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ASYMPTOTICS III: STATIONARY PHASE FOR TWO PARAMETERS WITH AN APPLICATION TO BESSEL FUNCTIONS

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1. Introduction. The method of stationary phase has long been a valuable analytical tool for investigating the asymptotic behavior as $p \rightarrow \infty$ of integrals of the form

$$I(p) = \int_{0}^{a} Q(t) \exp{(ipF(t))} dt$$
.

As a natural generalization of the method of stationary phase involving one parameter we will investigate the asymptotic behavior of an integral of the form

$$I(h, k) = \int_0^a t^{\gamma-1}q(t) \exp\left[i(ht^{\lambda}f(t) + kt^{\gamma}g(t))\right]dt$$

where h and k tend to infinity independently.

It will be shown that under certain restrictions between the real numbers λ , ν and γ that the asymptotic form of I(h, k) is determined by the behavior of the ratio $kh^{-\nu/\lambda}$ as $h, k \to \infty$ and by the character of f and g in a neighborhood of t = 0. For example, if $\gamma < \nu < \lambda$, $\gamma > 0, f(0) > 0, g(0) > 0$ and $kh^{-\nu/\lambda} \to \infty$ then

$$I(h, k) \sim \frac{q(0)\Gamma\left(\frac{\gamma}{\nu}\right)\exp\left(\frac{i\pi\gamma}{2\nu}\right)}{\nu [kg(0)]^{\gamma/\nu}}$$

As an immediate application of our results we will determine the asymptotic behavior of the Bessel function $J_{\nu}(x)$ in Watson's transition region, i.e. when ν , x and $|\nu - x|$ are large and ν/x is nearly equal to 1. In particular, we will obtain a simple rigorous proof of Nicholson's formulas under the restriction that $0 < \limsup x^{-1/3} |\nu - x| < \infty$.

2. General assumptions. Throughout the paper we shall use $A \sim B$ to mean $\lim A/B = 1$, and all limits will mean the limit as h and k tend to infinity. A similar remark applies to order symbols.

We shall consider I(h, k) under the following general assumptions:

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(i) k = o(h),

(ii) $\lambda > 0$, $\nu > 0$, and $\gamma > 0$,

(iii) $q(0) \neq 0$ and $g(0) \neq 0$,

- (iv) f, g and q are real valued functions such that $f \in C^2$, $g \in C^2$ and $q \in C$ on [0, a],
- (v) $\lambda f(t) + tf'(t) > 0$ on [0, a].

For convenience we shall consider here only the case f(0) > 0. If f(0) < 0 and -f satisfies certain obvious conditions one obtains analogous results with -i and -g replacing i and g, respectively.

3. Preliminary lemmas. We shall first establish the following lemmas.

LEMMA I. Consider $I(p) = \int_{0}^{b} \omega(t) \Psi(t) \exp(ip \Phi(t)) dt$. Suppose d is a nonnegative constant and p, α and μ are functions of h and k such that $p \to \infty$, $\mu \to 0$ and α is bounded as $h, k \to \infty$.

(i) $\Phi(t) = t^r \phi(\alpha t)$, $\Psi(t) = t^{s-1} \psi(\mu(t+d))$, the functions ϕ and ψ are real with $\psi(0) \neq 0$, r > 0, 0 < s < r, $\phi(\alpha t) > 0$ for $0 \leq t \leq c'$, c' > 0 and $\phi \in C^{n+3}$ and $\psi \in C^n$ for $0 \leq t \leq c'$ where m and n are the least integers such that mr > 1 and $n \geq m(r-s) + 1$, respectively.

(ii) b is a constant such that 0 < b < c' and $bMK/m_0r < 1$ where $M = \max_{0 \le t \le c'} |\phi'(t)|$, $m_0 = \min_{0 \le t \le c'} \phi(t)$ and $K \ge \alpha$ when h, k are sufficiently large.

(iii) $\omega = u + iv$ is a complex valued function such that u(0) = 1, v(0) = 0 and $u, v \in C^n$ for $0 \leq t \leq c'$. Then

$$I(p) \backsim rac{\psi(0) \varGamma \left(rac{s}{r}
ight) \exp \left(rac{i \pi s}{2r}
ight)}{r [p \phi(0)]^{s/r}}$$

Proof. We may set $v \equiv 0$ since it will be seen that the contribution from v to I(p) is negligible because v(0) = 0. Let $x = t[\phi(\alpha t)]^{1/r}$. Since x'(t) > 0 for $0 \leq t \leq b$ and $x \in C^{n+2}$ there exists a unique inverse function, say t(x), such that $t \in C^{n+2}$ for $0 \leq x \leq b[\phi(\alpha b)]^{1/r} = a$, t(0) = 0and $t'(0) = a_1 = [\phi(0)]^{-1/r}$. Hence we may write $t(x) = a_1x + a_2x^2 + \cdots + a_{n-1}x^{n-1} + A(x)x^n$ where $A \in C^2$ and a_i is bounded as $h, k \to \infty$ for $2 \leq l \leq n-1$. We may assume that c' is sufficiently small such that if $t(x) = a_1x(1 + w(x))$ then |w(x)| < 1 for $0 \leq x \leq a$. This implies that

$$(t(x))^{s-1} = a_1^{s-1}(1 + b_1x + \cdots + b_{n-2}x^{n-2} + z(x)x^{n-1})x^{s-1}$$

where $z \in C$ and b_i is independent of x for $1 \leq l \leq n-2$. If we now expand ψ and ω about t = 0 and substitute t(x), and let $B(x) = \omega(t(x))$ $(t(x))^{s-1}\psi(\mu(t(x) + d))$ we have

$$B(x) \ rac{dt(x)}{dx} = a_1^s u(\mu d) x^{s-1} + c_0 x^s + \cdots + c_{n-3} x^{s+n-3} + D(x) x^{n+s-2}$$

where h and k are sufficiently large such that $\mu d < b$, $D \in C$ and c_l is bounded as $h, k \to \infty$ and independent of x for $0 \leq l \leq n-3$. Therefore,

$$I(p) = a_1^s \psi(\mu d) \int_0^a x^{s-1} \exp{(ipx^r)} dx + J(p) + \sum_{l=0}^{n-3} c_l \int_0^a x^{s+l} \exp{(ipx^r)} dx$$

where $J(p) = \int_{0}^{a} D(x) x^{n+s-2} \exp(ipx^{r}) dx$. Since $\int_{0}^{\infty} e^{it} t^{\beta-1} dt = \exp(i\pi\beta/2) \Gamma(\beta)$ for $0 < \beta < 1$ and $r \le n+s-1 < r+1$ when r > 1 we have

$$egin{aligned} I(p) =& rac{a_1^s\psi(\mu d)\exp\left(rac{i\pi s}{2r}
ight)\Gamma\left(rac{s}{r}
ight)}{rp^{s/r}}+J(p)+o(p^{-s/r})\ &= rac{a_1^s\psi(0)\exp\left(rac{i\pi s}{2r}
ight)\Gamma\left(rac{s}{r}
ight)}{rp^{s/r}}+J(p)+o(p^{-s/r})\ . \end{aligned}$$

Finally an integration by parts yields J(p) = 0(1/p) since $n - (r+1-s) \ge 0$ by the choice of n and $D \in C$. This completes the proof of Lemma I for the case r > 1. For $0 < r \le 1$ one makes the change of variable $t = x^m$ and the desired result follows from the case r > 1.

LEMMA II. Suppose that in addition to the assumptions of Lemma I that r is an even integer, s = 1, $\phi(\alpha t) > 0$ for $-c' \leq t \leq c'$, b satisfies the same conditions as in Lemma I except that M and m_0 are now determined for $-c' \leq t \leq c'$, and ω, ψ and ϕ are now in their respective differentiability classes given in Lemma I for $-c' \leq t \leq c'$. Then

$$\int_{-b}^{b} \omega(t) \varPsi(t) \exp{(ip arPsi(t))} dt \sim rac{2 \psi(0) \Gamma\left(rac{1}{r}
ight) \exp\left(rac{i\pi}{2r}
ight)}{r[p \phi(0)]^{1/r}} \; .$$

The proof follows immediately from Lemma I.

We will introduce the following functions which will be used throughout the remainder of the paper:

$$F(t) = t^{\lambda} f(t), \ G(t) = t^{\nu} g(t) \text{ and } Q(t) = t^{\gamma-1} q(t).$$

LEMMA III. Under the general assumptions on F, G and Q we have for each arbitrarily small but fixed positive constant c < a that

$$L(h, k) = \int_{a}^{a} Q(t) \exp [i(hF(t) + kG(t))]dt = 0(1/h)$$

Proof. Let H(t) = F(t) + (k/h)G(t). Then H'(t) > 0 for $c \le t \le a$ and h, k sufficiently large since $\lambda f(t) + tf'(t) > 0$ by hypothesis and k = o(h). Hence an integration by parts implies L(h, k) = O(1/h).

This completes the necessary lemmas and the main results of the paper will now be presented.

4. The asymptotic evaluation of I(h, k). We shall first consider the case where $kh^{-\nu/\lambda} \rightarrow 0$ so that I(h, k) is almost completely determined by the character of hf at the origin.

THEOREM I. Suppose that

1. $f \in C^{n+2}$ and $q \in C^n$ for $0 \leq t \leq c$, c > 0, where m and n are the least integers such that $m\lambda > 1$ and $n \geq m(\lambda - \gamma) + 1$, respectively, 2. if $0 \leq t \leq c$ then $t^{\nu}g(\beta t) = b_0 + b_1t + \cdots + b_{n-2}t^{n-2} + B(t)t^{n-1}$ where

 $B \in C \text{ and } b_l \text{ is bounded as } \beta \rightarrow 0 \text{ for } 0 \leq l \leq n-2,$ 3. $k^{\lambda} = o(h^{\nu}) \text{ and } \gamma < \lambda.$ Then

$$I(h, k) \sim rac{q(0) \Gamma\left(rac{\gamma}{\lambda}
ight) \exp\left(rac{i\pi\gamma}{2\lambda}
ight)}{\lambda [hf(0)]^{\gamma/\lambda}} .$$

Proof of Theorem I. For c as given we have

$$I(h,k) = \int_{0}^{c} + \int_{c}^{a} = I'(h, k) + 0(1/h)$$

by Lemma III. Let $t = xk^{-1/\nu}$, $\tilde{f}(x) = f(xk^{-1/\nu})$, $\tilde{g}(x) = g(xk^{-1/\nu})$, $\tilde{Q}(x) = Q(xk^{-1/\nu})$ and $p = hk^{-\lambda/\nu}$. For any b such that 0 < b < c we have

$$egin{aligned} I'(h,\,k) &= k^{-1/
u} \! \int_{0}^{b} [\widetilde{Q}(x) \exp{(i\widetilde{g}(x)x^{
u})}] \exp{(ip\widetilde{f}(x)x^{\lambda})} dx \ &+ k^{-1/
u} \! \int_{b}^{ck^{1/
u}} &= I''(h,\,k) + J(h,\,k), ext{ respectively.} \end{aligned}$$

Set $\mu = k^{-1/\nu}$, $\psi = q$, $\phi = f$, $\lambda = r$, $\gamma = s$, $\omega(x) = \exp(ix^{\nu}\tilde{g}(x))$ and note that f(0) > 0 implies that $\tilde{f}(x) > 0$ for $0 \le x \le c'$, c' > 0, so that b may be chosen to satisfy the requirements of Lemma I. Hence by Lemma I

$$I^{\prime\prime}(h,\,k) \backsim rac{q(0) \varGamma(\gamma/\lambda) \exp{(i\pi \gamma/2\lambda)}}{\lambda [hf(0)]^{\gamma/\lambda}} \,.$$

Therefore to complete the proof of Theorem I it is sufficient to show that $h^{\gamma/\lambda}J(h, k) = o(1)$. Let $d = bk^{-1/\nu}$, H(t) = F(t) = + (k/h)G(t)and $P(t) = \lambda f(t) + tf'(t) + kt^{\nu-\lambda}/h[\nu g(t) + g'(t)t]$. Note that $P(d) \to \lambda f(0) =$ 2B > 0 as $h, k \to \infty$ since $k^{\lambda} = 0(h^{\nu})$ and P(t) is continuous for $0 < d \le t \le a$. We may assume that c is such that for h, k sufficiently large, $P(t) \ge B$ for the entire closed interval $d \le t \le c$. This implies $H'(t) \ge Bt^{\lambda-1} > 0$ for $0 < d \le t \le c$ and hence we can integrate J(h, k) by parts as follows:

$$\begin{split} J(h,k) &= \int_{a}^{c} Q(t) \exp{(ihH(t))} dt = \frac{Q(c) \exp{(ihH(c))}}{ihH'(c)} - \frac{Q(d) \exp{(ihH(d))}}{ihH'(d)} \\ &- \frac{1}{ih} \int_{a}^{c} \frac{Q'(t) \exp{(ihH(t))} dt}{H'(t)} + \frac{1}{ih} \int_{a}^{c} \frac{Q(t)H''(t) \exp{(ihH(t))} dt}{(H'(t))^{2}} \\ &= 0(1/h) + A + J'(h, k), \text{ respectively.} \end{split}$$

Using the estimates $H'(t) \ge Bt^{\lambda-1}$ and $|H''(t)| \le Kt^{\lambda-2}$ for some K we see immediately that $J = O(k^{(\lambda-\gamma)\nu}/h)$. Since $k^{\lambda} = o(h^{\nu})$ this implies $h^{\gamma/\lambda}J(h, k) = o(1)$ which completes the proof of Theorem I.

We state the following corollary to Theorem I which may apply when $t^{\nu}g(\beta t)$ does not have the required smoothness at the origin but f, g and q are highly differentiable on [0, c], c > 0.

Corollary. Suppose that $\nu + \gamma > \lambda$ and

- 1. $f \in C^{n+2}$, $g \in C^n$ and $q \in C^n$ for $0 \leq t \leq c$, c > 0 where *m* and *n* are the least integers such that $m(\nu + \gamma \lambda) \geq 2$, $m\lambda > 1$ and $n \geq m(\lambda \gamma) + 1$,
- 2. $k^{\lambda} = o(h^{\nu})$ and $\gamma < \lambda$. Then

$$I(h, k) \sim rac{q(0) \Gamma \Big(rac{\gamma}{\lambda} \Big) \exp \Big(rac{i \pi \gamma}{2 \lambda} \Big)}{\lambda [h f(0)]^{\gamma/\lambda}} .$$

Proof. Note that $m\nu \ge m(\lambda - \gamma) + 2 > n$ by the definition of n and hence $x^{m\nu} \in C^n$. The change of variable $t = x^m$ and the use of Theorem I completes the proof.

We shall next consider the case where the behavior of kg at the origin becomes a significant factor in the asymptotic evaluation of I(h, k).

THEOREM II. Suppose that

- 1. $q \in C^n$ and $g \in C^{n+2}$ for $0 \leq t \leq c$, c > 0, where m and n are the least integers such that $m\nu > 1$ and $n \geq m(\nu \gamma) + 1$, respectively,
- 2. if $0 \leq t \leq c$ then $t^{\lambda}f(\beta t) = b_0 + b_1t + \cdots + b_{n-2}t^{n-2} + B(t)t^{n-1}$ where $B \in C$ and b_1 is bounded as $\beta \to 0$ for $0 \leq l \leq n-2$,
- 3. g(0) > 0, $h^{\nu} = o(k^{\lambda})$ and $\gamma < \nu < \lambda$. Then

$$I(h, k) \backsim rac{q(0) arGamma \left(rac{\gamma}{
u}
ight) \exp \left(rac{\imath \pi \gamma}{2
u}
ight)}{
u [kg(0)]^{\gamma/
u}} \; .$$

Proof of Theorem II. The proof of Theorem II follows from the proof of Theorem I with the roles of f and g, λ and ν , h and k intercharged.

COROLLARY. Suppose that

1. $f \in C^n$, $q \in C^n$ and $g \in C^{n+2}$ for $0 \le t \le c$, c > 0, where m and n are the least integers such that $m(\lambda + \gamma - \nu) \ge 2$, $m\nu > 1$ and $n \ge m(\nu - \gamma) + 1$, 2. g(0) > 0, $h^{\nu} = o(k^{\lambda})$ and $\gamma < \nu < \lambda$. Then

$$I(h, k) \sim rac{q(0) \Gamma\left(rac{\gamma}{
u}
ight) \exp\left(rac{i\pi\gamma}{2
u}
ight)}{
u [kg(0)]^{\gamma/
u}} \; .$$

When $kh^{-\nu/\lambda} \to \infty$ and g(0) < 0 the character of both F and G in a neighborhood of t = 0 becomes important since for h and k sufficiently large they determine uniquely in some $(0, c_0)$ a number τ such that $hF'(\tau) + kG'(\tau) = 0$ and in terms of which the asymptotic form of I(h, k) may be expressed.

THEOREM III. Suppose that g(0) < 0, $\nu < \lambda$, $\gamma < \lambda$, $h^{\nu} = o(k^{\lambda})$, $f \in C^{6}$, $g \in C^{6}$ and $q \in C^{2}$ for $0 \leq t \leq c$, c > 0, and hypothesis 1 and 2 Theorem II are satisfied when $\nu \geq \gamma$. A. If $\nu < 2\gamma$ then

$$I(h, k) \sim \frac{\sqrt{2} q(0) \Gamma\left(\frac{1}{2}\right) \exp\left(\frac{i\pi}{4}\right)}{(\lambda - \nu)^{1/2}} \left[\frac{(\lambda h f(0))^{\nu - 2\gamma}}{(-\nu k g(0))^{\lambda - 2\gamma}}\right]^{1/2(\lambda - \nu)}$$

$$imes \, \exp \left[i (h F(au) + \, k G\left(au)
ight)
ight] \, .$$

B. If $\nu = 2\gamma$ then

$$I(h,k) \backsim rac{q(0) arGamma igg(rac{1}{2}igg) ext{exp}igg(rac{i\pi}{4}igg)}{(-
u k g(0))^{1/2}} igg\{rac{\sqrt{2} ext{exp} \left[i(hF(au)+kG(au))
ight]}{(\lambda-
u)^{1/2}} - i
u^{-1/2}igg\} \,.$$

C. If $\nu > 2\gamma$ then

$$I(h, k) \sim \frac{q(0)\Gamma\left(\frac{\gamma}{\nu}\right) \exp\left(\frac{\gamma\pi}{2\nu i}\right)}{\nu [-kg(0)]^{\gamma/\nu}} .$$

Proof of Theorem III. We may assume that c is such that G'(t) < 0and f(t) > 0 for $0 < t \leq c$. For $0 < t \leq c$ let D(t) = F'(t)/-G'(t) with D(0) = 0. Then $D'(t) = t^{\lambda+\nu-t}/(G'(t))^2[\nu\lambda f(t)g(t)(\nu-\lambda) + tE(t)]$ for $0 < t \leq c$ where E is continuous on [0, c]. Hence there exists c_0 such that $0 < c_0 < c$, D'(t) > 0 for $0 < t \leq c_0$, $D(c_0) > 0$ and for h and k sufficiently large $k/h < D(c_0)$. This implies that there exists a unique $\tau \in (0, c_0)$ such that $D(\tau) = k/h$ which is equivalent to $hF'(\tau) + kG'(\tau) = 0$. Moreover from the definition of D we have

$$au = \left(rac{-
u kg(0)}{\lambda hf(0)}
ight)^{1/\lambda-
u} (1+o(1))$$

which implies that $\tau^{\lambda-\nu} = o(k/h) = o(1)$.

If we now let H(t) = F(t) + (k/h)G(t) and expand $h(H(t) - H(\tau))$ about $t = \tau$ we have using the integral form of the remainder

$$egin{aligned} h(H(t)-H(au)) &= h {\int_{ au}^t {(t-y)F^{\prime\prime}(y)dy} + k {\int_{ au}^t {(t-y)G^{\prime\prime}(y)dy} \ &= hR(t, au) + kS(t, au), \ ext{respectively.} \end{aligned}$$

We may further assume that c_0 is so small that f, f', f'', g, g' and g'' are of constant sign for $0 \le t \le c_0$. If we apply the mean value theorem for integrals and substitute $t = \tau(x+1)$ we have for -1 < x < 1 that

$$egin{aligned} T(x,\, au) &= R(au(x+1),\, au) = rac{ au^{\lambda}x^2}{2} \left[\lambda(\lambda-1)\,f(au_{\scriptscriptstyle 0}(x))lpha_{\scriptscriptstyle 0}(x)
ight. \ &+ au(\lambda+1)lpha_{\scriptscriptstyle 1}(x)\,f'(au_{\scriptscriptstyle 1}(x)) + au^2 f^{\,\prime\prime}(au_{\scriptscriptstyle 2}(x))lpha_{\scriptscriptstyle 2}(x)
ight] = rac{ au^{\lambda}x^2}{2} \;P_1(x) \end{aligned}$$

where α_0 , α_1 , $\alpha_2 \in C^{\infty}$, $\alpha_0(0) = \alpha_2(0) = \alpha_1(0) = 1$, $P_1 \in C^4$ and $P_1(0) = \lambda(\lambda - 1)$ f(0) + o(1). Similarly

$$W(x, \tau) = S(\tau(x + 1), \tau) = \frac{1}{2} \tau^{\nu} x^2 P_2(x)$$

where $P_2 \in C^4$ and $P_2(0) = \nu(\nu - 1)g(0) + o(1)$. Let $d_0 = (c_0/\tau) - 1$, $I'(h, k) = \exp(-ihH(\tau))I(h, k)$ and choose b such that $\tau(b + 1) < c_0$ and 0 < b < 1.

$$egin{aligned} I'(h,\,k) = & \int_{0}^{c_{0}} \,+\, 0 \Big(rac{1}{h} \Big) = \,I''(h,\,k) \,+\, 0 \Big(rac{1}{h} \Big) \,. \ I''(h,\,k) = & au \int_{-1}^{-b} \,+\, au \int_{-b}^{b} \,+\, au \int_{b}^{a_{0}} Q(au(x\,+\,1)) \exp{[i(h\,T(x,\, au)\,+\,k\,W(x,\, au))]} dx \ &= L(h,\,k) \,+\, I'''(h,\,k) \,+\, J(h,\,k), \,\, ext{respectively}. \end{aligned}$$

Let

$$\mu= au,\ p=(1/2)h au^{\lambda},\ \psi=q,\ \omega(x)=(1+x)^{\gamma-1}\ ext{and}\ \phi(x)=P_1(x)+(k au^{
u-\lambda}/h)P_2(x).$$

Then $\phi(0) = \lambda(\lambda - \nu)f(0) + o(1)$ implies for *h*, *k* sufficiently large that $\phi(x) > 0$ for $-c' \leq x \leq c', c' > 0$. Hence *b* may be chosen small enough that the conditions on *b* in Lemma II are satisfied. Therefore

$$I^{\prime\prime\prime}(h,k) \sim \frac{\sqrt{2} q(0) \Gamma\left(\frac{1}{2}\right) \exp\left(\frac{i\pi}{4}\right)}{(\lambda-\nu)^{1/2}} \left[\frac{(\lambda hf(0))^{\nu-2\gamma}}{(-\nu kg(0))^{\lambda-2\gamma}}\right]^{1/2(\lambda-\nu)}$$

The contribution of L(h, k) to I(h, k) may be determined by considering

$$L'(h, k) = \int_0^{r(1-b)} Q(t) \exp{(ihH(t))} dt \; .$$

We note that the uniqueness of τ in $[\varepsilon, c_0]$, $\varepsilon > 0$, implies that $H'(t) \neq 0$ on $\varepsilon \leq t \leq \tau(1-b)$ for every $\varepsilon > 0$. In fact, there exists a number K > 0 which is independent of ε and for which we have $|H'(t)| \geq K(k/h)t^{\nu-1}$ for $\varepsilon \leq t \leq \tau(1-b)$.

(i) For $\nu < \gamma$ the usual integration by parts together with the above inequality for H'(t) yields that L'(h, k) = o(1/k). Hence $L'(h, k) = o(h^{\nu-2\gamma}/k^{\lambda-2\gamma})^{1/2(\lambda-\nu)}$ since $h^{\nu} = o(k^{\lambda})$.

(ii) For $\nu > \gamma$ we rewrite L'(h, k) as

$$L'(h, k) = \int_{0}^{\tau(1-b)} Q(t) \exp \{ -i[kt^{\nu}(-g(t)) + ht^{\lambda}(-f(t))] \} dt$$

and apply Theorem II with -g playing the role of f. Hence

$$L'(h, k) \sim rac{q(0) \Gamma\left(rac{\gamma}{
u}
ight) \exp\left(rac{-i\pi\gamma}{2
u}
ight)}{
u[-kg(0)]^{\gamma/
u}} \; .$$

(iii) Finally for $\nu = \gamma$ a closer examination of the proof of Lemma I together with the change of variable $t = xh^{-1/\lambda}$ implies for $p' = k/h^{\nu/\lambda}$ that $L'(h, k) = h^{-\gamma/\lambda} O(1/p') = O(1/k)$.

The given relation $h^{\nu} = o(k^{\lambda})$ and the calculation

$$k^{\gamma/
u} 0 \Big(\Big(rac{h^{
u-2\gamma}}{k^{\lambda-2\gamma}}\Big)^{1/2(\lambda-
u)} \Big) = 0 \, \Big(\Big(rac{h}{k^{\lambda/
u}}\Big)^{
u-2\gamma/2(\lambda-
u)} \Big)$$

then imply that $I'''(h, k) = o(k^{-\gamma/\nu})$ if $\nu > 2\gamma$ and $L'(h, k) = o((h^{\nu-2\gamma}/k^{\lambda-2\gamma})^{1/2(\lambda-\nu)})$ if $\gamma \leq \nu < 2\gamma$. When $\nu = 2\gamma$ we note that both L'(h, k) and I'''(h, k) are of the same order so that both terms contribute to I(h, k).

To complete the proof of Theorem III we need only show that J(h, k) is negligible compared to I'''(h, k). For P(t) defined as in the proof of Theorem I and $d = \tau(b+1)$ we have

$$P(d) = \lambda f(0)[1 - (b + 1)^{\nu - \lambda}](1 + o(1))$$
.

Then P(d) > 0 for h and k sufficiently large and hence proceeding exactly as in the proof of Theorem I we obtain $H'(t) \ge Bt^{\lambda-1} > 0$ for $0 < d \le t \le c_0$ and $2B = \lambda f(0)[1 - (1 + b)^{\nu-\lambda}]$. We now write

$$J(h, k) = \int_{d}^{c_0} Q(t) \exp{(ihH(t))} dt$$

and integrate by parts as in Theorem I to obtain $J(h, k) = 0((h^{\nu-\gamma}/k^{\lambda-\gamma})^{1/\lambda-\nu})$. Hence $h^{\nu} = o(k^{\lambda})$ implies that

$$\left(rac{k^{\lambda-2\gamma}}{h^{
u-2\gamma}}
ight)^{1/2(\lambda-
u)}J(h,\,k)=0\left(\left(rac{h^
u}{k^\lambda}
ight)^{1/2(\lambda-
u)}
ight)=o(1)\;.$$

To obtain the value of $\exp [i(hF(\tau) + kG(\tau))]$ in a more explicit form we need to know more about the exact relation between h and k. For example we shall state the following corollary under more stringent assumptions.

COROLLARY. If in addition to the above assumptions in Theorem III we have $k^{\lambda+1} = o(h^{\nu+1})$ then

$$\exp\left[i(hF(\tau)+kG(\tau))\right] \backsim \exp\left\{\frac{i(\nu-\lambda)}{\lambda\nu}\left[\frac{(-\nu kg(0))^{\lambda}}{(\lambda hf(0))^{\nu}}\right]^{1/\lambda-\nu}\right\}\,.$$

Proof of the Corollary to Theorem III. We will use the same notation as in the proof of Theorem III. If we expand $hH(\tau)$ about the origin and substitute for τ we have

$$\begin{split} \exp((ihH(\tau)) &= \exp\left\{i\left[hf(0)\left(\frac{-\nu kg(0)}{\lambda hf(0)}\right)^{\lambda/\lambda-\nu} \right. \\ &+ kg(0)\left(\frac{-\nu kg(0)}{\lambda hf(0)}\right)^{\nu/\lambda-\nu}\right]\left[1 + 0\left(\left(\frac{k}{h}\right)^{1/\lambda-\nu}\right)\right]\right\} \\ &= \exp\left\{\frac{i(\nu-\lambda)}{\lambda\nu}\left[\frac{(-\nu kg(0))^{\lambda}}{(\lambda hf(0))^{\nu}}\right]^{1/\lambda-\nu} + 0\left(\left(\frac{k^{\lambda+1}}{h^{\nu+1}}\right)^{1/\lambda-\nu}\right)\right\}. \end{split}$$

Hence if $k^{\lambda+1} = o(h^{\nu+1})$ the Corollary is established.

Finally, we shall consider the case where $\limsup kh^{-\nu/\lambda}$ is bounded away from both 0 and ∞ .

THEOREM IV. Suppose that $\gamma < \lambda$, $\nu < \lambda$ and $0 < \limsup p < \infty$ where $p = kh^{-\nu/\lambda}$. Then

Proof of Theorem IV. We will consider only values of c > 0 such that (i) $\nu g(t) + tg'(t)$ is of constant sign for $0 \leq t \leq c$ and (ii) for each $\varepsilon > 0$ we have $|q(t) - q(0)| < \varepsilon$, $|f(t) - f(0)| < \varepsilon$ and $|g(t) - g(0)| < \varepsilon$ for $0 \leq t \leq c$. Set H(t) = F(t) + (k/h) G(t) and $I'(h, k) = \int_{0}^{\varepsilon} as$ usual. Let $m = \min_{0 \leq t \leq a} \lambda f(t) + tf'(t) > 0$, $\omega = \limsup p$ and $M = \max_{0 \leq t \leq a} (1, |f^{(t)}|, |q|, |q'|, |g^{(t)}|, |\nu g(t) + tg'(t)|)$ for l = 0, 1, 2. Consider a number b > 1 chosen such that $b > N = (4M\omega/m)^{1/(\lambda-\nu)}$. If $d = bh^{-1/\lambda} < c$ then for $0 < d \leq t \leq c$ we have for g(0) < o

$$H'(t) \ge t^{\lambda-1} \Big(m - rac{kM}{h d^{\lambda-
u}} \Big) \ge t^{\lambda-1} \Big(m - rac{mk}{4\omega h^{
u/\lambda}} \Big) \ge rac{1}{2} m t^{\lambda-1}$$

since $2\omega > kh^{-\nu/\lambda}$ for h and k sufficiently large. Hence $H'(t) \ge (1/2)mt^{\lambda-1} > 0$ for $0 < d \le t \le c$. Let $t = xh^{-1/\lambda}$, $\tilde{q}(x) = q(xh^{-1/\lambda})$, $\tilde{f}(x) = f(xh^{-1/\lambda})$ and $\tilde{g}(x) = g(xh^{-1/\lambda})$. Then

$$\begin{split} h^{\gamma/\lambda}I'(h,\,k) = & \int_{0}^{ch^{1/\lambda}} x^{\gamma-1}\widetilde{q}(x) \exp\left[i(\widetilde{f}(x)x^{\lambda} + p\widetilde{g}(x)x^{\nu})\right] dx \\ = & \int_{0}^{b} + \int_{b}^{ch^{1/\lambda}} = I''(h,\,k) + J(h,\,k), \text{ respectively} \end{split}$$

We will first estimate J(h, k) in terms of the number b. Since $H'(t) \ge (1/2)mt^{\lambda-1} > 0$ for $0 < d \le t \le c$ we may integrate J(h, k) by parts as follows:

$$\begin{split} J(h, \, k) &= h^{\gamma/\lambda} \left\{ \frac{Q(c) \exp{(ihH(c))}}{ihH'(c)} - \frac{Q(d) \exp{(ihH(d))}}{ihH'(d)} \\ &- \frac{1}{ih} \int_{a}^{c} \frac{Q'(t) \exp{(ihH(t))}dt}{H'(t)} + \frac{1}{ih} \int_{a}^{c} \frac{Q(t)F''(t) \exp{(ihH(t))}dt}{[H'(t)]^2} \\ &+ \frac{k}{ih^2} \int_{a}^{c} \frac{Q(t)G''(t) \exp{(ihH(t))}dt}{[H'(t)]^2} \right\} \\ &= 0(h^{\gamma-\lambda/\lambda}) + A + J'(h, \, k) + J''(h, \, k) + J'''(h, \, k), \text{ respectively.} \end{split}$$

Hence $|A| \leq 2M/mb^{\lambda-\gamma} = Bb^{\gamma-\lambda}$,

$$egin{aligned} |J'(h,k)| &\leq rac{2Mh^{\gamma/\lambda}}{mh} \int_a^c t^{\gamma-\lambda-1} dt < rac{2M}{m(\lambda-\gamma)b^{\lambda-\gamma}} = B'b^{\gamma-\lambda} \ , \ &|J''(h,k)| &\leq rac{4M^2}{m^2(\lambda-\gamma)b^{\lambda-\gamma}} = B''b^{\gamma-\lambda} \ , \ ext{and} \ &|J'''(h,k)| &\leq rac{4M^2kh^{-
u/\lambda}}{m^2(2\lambda-\gamma-
u)b^{2\lambda-\gamma-
u}} &\leq rac{8M^2\omega}{m^2(2\lambda-\gamma-
u)b^{\lambda-\gamma}} = B'''b^{\gamma-\lambda} \ . \end{aligned}$$

Define

$$h^{\gamma/\lambda}I_{0}(h, k) = \int_{0}^{\infty} x^{\gamma-1}q(0) \exp\left[i(f(0)x^{\lambda} + pg(0)x^{\gamma})\right] dx = \int_{0}^{b} + R(b)$$

Then there exists a number K which is independent of h, k and ε and for which $|J(h, k)| \leq Kb^{\gamma-\lambda}$ and $|R(b)| \leq Kb^{\gamma-\lambda}$. Consider

where $P(x) = \exp \{i[\tilde{f}(x) - f(0))x^{\lambda} + p(\tilde{g}(x) - g(0))x^{\nu}]\}$. By the choice

of c for each $\varepsilon > 0$ we have $|L(h, k)| \leq M \varepsilon b^{\lambda+\gamma}$, $|L'(h, k)| < 2M \varepsilon b^{\lambda+\gamma} + 2M \varepsilon \omega b^{\nu+\gamma}$. If we take lim sup of both sides of (*) as $h, k \to \infty$ we obtain

$$0 \leq \limsup h^{\gamma/\lambda} |I_0 - I| \leq 3M arepsilon b^{\lambda+\gamma} + 2M arepsilon b^{
u+\gamma} + 2K b^{\gamma-\lambda}$$

which is true for $\varepsilon > 0$ and b > N. Since $\limsup h^{\gamma/\lambda} |I_0 - I|$ is independent of both ε and b we first let $\varepsilon \to 0$ and then $b \to \infty$. Hence $h^{\gamma/\lambda}(I_0 - I) = o(1)$ which implies that $I(h, k) \backsim I_0(h, k)$. To obtain the alternate form of $I_0(h, k)$ we let $x = h^{1/\lambda}(k/h)^{1/(\lambda-\nu)}t$.

5. Discussion of the suggested application. Consider for x > 0Schlafli's generalization of Bessel's integral:

$$egin{aligned} J_{
u}(x) &= rac{1}{\pi} \int_0^\pi \cos{(
u t - x \sin{t})} dt - rac{\sin{
u \pi}}{\pi} \int_0^\infty \exp{[-
u t - x \sin{ht}]} dt \ &= rac{1}{\pi} R \int_0^\pi \exp{[i(
u t - x \sin{t})]} dt + 0\left(rac{1}{
u}
ight). \end{aligned}$$

Let $F(t) = t - \sin t$ and |G(t)| = t. We rewrite F(t) as $F(t) = (1/6)t^3 \cos(r(t))$ and let h = x, $k = |\nu - x|$, $q(t) \equiv 1$ and $f(t) = 1 \setminus 6 \cos(r(t))$. It follows that the condition 3f(t) + tf'(t) > 0 for $0 \leq t \leq \pi$ is satisfied since $F'(t) = 1 - \cos t > 0$ for $0 < t \leq \pi$.

We note that our Theorem I and III yield the dominant terms of some well known complete asymptotic expansions for $J_{\nu}(x)$ with $\tau = \operatorname{Arccos} \nu/x$ in Theorem III¹. For the case $0 < \limsup x^{-1/3} | \nu - x | < \infty$ we have by Theorem IV with $p = x^{-1/3} | (\nu - x)x^{-1/3} |$ that

$$J_{
u}(x) \backsim rac{1}{\pi x^{1/3}} \int_{0}^{\infty} \cos \Big(rac{1}{6} t^3 + pt \Big) dt$$

where the expression on the right is one of Airy's integrals², whose evaluation for p > 0 and p < 0 yields precisely Nicholson's formulas when ν is an integer³.

INSTITUTE OF TECHNOLOGY UNIVERSITY OF MINNESOTA

¹ See W. Magnus and F. Oberhettinger, "Formeln und Satze fur die Speziellen Funktionen der Mathematischen Physik," Springer-Verlag, Berlin, 1948, pp. 33-34. Our theorems I and III give results which are equivalent to the dominant terms of the expansions (b_3) and (b_1) , respectively.

² See, for example, G. N. Watson, "Theory of Bessel Functions," Cambridge, 1944, pp. 188-190.

³ See G. N. Watson, op. cit., pp. 248-249.

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Pacific Journal of Mathematics Vol. 12, No. 4 April, 1962

Tsuyoshi Andô, On fundamental properties of a Banach space with a cone	1163		
Sterling K. Berberian, A note on hyponormal operators			
Errett Albert Bishop, Analytic functions with values in a Frechet space	1177		
(Sherman) Elwood Bohn, Equicontinuity of solutions of a quasi-linear			
equation	1193		
Andrew Michael Bruckner and E. Ostrow, <i>Some function classes related to the</i>			
class of convex functions	1203		
J. H. Curtiss, <i>Limits and bounds for divided differences on a Jordan curve in the complex domain</i>	1217		
P. H. Doyle, III and John Gilbert Hocking, <i>Dimensional invertibility</i>	1235		
David G. Feingold and Richard Steven Varga, <i>Block diagonally dominant matrices</i>			
and generalizations of the Gerschgorin circle theorem	1241		
Leonard Dubois Fountain and Llovd Kenneth Jackson. A generalized solution of the			
boundary value problem for $y'' = f(x, y, y')$	1251		
Robert William Gilmer, Jr., <i>Rings in which semi-primary ideals are primary</i>	1273		
Ruth Goodman <i>K</i> -nolar polynomials	1277		
Israel Halperin and Maria Wonenburger. On the additivity of lattice	1277		
completeness	1289		
Robert Winship Heath Arc-wise connectedness in semi-metric spaces	1301		
Isidore Heller and Alan Jerome Hoffman. On unimodular matrices	1321		
Pohert C. Haynomen, Duality in conord anodia theory	1220		
Charles Dev Hohm, <i>Duality in general ergoaic ineory</i>	1329		
Charles Ray Hobby, Abertan subgroups of p-groups	1545		
Kenneth Myron Hoffman and Hugo Rossi, <i>The minimum boundary for an analytic</i>	1247		
polynearon	1347		
Adam Koranyi, The Bergman kernel function for tubes over convex cones	1355		
Pesi Rustom Masani and Jack Max Robertson, The time-domain analysis of a	12(1		
continuous parameter weakly stationary stochastic process	1301		
William Schumacher Massey, Non-existence of almost-complex structures on	1270		
quaternionic projective spaces	13/9		
Deane Montgomery and Chung-Tao Yang, A theorem on the action of SO(3)	1385		
Ronald John Nunke, A note on Abelian group extensions	1401		
Carl Mark Pearcy, A complete set of unitary invariants for operators generating finite W*-algebras of type I	1405		
Edward C. Posner, Integral closure of rings of solutions of linear differential	1417		
Puere Sother Assure taking III. Station and the family a sure of the state of the s	141/		
Duane Sather, Asymptotics. III. Stationary phase for two parameters with an	1422		
application to Bessel functions	1423		
J. Sladkowska, Bounds of analytic functions of two complex variables in domains with the Bergman-Shilov boundary	1435		
Joseph Gail Stampfli, <i>Hyponormal operators</i>	1453		
George Gustave Weill, Some extremal properties of linear combinations of kernels			
on Riemann surfaces	1459		
Edward Takashi Kobayashi, Errata: "A remark on the Nijenhuis tensor"	1467		