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AN EXTREMAL LENGTH CRITERION FOR THE PARABOLICITY OF RIEMANNIAN SPACES

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AN EXTREMAL LENGTH CRITERION FOR THE PARABOLICITY OF RIEMANNIAN SPACES

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It is the purpose of this paper to show that a given Riemannian space satisfying a regularity condition is parabolic if and only if the extremal distance of a fixed ball in the space from the ideal boundary of the space is infinite.

We will also show that the harmonic modulus of a space bounded by two sets of boundary components coincides with the extremal distance between the two sets.

STATEMENTS OF MAIN RESULTS

1. Regularity condition. Throughout this paper we denote by R a noncompact C^{∞} Riemannian space with the ideal boundary β . We always assume that R is orientable and connected. Let A be the complement of a regular subregion of R with the relative boundary α . We also assume that $A \cdot \alpha$ is connected. We consider the following regularity condition for R (more precisely, for A):

For any nonconstant harmonic function u defined on a region $\Omega \subset A$, the set $\{x \in \Omega \mid | \nabla u(x) | = 0\}$ has zero capacity.

This condition is always satisfied if the dimension of R is two. This is also true, for example, when the metric tensor g_{ij} is real analytic on A- α . A typical case is furnished by a locally flat A- α .

In this paper we only consider those spaces R for which the above regularity condition is met.

2. Extremal length. Let ρ be a density, i.e. a nonnegative Borel function on A, and let Γ be a family of curves γ which issue from a point in α and lie in A- α . We define the harmonic extremal length, or simply the extremal length of Γ , by

(1)
$$\lambda(\Gamma) = \sup_{\rho \neq 0} \frac{L(\Gamma, \rho)^2}{V(A, \rho)},$$

where $V(A, \rho) = \int_{A} \rho^{2} dV$ and $L(\Gamma, \rho) = \inf_{\Gamma} \int_{\gamma} \rho ds$. Here dV and ds are the volume and the line element

We are particularly interested in the family $\Gamma_{\beta} \subset \Gamma$ of all curves $\gamma \in \Gamma$ terminating at β .

3. Parabolicity. We call R parabolic, $R \in O_G$, if R carries no nonconstant positive superharmonic function. The main object of this

paper is to prove:

THEOREM 1. The space R is parabolic if and only if $\lambda(\Gamma_{\beta}) = \infty$.

4. Moduli. Let Ω be a regular subregion of R with relative boundary $\beta_{g} \subset A - \alpha$, and let u_{g} be the continuous function on $\overline{\Omega} \cap A$ which is harmonic in the interior of $\overline{\Omega} \cap A$ with $u_{g}| = 0$ and $u_{g}|\beta_{g} = 1$. The constant μ_{g} given by

$$(2) \hspace{1cm} \log \mu_{\scriptscriptstyle \mathcal{G}} = 1/\int_{\overline{a} \cap A} du_{\scriptscriptstyle \mathcal{G}} \wedge * du_{\scriptscriptstyle \mathcal{G}}$$

is called the *harmonic modulus*, or simply the modulus of $\overline{\Omega} \cap A$ with respect to α . It is easy to see that

$$(3) \qquad \qquad \mu_{\mathfrak{g}} \leq \mu_{\mathfrak{g}'}$$

for $\Omega \subset \Omega'$. Therefore, we can define μ_R , the harmonic modulus of A with respect to α , as the directed limit

$$(4) \qquad \qquad \mu_{R} = \lim_{g \to R} \mu_{g} .$$

It is again easy to see that $u_R = \lim_{\rho \to R} u_\rho$ exists and is continuous on A, harmonic on $A - \alpha$ with $u_R | \alpha = 0$. Moreover,

(5)
$$\log \mu_R = 1/\int_{\mathcal{A}} du_R \wedge * du_R$$
.

It can be seen that $R \in 0_G$ if and only if $\mu_R = \infty$ (Glasner [3]). Thus Theorem 1 may be considered as a special case of

THEOREM 2. The following identity is valid:

$$\lambda(\Gamma_{\beta}) = \log \mu_{R}.$$

The proof will be given in 5-9.

PARABOLIC CASE

5. A general inequality. We start with proving

$$(7) \qquad \qquad \lambda(\Gamma_{\beta}) \ge \log \mu_{R} .$$

Let $\Gamma_{B_{\rho}}$ be the family of curves $\gamma \in \Gamma$ which lie in $\overline{\Omega} \cap A$ and terminate at a point of β_{ρ} . Define ρ as $(\log \mu_{\rho}) | \mathcal{F}u_{\rho} |$ in the interior of $\overline{\Omega} \cap A$ and as zero elsewhere in R. For $\gamma \in \Gamma_{\beta_{\rho}}$,

$$\int_{\mathbf{y}}
ho \ ds = \int_{\mathbf{y}} (\log \mu_{a}) \ | \ arpsi u_{a} \ | \ ds \ge (\log \mu_{a}) \int_{\mathbf{y}} rac{dh}{ds} \ ds = \log \mu_{a} \ .$$

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Therefore

$$L({\Gamma}_{{}^{eta}{}_{\mathcal{G}}},
ho)=\inf_{{}_{\Gamma}}\int_{\gamma}\!\!
ho ds\geq\log\mu_{a}$$
 .

By (2) we also obtain

and infer by (1) that

$$\lambda(\Gamma_{\beta g}) \ge \log \mu_g \,.$$

Since every $\gamma \in \Gamma_{\beta}$ contains a $\gamma' \in \Gamma_{\beta_{\alpha}}$, we can easily see that $\lambda(\Gamma_{\beta}) \geq \gamma(\Gamma_{\beta_{\alpha}})$ (cf. Ahlfors-Sario [1, p. 222]). Thus (8) implies that $\lambda(\Gamma_{\beta}) \geq \log \mu_{\alpha}$ for every Ω . On letting $\Omega \to R$ we obtain (7).

6. Now suppose that $R \in 0_G$. Then since $\mu_R = \infty$, (7) implies that

$$\lambda(\Gamma_{\beta}) = \log \mu_{R} = \infty .$$

In order to complete the proofs of Theorems 1 and 2, we have only to show the validity of (6) under the assumption $R \notin 0_c$. Note that in our discussion thus far we have not made any use of the regularity condition.

HYPERBOLIC CASE

7. *u*-lines. Hereafter we assume that $R \notin 0_G$. Then u_R , to be denoted simply by u, is not constant on A. Since $u \mid \alpha = 0$ and $u \mid A - \alpha > 0$, we infer that $|\Gamma u|$ can be extended continuously to all of A and that $|\Gamma u \mid \alpha \neq 0$.

For each $x \in \alpha$ we consider the unique curve l_x issuing from xand such that $l_x - x \subset A - \alpha$, *du = 0 on l_x , $|\nabla u| \neq 0$ on l_x . Moreover we require that l_x either terminates at β or at a point of A at which $|\nabla u| = 0$. Such an l_x will be called a *u*-line. As y traces l_x , u(y) increases. Thus we can classify points of α as follows:

$$egin{aligned} lpha_{\scriptscriptstyle 0} &= \{x \in lpha \mid \lim_{y o eta, y \in \iota_x} u(y) < 1\} \;, \ lpha_{\scriptscriptstyle 1} &= \{x \in lpha \mid \lim_{y o eta, y \in \iota_x} u(y) = 1\} \;, \end{aligned}$$

with

$$(10) \qquad \qquad \alpha = \alpha_0 \cup \alpha_1 .$$

8. Vanishing surface area. We denote by dS the surface element of α . We wish to show that

(11)
$$S(lpha_{\scriptscriptstyle 0})=\int_{lpha_{\scriptscriptstyle 0}} dS=0$$
 .

Let F_{-1} be the set of points $x \in \alpha$ such that l_x terminates at some point of R. Clearly $F_{-1} \subset \alpha_0$, and we set $F_0 = \alpha_0 - F_{-1}$. By the regularity condition in §1, we see that $S(F_{-1}) = 0$ (cf. Brelot-Choquet [2]). Therefore we only have to show that $S(F_0) = 0$. Let

$$F_n = \left\{ x \in F_0 \left| \lim_{y \to eta, y \in I_x} (1 - u(y)) \ge \frac{1}{n} \right\} \quad (n = 1, 2, \cdots) .$$

Since $F_0 = \bigcup_{i=1}^{\infty} F_n$, it is sufficient to show that $S(F_n) = 0$.

We can find a positive harmonic function ω in the interior of A with the following properties (cf. Nakai [4]): (a) ω has the boundary values 0 on α , (b) $\lim_{y\to\beta,y\in l_x}\omega(y) = \infty$ for $x\in F_0$, (c) $\int_A |\nabla \omega_c|^2 dV \leq c$, with $\omega_c = \min(\omega, c)$ for every positive number c.

Fix a c > 0 arbitrarily and a point $y_x \in l_x$ with $w_c(y_x) = c$ for each $x \in F_n$.

Set v = 1 - u on A. In a neighborhood of a point in α with respect to A we may incorporate v into a coordinate system, say $v = x^1$, while x^2, \dots, x^m are m - 1 linearly independent parameters for α . Then

$$\mid {\it _d} v \mid^{\scriptscriptstyle 2} = g^{\scriptscriptstyle 11} \Bigl({\partial v \over \partial x^{\scriptscriptstyle 1}} \Bigr)^{\scriptscriptstyle 2} = g^{\scriptscriptstyle 11} \; .$$

Since $*dv = |\nabla v| dS = \sqrt{g^{11}} dS$ on $\alpha, S(F_n) = 0$ is equivalent to $\int_{F_n} *dv = 0$. Observe that

$$egin{aligned} &c \int_{F_n} * dv &\leq \int_{F_n} \Bigl(\int_x^{y_x} rac{\partial \omega_o}{\partial v} \, dv \Bigr) * dv \ &= \int_{F_n} \int_x^{y_x} \Bigl| rac{\partial \omega_o}{\partial v} \Bigr| \, dv \wedge * dv \ &= \int_{F_n} \int_x^{y_x} \Bigl| rac{\partial \omega_o}{\partial v^1} \Bigr| \, g^{11} dV. \end{aligned}$$

By the Schwarz inequality we have

$$\int_{F_n} \int_x^{y_x} \left| \frac{\partial \omega_{\scriptscriptstyle c}}{\partial x^{\scriptscriptstyle 1}} \right| g^{\scriptscriptstyle 11} d\, V \leq \left(\int_{F_n} \int_x^{y_x} \left| \frac{\partial \omega_{\scriptscriptstyle c}}{\partial x^{\scriptscriptstyle 1}} \right|^2 g^{\scriptscriptstyle 11} d\, V \right)^{1/2} \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_{F_n} \int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} d\, V \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} dV \right)^{1/2} dV = \left(\int_x^{y_x} g^{\scriptscriptstyle 11} dV \right)^{1/2} dV = \left(\int_x^{y_x} g^{$$

$$egin{aligned} &\leq \left(\int_A | \, arphi oldsymbol{\omega}_{oldsymbol{s}} \, |^2 d \, V
ight)^{1/2} & \left(\int_A | \, arphi \, v \, |^2 d \, V
ight)^{1/2} \ &\leq \sqrt{-c} \left(\int_A d u \, \wedge \, st \, d u
ight)^{1/2} \, . \end{aligned}$$

From this we infer that

$$\left|\int_{{}^{F}{}_{n}}{}^{*}dv
ight|\leq 1/\sqrt{\mu_{{}_{R}}c}\;\;.$$

Since the number c can be arbitrarily large, we have $\int_{F_n} * dv = 0$, and (11) follows.

9. Let ρ be a density with $\rho \neq 0$ on A. Since

$$du \wedge *du = |ec vu| \, dV$$
 ,

we can compute

$$egin{aligned} V(A,\,
ho)&=\int_{A}
ho^{2}d\,V=\int_{A}rac{
ho^{2}}{|ec{
aligned} u||^{2}}du\,\wedge\,st du\ &\geqq\int_{lpha_{1}}\Bigl(\int_{lx}rac{
ho^{2}}{|ec{
aligned} u||^{2}}du\Bigr)st du\ &=\int_{lpha_{1}}\Bigl(\int_{lx}rac{
ho^{2}}{|ec{
aligned} u||^{2}}du\cdot\int_{lx}1^{2}du\Bigr)st du\ &\geqq\int_{lpha_{1}}\Bigl(\int_{lx}rac{
ho}{|ec{
aligned} u||^{2}}du\,
ight)^{2}st du\ . \end{aligned}$$

On $l_x(x \in \alpha_1)$ we have $du = |\nabla u| ds$, and thus

$$V(A, \rho) \ge \int_{\alpha_1} \left(\int_{I_x} \rho ds \right)^2 * du$$
.

From $l_x \in \Gamma_\beta$ for $x \in \alpha_1$ we obtain $\int_{l_x} \rho ds \ge L(\Gamma_\beta, \rho)$, and therefore

(12)
$$V(A, \rho) \ge L(\Gamma_{\beta}, \rho)^2 \int_{\alpha_1} * du .$$

On the other hand, by (11), we have $\int_{\alpha_1} * du = \int_{\alpha} * du$. Take an arbitrary regular region Ω with $\beta_{\rho} \subset A - \alpha$. Then

$$\int_{lpha} * du = \lim_{arrho o R} \int_{lpha} * du_{arrho} \; .$$

Here we see that

$$\int_{\mathfrak{a}} * du_{\mathfrak{g}} = \int_{\beta_{\mathcal{G}}} * du_{\mathfrak{g}} = \int_{\beta_{\mathcal{G}} - \mathfrak{a}} u_{\mathfrak{g}} * du_{\mathfrak{g}} = \int_{\overline{\mathfrak{g}} \cap A} du_{\mathfrak{g}} \wedge * du_{\mathfrak{g}} ,$$

and infer that

$$\int_{lpha} * du = \lim_{{m{g}} o R} \int_{\overline{m{g}} \cap A} du_{m{g}} \wedge * du_{m{g}} = \int_{A} du \wedge * du \; .$$

This together with (5) and (12) implies the inequality

$$\log \mu_{\scriptscriptstyle R} \geq rac{L(\varGamma_eta,
ho)^2}{V(A,
ho)} \; .$$

Since ρ was arbitrary, we now conclude that

 $\log \mu_R \geq \lambda(\Gamma_\beta)$.

We combine this with (7) and obtain (6).

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