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ON THE BOUNDARY VALUES OF SOLUTIONS OF THE HEAT EQUATION

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ON THE BOUNDARY VALUES OF SOLUTIONS OF THE HEAT EQUATION

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1. Introduction. In a recent paper Hartman and Wintner [3] consider solutions of the heat equation

(1)
$$u_{xx}(x, t) - u_t(x, t) = 0$$

in a rectangle R: 0 < x < 1 ($0 \le t < k \le \infty$). There they obtain necessary and sufficient conditions for a solution of (1) in R to be representable in the form

(2)
$$u(x, t) = \int_{0+}^{1-0} G(x, t; y, s) dA(y)$$

+ $\int_{0}^{t} G_{y}(x, t; 0, s) dB(s) - \int_{0}^{t} G_{y}(x, t; 1, s) dC(s),$

the Green's function G being defined by

(3)
$$G(x, t; y, s) = \frac{1}{2} \left[\vartheta_3 \left(\frac{x - y}{2}, t - s \right) - \vartheta_3 \left(\frac{x + y}{2}, t - s \right) \right]$$

where ϑ_3 is the well known Jacobi theta function. (The first integral in (2) is an absolutely convergent improper Riemann-Stieltjes integral.) They proceed to show that the functions representable in the form (2) exhibit the following behavior at the boundary of R:

(4)
$$\lim_{t \to 0+} u(x, t) = A'(x),$$

(5)
$$\lim_{x \to 0^+} u(x, t) = B'(t), \lim_{x \to 1^{-0}} u(x, t) = C'(t)$$

wherever the derivatives in question exist.

In the present note we present an improvement of (5) first given in the author's thesis [2]. The admittedly slight mathematical improvement is physically significant. A solution of (1) which admits the representation (2) gives the

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temperature at time t and position x in an insulated rod of length unity and with a certain initial temperature distribution, given essentially by (4), and imposed end temperatures, given essentially by (5). We note that such solutions are not uniquely determined by (4) and (5).

As x approaches the boundary of R along a line $t = t_0$, it seems intuitively clear that the limit should be independent of values of B (or C) for $t \ge t_0$. Hence the expected result (for the left side of R) would be

$$\lim_{x \to 0^+} u(x, t) = B'(t - 0) = \lim_{h \to 0^+} \frac{B(t - h) - B(t - 0)}{-h}$$

wherever this derivative exists.

2. Theorem. For the above improvement it is sufficient to establish the following result.

THEOREM. If B(s) is of bounded variation on every closed interval of $0 \le s < k \le \infty$, then

$$\lim_{x \to 0+} \int_0^t G_y(x, t; 0, s) \, dB(s) = B'(t-0)$$

wherever this derivative exists.

Proof. Let

$$u(x, t) = \int_0^t G_y(x, t; 0, s) dB(s).$$

Then since

$$\vartheta_3\left(\frac{x}{2},t\right) = (\pi t)^{-1/2} \sum_{n=-\infty}^{\infty} \exp\left[\frac{-(x+2n)^2}{4t}\right]$$

(see, for example, [1, p. 307]), we can write

$$u(x, t) = \frac{1}{2} x \pi^{-1/2} \int_0^t (t-s)^{-3/2} \exp \left[\frac{-x^2}{4(t-s)}\right] dB(s)$$

$$+ \frac{1}{2} \pi^{-1/2} \int_0^t (t-s)^{-3/2} \sum_{\substack{n = -\infty \\ n \neq 0}}^\infty (x+2n) \exp\left[\frac{-(x+2n)^2}{4(t-s)}\right] dB(s).$$

Clearly the latter integral vanishes with x. Then denoting the first integral on

the right by *l* and by setting $z = x^2/4$ and t - s = 1/v, we get

$$I = \left(\frac{z}{\pi}\right)^{1/2} \int_{v=1/t}^{\infty} e^{-zv} v^{3/2} dB(t-1/v).$$

If we define

$$\alpha(v) = \begin{cases} \int_{r=a}^{v} r^{3/2} dB(t-1/r) & (v \ge 1/t), \\ \\ \alpha(1/t) & (v < 1/t), \end{cases}$$

where a is a suitable constant, then we have

$$I = \left(\frac{z}{\pi}\right)^{1/2} \int_0^\infty e^{-zv} d\alpha(v).$$

To evaluate $\lim_{z\to\infty} I$ we apply [4, Theorem 1, p. 181], which states: If

$$f(s) = \int_0^\infty e^{-st} d\alpha(t),$$

then for any $\gamma \geq 0$ any constant A we have

$$\lim_{s\to 0^+} |S^{\gamma}f(s) - A| \leq \lim_{t\to\infty} |\alpha(t) t^{-\gamma} \Gamma(\gamma + 1) - A|.$$

To this end we evaluate $\lim_{v \to \infty} v^{-1/2} \alpha(v)$. Now

$$v^{-1/2} \ \alpha(v) = v^{-1/2} \int_{r=a}^{v} r^{3/2} \ dB(t-1/r)$$

= $v^{-1/2} \int_{a}^{v} r^{3/2} \ d[B(t-1/r) - B(t-0)]$
= $r^{3/2} \ v^{-1/2} [B(t-1/r) - B(t-0)] \Big|_{a}^{v}$
+ $\frac{3}{2} \ v^{-1/2} \int_{a}^{v} [B(t-0) - B(t-1/r)] r^{1/2} \ dr$
= $\frac{B(t-1/v) - B(t-0)}{1/v} - \frac{B(t-1/a) - B(t-0)}{v^{1/2}} a^{3/2}$
+ $\frac{3}{2} \ v^{-1/2} \ \int_{a}^{v} [B(t-0) - B(t-1/r)] r^{1/2} \ dr.$

As $v \to \infty$ the first expression on the right tends to -B'(t-0), if this derivative exists, and the second vanishes. Now consider the integral term: given $\epsilon > 0$, choose T so large that

$$\left|B'(t-0) - \frac{B(t-0) - B(t-1/r)}{1/r}\right| < \epsilon \text{ if } r > T.$$

Then

$$\frac{3}{2} v^{-1/2} \int_{r=a}^{v} \left[B(t-0) - B(t-1/r) \right] r^{1/2} dr$$

$$= \frac{3}{2} v^{-1/2} \int_{r=a}^{T} \left[B(t-0) - B(t-1/r) \right] r^{1/2} dr$$

$$+ \frac{3}{2} v^{-1/2} \int_{r=T}^{v} \frac{B(t-0) - B(t-1/r)}{1/r} r^{-1/2} dr.$$

The first integral on the right $\rightarrow 0$ as $v \rightarrow \infty$, and

$$\frac{3}{2} v^{-1/2} \int_{T}^{v} \frac{B(t-0) - B(t-1/r)}{1/r} r^{-1/2} dr$$
$$= 3[B'(t-0) + \eta (T, v)] (v^{1/2} - T^{1/2}) v^{-1/2}$$

where $|\eta| < \epsilon$ for all values of v > T. Let $v \to \infty$, then let $\epsilon \to 0$; the right side of the above equation approaches 3B'(t-0). Consequently we now have

,

$$\lim_{v \to \infty} v^{-1/2} \ \alpha(v) = 2 B'(t-0).$$

By applying the above-mentioned theorem with $\gamma = 1/2$, $A = \pi^{1/2} B'(t-0)$, we now obtain

$$\frac{\lim_{z \to 0}}{\lim_{v \to \infty}} \left| z^{1/2} \int_0^\infty e^{-zv} d\alpha(v) - \pi^{1/2} B'(t-0) \right|$$

$$\leq \frac{\lim_{v \to \infty}}{\lim_{v \to \infty}} \left| \frac{1}{2} \pi^{1/2} v^{-1/2} B(v) - \pi^{1/2} B'(t-0) \right| = 0.$$

Hence

$$\lim_{x \to 0+} u(x, t) = \lim_{z \to 0} l = B'(t-0).$$

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