\begin{aligned}
 \left\| \sum_{S \in \mathcal{S}^a} \mathbb{I}_{\mathcal{H}^a(S)}(w\chi_Q) \right\|_{L^2(\sigma)} &\leq \sum_{j=0}^{\infty} (j+1) \left\| \sum_{S \in \mathcal{S}^a} \langle w \rangle_S \cdot \chi_{E_j(S)} \right\|_{L^2(\sigma)} \\
 &= \sum_{j=0}^{\infty} (j+1) \left(\int \left[\sum_{S \in \mathcal{S}^a} \langle w \rangle_S \cdot \chi_{E_j(S)}(x) \right]^2 \sigma(x) \, dx \right)^{1/2} \\
 &\stackrel{(*)}{\leq} C \sum_{j=0}^{\infty} (j+1) \left(\int \sum_{S \in \mathcal{S}^a} \langle w \rangle_S^2 \cdot \chi_{E_j(S)}(x) \sigma(x) \, dx \right)^{1/2} \\
 &= C \sum_{j=0}^{\infty} (j+1) \left(\sum_{S \in \mathcal{S}^a} \langle w \rangle_S^2 \cdot \sigma(E_j(S)) \right)^{1/2} \\
 &\leq C \sum_{j=0}^{\infty} (j+1) \left(\sum_{S \in \mathcal{S}^a} \langle w \rangle_S^2 \cdot C e^{-cj} \sigma(S) \right)^{1/2} \quad (\text{by (3.10)}) \\
 &\leq C \sum_{j=0}^{\infty} e^{-cj} (j+1) \left(2^a \sum_{S \in \mathcal{S}^a} w(S) \right)^{1/2} \quad (\text{by (3.7) for } S \in \mathcal{S}^a \subset \mathcal{H}^a) \\
 &\leq C \cdot 2^{a/2} \left(\sum_{S \in \mathcal{S}^a} w(S) \right)^{1/2}.
 \end{aligned}$$

In (*) we used the fact that at a fixed x , the numbers $\langle w \rangle_S$ for the principal cubes $S \supset E_j(S) \ni x$ increase at least geometrically, so their ℓ^1 and ℓ^2 norms are comparable.

We now come to the crucial point, where we can improve the earlier A_2 bounds to A_∞ :

Lemma 3.11. *For the principal cubes as defined above, we have*

$$\sum_{S \in \mathcal{S}^a} w(S) \leq 2 \cdot [w]_{A_\infty}' \cdot w(Q).$$

Proof. Let

$$E(S) := S \setminus \bigcup_{S' \subsetneq S} S'. \tag{3.12}$$

The union is the union of its maximal members S' , which satisfy

$$|S'| = |S'|/w(S') \cdot w(S') \leq \frac{1}{2} |S|/w(S) \cdot w(S');$$

hence $\sum |S'| \leq \frac{1}{2} |S|$, and thus

$$|E(S)| \geq \frac{1}{2} |S|. \tag{3.13}$$

Therefore

$$\sum_{S \in \mathcal{S}^a} w(S) = \sum_{S \in \mathcal{S}^a} \frac{w(S)}{|S|} |S| \leq \sum_{S \in \mathcal{S}^a} \frac{w(S)}{|S|} 2|E(S)| \leq 2 \sum_{S \in \mathcal{S}^a} \int_{E(S)} M(w\chi_Q) = 2 \int_Q M(w\chi_Q) \leq 2 [w]_{A_\infty}' w(Q),$$

where the last step was the definition of $[w]_{A_\infty}'$. □

Substituting the estimates obtained back into (3.8), we conclude that

$$\begin{aligned} \|\mathbb{M}_{\mathcal{H}}(w\chi_Q)\|_{L^2(\sigma)} &\leq \sum_{a \leq \log_2[w]_{A_2}} \left\| \sum_{S \in \mathcal{F}^a} \mathbb{M}_{\mathcal{H}^a(S)}(w\chi_Q) \right\|_{L^2(\sigma)} \leq C \sum_{a \leq \log_2[w]_{A_2}} 2^{a/2} \left(\sum_{S \in \mathcal{F}^a} w(S) \right)^{1/2} \\ &\leq C \sum_{a \leq \log_2[w]_{A_2}} 2^{a/2} ([w]_{A_\infty}' \cdot w(Q))^{1/2} \leq C [w]_{A_2}^{1/2} ([w]_{A_\infty}')^{1/2} w(Q)^{1/2}. \end{aligned}$$

Recalling the initial splitting of $\{K \in \mathcal{D} : K \subseteq Q\}$ into $r + 1$ subcollections of the same form as \mathcal{H} , this concludes the proof of (3.6), and hence the proof of Theorem 3.3.

4. Two-weight theory for the maximal function

4A. Background. The two-weight problem was studied in the 1970s by Muckenhoupt and Wheeden and fully solved by E. Sawyer [1982]. The general question is to find a necessary and sufficient condition for a pair of unrelated weights w and σ for which the estimate

$$\|M(f\sigma)\|_{L^p(w)} \leq B \|f\|_{L^p(\sigma)} \tag{4.1}$$

holds for a finite constant B . Then the main result of E. Sawyer shows that this is the case if and only if there exists a finite c such that

$$\int_Q M(\sigma\chi_Q)(y)^p w(y) \, dy \leq c \sigma(Q)$$

for all cubes Q . Furthermore, it is shown in [Moen 2009] that if B denotes the best constant, then

$$B \approx \sup_Q \left(\frac{\int_Q M(\sigma\chi_Q)^p w \, dx}{\sigma(Q)} \right)^{1/p}$$

Since this condition is hard to verify in practice, the second author considered in [Pérez 1995] conditions closer in spirit to the classical two-weight A_p condition,

$$A_p[w, \sigma] := \sup_Q \left(\int_Q w \right) \left(\int_Q \sigma \right)^{p-1},$$

which reduces to $[w]_{A_p}$ if $\sigma = w^{-1/(p-1)}$. As a consequence of the main result in that work, if $\delta > 0$ and

$$\sup_Q \left(\int_Q w \right) \|\sigma\|_{L(\log L)^{p-1+\delta}, Q}^{p-1} < \infty, \tag{4.2}$$

then the two-weight norm inequality (4.1) holds. Recent advances in collaboration with M. Mastyło [Mastyło and Pérez 2013] allow one to go beyond condition (4.2) and improve the main results from [Pérez 1995].

In this paper, we consider a different new quantity, namely

$$B_p[w, \sigma] := \sup_Q \left(\int_Q w \right) \left(\int_Q \sigma \right)^p \exp \left(\int_Q \log \sigma^{-1} \right).$$

To understand this new quantity, we observe that it is simply the functional on Q defining the $A_p[w, \sigma]$ condition multiplied by $\int_Q \sigma \exp(f_Q \log \sigma^{-1}) \geq 1$. Then it is immediate that

$$A_p[w, \sigma] \leq B_p[w, \sigma] \leq A_p[w, \sigma]A_\infty[\sigma],$$

the difference between the last two being that $A_p[w, \sigma]A_\infty[\sigma]$ involves two independent suprema, as opposed to just one in $B_p[u, v]$.

We will consider first the dyadic maximal operator M_d , for which we can prove a dimension-free bound. Let us also introduce the weighted dyadic maximal function

$$M_{d,\sigma} f := \sup_{Q \in \mathfrak{D}} \frac{\chi_Q}{\sigma(Q)} \int_Q |f(y)| \sigma(y) dy,$$

which controls $M_d(f\sigma)$ as follows:

Theorem 4.3. *Let $p \in (1, \infty)$; then*

$$\begin{aligned} \|M_d(f\sigma)\|_{L^p(w)} &\leq 4e \cdot (B_p[w, \sigma])^{1/p} \|M_{d,\sigma} f\|_{L^p(\sigma)} \\ &\leq 4e \cdot p' \cdot (B_p[w, \sigma])^{1/p} \|f\|_{L^p(\sigma)}, \end{aligned}$$

and also

$$\begin{aligned} \|M_d(f\sigma)\|_{L^p(w)} &\leq 4e \cdot ([w]_{A_p} [\sigma]_{A_\infty}')^{1/p} \|M_{d,\sigma} f\|_{L^p(\sigma)} \\ &\leq 4e \cdot p' \cdot ([w]_{A_p} [\sigma]_{A_\infty}')^{1/p} \|f\|_{L^p(\sigma)}. \end{aligned}$$

The main estimate in both chains of inequalities is of course the first one, since the second is simply the universal estimate for the weighted dyadic maximal function on the weighted L^p space with the same weight:

$$\|M_{d,\sigma}\|_{\mathfrak{B}(L^p(\sigma))} \leq p'.$$

Obviously, in this dyadic version, it suffices to have the supremum in the weight constants over dyadic cubes only, and to only use the dyadic square function in the definition of $[\sigma]_{A_\infty}'$. And specializing to the case $\sigma = w^{-1/(p-1)}$, by the standard dual weight trick, we also get the bounds

$$\|M_d f\|_{L^p(w)} \leq \begin{cases} 4e \cdot p' \cdot (B_p[w, w^{-1/(p-1)}])^{1/p} \|f\|_{L^p(w)}, \\ 4e \cdot p' \cdot ([w]_{A_p} [w^{-1/(p-1)}]_{A_\infty}')^{1/p} \|f\|_{L^p(w)}. \end{cases}$$

Let us also recall how such dyadic bounds yield corresponding results for the Hardy–Littlewood maximal operator by a standard argument.

Proof of Theorem 1.7. Consider the 2^d shifted dyadic systems

$$\mathfrak{D}^\alpha := \{2^{-k}([0, 1]^d + m + (-1)^k \alpha) : k \in \mathbb{Z}, m \in \mathbb{Z}^d\}, \quad \alpha \in \{0, \frac{1}{3}\}^d.$$

One can check (perhaps best in dimension $d = 1$ first) that any cube Q is contained in a shifted dyadic cube $Q^\alpha \in \mathfrak{D}^\alpha$ with $\ell(Q^\alpha) \leq 6\ell(Q)$, for some α . Hence

$$\int_Q |f| \leq 6^d \int_{Q^\alpha} |f| \leq 6^d M_d^\alpha f,$$

and therefore

$$Mf \leq 6^d \sum_{\alpha \in \{0, \frac{1}{3}\}^d} M_d^\alpha f.$$

Thus, the norm bound for M_d may be multiplied by 12^d to give a bound for M . □

Remark 4.4. A recent result in collaboration with A. Kairema [Hytönen and Kairema 2012] allows one to perform a similar trick with adjacent dyadic systems even in an abstract space of homogeneous type. Thus, Theorem 1.7 readily extends to this generality as well.

4B. Proof of Theorem 4.3. We start by observing that it suffices to have a uniform bound over all linearizations

$$\tilde{M}(f\sigma) = \sum_{Q \in \mathfrak{D}} \chi_{E(Q)} \langle f\sigma \rangle_Q,$$

where the sets $E(Q) \subseteq Q$ are pairwise disjoint. Here we use the notation

$$\langle f \rangle_Q = \int_Q f = \int_Q f(x) \, dx$$

and

$$\langle f \rangle_Q^\sigma = \frac{1}{\sigma(Q)} \int_Q f(x)\sigma(x) \, dx,$$

where, as usual, $\sigma(E) = \int_Q \sigma(x) \, dx$.

By this disjointness,

$$\|\tilde{M}(f\sigma)\|_{L^p(w)} = \left(\sum_{Q \in \mathfrak{D}} w(E(Q)) \langle f\sigma \rangle_Q^p \right)^{1/p} = \left(\sum_{Q \in \mathfrak{D}} w(E(Q)) \left(\frac{\sigma(Q)}{|Q|} \right)^p (\langle f \rangle_Q^\sigma)^p \right)^{1/p}.$$

Now recall:

Theorem 4.5 (dyadic Carleson embedding theorem). *Suppose that the nonnegative numbers a_Q satisfy*

$$\sum_{Q \subseteq R} a_Q \leq A\sigma(R) \quad \text{for all } R \in \mathfrak{D}.$$

Then, for all $p \in [1, \infty)$ and $f \in L^p(\sigma)$,

$$\left(\sum_{Q \in \mathfrak{D}} a_Q (\langle f \rangle_Q^\sigma)^p \right)^{1/p} \leq A^{1/p} \|M_{d,\sigma} f\|_{L^p(\sigma)} \leq A^{1/p} \cdot p' \cdot \|f\|_{L^p(\sigma)} \quad \text{if } p > 1.$$

Since this is a slightly nonstandard formulation, although immediate by inspection of the usual argument, we provide a proof for completeness:

Proof. We view the sum $\sum_Q a_Q (\langle f \rangle_Q^\sigma)^p$ as an integral on a measure space (\mathfrak{D}, μ) built over the set of dyadic cubes \mathfrak{D} , assigning to each $Q \in \mathfrak{D}$ the measure a_Q . Thus

$$\sum_{Q \in \mathfrak{D}} a_Q (\langle f \rangle_Q^\sigma)^p = \int_0^\infty p\lambda^{p-1} \mu(\{Q \in \mathfrak{D} : \langle f \rangle_Q^\sigma > \lambda\}) \, d\lambda =: \int_0^\infty p\lambda^{p-1} \mu(\mathfrak{Q}_\lambda) \, d\lambda.$$

Let \mathfrak{Q}_λ^* be the set of maximal dyadic cubes R with the property that $\langle f \rangle_R^\sigma > \lambda$. The cubes $R \in \mathfrak{Q}_\lambda^*$ are

disjoint, and their union is equal to the set $\{M_{d,\sigma} f > \lambda\}$. Thus

$$\mu(\mathcal{Q}_\lambda) = \sum_{Q \in \mathcal{Q}_\lambda} a_Q \leq \sum_{R \in \mathcal{Q}_\lambda^*} \sum_{Q \subseteq R} a_Q \leq \sum_{R \in \mathcal{Q}_\lambda^*} A\sigma(R) = A\sigma(M_{d,\sigma} f > \lambda),$$

and hence

$$\sum_{Q \in \mathcal{Q}} a_Q ((f)_Q)^p \leq A \int_0^\infty p\lambda^{p-1} \sigma(M_{d,\sigma} f > \lambda) d\lambda = A \|M_{d,\sigma} f\|_{L^p(\sigma)}^p. \quad \square$$

If we apply the Carleson embedding with $a_Q = w(E(Q))(\sigma(Q)/|Q|)^p$, we find that

$$\|\tilde{M}(f\sigma)\|_{L^p(w)} \leq A^{1/p} \|M_{d,\sigma} f\|_{L^p(\sigma)}, \tag{4.6}$$

provided that

$$\sum_{Q \subseteq R} w(E(Q)) \left(\frac{\sigma(Q)}{|Q|}\right)^p \leq A\sigma(R) \quad \text{for all } R \in \mathcal{Q}. \tag{4.7}$$

Note that on $E(Q) \subseteq Q \subseteq R$, we have $\sigma(Q)/|Q| \leq M(\sigma \chi_R)$, and hence

$$\begin{aligned} \sum_{Q \subseteq R} w(E(Q)) \left(\frac{\sigma(Q)}{|Q|}\right)^p &= \int \sum_{Q \subseteq R} \chi_{E(Q)} \left(\frac{\sigma(Q)}{|Q|}\right)^p w \\ &\leq \int \sum_{Q \subseteq R} \chi_{E(Q)} M(\chi_R \sigma)^p w \leq \int_R M(\chi_R \sigma)^p w. \end{aligned}$$

So if $\|\chi_R M(\chi_R \sigma)\|_{L^p(w)} \leq A^{1/p} \sigma(R)^{1/p}$, then (4.7) holds, and hence by Carleson’s embedding also (4.6), and therefore the original two-weight inequality

$$\|M(f\sigma)\|_{L^p(w)} \leq A^{1/p} \|M_{d,\sigma} f\|_{L^p(\sigma)}.$$

Hence, we are reduced to proving that

$$\|\chi_R M(\chi_R \sigma)\|_{L^p(w)}^p \leq A\sigma(R), \quad A = (4e)^{1/p} \cdot B_p[w, \sigma]. \tag{4.8}$$

(In fact, the argument up to this point was essentially reproving Sawyer’s two-weight characterization for the maximal function, paying attention to the constants.)

To prove (4.8), we exploit another linearization of M involving the *principal cubes*, as in the proof of the A_2 theorem: let $\mathcal{S}_0 := \{R\}$ and recursively let

$$\mathcal{S}_k := \bigcup_{S \in \mathcal{S}_{k-1}} \{Q \subset S : \langle \sigma \rangle_Q > 2\langle \sigma \rangle_S, Q \text{ is a maximal such cube}\},$$

and then $\mathcal{S} := \bigcup_{k=0}^\infty \mathcal{S}_k$. The pairwise disjoint subsets $E(S) \subseteq S$, defined in (3.12), satisfy $|E(S)| \geq \frac{1}{2}|S|$ by (3.13), and they partition R .

If $x \in E(S)$ and $Q \ni x$, then $\langle \sigma \rangle_Q \leq 2\langle \sigma \rangle_S$, and hence $\chi_R M(\chi_R \sigma) \leq 2\langle \sigma \rangle_S$ on $\chi_{E(S)}$. So altogether,

$$\begin{aligned}
 \|\chi_R M(\chi_R \sigma)\|_{L^p(w)}^p &\leq 2^p \left\| \sum_{S \in \mathcal{S}} \chi_{E(S)} \langle \sigma \rangle_S \right\|_{L^p(w)}^p = 2^p \sum_{S \in \mathcal{S}} w(E(S)) \left(\frac{\sigma(S)}{|S|} \right)^p \\
 &\leq 2^p \sum_{S \in \mathcal{S}} \frac{w(S)}{|S|} \left(\frac{\sigma(S)}{|S|} \right)^p |S| \\
 &\leq 2^{p+1} \sum_{S \in \mathcal{S}} B_p[w, \sigma] \exp\left(\int_S \log \sigma \right) |E(S)| \\
 &\leq 2^{p+1} B_p[w, \sigma] \int_R \sum_{S \in \mathcal{S}} \exp\left(\int_S \log \sigma \right) \chi_{E(S)} \\
 &\leq 2^{p+1} B_p[w, \sigma] \int_R \sup_{Q \in \mathcal{Q}} \chi_Q \exp\left(\int_Q \log \sigma \chi_R \right) \\
 &= 2^{p+1} B_p[w, \sigma] \int_R M_0(\chi_R \sigma),
 \end{aligned} \tag{4.9}$$

where M_0 is the (dyadic) logarithmic maximal function introduced in [Lemma 2.1](#). By this lemma, we then have

$$\|\chi_R M(\chi_R \sigma)\|_{L^p(w)}^p \leq 4^p B_p[w, \sigma] \cdot e \cdot \sigma(R),$$

which proves [\(4.8\)](#), and hence [Theorem 4.3](#), upon taking the p -th root.

In order to prove the second version of [Theorem 4.3](#), we only need to make a slight modification in the estimate [\(4.9\)](#). We then compute:

$$\begin{aligned}
 \|\chi_R M(\chi_R \sigma)\|_{L^p(w)}^p &\leq 2^p \sum_{S \in \mathcal{S}} \frac{w(S)}{|S|} \left(\frac{\sigma(S)}{|S|} \right)^p |S| \leq 2^{p+1} \sum_{S \in \mathcal{S}} [w]_{A_p} \frac{\sigma(S)}{|S|} |E(S)| \\
 &\leq 2^{p+1} [w]_{A_p} \sum_{S \in \mathcal{S}} \int_{E(S)} M(\sigma \chi_Q) = 2^{p+1} [w]_{A_p} \int_Q M(\sigma \chi_Q) \\
 &= 2^{p+1} [w]_{A_p} [\sigma]'_{A_\infty} \sigma(Q),
 \end{aligned}$$

by a direct application of the definition of $[\sigma]'_{A_\infty}$ in the last step, and this completes the alternative argument.

4C. Another proof of [Theorem 4.3](#). We finish this section by providing yet another proof variant for [Theorem 4.3](#). This proof is more elementary, since it does not need the reduction to the testing condition [\(4.8\)](#), and it uses the more standard Calderón–Zygmund-type stopping cubes instead of the principal cubes. Its disadvantage is the fact that we cannot recover the dimension-independence by this argument. On the other hand, the proof may be extended to maximal functions defined in terms of a general basis; see [\[García-Cuerva and Rubio de Francia 1985, Section IV.4\]](#).

A simpler proof of [Theorem 4.3](#) with a dimension-dependent bound. Fix $a > 2^d$. For each integer k , let

$$\Omega_k = \{x \in \mathbb{R}^d : M_d(f \sigma)(x) > a^k\}.$$

By standard arguments, we consider the Calderón–Zygmund decomposition, and there is a family of maximal nonoverlapping dyadic cubes $\{Q_{k,j}\}$ for which $\Omega_k = \bigcup_j Q_{k,j}$ and

$$a^k < \frac{1}{|Q_{k,j}|} \int_{Q_{k,j}} |f(y)|\sigma(y) \, dy \leq 2^d a^k. \tag{4.10}$$

Now,

$$\begin{aligned} \int_{\mathbb{R}^d} M_d(f\sigma)^p w \, dx &= \sum_k \int_{\Omega_k \setminus \Omega_{k+1}} M_d(f\sigma)^p w \, dx \\ &\leq a^p \sum_k a^{kp} w(\Omega_k) = a^p \sum_{k,j} a^{kp} w(Q_{k,j}) \\ &\leq a^p \sum_{k,j} \left(\frac{1}{|Q_{k,j}|} \int_{Q_{k,j}} |f(y)|\sigma(y) \, dy \right)^p w(Q_{k,j}) \\ &= a^p \sum_{k,j} (\langle |f| \rangle_{Q_{k,j}}^\sigma)^p \left(\frac{\sigma(Q_{k,j})}{|Q_{k,j}|} \right)^p w(Q_{k,j}) \\ &\leq a^p B_p[w, \sigma] \sum_{k,j} (\langle |f| \rangle_{Q_{k,j}}^\sigma)^p |Q_{k,j}| \exp\left(\int_{Q_{k,j}} \log \sigma(t) \, dt \right) \\ &= a^p B_p[w, \sigma] \sum_{Q \in \mathfrak{D}} (\langle |f| \rangle_Q^\sigma)^p a_Q, \end{aligned}$$

where

$$a_Q = \begin{cases} |Q| \exp\left(\int_Q \log \sigma \right) & \text{if } Q = Q_{k,j} \text{ for some } (k, j), \\ 0 & \text{else.} \end{cases}$$

By the dyadic Carleson embedding theorem, we can hence conclude that

$$\int_{\mathbb{R}^d} M^d(f\sigma)^p w \, dx \leq a^p B_p[w, \sigma] A \int_{\mathbb{R}^d} (M_{d,\sigma} f)^p \sigma \, dx,$$

provided that we check the condition

$$\sum_{Q \subseteq R} a_Q = \sum_{k,j: Q_{k,j} \subseteq R} |Q_{k,j}| \exp\left(\int_{Q_{k,j}} \log \sigma \right) \leq A |R|. \tag{4.11}$$

To estimate the left side of (4.11), we first do the following: for each (k, j) , we set $E_{k,j} = Q_{k,j} \setminus \Omega_{k+1}$. Observe that the sets of the family $E_{k,j}$ are pairwise disjoint. We claim that

$$|Q_{k,j}| < \frac{a}{a-2^d} |E_{k,j}| \tag{4.12}$$

for each k, j . Indeed, by (4.10) and Hölder’s inequality,

$$|Q_{k,j} \cap \Omega_{k+1}| = \sum_{Q_{k+1,l} \subset Q_{k,j}} |Q_{k+1,l}| < \frac{1}{a^{k+1}} \sum_{Q_{k+1,l} \subset Q_{k,j}} \int_{Q_{k+1,l}} |f| \sigma \leq \frac{1}{a^{k+1}} \int_{Q_{k,j}} |f| \sigma \leq \frac{2^d}{a} |Q_{k,j}|,$$

which proves (4.12). With $\beta = a/(a - 2^d)$, we can estimate the left side of (4.11) as follows:

$$\begin{aligned} \sum_{Q \subseteq R} a_Q &\leq \beta \sum_{(k,j): Q_{k,j} \subseteq R} |E_{k,j}| \exp\left(\int_{Q_{k,j}} \log \sigma(t) dt\right) \\ &\leq \beta \sum_{(k,j): Q_{k,j} \subseteq R} \int_{E_{k,j}} M_0(\sigma 1_R)(x) dx \\ &\leq \beta \int_R M_0(\sigma 1_R)(x) dx \leq \beta e \sigma(R), \end{aligned}$$

where we used the definition and the L^1 boundedness of the logarithmic dyadic maximal function. This proves (4.11) with $A = \beta e$, concluding the proof. \square

5. Proof of the extrapolation theorems

We will prove in this section the upper and lower extrapolation theorems 1.24 and 1.25. Recall that the initial hypothesis is given by the expression

$$\|Tf\|_{L^r(w)} \leq \varphi([w]_{A_r}, [w]_{A_\infty}, [w^{-1/(r-1)}]_{A_\infty}^{(r-1)}) \|f\|_{L^r(w)},$$

for some $r \in (1, \infty)$.

Proof of Theorem 1.24. Our argument is modeled after a simplified proof by Duoandikoetxea [2011] of the already cited result from [Dragičević et al. 2005] (see also [Cruz-Uribe et al. 2011]).

Fix some $p \in (1, r)$, $w \in A_p$, $f \in L^p(w)$ and $g := |f|/\|f\|_{L^p(w)}$. Let

$$Rg := \sum_{k=0}^{\infty} \frac{2^{-k} M^k g}{\|M\|_{\mathcal{B}(L^p(w))}^k},$$

so that

$$|g| \leq Rg, \quad \|Rg\|_{L^p(w)} \leq 2\|g\|_{L^p(w)} = 2, \quad [Rg]_{A_1} \leq 2\|M\|_{L^p(w)}.$$

Then by Hölder's inequality,

$$\begin{aligned} \|Tf\|_{L^p(w)} &= \left(\int |Tf|^p (Rg)^{-(r-p)p/r} (Rg)^{(r-p)p/r} w \right)^{1/p} \\ &\leq \left(\int |Tf|^r (Rg)^{-(r-p)} w \right)^{1/r} \left(\int (Rg)^p w \right)^{1/p-1/r} \\ &\leq \|Tf\|_{L^r(W)} (2^p)^{1/p-1/r} \leq 2\|Tf\|_{L^r(W)}, \end{aligned}$$

where

$$W := (Rg)^{-(r-p)} w.$$

By assumption, we have

$$\|Tf\|_{L^r(W)} \leq \varphi([W]_{A_r}, [W]_{A_\infty}, [W^{-1/(r-1)}]_{A_\infty}^{(r-1)}) \|f\|_{L^r(W)},$$

where

$$\|f\|_{L^r(W)} = \left(\int |f|^r (Rf)^{-(r-p)} w \right)^{1/r} \|f\|_{L^p(w)}^{(r-p)/r} \leq \left(\int |f|^r |f|^{-(r-p)} w \right)^{1/r} \|f\|_{L^p(w)}^{(r-p)/r} = \|f\|_{L^p(w)},$$

so it remains to estimate the weight constants

$$[W]_{A_r}, \quad [W]_{A_\infty}, \quad [W^{-1/(r-1)}]_{A_\infty}.$$

Using $\sup_Q (Rg)^{-1} \leq [Rg]_{A_1} \langle Rg \rangle_Q^{-1}$ or Hölder’s or Jensen’s inequality where appropriate, we compute

$$\begin{aligned} \langle W \rangle_Q &= \langle (Rg)^{-(r-p)} w \rangle_Q \leq [Rg]_{A_1}^{r-p} \langle Rg \rangle_Q^{-(r-p)} \langle w \rangle_Q, \\ \langle W^{-1/(r-1)} \rangle_Q^{r-1} &= \langle (Rg)^{(r-p)/(r-1)} w^{-1/(r-1)} \rangle_Q^{r-1} \leq \langle Rg \rangle_Q^{r-p} \langle w^{-1/(p-1)} \rangle_Q^{p-1}, \\ \exp\langle -\log W \rangle_Q &= (\exp\langle \log(Rg) \rangle_Q)^{r-p} \exp\langle -\log w \rangle_Q \leq \langle Rg \rangle_Q^{r-p} \exp\langle -\log w \rangle_Q, \end{aligned}$$

and

$$\begin{aligned} (\exp\langle -\log W^{-1/(r-1)} \rangle_Q)^{r-1} &= (\exp\langle \log(Rg)^{-1} \rangle_Q)^{r-p} (\exp\langle -\log w^{-1/(r-1)} \rangle_Q)^{r-1} \\ &\leq [Rg]_{A_1}^{r-p} \langle Rg \rangle_Q^{-(r-p)} (\exp\langle -\log w^{-1/(p-1)} \rangle_Q)^{p-1}. \end{aligned}$$

Multiplying the appropriate estimates and using the definition, we then have

$$[W]_{A_r} \leq [Rg]_{A_1}^{r-p} [w]_{A_p}, \quad [W]_{A_\infty} \leq [Rg]_{A_1}^{r-p} [w]_{A_\infty}, \quad [W^{-1/(r-1)}]_{A_\infty}^{r-1} \leq [Rg]_{A_1}^{r-p} [w^{-1/(p-1)}]_{A_\infty}^{p-1}.$$

(We do not know whether it is possible to make similar estimates for $[W]_{A_\infty}'$ in terms of $[w]_{A_\infty}'$; this is the reason why we need to use the $[]_{A_\infty}$ constants in this proof.)

Next, recall that

$$[Rg]_{A_1} \leq 2 \|M\|_{\mathfrak{B}(L^p(w))} \leq c_d \cdot p' \cdot [w]_{A_p}^{1/p} ([w^{-1/(p-1)}]_{A_\infty}')^{1/p}.$$

Thus we conclude the proof with

$$\begin{aligned} \|Tf\|_{L^p(w)} &\leq 2 \|Tf\|_{L^r(W)} \leq 2\varphi([W]_{A_r}, [W]_{A_\infty}, [W^{-1/(r-1)}]_{A_\infty}^{(r-1)}) \|f\|_{L^r(W)} \\ &\leq 2\varphi\left([Rg]_{A_1}^{r-p} ([w]_{A_p}, [w]_{A_\infty}, [w^{-1/(p-1)}]_{A_\infty}^{(p-1)})\right) \|f\|_{L^p(w)} \\ &\leq 2\varphi\left(2^{r-p} \|M\|_{\mathfrak{B}(L^p(w))}^{r-p} ([w]_{A_p}, [w]_{A_\infty}, [w^{-1/(p-1)}]_{A_\infty}^{(p-1)})\right) \|f\|_{L^p(w)}. \quad \square \end{aligned}$$

Proof of Theorem 1.25. Again, our argument is inspired by Duoandikoetxea’s simplification [2011] of the proof of a result in [Dragičević et al. 2005] (see also [Cruz-Uribe et al. 2011]).

Fix some $p \in (r, \infty)$, $w \in A_p$, $f \in L^p(w)$. By duality, we have

$$\|Tf\|_{L^p(w)} = \sup_{\substack{h \geq 0 \\ \|h\|_{L^{p'}(w)} = 1}} \int |Tf| h w.$$

We fix one such h , and try to bound the expression on the right.

Observe that the pointwise multiplication operators

$$h \mapsto wh : L^{p'}(w) \rightarrow L^{p'}(w^{1-p'}), \quad g \mapsto \frac{1}{w}g : L^{p'}(w^{1-p'}) \rightarrow L^{p'}(w)$$

are isometric. Let R be as in the previous proof, except with p' and $\sigma = w^{1-p'}$ in place of p and w :

$$Rg := \sum_{k=0}^{\infty} \frac{2^{-k} M^k g}{\|M\|_{\mathfrak{B}(L^{p'}(\sigma))}^k},$$

and $R'h := w^{-1}R(wh)$. Then

$$h \leq R'h, \quad \|R'h\|_{L^{p'}(w)} \leq 2\|h\|_{L^{p'}(w)} = 2, \quad [wR'h]_{A_1} \leq 2\|M\|_{\mathfrak{B}(L^{p'}(\sigma))}.$$

Then, by Hölder's inequality,

$$\begin{aligned} \int |Tf| h w &\leq \int |Tf| (R'h) w = \int |Tf| (R'h)^{(p-r)/[r(p-1)]} (R'h)^{(r-1)p/[r(p-1)]} w \\ &\leq \left(\int |Tf|^r (R'h)^{(p-r)/(p-1)} w \right)^{1/r} \left(\int (R'h)^{p/(p-1)} w \right)^{1/r'} \\ &\leq \|Tf\|_{L^r(W)} 2^{p'/r'}, \end{aligned}$$

where

$$W := (R'h)^{(p-r)/(p-1)} w.$$

By assumption,

$$\|Tf\|_{L^r(W)} \leq \varphi([W]_{A_r}, [W]_{A_\infty}, [W^{-1/(r-1)}]_{A_\infty}^{(r-1)}) \|f\|_{L^r(W)}, \quad (5.1)$$

where, by Hölder's inequality with exponents p/r and $p/(p-r)$,

$$\begin{aligned} \|f\|_{L^r(W)} &= \left(\int |f|^r w^{r/p} \cdot (R'h)^{(p-r)/(p-1)} w^{(p-r)/p} \right)^{1/r} \\ &\leq \left(\int |f|^p w \right)^{1/p} \left(\int (R'h)^{p/(p-1)} w \right)^{1/r-1/p} \leq \|f\|_{L^p(w)} (2^{p'})^{1/r-1/p}, \end{aligned}$$

so altogether, suppressing the arguments of φ from (5.1),

$$\begin{aligned} \int |Tf| h w &\leq \|Tf\|_{L^r(W)} 2^{p'/r'} \leq \varphi(\dots) \|f\|_{L^r(W)} 2^{p'/r'} \\ &\leq \varphi(\dots) (2^{p'})^{1/r-1/p} \|f\|_{L^p(w)} 2^{p'/r'} = 2\varphi(\dots) \|f\|_{L^p(w)}. \end{aligned}$$

It remains to estimate

$$[W]_{A_r}, \quad [W]_{A_\infty}, \quad [W^{-1/(r-1)}]_{A_\infty}^{(r-1)}$$

for

$$W = (R'h)^{(p-r)/(p-1)} w = [(R'h)w]^{(p-r)/(p-1)} w^{(r-1)/(p-1)}.$$

We thus compute

$$\begin{aligned} \langle W \rangle_Q &= \langle (R'h)^{(p-r)/(p-1)} w \rangle_Q \leq \langle (R'h)w \rangle_Q^{(p-r)/(p-1)} \langle w \rangle_Q^{(r-1)/(p-1)}, \\ \langle W^{-1/(r-1)} \rangle_Q^{r-1} &= \langle (wR'h)^{-(p-r)/[(p-1)(r-1)]} w^{-1/(p-1)} \rangle_Q^{r-1} \\ &\leq [wR'h]_{A_1}^{(p-r)/(p-1)} \langle (R'h)w \rangle_Q^{-(p-r)/(p-1)} \langle w^{-1/(p-1)} \rangle_Q^{r-1}, \\ \exp(-\langle \log W \rangle_Q) &= (\exp \langle \log (wR'h)^{-1} \rangle_Q)^{(p-r)/(r-1)} (\exp \langle -\log w \rangle_Q)^{(r-1)/(p-1)} \\ &\leq [(R'h)w]_{A_1}^{(p-r)/(r-1)} \langle (R'h)w \rangle_Q^{-(p-r)/(r-1)} (\exp \langle -\log w \rangle_Q)^{(r-1)/(p-1)}, \end{aligned}$$

and

$$\begin{aligned} (\exp(-\langle \log W^{-1/(r-1)} \rangle_Q))^{r-1} &= (\exp(\langle \log (wR'h) \rangle_Q))^{(p-r)/(p-1)} (\exp \langle -\log w^{-1/(p-1)} \rangle_Q)^{r-1} \\ &\leq \langle (R'h)w \rangle_Q^{(p-r)/(r-1)} (\exp \langle -\log w^{-1/(p-1)} \rangle_Q)^{r-1}. \end{aligned}$$

Multiplying the relevant quantities, it follows that

$$\begin{aligned} [W]_{A_r} &\leq [(R'h)w]_{A_1}^{(p-r)/(p-1)} [w]_{A_p}^{(r-1)/(p-1)}, \\ [W]_{A_\infty} &\leq [(R'h)w]_{A_1}^{(p-r)/(p-1)} [w]_{A_\infty}^{(r-1)/(p-1)}, \\ [W^{-1/(r-1)}]_{A_\infty}^{r-1} &\leq [(R'h)w]_{A_1}^{(p-r)/(p-1)} [w^{-1/(p-1)}]_{A_\infty}^{(r-1)}. \end{aligned}$$

Also recall that

$$[(R'h)w]_{A_1} \leq 2 \|M\|_{\mathfrak{B}(L^{p'}(w^{1-p'}))} \leq c_d [w^{1-p'}]_{A_{p'}}^{1/p'} [w]_{A_\infty}^{1/p'} = c_d [w]_{A_p}^{1/p} [w]_{A_\infty}^{1/p'},$$

and thus we conclude with

$$\begin{aligned} \|Tf\|_{L^p(w)} &\leq \int |Tf| h w \leq 2\varphi([W]_{A_r}, [W]_{A_\infty}, [W^{-1/(r-1)}]_{A_\infty}^{(r-1)}) \|f\|_{L^p(w)} \\ &\leq 2\varphi([(R'h)w]_{A_1}^{(p-r)/(p-1)} ([W]_{A_r}, [W]_{A_\infty}, [W^{-1/(r-1)}]_{A_\infty}^{(r-1)})) \|f\|_{L^p(w)} \\ &\leq 2\varphi\left((2\|M\|_{\mathfrak{B}(L^{p'}(w^{1-p'}))})^{(p-r)/(p-1)} \right. \\ &\quad \left. \times ([w]_{A_r}^{(r-1)/(p-1)}, [w]_{A_\infty}^{(r-1)/(p-1)}, [w^{-1/(p-1)}]_{A_\infty}^{(r-1)}) \right) \|f\|_{L^p(w)}. \quad \square \end{aligned}$$

6. The A_1 theory, proof of Theorem 1.14 and its consequences

6A. The main lemma. The proofs of the theorems will be based on the following lemma.

Lemma 6.1. *Let T be any Calderón–Zygmund singular integral operator and let w be any weight. Also let $p, r \in (1, \infty)$. Then there is a constant $c = c_{d,T}$ such that*

$$\|Tf\|_{L^p(w)} \leq c p p'(r')^{1/p'} \|f\|_{L^p(M_r w)},$$

where, as usual, we denote $M_r w = M(w^r)^{1/r}$.

This is a consequence of the estimate

$$\|Tf\|_{L^p(w)} \leq c p p' \left(\frac{1}{r-1}\right)^{1-1/pr} \|f\|_{L^p(M_r w)}$$

(which can be found in [Lerner et al. 2009a] when $r \in (1, 2)$), since

$$\left(\frac{1}{r-1}\right)^{1-1/pr} \leq (r')^{1-1/p+1/pr'} \leq 2(r')^{1/p'}$$

where we used $t^{1/t} \leq 2$, $t \geq 1$.

6B. Proof of the sharp reverse Hölder's inequality.

Lemma 6.2. For any cube Q and any measurable function w ,

$$\int_Q w \log\left(e + \frac{w}{\langle w \rangle_Q}\right) dx \leq 2^{d+1} \int_Q M(w \chi_Q) dx. \quad (6.3)$$

Hence, if $w \in A_\infty$,

$$\sup_Q \frac{1}{w(Q)} \int_Q w(y) \log\left(e + \frac{w(y)}{\langle w \rangle_Q}\right) dy \leq 2^{d+1} [w]'_{A_\infty}. \quad (6.4)$$

The essential idea of the proof can be traced back to the well-known $L \log L$ estimate for M in [Stein 1969]. However, these estimates are not homogeneous. A proof of this lemma within the context of spaces of homogeneous type can essentially be found in [Pérez and Wheeden 2001, Lemma 8.5] (see also [Wilson 2008, p. 17, inequality (2.15)] for a different proof).

Proof of Lemma 6.2. Fix a cube Q . By homogeneity, we assume that $\langle w \rangle_Q = 1$. The key estimate follows from the “reverse weak type (1, 1) estimate”: if w is nonnegative and $t > \langle w \rangle_Q$,

$$\frac{1}{t} \int_{\{x \in Q : w(x) > t\}} w dx \leq 2^d |\{x \in Q : M(w \chi_Q)(x) > t\}|. \quad (6.5)$$

Now,

$$\frac{1}{|Q|} \int_Q w \log(e + w) dx = \frac{1}{|Q|} \int_0^\infty \frac{1}{e+t} w(\{x \in Q : w(x) > t\}) dt = I + II.$$

Here

$$I := \frac{1}{|Q|} \int_0^1 \frac{1}{e+t} w(\{x \in Q : w(x) > t\}) dt \leq 1 \leq \frac{1}{|Q|} \int_Q M(w \chi_Q) dx,$$

while for the complementary term II we use the estimate (6.5):

$$\begin{aligned} II &= \frac{1}{|Q|} \int_1^\infty \frac{1}{e+t} w(\{x \in Q : w(x) > t\}) dt \\ &\leq \frac{2^d}{|Q|} \int_1^\infty \frac{t}{e+t} |\{x \in Q : M(w \chi_Q)(x) > t\}| dt \\ &\leq \frac{2^d}{|Q|} \int_0^\infty |\{x \in Q : M(w \chi_Q)(x) > t\}| dt \\ &= \frac{2^d}{|Q|} \int_Q M(w \chi_Q)(x) dx. \end{aligned}$$

This gives (6.3), and (6.4) follows from the definition of $[w]_{A_\infty}'$. □

The main use of the lemma is the following key observation.

Lemma 6.6 [Wilson 2008, p. 45]. *Let $S \subset Q$ and let $\lambda > 0$; then*

$$\frac{|S|}{|Q|} < e^{-\lambda} \quad \text{implies} \quad \frac{w(S)}{w(Q)} < \frac{2^{d+2}[w]_{A_\infty}'}{\lambda} + e^{-\lambda/2}. \tag{6.7}$$

Proof. Indeed, if $E_\lambda = \{x \in Q : w(x) > e^\lambda \langle w \rangle_Q\}$, then $w(E_\lambda) \leq (2^{d+1}/\lambda)[w]_{A_\infty}' w(Q)$ by (6.4). Therefore

$$\begin{aligned} w(S) &\leq w(S \cap E_{\lambda/2}) + w(S \setminus E_{\lambda/2}) \leq \frac{2^{d+2}[w]_{A_\infty}'}{\lambda} w(Q) + e^{\lambda/2} \langle w \rangle_Q |S| \\ &\leq \frac{2^{d+2}[w]_{A_\infty}'}{\lambda} w(Q) + e^{\lambda/2} e^{-\lambda} w(Q) \quad \text{by the hypothesis in (6.7)} \\ &= \frac{2^{d+2}[w]_{A_\infty}'}{\lambda} w(Q) + e^{-\lambda/2} w(Q), \end{aligned}$$

and this proves the claim (6.7). □

Proof of Theorem 2.3. Recall that we have to prove that

$$\left(\int_Q w^{r(w)} \right)^{1/r(w)} \leq 2 \int_Q w,$$

where

$$r(w) := 1 + \frac{1}{\tau_d [w]_{A_\infty}'},$$

and where τ_d is a large dimensional constant.

Observe that by homogeneity, we can assume that $\int_Q w = 1$. We use the dyadic maximal function on the dyadic subcubes of a given Q :

$$\begin{aligned} \int_Q w^{1+\varepsilon} &\leq \int_Q M_d(w \chi_Q)^\varepsilon w = \int_0^\infty \varepsilon t^{\varepsilon-1} w(\{x \in Q : M_d(w \chi_Q) > t\}) dt \\ &\leq \int_0^1 \varepsilon t^{\varepsilon-1} w(Q) dt + \varepsilon \int_1^\infty \varepsilon t^\varepsilon w(\{x \in Q : M_d(w \chi_Q) > t\}) \frac{dt}{t} \\ &\leq |Q| + \varepsilon \sum_{k \geq 0} \int_{a^k}^{a^{k+1}} t^\varepsilon w(\{x \in Q : M_d(w \chi_Q) > t\}) \frac{dt}{t} \\ &\leq |Q| + \varepsilon a^\varepsilon \sum_{k \geq 0} a^{k\varepsilon} \int_{a^k}^{a^{k+1}} w(\{x \in Q : M_d(w \chi_Q) > a^k\}) \frac{dt}{t} \quad \text{for } a \gg 1, \\ &= |Q| + \varepsilon a^\varepsilon \log a \sum_{k \geq 0} a^{k\varepsilon} w(\Omega_k), \end{aligned}$$

where

$$\Omega_k = \{x \in Q : M_d(w \chi_Q)(x) > a^k\}.$$

Since $a^k \geq 1 = \int_Q w$, we can consider the Calderón–Zygmund decomposition of w adapted to Q . There is a family of maximal nonoverlapping dyadic cubes $\{Q_{k,j}\}$ strictly contained in Q for which $\Omega_k = \bigcup_j Q_{k,j}$ and

$$a^k < \int_{Q_{k,j}} w \leq 2^d a^k. \tag{6.8}$$

Now,

$$\sum_{k \geq 0} a^{k\varepsilon} w(\Omega_k) = \sum_{k,j} a^{k\varepsilon} w(Q_{k,j}) \leq \sum_{k,j} \left(\frac{1}{|Q_{k,j}|} \int_{Q_{k,j}} w(y) dy \right)^\varepsilon w(Q_{k,j}).$$

We need to estimate $w(Q_{k,j})$, which we pursue similarly to Section 4C; see in particular (4.12). For each (k, j) we set $E_{k,j} = Q_{k,j} \setminus \Omega_{k+1}$. Observe that the sets of the family $E_{k,j}$ are pairwise disjoint. But exactly as in (4.12), we have that for $a > 2^d$ and for each k, j ,

$$|Q_{k,j}| < \frac{a}{a-2^d} |E_{k,j}|. \tag{6.9}$$

We now apply (6.7) with $Q = Q_{k,j}$ and $S = Q_{k,j} \cap \Omega_{k+1}$. Choose λ such that $e^{-\lambda} = 2^d/a$, namely $\lambda = \log(a/2^d)$. Then applying (6.7), we have that

$$\frac{w(Q_{k,j} \cap \Omega_{k+1})}{w(Q_{k,j})} < \frac{2^{d+2} [w]_{A_\infty}'}{\log(a/2^d)} + \left(\frac{2^d}{a}\right)^{1/2}.$$

Since $a > 2^d$ is available, we choose $a = 2^d e^{L[w]_{A_\infty}'}$, with L a large dimensional constant to be chosen. If in particular $L \geq 2^{d+4}$, we have

$$\frac{w(Q_{k,j} \cap \Omega_{k+1})}{w(Q_{k,j})} < \frac{2^{d+2}}{L} + e^{-[w]_{A_\infty}' L/2} < \frac{1}{4} + \frac{1}{4} = \frac{1}{2}.$$

This yields that $w(Q_{k,j}) \leq 2w(E_{k,j})$, and we can continue with the sum estimate:

$$\begin{aligned} \sum_{k \geq 0} a^{k\varepsilon} w(\Omega_k) &\leq 2 \sum_{k,j} \left(\frac{1}{|Q_{k,j}|} \int_{Q_{k,j}} w(y) dy \right)^\varepsilon w(E_{k,j}) \\ &\leq 2 \sum_{k,j} \int_{E_{k,j}} M_d(w \chi_Q)^\varepsilon w dx \leq 2 \int_Q M_d(w \chi_Q)^\varepsilon w dx. \end{aligned}$$

Combining estimates, we end up with

$$\int_Q M_d(w \chi_Q)^\varepsilon w \leq 1 + 2\varepsilon a^\varepsilon \log a \int_Q M_d(w \chi_Q)^\varepsilon w dx,$$

for any $\varepsilon > 0$. Recall that $a = 2^d e^{L[w]_{A_\infty}'}$. Hence, if we choose

$$L = 2^{d+4}, \quad \varepsilon = \frac{1}{2^7 L [w]_{A_\infty}'} = \frac{1}{2^{11+d} [w]_{A_\infty}'},$$

we can compute

$$2\varepsilon a^\varepsilon \log a < \frac{1}{2}, \quad \int_Q M_d(w \chi_Q)^\varepsilon w \leq 2,$$

concluding the proof of the theorem. □

6C. Proof of Theorem 1.14: the strong case. The proof is, as in [Lerner et al. 2009a], just an application of Lemma 6.1 with a specific parameter r coming from the sharp reverse Hölder inequality given by Theorem 2.3. Indeed, since $w \in A_1 \subset A_\infty$, and if we write

$$r(w) := 1 + \frac{1}{\tau_d [w]_{A_\infty}'},$$

we have

$$\left(\int_Q w^{r(w)} \right)^{1/r(w)} \leq 2 \int_Q w. \tag{6.10}$$

Now, by Lemma 6.1 with $r = r(w)$, we have

$$\begin{aligned} \|Tf\|_{L^p(w)} &\leq c p p'(r')^{1/p'} \|f\|_{L^p(M_r w)} \leq c p p' ([w]_{A_\infty}')^{1/p'} \|f\|_{L^p(2Mw)} \\ &\leq c p p' ([w]_{A_\infty}')^{1/p'} [w]_{A_1}^{1/p} \|f\|_{L^p(w)}, \end{aligned}$$

using the standard notation $M_r w = M(w^r)^{1/r}$. This concludes the proof of the theorem.

6D. Proof of Theorem 1.15: the weak case. We follow here the classical method of Calderón and Zygmund, with the modifications considered in [Pérez 1994]. Applying the Calderón–Zygmund decomposition to f at level λ , we get a family of pairwise disjoint cubes $\{Q_j\}$ such that

$$\lambda < \frac{1}{|Q_j|} \int_{Q_j} |f| \leq 2^d \lambda.$$

Let $\Omega = \bigcup_j Q_j$ and $\tilde{\Omega} = \bigcup_j 2Q_j$. The “good part” is defined by

$$g = \sum_j f_{Q_j} \chi_{Q_j}(x) + f(x) \chi_{\Omega^c}(x),$$

and the “bad part” b as

$$b = \sum_j b_j,$$

where

$$b_j(x) = (f(x) - f_{Q_j}) \chi_{Q_j}(x).$$

Then $f = g + b$. We split the level set as

$$\begin{aligned} w \{x \in \mathbb{R}^d : |Tf(x)| > \lambda\} &\leq w(\tilde{\Omega}) + w \left\{ x \in (\tilde{\Omega})^c : |Tb(x)| > \frac{\lambda}{2} \right\} + w \left\{ x \in (\tilde{\Omega})^c : |Tg(x)| > \frac{\lambda}{2} \right\} \\ &=: I + II + III. \end{aligned}$$

Exactly as in [Pérez 1994], the main term is III . We first deal with the easy terms I and II , which actually satisfy the better bound

$$I + II \lesssim \frac{1}{\lambda} [w]_{A_1} \|f\|_{L^1(w)}.$$

Indeed, the first term is essentially the level set of Mf :

$$I = w \{x \in \mathbb{R}^d : Mf(x) > c_d \lambda\},$$

and the result follows by the classical Fefferman–Stein inequality:

$$\|Mf\|_{L^{1,\infty}(w)} \leq c_d \|f\|_{L^1(Mw)}.$$

For the second term, we use the following estimate: there is a dimensional constant c such that for any cube Q and any function b supported on Q such that $\int_Q b(x) dx = 0$ and any weight w , we have

$$\int_{\mathbb{R}^n \setminus 2Q} |Tb(y)|w(y) dy \leq c_d \int_Q |b(y)|Mw(y) dy. \quad (6.11)$$

This can be found in Lemma 3.3 of [García-Cuerva and Rubio de Francia 1985, p. 413]. Now, using this estimate with w replaced by $w\chi_{\mathbb{R}^n \setminus 2Q_j}$, we have

$$\begin{aligned} II &\leq \frac{c}{\lambda} \int_{\mathbb{R}^n \setminus \tilde{\Omega}} |Tb(y)|w(y) dy \\ &\leq \frac{c}{\lambda} \sum_j \int_{\mathbb{R}^n \setminus 2Q_j} |Tb_j(y)|w(y) dy \leq \frac{c}{\lambda} \sum_j \int_{Q_j} |b_j(y)|M(w\chi_{\mathbb{R}^n \setminus 2Q_j})(y) dy \\ &\leq \frac{c}{\lambda} \int_{\mathbb{R}^d} |f(y)|Mw(y) dy + \frac{c}{\lambda} \sum_j \frac{1}{|Q_j|} \int_{Q_j} M(w\chi_{\mathbb{R}^n \setminus 2Q_j})(x) dx \int_{Q_j} |f(x)| dx. \end{aligned}$$

To estimate the inner sum, we use that $M(\chi_{\mathbb{R}^n \setminus 2Q}\mu)$ is essentially constant on Q :

$$M(\chi_{\mathbb{R}^n \setminus 2Q}\mu)(y) \approx M(\chi_{\mathbb{R}^n \setminus 2Q}\mu)(z), \quad y, z \in Q, \quad (6.12)$$

where the constants are dimensional. This fact that can be found in [ibid., p. 159]. Hence, the sum is controlled by

$$c_d \sum_j \inf_{x \in Q} M(w\chi_{\mathbb{R}^n \setminus 2Q_j})(x) \int_{Q_j} |f(x)| dx \leq c_d \int_{\mathbb{R}^n} |f(x)|Mw(x) dx.$$

This gives the required estimate.

We now consider the singular term III , to which we apply the Chebyshev inequality and Lemma 6.1 with exponents $p, r \in (1, \infty)$ to be chosen soon:

$$\begin{aligned} III &= w \left\{ x \in (\tilde{\Omega})^c : |Tg(x)| > \frac{\lambda}{2} \right\} \leq \frac{2^p}{\lambda^p} \|T(g)\|_{L^p(w\chi_{(\tilde{\Omega})^c})}^p \\ &\leq c(pp')^p (r')^{p/p'} \frac{1}{\lambda^p} \int_{\mathbb{R}^d} |g|^p M_r(w\chi_{(\tilde{\Omega})^c}) dx = c(pp')^p (r')^{p-1} \frac{1}{\lambda} \int_{\mathbb{R}^d} |g| M_r(w\chi_{(\tilde{\Omega})^c}) dx. \end{aligned}$$

Now, after using the definition of g , we apply the same argument as above, using (6.12) with M replaced by M_r . Then we have

$$\begin{aligned} \int_{\Omega} |g| M_r(w\chi_{(\tilde{\Omega})^c}) dx &\leq \sum_j \frac{1}{|Q_j|} \int_{Q_j} |f(x)| dx \int_{Q_j} M_r(w\chi_{\mathbb{R}^n \setminus 2Q_j})(x) dx \\ &\leq c_d \sum_j \inf_{x \in Q} M_r(w\chi_{\mathbb{R}^n \setminus 2Q_j})(x) \int_{Q_j} |f(x)| dx \leq c_d \int_{\mathbb{R}^d} |f(x)| M_r w(x) dx, \end{aligned}$$

and of course

$$\int_{\Omega^c} |g| M_r(w \chi_{(\tilde{\Omega})^c}) \, dx \leq \int_{\Omega^c} |f| M_r w \, dx.$$

Note that r is not chosen yet, and we conclude by choosing as above the exponent from [Theorem 2.3](#),

$$r = r(w) := 1 + \frac{1}{\tau_d [w]_{A_\infty}'},$$

namely the sharp A_∞ reverse Hölder’s exponent. We also choose

$$p = 1 + \frac{1}{\log(e + [w]_{A_\infty}')},$$

where $p < 2$ and $p' \approx \log(e + [w]_{A_\infty}')$. Then we continue with

$$\begin{aligned} w \left\{ x \in (\tilde{\Omega})^c : |Tg(x)| > \frac{\lambda}{2} \right\} &\leq c \log(e + [w]_{A_\infty}') \frac{[w]_{A_\infty}'^{(p-1)}}{\lambda} \int_{\mathbb{R}^d} |f| 2Mw \, dx. \\ &\leq \frac{c [w]_{A_1} (e + \log [w]_{A_\infty}')}{\lambda} \int_{\mathbb{R}^d} |f| w \, dx. \end{aligned}$$

This estimate combined with the previous ones for I and II completes the proof.

6E. Proof of [Theorem 1.16](#): the dual weak case. We adapt here the method from [\[Lerner et al. 2009b\]](#), where a variant of the Calderón–Zygmund decomposition is used — namely, the Calderón–Zygmund cubes are replaced by Whitney cubes. Fix $\lambda > 0$, and set

$$\Omega_\lambda = \{x \in \mathbb{R}^d : M_w^c(f/w)(x) > \lambda\},$$

where M_w^c denotes the weighted centered maximal function. Let $\bigcup_j Q_j$ be the Whitney covering of Ω_λ , and set the Calderón–Zygmund decomposition $f = g + b$ with respect to these cubes: the “good part” is defined by

$$g = \sum_j f_{Q_j} \chi_{Q_j}(x) + f(x) \chi_{\Omega^c}(x),$$

and then the “bad part” b is given by

$$b = \sum_j b_j,$$

where

$$b_j(x) = (f(x) - \langle f \rangle_{Q_j}) \chi_{Q_j}(x).$$

By the classical Besicovitch lemma, we have

$$w(\Omega_\lambda) \leq \frac{c_n}{\lambda} \|f\|_{L^1(\mathbb{R}^d)}.$$

Hence, we have to estimate

$$w \left\{ x \notin \Omega_\lambda : \frac{|Tf(x)|}{w(x)} > \lambda \right\} \leq w \left\{ x \notin \Omega_\lambda : \frac{|Tb(x)|}{w(x)} > \frac{\lambda}{2} \right\} + w \left\{ x \notin \Omega_\lambda : \frac{|Tg(x)|}{w(x)} > \frac{\lambda}{2} \right\} =: I_1 + I_2.$$

By using (6.11) again, with $w = 1$, we obtain

$$I_1 \leq \frac{2}{\lambda} \int_{\mathbb{R}^d \setminus \Omega_\lambda} |Tb(x)| \, dx \leq \frac{c}{\lambda} \sum_j \int_{Q_j} |f - \langle f \rangle_{Q_j}| \, dx \leq \frac{c}{\lambda} \|f\|_{L^1(\mathbb{R}^d)},$$

where $c = c_{d,T}$.

To estimate I_2 , we will use the dual version of Lemma 6.1, namely

$$\|Tf\|_{L^{p'}((M_r w)^{1-p'})} \leq c p p'(r')^{1/p'} \|f\|_{L^{p'}((w)^{1-p'})}. \tag{6.13}$$

As before, we use Theorem 2.3 with

$$r = r(w) := 1 + \frac{1}{\tau_d [w]_{A_\infty}}$$

such that

$$\left(\int_Q w^r \right)^{1/r} \leq 2 \int_Q w.$$

Then $M_r w \leq 2Mw \leq 2[w]_{A_1} w$, where, as usual, $M_r w = M(w^r)^{1/r}$. Combining the Chebyshev inequality with (6.13) for a value of $p \in (1, \infty)$ to be chosen soon, we have

$$\begin{aligned} I_2 &\leq \frac{2^{p'}}{\lambda^{p'}} \int_{\mathbb{R}^d} |Tg|^{p'} w^{1-p'} \, dx \leq \frac{4^{p'} [w]_{A_1}^{p'-1}}{\lambda^{p'}} \int_{\mathbb{R}^d} |Tg|^{p'} M_r w^{1-p'} \, dx \\ &\leq (c p p')^{p'} r' \frac{[w]_{A_1}^{p'-1}}{\lambda^{p'}} \int_{\mathbb{R}^d} |g|^{p'} w^{1-p'} \, dx \\ &\leq (c p' p)^{p'} r' \frac{[w]_{A_1}^{p'-1}}{\lambda^{p'}} \left(\int_{\mathbb{R}^d \setminus \Omega_\lambda} |f|^{p'} w^{1-p'} \, dx + \sum_j (\langle |f| \rangle_{Q_j})^{p'} \int_{Q_j} w^{1-p'} \, dx \right). \end{aligned}$$

We have that $|f| \leq \lambda w$ almost everywhere in $\mathbb{R}^d \setminus \Omega_\lambda$, and hence

$$\int_{\mathbb{R}^d \setminus \Omega_\lambda} |f|^{p'} w^{1-p'} \, dx \leq \lambda^{p'-1} \|f\|_{L^1(\mathbb{R}^d)}.$$

Next, following again [Lerner et al. 2009b], by properties of the Whitney covering, it is easy to see that for any cube Q_j there exists a cube Q_j^* such that $Q_j \subset Q_j^*$, $|Q_j^*| \leq c_n |Q_j|$, and the center of Q_j^* lies outside of Ω_λ . Therefore,

$$\begin{aligned} (\langle |f| \rangle_{Q_j})^{p'-1} \int_{Q_j} w^{1-p'} \, dx &\leq [w]_{A_1}^{p'-1} (\langle |f| \rangle_{Q_j})^{p'-1} \int_{Q_j} (Mw)^{1-p'} \, dx \\ &\leq [w]_{A_1}^{p'-1} |Q_j| \left(\frac{c \langle |f| \rangle_{Q_j^*}}{\langle w \rangle_{Q_j^*}} \right)^{p'-1} \leq (c \lambda [w]_{A_1})^{p'-1} |Q_j|, \end{aligned}$$

which gives

$$\sum_j (\langle |f| \rangle_{Q_j})^{p'} \int_{Q_j} w^{1-p'} \, dx \leq (c \lambda [w]_{A_1})^{p'-1} \sum_j \langle |f| \rangle_{Q_j} |Q_j| \leq (c \lambda [w]_{A_1})^{p'-1} \|f\|_{L^1(\mathbb{R}^d)}.$$

Combining the previous estimates and recalling that $r' \approx [w]_{A_\infty}'$, we obtain

$$I_2 \leq c^{p'} [w]'_{A_\infty} p(p')^{p'} \frac{p^{p'-1} [w]_{A_1}^{2(p'-1)}}{\lambda} \|f\|_{L^1(\mathbb{R}^d)},$$

and choosing now p such that $p' = 1 + \frac{1}{\log(e + [w]_{A_1})} \leq 2$, we get

$$I_2 \leq \frac{c [w]'_{A_\infty} \log(e + [w]_{A_1})}{\lambda} \|f\|_{L^1(\mathbb{R}^d)}.$$

This, along with estimates for I_1 and for $w(\Omega_\lambda)$, completes the proof of [Theorem 1.16](#).

7. Commutators, proof of [Theorem 1.17](#) and its consequences

For the proof, we need a sharp version of the John–Nirenberg theorem, which can be essentially found in [\[Journé 1983, pp. 31–32\]](#).

Lemma 7.1 (sharp John–Nirenberg theorem). *There are dimensional constants $0 \leq \alpha_d < 1 < \beta_d$ such that*

$$\sup_Q \frac{1}{|Q|} \int_Q \exp\left(\frac{\alpha_d}{\|b\|_{\text{BMO}}} |b(y) - \langle b \rangle_Q|\right) dy \leq \beta_d. \tag{7.2}$$

In fact, we can take $\alpha_d = 1/2^{d+2}$.

A key consequence of this lemma for the present purposes is that $e^{\text{Re } z b} w$ inherits the good weight properties of w when the complex number z is small enough. More precisely, for the A_2 constant, we have:

Lemma 7.3. *There are dimensional constants ϵ_d and c_d such that*

$$[e^{\text{Re } z b} w]_{A_2} \leq c_d [w]_{A_2} \quad \text{if } |z| \leq \frac{\epsilon_d}{\|b\|_{\text{BMO}}([w]'_{A_\infty} + [w^{-1}]'_{A_\infty})}.$$

Proof. From the reverse Hölder inequality with exponent $r = 1 + 1/(\tau_d [w]'_{A_\infty})$, and the John–Nirenberg inequality, we have for an arbitrary Q :

$$\begin{aligned} \int_Q w e^{\text{Re } z b} &\leq \left(\int_Q w^r\right)^{1/r} \left(\int_Q e^{r' \text{Re } z (b - \langle b \rangle_Q)}\right)^{1/r'} e^{\text{Re } z \langle b \rangle_Q} \\ &\leq \left(2 \int_Q w\right) \cdot \beta_d \cdot e^{\text{Re } z \langle b \rangle_Q}, \quad \text{if } |z| \leq \frac{\epsilon_d}{\|b\|_{\text{BMO}} [w]'_{A_\infty}}. \end{aligned}$$

By symmetry, we also have

$$\int_Q w^{-1} e^{-\text{Re } z b} \leq 2\beta_d \left(\int_Q w^{-1}\right) e^{-\text{Re } z \langle b \rangle_Q} \quad \text{if } |z| \leq \frac{\epsilon_d}{\|b\|_{\text{BMO}} [w^{-1}]'_{A_\infty}}.$$

Multiplication of the two estimates gives

$$\left(\int_Q w e^{\text{Re } z b}\right) \left(\int_Q w^{-1} e^{-\text{Re } z b}\right) \leq 4\beta_d^2 [w]_{A_2},$$

for all z as in the assertion, and completes the proof. □

There is an analogous statement for the A_∞ constant $[\]'_{A_\infty}$. (A similar result for $[\]_{A_\infty}$ is also true, and easier, but we will have no need for it, and it is therefore left as an exercise for the reader.)

Lemma 7.4. *There are dimensional constants ϵ_d and c_d such that*

$$[e^{\operatorname{Re} z b} w]'_{A_\infty} \leq c_d [w]'_{A_\infty} \quad \text{if } |z| \leq \frac{\epsilon_d}{\|b\|_{BMO} [w]'_{A_\infty}}.$$

Proof. We know that w satisfies the reverse Hölder inequality $(\int_Q w^{1+3\delta})^{1/(1+3\delta)} \leq 2 \int_Q w$ with a constant $\delta = c_d/[w]'_{A_\infty} < 2^{-1}$, where c_d is a small dimensional constant. We will prove that $e^{\operatorname{Re} z b} w$ satisfies a reverse Hölder estimate

$$\left(\int_Q (e^{\operatorname{Re} z b} w)^{1+\delta} \right)^{1/(1+\delta)} \leq C_d \int_Q e^{\operatorname{Re} z b} w, \tag{7.5}$$

for all z as in the assertion. By part (b) of [Theorem 2.3](#), this shows that

$$[e^{\operatorname{Re} z b} w]'_{A_\infty} \leq 2C_d/\delta \leq c_d [w]'_{A_\infty},$$

To prove (7.5), we first have

$$\begin{aligned} \left(\int_Q (e^{\operatorname{Re} z b} w)^{1+\delta} \right)^{1/(1+\delta)} &= e^{\operatorname{Re} z \langle b \rangle_Q} \left(\int_Q (e^{\operatorname{Re} z (b - \langle b \rangle_Q)} w)^{1+\delta} \right)^{1/(1+\delta)} \\ &\leq e^{\operatorname{Re} z \langle b \rangle_Q} \left(\int_Q e^{\operatorname{Re} z (b - \langle b \rangle_Q)(1+\delta)^2/\delta} \right)^{\delta/(1+\delta)^2} \left(\int_Q w^{(1+\delta)^2} \right)^{1/(1+\delta)^2}, \end{aligned}$$

where we applied Hölder’s inequality with exponents $(1 + \delta)/\delta$ and $1 + \delta$. Now

$$(1 + \delta)^2 = 1 + 2\delta + \delta^2 \leq 1 + 3\delta,$$

and hence the last factor is bounded by $2 \int_Q w$. Moreover, by [Lemma 7.1](#), we have

$$\int_Q e^{\operatorname{Re} z (b - \langle b \rangle_Q)(1+\delta)^2/\delta} \leq \beta_d \quad \text{if } |z| \leq \frac{\alpha_d \delta}{4\|b\|_{BMO}}.$$

So altogether,

$$\left(\int_Q (e^{\operatorname{Re} z b} w)^{1+\delta} \right)^{1/(1+\delta)} \leq e^{\operatorname{Re} z \langle b \rangle_Q} \cdot \beta_d \cdot 2 \int_Q w, \tag{7.6}$$

and we concentrate on the last factor. We observe that

$$\left(\int_Q w \right)^2 = \left(\int_Q w^{(1+\delta)/2} w^{(1-\delta)/2} \right)^2 \leq \left(\int_Q w^{1+\delta} \right) \left(\int_Q w^{1-\delta} \right) \leq \left(2 \int_Q w \right)^{1+\delta} \left(\int_Q w^{1-\delta} \right),$$

and hence

$$\begin{aligned} \int_Q w &\leq 2^{(1+\delta)/(1-\delta)} \left(\int_Q w^{1-\delta} \right)^{1/(1-\delta)} \leq 8 \left(\int_Q w^{1-\delta} e^{\operatorname{Re} z b(1-\delta)} e^{-\operatorname{Re} z b(1-\delta)} \right)^{1/(1-\delta)} \\ &\leq 8 \left(\int_Q w e^{\operatorname{Re} z b} \right) \left(\int_Q e^{-\operatorname{Re} z b(1-\delta)/\delta} \right)^{\delta/(1-\delta)}, \end{aligned}$$

where we used Hölder’s inequality with exponents $1/(1 - \delta)$ and $1/\delta$.

Combining with (7.6), we have shown that

$$\begin{aligned} \left(\int_Q (e^{\operatorname{Re} z b} w)^{1+\delta}\right)^{1/(1+\delta)} &\leq e^{\operatorname{Re} z \langle b \rangle_Q} \cdot \beta_d \cdot 16 \left(\int_Q w e^{\operatorname{Re} z b}\right) \left(\int_Q e^{-\operatorname{Re} z b(1-\delta)/\delta}\right)^{\delta/(1-\delta)} \\ &= 16\beta_d \cdot \left(\int_Q w e^{\operatorname{Re} z b}\right) \left(\int_Q e^{-\operatorname{Re} z (b-\langle b \rangle_Q)(1-\delta)/\delta}\right)^{\delta/(1-\delta)} \\ &\leq 16\beta_d \cdot \left(\int_Q w e^{\operatorname{Re} z b}\right) \cdot \beta_d, \end{aligned}$$

provided that $|z| \leq \alpha_d \delta / \|b\|_{\text{BMO}}$ in the last step. Altogether, we have proven (7.5) with $C_d = 16\beta_d^2$, under the condition that $|z| \leq \alpha_d \delta / (4\|b\|_{\text{BMO}})$, and this completes the proof. \square

Proof of Theorem 1.17. The proof is a revised version of that of [Chung et al. 2012], following the second proof in the classical L^p theorem for commutators that can be found in [Coifman et al. 1976]. Indeed, we begin by considering the “conjugate” of the operator given by

$$T_z(f) = e^{zb} T(e^{-zb} f),$$

where z is any complex number. Then a computation gives (for instance for “nice” functions)

$$[b, T](f) = \frac{d}{dz} T_z(f)|_{z=0} = \frac{1}{2\pi i} \int_{|z|=\epsilon} \frac{T_z(f)}{z^2} dz, \quad \epsilon > 0,$$

by the Cauchy integral theorem. Now, by Minkowski’s inequality

$$\|[b, T](f)\|_{L^2(w)} \leq \frac{1}{2\pi\epsilon^2} \int_{|z|=\epsilon} \|T_z(f)\|_{L^2(w)} |dz|, \quad \epsilon > 0, \tag{7.7}$$

all we need to do is estimate $\|T_z(f)\|_{L^2(w)} = \|T(e^{-zb} f)\|_{L^2(e^{2\operatorname{Re} z b} w)}$, for $|z| = \epsilon$ with appropriate ϵ . By the main hypothesis of the theorem, we have

$$\|T(e^{-zb} f)\|_{L^2(w)} \leq \varphi\left([e^{2\operatorname{Re} z b} w]_{A_2}, [e^{2\operatorname{Re} z b} w]_{A_\infty}', [e^{2\operatorname{Re} z b} \sigma]_{A_\infty}'\right) \|e^{-zb} f\|_{L^2(e^{2\operatorname{Re} z b} w)},$$

where $\|e^{-zb} f\|_{L^2(e^{2\operatorname{Re} z b} w)} = \|f\|_{L^2(w)}$.

By Lemmas 7.3 and 7.4 (the latter applied to both w and w^{-1}), we have

$$[we^{2\operatorname{Re} b z}]_{A_2} \leq C_d [w]_{A_2}, \quad [we^{2\operatorname{Re} b z}]_{A_\infty}' \leq C_d [w]_{A_\infty}', \quad [w^{-1}e^{-2\operatorname{Re} b z}]_{A_\infty}' \leq C_d [w^{-1}]_{A_\infty}',$$

provided that

$$|z| = \epsilon \leq \frac{\epsilon_d}{\|b\|_{\text{BMO}}([w]_{A_\infty}' + [w^{-1}]_{A_\infty}')}. \tag{7.7}$$

Using this radius and the above estimates in (7.7), we obtain

$$\begin{aligned} \|[b, T](f)\|_{L^2(w)} &\leq \frac{1}{2\pi\epsilon^2} \int_{|z|=\epsilon} \varphi(C_d [w]_{A_2}, C_d [w]_{A_\infty}', C_d [w^{-1}]_{A_\infty}') \|f\|_{L^2(w)} |dz| \\ &\leq C_d \|b\|_{\text{BMO}}([w]_{A_\infty}' + [w^{-1}]_{A_\infty}') \times \varphi(C_d [w]_{A_2}, C_d [w]_{A_\infty}', C_d [w^{-1}]_{A_\infty}') \|f\|_{L^2(w)}. \end{aligned}$$

This concludes the proof of the main part of the theorem. The estimate for T_b^k is deduced by iterating from the case $k = 1$. □

8. Examples

We compare our new estimates with earlier quantitative results by means of some examples.

8A. Power weights and the maximal inequality. Let $d = 1$ and $p \in (1, \infty)$ be fixed; we do not pay attention to the dependence of multiplicative constants on p . For $w(x) = |x|^\alpha$ and $-1 < \alpha < p - 1$, one easily checks that

$$[w]_{A_p} \approx \frac{1}{1+\alpha} \cdot \frac{1}{((p-1)-\alpha)^{p-1}}, \quad [w]_{A_\infty} \approx \frac{1}{1+\alpha}, \quad [w^{-1/(p-1)}]_{A_\infty} \approx \frac{1}{(p-1)-\alpha};$$

moreover, the functionals $[\]_{A_\infty}$ and $[\]'_{A_\infty}$ are comparable for these weights.

Letting $\alpha \rightarrow -1$ or $\alpha \rightarrow p - 1$, this shows that we have power weights with $[w]_{A_p} = t \gg 1$ and either $[w]_{A_\infty} \approx t$ and $[w^{-1/(p-1)}]_{A_\infty} \approx 1$, or $[w]_{A_\infty} \approx 1$ and $[w^{-1/(p-1)}]_{A_\infty} \approx t^{1/(p-1)}$.

With $[w]_{A_p} \approx [w]_{A_\infty} \approx t \gg 1$ and $[w^{-1/(p-1)}]_{A_\infty} \approx 1$, our maximal estimate

$$\|M\|_{\mathcal{B}(L^p(w))} \lesssim ([w]_{A_p} [w^{-1/(p-1)}]_{A_\infty})^{1/p} \approx t^{1/p}$$

clearly improves on Buckley’s bound

$$\|M\|_{\mathcal{B}(L^p(w))} \lesssim [w]_{A_p}^{1/(p-1)} \approx t^{1/(p-1)}.$$

Despite this improvement over earlier estimates, our bounds fail to provide a two-sided estimate for the norm of the maximal operator: A. Lerner and S. Ombrosi (personal communication, 2008) have constructed a family of weights which shows that

$$\inf_{w \in A_2} \frac{\|M\|_{\mathcal{B}(L^2(w))}}{([w]_{A_2} [w^{-1}]'_{A_\infty})^{1/2}} = 0.$$

The weights of their example are products of power weights and the two-valued weights considered in the next subsection.

8B. Two-valued weights and Calderón–Zygmund operators. The estimates for the Muckenhoupt constants of power weights in the previous subsection show that

$$[w]_{A_2} \approx [w]_{A_\infty} + [w^{-1}]_{A_\infty} \approx [w]'_{A_\infty} + [w^{-1}]'_{A_\infty} \quad \text{for } w(x) = |x|^\alpha \text{ and } d = 1,$$

so the improvement of our bound

$$\|T\|_{\mathcal{B}(L^2(w))} \lesssim [w]_{A_2}^{1/2} ([w]'_{A_\infty} + [\sigma]'_{A_\infty})^{1/2}$$

over $\|T\|_{\mathcal{B}(L^2(w))} \lesssim [w]_{A_2}$ is invisible to such weights.

However, the difference can be observed with weights of the form $w = t \cdot \chi_E + \chi_{\mathbb{R} \setminus E}$, where $t > 0$ and $E \subset \mathbb{R}$ is a measurable set, so that both E and $\mathbb{R} \setminus E$ have positive Lebesgue measure. As I ranges over all intervals of \mathbb{R} , the ratio $|E \cap I|/|I|$ ranges (at least) over all values $\alpha \in (0, 1)$, and hence

$$[w]_{A_2} = \sup_{\alpha \in (0,1)} (\alpha t + 1 - \alpha)(\alpha t^{-1} + 1 - \alpha) = \frac{(t+1)^2}{4t},$$

$$[w]_{A_\infty} = \sup_{\alpha \in (0,1)} f(\alpha), \quad \text{with } f(\alpha) := (\alpha t + 1 - \alpha)e^{-\alpha \log t}.$$

Now $f'(\alpha) = 0$ at the unique point $\hat{\alpha} = \frac{1}{\log t} - \frac{1}{t-1} \in (0, 1)$, and so

$$[w]_{A_\infty} = f(\hat{\alpha}) = e^{-1} \frac{t-1}{\log t} \exp \frac{\log t}{t-1} \approx \begin{cases} \frac{t}{\log t} & \text{if } t \gg 1, \\ \frac{t^{-1}}{\log t^{-1}} & \text{if } 0 < t \ll 1. \end{cases}$$

Assume that $t \gg 1$, so $[w]_{A_\infty} \approx t / \log t$. Since σ is a weight of the same form with $t^{-1} \ll 1$ in place of t , we also have $[\sigma]_{A_\infty} \approx t / \log t$. Thus

$$[w]_{A_2} \approx t, \quad \|T\|_{\mathcal{B}(L^2(w))} \lesssim [w]_{A_2}^{1/2} ([w]_{A_\infty} + [\sigma]_{A_\infty})^{1/2} \approx \frac{t}{\sqrt{\log t}}.$$

In particular, these estimates already show that

$$\inf_{w \in A_2} \frac{\|T\|_{\mathcal{B}(L^2(w))}}{[w]_{A_2}} = 0.$$

If we use the sharper version of our A_2 theorem with the weight constants $[\]'_{A_\infty}$ instead, we find that $\|T\|_{\mathcal{B}(L^2(w))}$ can actually grow much slower than $[w]_{A_2}$:

Lemma 8.1. *For $w = t \cdot \chi_E + \chi_{\mathbb{R} \setminus E}$ and $t \geq 3$, we have $[w]'_{A_\infty} \leq 4 \log t$.*

With the earlier estimate for $[w]_{A_\infty}$, this shows that $[w]_{A_\infty}$ can be exponentially larger than $[w]'_{A_\infty}$. In fact, Lemma 3.11 of [Beznosova and Reznikov 2011] implies the even more surprising possibility that $[w]'_{A_\infty} \lesssim \log \log [w]_{A_\infty}$, which is also sharp, in that the converse always holds [ibid., Theorem 1.2]; however, the example there consists of the power weights $w(x) = |x|^t$ with $t \rightarrow \infty$, which fall outside A_2 as soon as $t \geq 1$, so they are not directly relevant for the present discussion of sharp A_2 bounds.

Proof. Note that

$$\begin{aligned} \chi_I M(w \chi_I) &= \chi_I \sup_{J \subseteq I} \chi_J \int_J w = \chi_I \sup_{J \subseteq I} \chi_J \frac{1}{|J|} (|J \setminus E| + t|J \cap E|) \\ &= \chi_I \sup_{J \subseteq I} \chi_J \left(1 + (t-1) \frac{|J \cap E|}{|J|} \right) = \chi_I (1 + (t-1) M(\chi_{I \cap E})), \end{aligned}$$

and hence, with the abbreviations $\tau := t - 1$ and $a := \frac{|I \cap E|}{|I|}$,

$$\begin{aligned} \int_I M(w \chi_I) &= |I| + \tau \int_I M(\chi_{I \cap E}) = |I| + \tau \int_0^1 |I \cap \{M(\chi_{I \cap E}) > \lambda\}| \, d\lambda \\ &\leq |I| + \tau \left(\int_0^a |I| \, d\lambda + \int_a^1 \frac{2}{\lambda} |I \cap E| \, d\lambda \right) = |I| + \tau \left(a|I| + 2|I \cap E| \log \frac{1}{a} \right) \\ &= |I| + \tau |I \cap E| \left(1 + 2 \log \frac{|I|}{|I \cap E|} \right), \end{aligned}$$

where the factor 2 is the weak-type $(1, 1)$ norm of the maximal operator on the real line. Since $w(I) = |I| + \tau|I \cap E|$, we have

$$[w]'_{A_\infty} = \sup_I \frac{1}{w(I)} \int_I M(w\chi_I) \leq \sup_{\alpha \in (0,1)} \frac{1 + \tau\alpha(1 + 2 \log \alpha^{-1})}{1 + \tau\alpha} = 1 + 2 \sup_{\alpha \in (0,1)} \frac{\tau\alpha}{1 + \tau\alpha} \log \frac{1}{\alpha}, \tag{8.2}$$

recalling that the ratio $|I \cap E|/|I|$ attains at least all values $\alpha \in (0, 1)$ as I ranges over all intervals.

If $\alpha \geq \tau^{-1}$, then $\log \alpha^{-1} \leq \log \tau$, while $\tau\alpha/(1 + \tau\alpha) \leq 1$. If $\alpha \leq \tau^{-1}$, then

$$\tau\alpha \log \frac{1}{\alpha} = \tau\alpha \log \frac{1}{\tau\alpha} + \tau\alpha \log \tau \leq \frac{1}{e} + \log \tau,$$

as $x \log x^{-1} \leq e^{-1}$ and $x \leq 1$ for $x = \tau\alpha \in (0, 1)$. Altogether, recalling that $t = \tau + 1 \geq 3$, we have

$$[w]'_{A_\infty} \leq 1 + 2\left(\frac{1}{e} + \log \tau\right) \leq \left(1 + \frac{2}{e}\right) + 2 \log t \leq 4 \log t. \quad \square$$

Since $\sigma = w^{-1}$ is a weight of the same form, we find that for these particular weights,

$$[w]_{A_2} \approx t, \quad \|T\|_{\mathcal{B}(L^2(w))} \lesssim [w]_{A_2}^{1/2} ([w]'_{A_\infty} + [\sigma]'_{A_\infty})^{1/2} \lesssim (t \log t)^{1/2},$$

so indeed $\|T\|_{\mathcal{B}(L^2(w))}$ can grow much slower than $[w]_{A_2}$ for such particular families of weights. This example also motivates the use of the A_∞ constants $[w]'_{A_\infty}$, rather than $[w]_{A_\infty}$, whenever this is possible.

In a similar way, we can show that the main result from [Theorem 1.14](#) strictly improves on the earlier estimate [\(1.12\)](#). Indeed, if we let w be the previous weight with $t \gg 1$ so that $[w]_{A_1} \approx t$ and $[w]'_{A_\infty} \approx \log t$, then

$$[w]_{A_1}^{1/p} [w]'_{A_\infty}^{1/p'} \approx t^{1/p} (\log t)^{1/p'}.$$

As above, this family of weights shows that

$$\inf_{w \in A_1} \frac{\|T\|_{\mathcal{B}(L^p(w))}}{[w]_{A_1}} = 0, \quad 1 < p < \infty.$$

8C. Two-valued weights and dyadic shifts. Although it was not stated explicitly above, from the proof it is clear that our weighted bound for the dyadic shifts only depends on the dyadic Muckenhoupt constants, where the supremum is over dyadic cubes only, instead of all cubes. This makes a difference for the two-valued weights $w = t \cdot \chi_E + \chi_{\mathbb{R} \setminus E}$ considered above, when the set E is appropriately chosen. Indeed, with $E := \bigcup_{k \in \mathbb{Z}} [2k, 2k + 1)$, one observes that the ratio $|E \cap I|/|I|$ only attains the values $0, \frac{1}{2}, 1$ as I ranges over the dyadic intervals. Consequently, the dyadic A_∞ constant has a different expression:

$$[w]_{A_\infty}^d = \max_{\alpha \in \{0, 1/2, 1\}} (\alpha t + 1 - \alpha)e^{-\alpha t} = \frac{t+1}{2\sqrt{t}} = ([w]_{A_2}^d)^{1/2},$$

where $[w]_{A_2}^d = [w]_{A_2}$, as one easily observes. Repeating the proof of [Lemma 8.1](#) in the dyadic case (recalling that the weak-type $(1, 1)$ norm is $C_d = 1$ for the dyadic maximal operator), we get in place of [\(8.2\)](#) that

$$[w]'_{A_\infty}{}^d \leq 1 + \sup_{\alpha \in \{0, 1/2, 1\}} \frac{\tau\alpha}{1 + \tau\alpha} \log \frac{1}{\alpha} = 1 + \frac{\frac{1}{2}\tau}{1 + \frac{1}{2}\tau} \log 2 \leq 1 + \log 2.$$

So these constants are actually uniformly bounded over the choice of the parameter t .

By symmetry, we also have $[w^{-1}]_{A_\infty}^d = [w]_{A_\infty}^d$ and $[w^{-1}]_{A_\infty}'^d = [w]_{A_\infty}'^d$, and hence, for this particular E and $w = t \cdot \chi_E + \chi_{\mathbb{R} \setminus E}$,

$$\|\text{III}\|_{\mathfrak{B}(L^2(w))} \lesssim (r + 1)^2 ([w]_{A_2}^d)^{1/2} ([w]_{A_\infty}'^d + [w^{-1}]_{A_\infty}'^d)^{1/2} \lesssim (r + 1)^2 ([w]_{A_2}^d)^{1/2}.$$

Hruščev’s A_∞ constants $[]_{A_\infty}^d$ would have given the weaker bound $\|\text{III}\|_{\mathfrak{B}(L^2(w))} \lesssim (r + 1)^2 ([w]_{A_2}^d)^{3/4}$, instead.

8D. The extrapolated bounds for Calderón–Zygmund operators. It is interesting to compare our estimate (1.28), namely

$$\begin{aligned} \|T\|_{\mathfrak{B}(L^p(w))} &\lesssim [w]_{A_p}^{2/p-1/[2(p-1)]} ([w]_{A_\infty}^{1/[2(p-1)]} + [\sigma]_{A_\infty}^{1/2}) ([w]_{A_\infty}')^{1-2/p} \\ &\lesssim [w]_{A_p}^{2/p} ([w]_{A_\infty}')^{1-2/p}, \end{aligned} \tag{8.3}$$

which is valid for any Calderón–Zygmund operator and for all $p \geq 2$, with an estimate implicitly contained in the proof of a related result by Lerner [2011, Theorem 1.2]. He considers maximal truncations T_* of convolution-type Calderón–Zygmund operators, and obtains the bound

$$\|T_*\|_{\mathfrak{B}(L^p(w))} \lesssim [w]_{A_p}^{1/2} ([w]_{A_\infty}')^{1/2} + \|M\|_{\mathfrak{B}(L^p(w))} \lesssim [w]_{A_p}^{1/2} ([w]_{A_\infty}')^{1/2}, \quad p \in [3, \infty), \tag{8.4}$$

where the second estimate is an application of Buckley’s result (we do not even need our improvement at this point),

$$\|M\|_{\mathfrak{B}(L^p(w))} \lesssim [w]_{A_p}^{1/(p-1)} \leq [w]_{A_p}^{1/2}, \quad p \in [3, \infty).$$

In (8.4), the factor $([w]_{A_\infty}')^{1/2}$ comes from an estimate of Wilson [1989] relating the weighted norms of the grand maximal function and a certain square function, while $[w]_{A_p}^{1/2}$ is Lerner’s bound for the weighted norm of such square functions (whose exponent is optimal by [Cruz-Uribe et al. 2012]).

To simplify comparison, let us only consider the simpler form of our bound (8.3). Then the sum of the powers of $[w]_{A_p}$ and $[w]_{A_\infty}'$ in both (8.3) and (8.4) is $2/p + (1 - 2/p) = \frac{1}{2} + \frac{1}{2} = 1$, and the sharper bound is the one where the larger weight constant $[w]_{A_p}$ has the smaller power. We have $2/p \leq \frac{1}{2}$ if and only if $p \geq 4$, and hence Lerner’s bound is sharper for $p \in [3, 4)$ and ours for $p \in (4, \infty)$. This indicates that the present results are not the last word on joint A_p – A_∞ -control, but there is place for further investigation (which indeed has already taken place since the first public distribution of this paper; see [Hytönen et al. 2011, Section 12; Lacey 2012; Hytönen and Lacey 2011]).

9. Proof of the end-point theory at $p = \infty$

The proof again relies on the sharp reverse Hölder inequality in Theorem 2.3: if $w \in A_\infty$ and if we let

$$r = r(w) := 1 + \frac{1}{c_d [w]_{A_\infty}'},$$

then

$$\left(\int_Q w^r \, dx \right)^{1/r} \leq \frac{2}{|Q|} \int_Q w.$$

Proof of Theorem 1.19. For $c = \langle f \rangle_Q$,

$$\begin{aligned} \frac{1}{w(Q)} \int_Q |f - c|w &= \frac{|Q|}{w(Q)} \int_Q |f - c|w \leq \frac{|Q|}{w(Q)} \left(\int_Q |f - c|^{r(w)'} \right)^{1/r(w)'} \left(\int_Q w^{r(w)} \right)^{1/r(w)} \\ &\leq \frac{|Q|}{w(Q)} (C_d r(w)') \|f\|_{\text{BMO}} \left(2 \int_Q w \right) = C_d r(w)' \|f\|_{\text{BMO}} \leq C_d [w]'_{A_\infty} \|f\|_{\text{BMO}}, \end{aligned}$$

which shows that $\|f\|_{\text{BMO}(w)} \leq C_d [w]'_{A_\infty} \|f\|_{\text{BMO}}$. Note that we used the sharp order of growth of the local L^p norms of BMO functions as $p \rightarrow \infty$, which follows easily from the exponential integrability.

To see the sharpness for $d = 1$, consider $w(x) = |x|^{-1+\varepsilon}$, which has $[w]_{A_\infty} \approx [w]'_{A_\infty} \approx 1/\varepsilon$ and $f(x) = \log|x|$. We check that

$$\|f\|_{\text{BMO}(w)} \geq \inf_a \frac{1}{w([0, 1])} \int_0^1 \left| \log \frac{1}{x} - a \right| w(x) dx \geq \frac{c}{\varepsilon} \geq c[w]_{A_\infty} \geq c[w]'_{A_\infty},$$

which proves the claim. It is immediate that $w([0, 1]) = \int_0^1 x^{-1+\varepsilon} dx = \frac{1}{\varepsilon}$. It remains to compute

$$\int_0^1 \left| \log \frac{1}{x} - a \right| x^{-1+\varepsilon} dx = \int_0^\infty |t - a| e^{-\varepsilon t} dt = \frac{1}{\varepsilon^2} \int_0^\infty |u - \varepsilon a| e^{-u} du.$$

It suffices to check that $\psi(\alpha) := \int_0^\infty |u - \alpha| e^{-u} du \geq c > 0$ for all $\alpha \in \mathbb{R}$. But this is an easy calculus exercise. \square

We now prove [Corollary 1.20](#) on end-point estimates for Calderón–Zygmund operators.

Proof of Corollary 1.20. For the positive estimate, it suffices to factorize $T = I \circ T$, where $T : L^\infty \rightarrow \text{BMO}$ and $I : \text{BMO} \rightarrow \text{BMO}(w)$ have norm bounds c_T and $c_d [w]'_{A_\infty}$, respectively. Concerning sharpness, note that the Hilbert transform of $\chi_{(-1,0)}$ is $\log(x+1) - \log x$ for $x > 0$. Since $\log(x+1)$ is bounded on $[0, 1]$, the computation proving the sharpness of the embedding $\text{BMO} \hookrightarrow \text{BMO}(w)$ also gives the lower bound

$$\|H\chi_{(-1,0)}\|_{\text{BMO}(|x|^{-1+\varepsilon})} \geq \frac{c}{\varepsilon} = c[x^{-1+\varepsilon}]_{A_\infty} \|\chi_{(-1,0)}\|_{L^\infty} \geq c[x^{-1+\varepsilon}]'_{A_\infty} \|\chi_{(-1,0)}\|_{L^\infty}. \quad \square$$

We conclude with the proof of [Proposition 1.21](#) on the sharp relation of A_∞ and BMO. Note that here we use the larger constant $[w]_{A_\infty}$, not $[w]'_{A_\infty}$.

Proof of Proposition 1.21. Let Q be a cube. We estimate

$$\begin{aligned} \int_Q |\log w - \log c| &= \int_{Q \cap \{w \geq c\}} \log \frac{w}{c} + \int_{Q \cap \{w < c\}} \log \frac{c}{w} \\ &= \int_{Q \cap \{w \geq c\}} \log \frac{w}{c} + \left(\int_Q - \int_{Q \cap \{w \geq c\}} \right) \log \frac{c}{w} \\ &= 2 \int_{Q \cap \{w \geq c\}} \log \frac{w}{c} + \int_Q \log c + \int_Q \log \frac{1}{w} \\ &\leq 2 \int_{Q \cap \{w \geq c\}} \frac{w}{c} + |Q| \log c + |Q| \log \left([w]_{A_\infty} / \int_Q w \right). \end{aligned}$$

Hence

$$\int_Q |\log w - \log c| \leq \frac{2}{c} \int_Q w + \log c + \log[w]_{A_\infty} - \log\left(\int_Q w\right).$$

Choosing $c = c_Q = 2 \int_Q w$, we get

$$\int_Q |\log w - \log c_Q| \leq 1 + \log 2 + \log\left(\int_Q w\right) + \log[w]_{A_\infty} - \log\left(\int_Q w\right) = \log(2e[w]_{A_\infty}),$$

and this proves that

$$\|\log w\|_{\text{BMO}} \leq \log(2e[w]_{A_\infty}). \quad \square$$

Remark 9.1. In the last estimate, we cannot replace $[w]_{A_\infty}$ by $[w]'_{A_\infty}$. Indeed, for the two-valued weight $w = t \cdot 1_E + 1_{\mathbb{R} \setminus E}$, one readily checks that $\|\log w\|_{\text{BMO}} \approx \log t$, whereas [Lemma 8.1](#) shows that also $[w]'_{A_\infty} \lesssim \log t$. Thus $\|\log w\|_{\text{BMO}} \leq \log(c[w]'_{A_\infty})$ would lead to the obvious contradiction that $\log t \leq c + \log \log t$.

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Cauchy problem for ultrasound-modulated EIT GUILLAUME BAL	751
Sharp weighted bounds involving A_∞ TUOMAS HYTÖNEN and CARLOS PÉREZ	777
Periodicity of the spectrum in dimension one ALEX IOSEVICH and MIHAL N. KOLOUNTZAKIS	819
A codimension-two stable manifold of near soliton equivariant wave maps IOAN BEJENARU, JOACHIM KRIEGER and DANIEL TATARU	829
Discrete Fourier restriction associated with KdV equations YI HU and XIAOCHUN LI	859
Restriction and spectral multiplier theorems on asymptotically conic manifolds COLIN GUILLARMOU, ANDREW HASSELL and ADAM SIKORA	893
Homogenization of Neumann boundary data with fully nonlinear operator SUNHI CHOI, INWON C. KIM and KI-AHM LEE	951
Long-time asymptotics for two-dimensional exterior flows with small circulation at infinity THIERRY GALLAY and YASUNORI MAEKAWA	973
Second order stability for the Monge–Ampère equation and strong Sobolev convergence of optimal transport maps GUIDO DE PHILIPPIS and ALESSIO FIGALLI	993