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**ASYMPTOTIC LOWER BOUNDS FOR THE FREQUENCIES OF  
CERTAIN POLYGONAL MEMBRANES**

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**1. Background.** Let the bounded, simply connected, open region  $R$  of the  $(x, y)$  plane have the boundary curve  $C$ . If a uniform elastic membrane of unit density is uniformly stretched upon  $C$  with unit tension across each unit length, the square  $\lambda = \lambda(R)$  of the fundamental frequency satisfies the conditions (subscripts denote differentiation)

$$(1a) \quad \begin{cases} \Delta u \equiv u_{xx} + u_{yy} = -\lambda u & \text{in } R, \\ \lambda = \text{minimum}, \end{cases}$$

with the boundary condition

$$(1b) \quad u(x, y) = 0 \quad \text{on } C.$$

The solution  $u$  of problem (1) is unique up to a constant factor. It is known [13, p. 24] that  $\lambda$  is the minimum over all piecewise smooth functions  $u$  satisfying (1b) of the Rayleigh quotient

$$(2) \quad \rho(u) = \iint_R |\nabla u|^2 dx dy / \iint_R u^2 dx dy,$$

where  $|\nabla u|^2 = u_x^2 + u_y^2$ . In many practical methods for approximating  $\lambda$  one essentially determines  $\rho(u)$  for functions  $u$  satisfying (1b) which are close to a solution of the boundary value problem (1). See [9, p. 112; 6, p. 276; 11, and 12]. By (2) these approximations are known to be *upper bounds* for  $\lambda$ ; they can be made arbitrarily good with sufficient labor. It is obviously of equal importance to obtain close lower bounds for  $\lambda$ ; cf. [14].

The lower bounds for  $\lambda$  given by Pólya and Szegő [13] are ordinarily far

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from close. Those obtainable from  $\rho(u)$ ,  $\iint_R u^2 dx dy$ , and  $\iint_R |\Delta u|^2 dx dy$  by methods due to Temple [15], D. H. Weinstein [17], Wielandt [18], and Kato [8] (for expositions see [3] and [16]) are arbitrarily good, but presuppose knowledge of a lower bound for the second eigenvalue  $\lambda_2$  of the problem (1). The same is true of Davis's proposals in [4]. It is possible, following Aronszajn and Zeichner [1], to get close lower bounds for  $\lambda$  by minimizing  $\rho(u)$  over a class of functions  $u$  permitted some discontinuity in  $R$  (method of A. Weinstein); the author has no knowledge of the practicability of the method.

A common method of approximating  $\lambda$  is to replace the boundary value problem (1) by a similar problem in finite differences. Divide the plane into squares of side  $h$  by the network of lines  $x = \mu h$ ,  $y = \nu h$  ( $\mu, \nu = 0, \pm 1, \pm 2, \dots$ ). The points  $(\mu h, \nu h)$  are the nodes of the net. A half-square is an isosceles right triangle whose vertices are three nodes of one square of the net. Assume that

$$(3) \quad R \text{ is the union of a finite number of squares and half-squares.}$$

Then every interior node of  $R$  has four neighboring nodes in  $R \cup C$ .

Define  $\Delta_h$ , a finite-difference approximation to  $\Delta$ , by the relation

$$h^2 \Delta_h v(x, y) = v(x+h, y) + v(x-h, y) + v(x, y+h) + v(x, y-h) - 4v(x, y).$$

Let  $\lambda_h$  be the least number satisfying the following difference equation for a net function  $v$  defined on the nodes  $(x, y)$  of the net:

$$(4a) \quad \Delta_h v = -\lambda_h v \text{ at the nodes in } R,$$

with the boundary condition

$$(4b) \quad v = 0 \text{ at the nodes on } C.$$

One can interpret  $\lambda_h$  as the square of the fundamental frequency of a network of massless strings with uniform tension  $h$ , fastened to  $C$ , and supporting a particle of mass  $h^2$  at each node. That is, a certain lumping of the distributed masses and tensions of problem (1) yields problem (4).

It is easily verified for a rectangular region of commensurable sides  $\pi/p$ ,  $\pi/q$ , and for  $h$  such that (3) holds, that one has  $u = v = \sin px \sin qy$ , and that

$$(5) \quad \frac{\lambda_h}{\lambda} = \frac{\sin^2(ph/2) + \sin^2(qh/2)}{(ph/2)^2 + (qh/2)^2} = 1 - \frac{p^4 + q^4}{p^2 + q^2} \frac{h^2}{12} + o(h^2) \quad (h \rightarrow 0).$$

Hence  $\lambda_h < \lambda$  for all  $h$ , and one can use  $\lambda_h$  as a lower bound for  $\lambda$ . However,

since  $\lambda$  is known exactly for rectangular regions, relation (5) contributes nothing to its computation. For general regions  $R$ , it was stated [3, p.405] in 1949 that nothing could be said about the relation of  $\lambda_h$  to  $\lambda$ .

**2. A new result.** An asymptotic relation resembling (5) will now be established for any *convex* polygonal region  $R$  satisfying (3). Such regions are polygons of at most eight sides, having interior vertex angles of  $45^\circ$ ,  $90^\circ$ , or  $135^\circ$ . The following theorem<sup>1</sup> will be proved in § 3 by use of the lemmas of § 4:

**THEOREM.** *Let  $R$  be a convex region which is a finite union of squares and half-squares for all  $h$  under consideration. Let  $u$  solve problem (1) for  $R$ , and let*

$$a = a(R) = \frac{\iint_R (u_{xx}^2 + u_{yy}^2) dx dy}{\iint_R (u_x^2 + u_y^2) dx dy}.$$

*Then, as  $h \rightarrow 0$ , one has*

$$(6) \quad \frac{\lambda_h}{\lambda} \leq 1 - \frac{a}{12} h^2 + o(h^2) \quad (h \rightarrow 0).$$

It is a consequence of the theorem that, for all sufficiently small  $h$ , say for  $h \leq h_0$ ,  $\lambda_h$  is a *lower bound* for  $\lambda$ . The ordinary finite-difference method thus complements any method based on Rayleigh quotients; and, since  $\lambda_h \rightarrow \lambda$  as  $h \rightarrow 0$ , together two such methods can confine  $\lambda$  to an arbitrarily short interval. In particular, Pólya [11 and 12] devises modified finite-difference approximations to problem (1) which furnish upper bounds to  $\lambda$  for all  $h$ . Hence arbitrarily good two-sided bounds to  $\lambda$  can be found by finite-difference methods alone.

The constant  $a$  of the theorem is the best possible for a rectangle  $R$  of sides  $\pi/p$ ,  $\pi/q$ . For this region, we have  $a = (p^4 + q^4) \cdot (p^2 + q^2)^{-1}$ , and (6) is seen by (5) to be actually an equality up to terms  $o(h^2)$ .

Using heuristic reasoning, Milne [9, p.238, (97.5)] finds an approximate formula which, specialized to the fundamental eigenvalue and set in our notation, says

$$(7) \quad \frac{\lambda_h}{\lambda} \doteq 1 - \frac{\lambda h^2}{24} + o(h^2) \quad (h \rightarrow 0).$$

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<sup>1</sup>The author gratefully acknowledges many helpful conversations with his colleague Dr. Wolfgang Wasow on the subject of this paper.

For a rectangle of sides  $\pi/p$ ,  $\pi/q$ , the coefficient of  $-h^2/12$  in (7) is  $(p^2+q^2)/2$ . Since

$$\frac{p^2+q^2}{2} + \frac{(p^2-q^2)^2}{p^2+q^2} = \frac{p^4+q^4}{p^2+q^2},$$

the coefficient of  $h^2$  in (7) is low for all rectangles with  $p \neq q$ , and exact for squares. Hence (7) cannot ordinarily be expected to be exact in its  $h^2$  term.

The use of the theorem to bound  $\lambda$  is limited by our lack of knowledge of  $h_0$ . However, it is the author's conjecture that, for the regions  $R$  of the theorem,  $\lambda_h < \lambda$  for all  $h$ .

The convexity of  $R$  is vital to the statement and proof of the theorem; in fact, by the remark after Lemma 4,  $a = \infty$  for nonconvex polygons. A heuristic argument, supported by the numerical example of § 5, has in fact convinced the author that, for nonconvex polygons,  $\lambda_h > \lambda$  for all sufficiently small  $h$ .

The restriction of  $R$  and  $h$  to satisfy (3) is less essential, but is used in two ways: (i) to be sure that no interior node has a neighboring node outside  $R$ ; (ii) to prove that  $\Gamma = 0$  in Lemma 7. With an appropriate alteration of  $\Delta_h$  near  $C$ , and with a modification of Lemma 7, one can extend the present method to obtain formulas of type (6) without assuming (3)—and even for convex regions  $R$  bounded by piecewise analytic curves  $C$ . See [5]. Analogous results can be expected in  $n$  dimensions.

**3. Proof of the theorem.** Let  $K$  be the class of functions  $u$  which vanish on  $C$ , such that  $(uu_x)_x$  and  $(uu_y)_y$  are continuous in  $R \cup C$ . Applying Gauss's divergence formula (27) with  $p = uu_x$ ,  $q = uu_y$ , one finds that, for all  $u$  in  $K$ , Green's formula is valid in the form

$$\iint_R |\nabla u|^2 dx dy = - \iint_R u \Delta u dx dy.$$

Hence, for all  $u \in K$ ,  $\rho(u)$  in (2) can be rewritten with  $-\iint_R u \Delta u dx dy$  in the numerator.

Since, by Lemma 1, the function  $u$  which minimizes (2) and solves (1) belongs to  $K$ , and since any function in  $K$  is piecewise smooth, one may alternatively define  $\lambda$  as the minimum, over all functions in  $K$ , of the quotient

$$\rho(u) = - \iint_R u \Delta u dx dy / \iint_R u^2 dx dy.$$

Analogously, without having to worry about function classes, one can show that  $\lambda_h$  is the minimum, over all net functions  $v$  satisfying (4b), of the quotient

$$(8) \quad \rho_h(v) = -h^2 \sum_{N_h} \sum v \Delta_h v / h^2 \sum_{N_h} \sum v^2,$$

where the sums are extended over all nodes  $N_h$  of the net inside  $R$ .

The key to proving the theorem is to set the solution  $u$  of problem (1) into the Rayleigh quotient (8) of problem (4). It will be shown that

$$(9) \quad \frac{\rho_h(u)}{\lambda} = 1 - \frac{1}{12} ah^2 + o(h^2) \quad (h \rightarrow 0).$$

Since  $\lambda_h \leq \rho_h(u)$ , the theorem follows from (9). Henceforth  $u$  will always denote a solution of problem (1).

The denominator of  $\rho_h(u)$  is a Riemann sum for  $\iint_R u^2 dx dy$ . Since  $u^2$  is continuous and hence Riemann integrable over  $R$ ,

$$(10) \quad h^2 \sum_{N_h} \sum u^2 = \iint_R u^2 dx dy + o(1) \quad (h \rightarrow 0).$$

(It can be shown that one can replace  $o(1)$  by  $o(h^2)$  in (10), but we shall not need to do this.)

The nodes  $N_h$  inside  $R$  are divided into two classes:

$N'_h$ : those at a distance  $h$  from some  $135^\circ$  vertex of  $C$ ;

$N''_h$ : the other nodes of  $N_h$ .

Split the numerator of  $\rho_h(u)$  accordingly:

$$(11) \quad -h^2 \sum_{N_h} \sum u \Delta_h u = -h^2 \sum_{N'_h} \sum u \Delta_h u - h^2 \sum_{N''_h} \sum u \Delta_h u = S'_h(u) + S''_h(u).$$

To estimate  $S'_h(u)$  note that, since there are at most eight  $135^\circ$  vertices, the number of nodes in  $N'_h$  is at most 8, for any  $h$ . At any node in  $N'_h$ ,

$$h^2 |u \Delta_h u| \leq h^2 \left( \frac{u-0}{h} \right) \sum_{i=1}^4 \left| \frac{u-u_i}{h} \right| \leq 4h^2 \max |\nabla u|^2,$$

where the maximum of  $|\nabla u|^2$  is taken for all points  $(x, y)$  within a distance  $2h$  of some  $135^\circ$  vertex. Hence, by Lemma 2, as  $h \rightarrow 0$  through values such that (3) holds,

$$(12) \quad |S'_h(u)| \leq 32h^2 \max |\nabla u|^2 = o(h^2) \quad (h \rightarrow 0).$$

Now, using the notation and assertion of Lemma 5, one obtains

$$(13) \quad S''_h(u) = -h^2 \sum_{N''_h} \sum u \Delta u - \frac{h^4}{12} \sum_{N''_h} \sum u (u'_{xxxx} + u''_{yyyy}).$$

Since  $u$  satisfies (1a),

$$(14) \quad -h^2 \sum_{N''_h} \sum u \Delta u = \lambda h^2 \sum_{N''_h} \sum u^2 = \lambda h^2 \sum_{N_h} \sum u^2 + o(h^2) \quad (h \rightarrow 0);$$

the last step is correct because  $u(x, y) \rightarrow 0$  as  $(x, y) \rightarrow C$ .

Combining (13) and (14), one finds that, as  $h \rightarrow 0$ ,

$$(15) \quad \begin{aligned} S''_h(u) &= \lambda h^2 \sum_{N_h} \sum u^2 - \frac{h^4}{12} \sum_{N''_h} \sum u (u'_{xxxx} + u''_{yyyy}) + o(h^2) \\ &= \lambda h^2 \sum_{N_h} \sum u^2 - \frac{h^2}{12} \iint_R u (u_{xxxx} + u_{yyyy}) dx dy + o(h^2), \end{aligned}$$

by Lemma 6. The integrals used in this proof exist, by Lemma 3. Using (11), (12), (15), and Lemma 7, one finds that

$$(16) \quad \begin{aligned} -h^2 \sum_{N_h} \sum u \Delta_h u \\ = \lambda h^2 \sum_{N_h} \sum u^2 - \frac{h^2}{12} \iint_R (u_{xx}^2 + u_{yy}^2) dx dy + o(h^2) \quad (h \rightarrow 0). \end{aligned}$$

Dividing (16) by the denominator of  $\rho_h(u)$ , one gets

$$\rho_h(u) = \lambda - \frac{h^2}{12} \frac{\iint_R (u_{xx}^2 + u_{yy}^2) dx dy}{h^2 \sum_{N_h} u^2} + o(h^2).$$

Hence, by (10),

$$(17) \quad \rho_h(u) = \lambda - \frac{h^2}{12} \frac{\iint_R (u_{xx}^2 + u_{yy}^2) dx dy}{\iint_R u^2 dx dy} + o(h^2) \quad (h \rightarrow 0).$$

If one divides (17) by  $\lambda$ , and notes from (2) that  $\lambda \iint_R u^2 dx dy = \iint_R |\nabla u|^2 dx dy$ , it is seen that

$$\frac{\rho_h(u)}{\lambda} = 1 - \frac{h^2}{12} \frac{\iint_R (u_{xx}^2 + u_{yy}^2) dx dy}{\iint_R |\nabla u|^2 dx dy} + o(h^2) \quad (h \rightarrow 0).$$

By the definition of  $a$  we have proved (9) and hence the theorem.

**4. Some lemmas.** Lemma 1, suggested to the author by Professor Max Shiffman, is used to establish Lemmas 2 to 7, which were applied to prove the theorem. In all the lemmas  $R$  is the convex union of squares and half-squares of the network, while  $u = u(x, y)$  is a function solving problem (1) in  $R$ .

LEMMA 1. *The function  $u$  is an analytic function of  $x$  and  $y$  in  $R \cup C$ , except at the  $135^\circ$  vertices of  $C$ . Let  $r, \theta$  be local polar coordinates centered at a  $135^\circ$  vertex  $P_k$ , with  $0 < \theta < 3\pi/4$  in  $R$ . Then*

$$(18) \quad u = \gamma_k r^{4/3} \sin(4\theta/3) + r^{7/3} E_k(r, \theta),$$

where  $\gamma_k$  is a constant, and where  $E_k(r, \theta)$ , together with all its derivatives, is bounded in a neighborhood of  $P_k$ .

*Proof.* By reflection one can continue  $u$  antisymmetrically across each straight segment of  $C$ , and (1a) is satisfied by the extended  $u$  at all points of  $R \cup C$  except the  $135^\circ$  vertices. The first sentence of the lemma then follows from [2, p. 179].

For  $(\xi, \eta) \in R$ , write  $t = \xi + i\eta$ . For each  $t$ , let  $w = f(z, t)$  be an analytic function of the complex variable  $z = x + iy$  which maps  $R$  into the unit circle  $|w| < 1$ , with  $f(t, t) = 0$ . To study  $f$  near a vertex  $z_k$  of  $C$ , one may assume



that  $f(z_k, t) = 1$ . Let the interior vertex angle of  $C$  at  $z_k$  be  $\pi/\alpha_k$  ( $\alpha_k = 4, 2,$  or  $4/3$ ). It is a property of the Schwarz-Christoffel transformation [10, p. 189] that

$$(19) \quad f(z, t) = 1 + (z - z_k)^{\alpha_k} g_k(z, t),$$

where  $g_k$  is an analytic function of  $z$  regular at  $z_k$ .

Let  $G(z, t) = G(x, y; \xi, \eta)$  be Green's function for  $\Delta u$  in  $R$ . Now  $G(z, t) = -(2\pi)^{-1} \log |f(z, t)|$ ; see [10, p. 181]. It then follows from (19) that, in the notation of the lemma, when  $\alpha_k = 4/3$ ,

$$(20) \quad G(z, t) = \gamma_k(t)r^{4/3} \sin(4\theta/3) + r^{7/3} E_k(r, \theta, t).$$

Moreover,  $\gamma_k(t)$  and  $E_k(r, \theta, t)$  are integrable over  $R$ , since the only discontinuity of  $G(z, t)$  is a logarithmic one at  $t = z$ .

The function  $u$  is representable by the integral [2, pp. 182-3]

$$(21) \quad u(x, y) = \lambda \iint_R G(x, y; \xi, \eta) u(\xi, \eta) d\xi d\eta.$$

Substituting (20) into (21) proves (18) and the lemma.

LEMMA 2.  $|\nabla u(x, y)| \rightarrow 0$  as  $(x, y) \rightarrow$  any  $135^\circ$  vertex of  $C$ .

*Proof.* By (18),  $|\nabla u| = O(r^{1/3})$ , as  $(x, y) \rightarrow$  any  $135^\circ$  vertex of  $C$ .

LEMMA 3. The functions  $u_{xx}^2, u_x u_{xxx}, uu_{xxxx}, u_{yy}^2, u_y u_{yyy}$ , and  $uu_{yyyy}$  are Lebesgue-integrable in  $R$ .

*Proof.* By Lemma 1 these functions are continuous in  $R \cup C$ , except at the  $135^\circ$  vertices  $P_k$ . At these vertices (18) implies that they are  $O(r^{-4/3})$  and are hence integrable.

LEMMA 4. The Lebesgue integrals  $\int_C u_y u_{yy} dx$  and  $\int_C u_x u_{xx} dy$  exist.

*Proof.* Analogous to that of Lemma 3.

REMARK. Lemmas 2, 3, and 4 are false for polygonal regions  $R$  which are not convex, since in general the exponent in (18) is  $\alpha_k$ , where  $\pi/\alpha_k$  is the interior angle at the vertex  $P_k$ .

LEMMA 5. At each node  $(x, y)$  in  $R$  of the network of section 1, one has

$$(22) \quad \Delta_h u = \Delta u + \frac{1}{12} h^2 (u'_{xxxx} + u''_{yyyy}),$$

where

$$(23) \quad \begin{cases} u'_{xxxx} = u_{xxxx}(x + \theta' h, y), & -1 < \theta' < 1; \\ u''_{yyyy} = u_{yyyy}(x, y + \theta'' h), & -1 < \theta'' < 1. \end{cases}$$

*Proof.* By Lemma 1,  $u_{xxxx}$  is continuous in the open line segment from  $(x - h, y)$  to  $(x + h, y)$  (though infinite at any 135° vertex). Since  $u$  is continuous in  $R \cup C$ , it follows from Taylor's formula [7, p. 357] that, if we fix  $y$  and set  $\phi(x) = u(x, y)$ ,

$$\begin{aligned} &\phi(x + h) + \phi(x - h) - 2\phi(x) \\ &= h^2 \phi''(x) + \frac{1}{24} h^4 [\phi''''(x + \theta_1 h) + \phi''''(x - \theta_2 h)], \end{aligned}$$

where  $0 < \theta_i < 1$  ( $i = 1, 2$ ). By the continuity of  $\phi''''$ , the last bracket equals  $2\phi''''(x + \theta' h)$ , where  $-1 < \theta' < 1$ .

A similar formula for  $\psi(y) = u(x, y)$ , when added to the above and divided by  $h^2$ , yields (22) and (23).

LEMMA 6. Define  $N_h''$  as in §3. For each node  $(x, y)$  in  $N_h''$ , use the notation of (23). Then, as  $h \rightarrow 0$  over values such that (3) holds, one has

$$(24) \quad h^2 \sum_{N_h''} (u'_{xxxx} + u''_{yyyy}) = \iint_R u (u_{xxxx} + u_{yyyy}) dx dy + o(1) \quad (h \rightarrow 0).$$

*Proof.* For all  $(x, y)$  in the entire plane  $E_2$  define

$$f(x, y) = \begin{cases} u(u_{xxxx} + u_{yyyy}), & \text{if } (x, y) \in R; \\ 0, & \text{elsewhere.} \end{cases}$$

By the proof of Lemma 3 one sees that  $f(x, y)$  is  $O(r^{-4/3})$  in the neighborhood of each 135° vertex  $P_k$  of  $C$ , and continuous elsewhere. Divide the nodes  $(x, y) = (\mu h, \nu h)$  of  $N_h'' \subset R$  into four classes  $K^{(i)}$  ( $i = 1, 2, 3, 4$ ) according to the parity of  $(\mu, \nu)$ . Fix any class  $K^{(i)}$ . For each vertex  $(x, y)$  in  $K^{(i)}$  let  $S(x, y)$  be the union of the four closed network squares of  $E_2$  which contain  $(x, y)$ . The area

of each  $S(x, y)$  is  $4h^2$ ; ordinarily certain of the  $S(x, y)$  contain points not in  $R$ . Define

$$f_h^{(i)}(\xi, \eta) = \begin{cases} u(x, y) (u'_{xxxx} + u''_{yyyy}), & \text{for } (\xi, \eta) \in S(x, y); \\ 0, & \text{for } (\xi, \eta) \notin \cup S(x, y). \end{cases}$$

Then  $f_h^{(i)}(\xi, \eta) \rightarrow f(\xi, \eta)$ , as  $h \rightarrow 0$ , for almost all  $(\xi, \eta)$  in the plane. Using the fact that no node of  $N_h''$  is adjacent to a  $135^\circ$  vertex of  $C$ , one can show that for all  $i$ , uniformly in  $h$ ,  $|f_h^{(i)}(\xi, \eta)| \leq F(\xi, \eta)$ , where  $F$  is an integrable function in  $E_2$ .

Each term of the sum (24) for which  $(x, y) \in K^{(i)}$  is equal to

$$\frac{1}{4} \iint_{S(x, y)} f_h^{(i)}(\xi, \eta) d\xi d\eta.$$

Hence, applying Lebesgue's convergence theorem, one sees that, as  $h \rightarrow 0$ , for each  $i$ ,

$$\begin{aligned} \sum_{N_h'' \cap K^{(i)}} u(u'_{xxxx} + u''_{yyyy}) &= \frac{1}{4} \iint_{E_2} f_h^{(i)}(\xi, \eta) d\xi d\eta \\ (25) \qquad \qquad \qquad &\rightarrow \frac{1}{4} \iint_{E_2} f(\xi, \eta) d\xi d\eta \qquad (h \rightarrow 0). \end{aligned}$$

Summing (25) over  $i = 1, 2, 3, 4$  proves (24) and the lemma.

LEMMA 7. *One has*

$$(26) \qquad \iint_R u(u_{xxxx} + u_{yyyy}) dx dy = \iint_R (u_{xx}^2 + u_{yy}^2) dx dy.$$

*Proof.* The following applications of Gauss's divergence theorem in the form

$$(27) \qquad \iint_R (p_x + q_y) dx dy = \int_C (p dy - q dx)$$

can be justified by integrating over the region  $R^*$  interior to a smooth convex curve  $C^*$  inside  $R$ , and then letting  $C^* \rightarrow C$  appropriately. The continuity of

the integrals in the limit follows from Lemmas 1, 3, and 4.

In the divergence theorem for  $p = uu_{xxx}$ ,  $q = uu_{yyy}$ , the line integral vanishes, and one finds

$$(28) \quad \iint_R u(u_{xxx} + u_{yyy}) dx dy = - \iint_R (u_x u_{xx} + u_y u_{yy}) dx dy.$$

A second application of the divergence theorem with  $p = u_x u_{xx}$ ,  $q = u_y u_{yy}$ , combined with (28), shows that

$$(29) \quad \iint_R u(u_{xxx} + u_{yyy}) dx dy = \iint_R (u_{xx}^2 + u_{yy}^2) dx dy + \Gamma,$$

where  $\Gamma = \int_C (u_y u_{yy} dx - u_x u_{xx} dy)$ .

By (1a),  $u_{xx} = -u_{yy}$  on  $C$ , whence  $\Gamma = \int_C u_{yy} (u_y dx + u_x dy)$ . On the segments of  $C$  parallel to the axes,  $u_{xx} = u_{yy} = 0$ , so that there the contribution to  $\Gamma$  is zero.

Now the vector  $\nabla u = (u_x, u_y)$  is perpendicular to  $C$ . On the segments of  $C$  making a  $45^\circ$  or  $135^\circ$  angle with the  $x$ -axis,  $(u_y, u_x)$  is parallel to  $(u_x, u_y)$ , whence  $(u_y, u_x)$  is perpendicular to  $C$ . Thus  $u_y dx + u_x dy \equiv 0$  when  $(dx, dy)$  is tangent to  $C$ , so that the contribution to  $\Gamma$  from these  $45^\circ$  and  $135^\circ$  segments of  $C$  is also zero.

Hence  $\Gamma = 0$ , and the lemma follows from (29).

**5. Numerical example.** Let  $R_1$  be the six-sided, nonconvex,  $L$ -shaped region whose closure is the union of the three unit squares

$$\begin{cases} -1 \leq x \leq 0, & 0 \leq y \leq 1; \\ 0 \leq x \leq 1, & 0 \leq y \leq 1; \\ 0 \leq x \leq 1, & -1 \leq y \leq 0. \end{cases}$$

The fundamental frequencies  $\lambda_h = \lambda_h(R_1)$  and corresponding net functions  $v$  were computed by B. F. Handy on the SWAC (National Bureau of Standards Western Automatic Computer) for  $1/h = 3, 4, \dots, 8$ . The computation used a *power method*; for some initial net function  $v_0$ ,  $(h^2 \Delta_h + 5I)^m v_0$  was determined for large positive integers  $m$ , where  $I$  is the identity operator. On the basis of Collatz's inclusion theorem [3, p. 289], the values in the accompanying table are believed to have errors less than  $5 \times 10^{-6}$ . Observe that  $\lambda_h(R_1)$  is less for  $h = 1/8$  than for  $h = 1/7$ .

TABLE

$h$	$\lambda_h(R_1)$	$\lambda_h(R_2)$
1/2	9.07180	12.00000
1/3	9.52514	13.73700
1/4	9.64143	14.37340
1/5	9.67860	14.67081
1/6	9.69083	14.83259
1/7	9.69384	14.93003
1/8	9.69316	14.99315

Since  $R_1$  is not convex, the theorem of § 2 does not apply, but a heuristic argument suggests that  $\lambda_h(R_1) - \lambda(R_1) = O(h^{4/3})$ . A least-squares fit to the values of  $\lambda_h(R_1)$  for  $1/8 \leq h \leq 1/4$  of a function of type

$$\lambda_h(R_1) \doteq \alpha_1 + \beta_1 h^{4/3} + \gamma_1 h^2 = \phi_1(h)$$

yielded the values

$$(30) \quad \alpha_1 = 9.63632, \quad \beta_1 = 2.40286, \quad \gamma_1 = -5.97212.$$

The maximum of  $|\lambda_h(R_1) - \phi_1(h)|$  for the five values of  $h$  is .00013. Hence  $\alpha_1$  is a working estimate of  $\lambda(R_1)$ .

The fact that  $\beta_1 > 0$  in (30) supports the author's conjecture that, for nonconvex polygonal domains satisfying (3),  $\lambda_h > \lambda$  for all sufficiently small  $h$ .

The table also gives Handy's values for the second eigenvalues of  $R_1$ , which are the fundamental eigenvalues  $\lambda_h(R_2)$  of the trapezoidal halfdomain  $R_2$  of  $R_1$  for which  $x > y$ . Since the theorem does apply to  $R_2$ , a least-squares fit to the values of  $\lambda_h(R_2)$  for  $1/8 \leq h \leq 1/4$  of a function of type

$$\lambda_h(R_2) \doteq \alpha_2 + \beta_2 h^2 = \phi_2(h)$$

seemed appropriate, and yielded the values

$$\alpha_2 = 15.19980, \quad \beta_2 = -13.22219.$$

The maximum of  $|\lambda_h(R_2) - \phi_2(h)|$  for the five values of  $h$  was .00010. Hence  $\alpha_2$  is a working estimate of  $\lambda(R_2)$ .

The value of  $\beta_2$  is negative, in agreement with (6), but the quantity

$-12\beta_2/\alpha_2 = 10.4387$  is something like one-fifth larger than an estimate of the corresponding quantity  $a(R_2)$  of the theorem. One therefore suspects that  $a$  is not the best possible constant in (6) for the region  $R_2$ .

In the table, note the relative closeness of the values of  $\lambda_h(R_2)$  to the working estimate,  $\alpha_2$ , of  $\lambda(R_2)$ , even for a coarse net. Thus the value 12 for  $\lambda_{1/2}(R_2)$ , which is obtained by pencil and paper from a simple quadratic equation, is comparable to the lower bounds 12.1 and  $5\pi^2/4$  obtained respectively by comparison with  $\lambda$  for the circular membrane of equal area [13, p. 8] and with  $\lambda$  for the rectangular region  $0 < x < 1$ ;  $-1 < y < 1$ . The value  $\lambda_{1/3}(R_2) = 13.737$  requires getting the least eigenvalue of a 7th-order matrix, a relatively easy procedure with a desk machine.

The monotonicity of  $\lambda_h(R_2)$  supports the author's conjecture<sup>2</sup> that, for the  $R$  of the theorem,  $\lambda_h < \lambda$  for all  $h$ .

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<sup>2</sup>See page 470.

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