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SOME DUAL SERIES EQUATIONS INVOLVING LAGUERRE POLYNOMIALS

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In this paper an exact solution is found for the dual series equations

$$(\ 1\) \quad \sum\limits_{n=0}^{\infty} C_n arGamma(lpha+eta+n) L_n(lpha\,;\,x) = f(x) \; ext{,} \qquad 0 \leqq x < d \; ext{,}$$

$$(\; 2\;) \quad \sum\limits_{n=0}^{\infty} C_n arGamma(lpha+1+n) L_n(lpha\,;\, x) = g(x)$$
 , $\qquad d < x < \infty$,

where $\alpha + \beta > 0$, $0 < \beta < 1$, $L_n(\alpha; x) = L_n^{\alpha}(x)$ is the Laguerre polynomial and f(x) and g(x) are known functions.

In a recent paper Srivastava [3] has solved the equations

$$(\ 3\) \qquad \qquad \sum\limits_{n=0}^{\infty}\left\{ A_{n}/\Gamma(lpha+1\,+\,n)
ight\} L_{n}(lpha\ ;\,x)=f(x)\ , \qquad 0 \leq x < d\ ,$$

(4)
$$\sum\limits_{n=0}^{\infty} \{A_n/\Gamma(lpha+1/2+n)\}L_n(lpha\,;\,x) = g(x)$$
 , $d < x < \infty$, $lpha > -1/2$,

by considering separately the equations when (a) $g(x) \equiv 0$, (b) $f(x) \equiv 0$, and reducing the problem in each case to that of solving an Abel integral equation. Srivastava's equations are a special case of (1) and (2) with $\beta = 1/2$ and $A_n = \Gamma(\alpha + 1 + n)\Gamma(\alpha + 1/2 + n)C_n$.

The solution presented in this paper employs a multiplying factor technique which is more direct than the method given in [3] and is similar to that used by Noble [2] to solve some dual series equations involving Jacobi polynomials.

2. In the course of the analysis we shall use the following results.

From [1, p. 293(5), p. 405(20)] it is readily shown that

$$(5) \int_0^y x^\alpha (y-x)^{\beta-1} L_n(\alpha\ ;\ x) dx = \frac{\Gamma(\beta)\Gamma(\alpha\ +\ 1\ +\ n)}{\Gamma(\alpha\ +\ \beta\ +\ 1\ +\ n)} \ y^{\alpha+\beta} \ L_n(\alpha\ +\ \beta\ ;\ y) \ ,$$

where $-1 < \alpha, \beta > 0$, and

$$(\ 6\) \quad \int_y^\infty (x-y)^{-\beta} e^{-x} \, L_{\scriptscriptstyle n}(\alpha\ ; \, x) dx = \varGamma (1-\beta) e^{-y} \, L_{\scriptscriptstyle n}(\alpha+\beta-1\ ; \, y) \ ,$$

where $1 > \beta$, $\alpha + \beta > 0$.

The orthogonality relation for the Laguerre polynomials is

(7)
$$\int_0^\infty x^\alpha e^{-x} L_n(\alpha \, ; x) L_m(\alpha \, ; x) dx = \frac{\Gamma(\alpha+1+n)}{\Gamma(n+1)} \, \delta_{mn}, \, \alpha > -1 \, ,$$

where δ_{mn} is the Kronecker delta.

3. Solution of the problem. Multiplying equation (1) by $x^{\alpha}(y-x)^{\beta-1}$, equation (2) by $(x-y)^{-\beta}e^{-x}$ and integrating with respect to x over (0, y) and (y, ∞) respectively we find on using the results (5) and (6)

$$(8) \quad \sum_{n=0}^{\infty} C_n \frac{\Gamma(\alpha+1+n)}{(\alpha+\beta+n)} \ L_n(\alpha+\beta \ ; \ y) = \frac{y^{-\alpha-\beta}}{\Gamma(\beta)} \int_0^y x^{\alpha} (y-x)^{\beta-1} f(x) dx \ ,$$

where $0 < y < d, \alpha > -1, \beta > 0$, and

$$egin{align} \sum_{n=0}^\infty C_n \, arGamma(lpha+1+n) L_n(lpha+eta-1\,;\,y) \ &= rac{e^y}{arGamma(1-eta)} \int_y^\infty (x-y)^{-eta} \, e^{-x} \, g(x) dx \; , \end{align}$$

for $d < y < \infty$, $1 > \beta$, $\alpha + \beta > 0$.

If we now multiply equation (8) by $y^{\alpha+\beta}$, differentiate with respect to y and use the formula

(10)
$$\frac{d}{dx} \{x^{\alpha} L_n(\alpha; x)\} = (n + \alpha) x^{\alpha-1} L_n(\alpha - 1; x) ,$$

we find

(11)
$$\begin{split} \sum_{n=0}^{\infty} C_n \, \varGamma(\alpha+1+n) L_n(\alpha+\beta-1\ ;\ y) \\ &= \frac{y^{1-\alpha-\beta}}{\varGamma(\beta)} \, \frac{d}{dy} \int_0^y x^{\alpha} (y-x)^{\beta-1} f(x) dx \ , \end{split}$$

where $0 < y < d, \beta > 0, \alpha > -1$.

The left hand sides of equations (9) and (11) are now identical and using the orthogonality relation (7) we see that the solution of equations (1) and (2) for $\alpha + \beta > 0$, $0 < \beta < 1$, is given by

(12)
$$C_n = \frac{\Gamma(n+1)}{\Gamma(\alpha+1+n)\Gamma(\alpha+\beta+n)} B_n(\alpha,\beta;d) ,$$

where

(13)
$$B_n(\alpha,\beta;d) = \frac{1}{\Gamma(\beta)} \int_0^d e^{-y} L_n(\alpha+\beta-1;y) F(y) dy + \frac{1}{\Gamma(1-\beta)} \int_d^\infty y^{\alpha+\beta-1} L_n(\alpha+\beta-1;y) G(y) dy,$$

and

(14)
$$F(y) = \frac{d}{dy} \int_0^y x^{\alpha} (y-x)^{\beta-1} f(x) dx,$$

(15)
$$G(y) = \int_{y}^{\infty} (x - y)^{-\beta} e^{-x} g(x) dx.$$

To obtain the solution of Srivastava's equations (3) and (4) we write $\beta=1/2$, $A_n=\Gamma(\alpha+1+n)\Gamma(\alpha+1/2+n)C_n$ in (12) and find that

(16)
$$A_{n} = \frac{\Gamma(n+1)}{\Gamma(1/2)} \left\{ \int_{d}^{d} e^{-y} L_{n}(\alpha-1/2\;;y) F_{1}(y) dy \right. \\ \left. + \int_{d}^{\infty} y^{\alpha-1/2} L_{n}(\alpha-1/2\;;y) G_{1}(y) dy \right\},$$

for $\alpha > -1/2$, and where $F_1(y)$ and $G_1(y)$ are given by equations (14) and (15) respectively with $\beta = 1/2$.

Comparing the above solution with that obtained in [3] it can be seen that they are in agreement except for the form of the function $G_1(y)$. The limits on the integrals of equations (4.7) and (4.8) in Srivastava's paper are wrong and should read (x, ∞) and (u, ∞) respectively. When these corrections have been made we find that his term corresponding to $G_1(y)$ can be written in the notation of the present paper as

(17)
$$-\frac{d}{dy} \int_{y}^{\infty} (x-y)^{-1/2} dx \int_{x}^{\infty} e^{-u} g(u) du .$$

After inverting the order of integration, carrying out the integration in x and performing the differentiation with respect to y it is found that (17) is equal to $G_1(y)$. Hence with this simplification Srivastava's solution reduces to that given by equation (16).

- 4. It is also possible without computing the coefficients C_n to find the values of series (1) and (2) in the regions where their values are not specified. We define (1) to have the value h(x), $d < x < \infty$, and (2) to have the value k(x), $0 \le x < d$.
- (a) Calculation of h(x). Substituting for C_n from equation (12) into (1) and interchanging the order of integration and summation we find

(18)
$$h(x) = \frac{1}{\Gamma(\beta)} \int_0^d e^{-y} F(y) S_1(x, y) dy + \frac{1}{\Gamma(1-\beta)} \int_d^\infty y^{\alpha+\beta-1} G(y) S_1(x, y) dy , \qquad d < x < \infty ,$$

where

(19)
$$S_1(x, y) = \sum_{n=0}^{\infty} \frac{\Gamma(n+1)}{\Gamma(\alpha+1+n)} L_n(\alpha; x) L_n(\alpha+\beta-1; y).$$

Using the results (6) and (7) it is easily shown that

(20)
$$S_{1}(x, y) = \frac{e^{y} x^{-\alpha}(x - y)^{-\beta}}{\Gamma(1 - \beta)} H(x - y) ,$$

where H(x) is the Heaviside unit function.

From equations (18) and (20) we see that h(x) is given by

(21)
$$\Gamma(1-\beta)x^{\alpha}h(x) = \frac{1}{\Gamma(\beta)} \int_0^t (x-y)^{-\beta} F(y) dy$$
$$+ \frac{1}{\Gamma(1-\beta)} \int_d^x e^y y^{\alpha+\beta-1} (x-y)^{-\beta} G(y) dy,$$

for $d < x < \infty$, where F(y) and G(y) are given by equations (14) and (15).

(b) Calculation of k(x). Using the differentiation formula

(22)
$$e^{-x} L_n(\alpha; x) = -\frac{d}{dx} \{e^{-x} L_n(\alpha - 1; x)\},$$

we may write equation (2) as

(23)
$$\frac{d}{dx} e^{-x} \sum_{n=0}^{\infty} C_n \Gamma(\alpha + 1 + n) L_n(\alpha - 1; x)$$

$$= -e^{-x} k(x), \quad 0 \leq x < d.$$

Substituting for C_n and interchanging the order of integration and summation we find

(24)
$$e^{-x} k(x) = -\frac{d}{dx} e^{-x} \left\{ \frac{1}{\Gamma(\beta)} \int_{0}^{d} e^{-y} F(y) S_{2}(x, y) dy + \frac{1}{\Gamma(1-\beta)} \int_{d}^{\infty} y^{\alpha+\beta-1} G(y) S_{2}(x, y) dy \right\},$$

for $0 \le x < d$, and

$$S_2(x,\,y) = \sum\limits_{n=0}^{\infty} rac{ \Gamma(n+1)}{\Gamma(lpha+eta+n)} L_n(lpha-1\,;\,x) L_n(lpha+eta-1\,;\,y) \ = rac{1}{\Gamma(eta)} e^z (y-x)^{eta-1} \, y^{1-lpha-eta} \, H(y-x) \; ,$$

where the series has been summed using the results (6) and (7). Substituting for $S_2(x, y)$ in (24) we see that k(x) is given by

(26)
$$\Gamma(\beta)e^{-x} k(x) = -\frac{d}{dx} \left\{ \frac{1}{\Gamma(\beta)} \int_{x}^{d} e^{-y} (y-x)^{\beta-1} y^{1-\alpha-\beta} F(y) dy + \frac{1}{\Gamma(1-\beta)} \int_{d}^{\infty} (y-x)^{\beta-1} G(y) dy \right\},$$

when $0 \le x < d$.

It is perhaps interesting to note that the expressions for the functions k(x) and h(x) do not involve Laguerre polynomials.

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