

Pacific Journal of Mathematics

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This paper is concerned with random series of the form $\sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t)$ where the X_n 's are random variables, the a_n 's are real numbers, and the v_n 's are heat polynomials as introduced by P. C. Rosenbloom and D. V. Widder. The sequences $\{a_n\}$ are assumed to satisfy $\limsup_{n \rightarrow \infty} |a_n|^{2/n} (2n/e) = 1$ which implies $\sum_{n=0}^{\infty} a_n v_n(x, t)$ has $|t| < 1$ as its strip of convergence, i.e., it converges to a C^2 -solution of the heat equation in $|t| < 1$ and does not converge everywhere in any larger open strip. Associated with each sequence $\{a_n\}$ is its classification number from $[0, 1]$ which indicates how rapidly a_n tends to zero. Assumptions are placed on the random variables which imply that for almost every ω the series $\sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t)$ has $|t| < 1$ as its strip of convergence.

The main results of the paper are two theorems. The first states that if $\{a_n\}$ has its classification number in $[0, 1/2]$, then for almost every ω the lines $t = 1$ and $t = -1$ form the natural boundary for $\sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t)$. The second is concerned with sequences having their classification numbers in $(1/2, 1]$. The conclusion implies that for almost every ω no interval of either of the lines $t = 1$ or $t = -1$ is part of the natural boundary for $\sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t)$.

The present work had its original motivation in the study of the boundary behavior of random power series. These are series of the form $\sum_{n=0}^{\infty} a_n(\omega) z^n$ where the a_n 's are complex valued random variables and z is a complex number. Reference [1] contains a history of results in this area. One of the early results which helped to motivate the first part of the proof of our Theorem 1 is due to Paley and Zygmund in a 1932 paper [see 6, p. 220]. In this theorem it is assumed that $\sum_{n=0}^{\infty} a_n z^n$ is an ordinary power series with a finite radius of convergence. Letting $\{\phi_n\}$ be the sequence of Rademacher functions, the conclusion is that for almost every ω in $[0, 1]$ the series $\sum_{n=0}^{\infty} \phi_n(\omega) a_n z^n$ has its circle of convergence as its natural boundary.

More recently [see 3] V. L. Shapiro has considered series of the form $\sum_{n=0}^{\infty} X_n(\omega) H_n(x)$ where the X_n 's are random variables and

$$\sum_{n=0}^{\infty} H_n(x)$$

is the spherical harmonic representation of a harmonic function in the unit ball. The harmonic continuability across the boundary of the unit ball of the functions $\sum_{n=0}^{\infty} X_n(\omega) H_n(x)$ was investigated. This

work further motivated the first part of the proof of our Theorem 1 and influenced our choice of the class of random variables to be considered.

2. Definitions and preliminary comments. For a point (x_0, t_0) in the plane and a number $\rho > 0$ we let

$$S(x_0, t_0; \rho) = \{(x, t): |x - x_0| < \rho \text{ and } |t - t_0| < \rho\}.$$

If $u(x, t)$ is a C^2 -solution to the heat equation in the strip $|t| < \sigma$ we say the line $t = -\sigma$ ($t = \sigma$) is part of the natural boundary for u in case for every x_0 and every $\rho > 0$ there is no C^2 -solution $v(x, t)$ in $S(x_0, -\sigma; \rho)$ ($S(x_0, \sigma; \rho)$) which agrees with $u(x, t)$ where u and v are both defined.

Let E_0 be the set of all sequences $\{a_n\}_{n=0}^\infty$ of real numbers. For $r > 0$ let

$$E_r = \{\{a_n\} \in E_0: |a_n| (2n/e)^{n/2} = O(e^{-n^r}) \text{ as } n \rightarrow \infty\}.$$

We call $\sup\{r: \{a_n\} \in E_r\}$ the classification number of $\{a_n\}$.

Explicitly, from [2, p. 222]

$$(2.1) \quad v_n(x, t) = n! \sum_{k=0}^{\lfloor n/2 \rfloor} \frac{x^{n-2k}}{(n-2k)!} \frac{t^k}{k!}, \quad n = 0, 1, \dots.$$

In [2, Th. 5.3, p. 231] it was shown that the series $\sum_{n=0}^\infty a_n v_n(x, t)$ converges to a C^2 -solution of the heat equation in the strip $|t| < \sigma$ where

$$(2.2) \quad \sigma = (\limsup |a_n|^{2/n} (2n/e))^{-1}$$

and that this strip is the largest open strip of convergence of the series. One easily shows that sequences $\{a_n\}$ satisfying

$$\limsup |a_n|^{2/n} (2n/e) = 1$$

have their classification numbers in $[0, 1]$.

We will make repeated use of the following bounds which appear in [4] by S. Täcklind. Assume $u(x, t)$ is continuous on the rectangle $R = \{(x, t): |x| \leq \mathcal{L}, 0 \leq t \leq T\}$, is a C^2 -solution to the heat equation in the interior of R , and μ is an upper bound for $|u(x, t)|$ on R ; then $u(x, t)$ is in class C^∞ on the interior of R and for $n = 2, 3, \dots, |x| < \mathcal{L}$, and $0 < t \leq T$

$$(2.3) \quad \left| \frac{\partial^n u}{\partial x^n}(x, t) \right| \leq \frac{\mu}{2\sqrt{\pi}} \frac{2^{(n+3)/2}}{t^{n/2}} \Gamma((n+1)/2) \\ + \frac{\mu}{\sqrt{\pi}} \left(\frac{\pi}{2} \right)^{5/2} \frac{2^{3n/2}}{(\mathcal{L} - |x|)^n} \Gamma(n+1).$$

3. THEOREM 1. Let $\{X_n\}_{n=0}^\infty$ be a sequence of symmetric independent random variables defined on the complete probability space (Ω, \mathcal{F}, P) and satisfying

(i) there exists a number M such that

$$\int_{\Omega} |X_n(\omega)|^2 dP(\omega) \leq M \text{ for } n = 0, 1, \dots, \text{ and}$$

(ii) there exists $N > 0$ such that

$$N \leq \int_{\Omega} |X_n(\omega)| dP(\omega), n = 0, 1, \dots.$$

Assume $\{a_n\}$ satisfies $\limsup |a_n|^{2/n} (2n/e) = 1$ and has its classification number in $[0, 1/2)$. Then for almost every ω in Ω the lines $t = 1$ and $t = -1$ form the natural boundary for

$$u_{\omega}(x, t) = \sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t).$$

Proof. Letting $\Omega' = \{\omega \in \Omega: \sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t) \text{ converges in the strip } |t| < 1\}$, we will first show $P(\Omega') = 1$. Clearly

$$[\limsup |X_n|^{2/n} \leq 1] \supseteq \bigcup_{k=1}^{\infty} \bigcap_{n=k}^{\infty} [|X_n| \leq nM^{1/2}]$$

and by the Borel-Cantelli Lemma the last set has probability 1 since $P[|X_n| > nM^{1/2}] \leq 1/n^2$ from (i). Hence

$$P\{\omega: \limsup |X_n(\omega) a_n|^{2/n} (2n/e) \leq 1\} = 1$$

which by (2.2) shows $P(\Omega') = 1$.

The following fact is essentially a merger of Lemma 1 from [3] and a special case of Lemma 2 from [3]. There exist numbers A in $(0, 1)$ and $B > 0$ with the following property: for $E \in \mathcal{F}$ with $P(E) > A$ there is a positive integer n_0 such that for $n \geq n_0$, every sequence $\{c_j\}_{j=0}^{\infty}$ of real numbers, and $k \geq 1$ we have

$$(3.1) \quad \sum_{j=n}^{n+k} c_j^2 \leq B \int_E \left\{ \sum_{j=n}^{n+k} c_j X_j(\omega) \right\}^2 dP(\omega).$$

We will show that for almost every ω the line $t = -1$ is part of the natural boundary for u_{ω} and will use this in the proof for the line $t = 1$.

Assume it is false that for a.e. ω in Ω the line $t = -1$ is part of the natural boundary for u_{ω} . The first part of the argument we give in order to obtain a contradiction is analogous to parts of the proof of Theorem 1 in [3] by V. L. Shapiro. We will employ (2.3), (3.1), and an asymptotic estimate for heat polynomials from [2] in

order to obtain conditions on the sequence $\{a_n\}$ which contradict the fact that the classification number of $\{a_n\}$ is in $[0, 1/2)$.

Let $E = \{\omega \in \Omega': t = -1 \text{ is not part of the natural boundary for } u_\omega\}$. Then either (i) $E \notin \mathcal{F}$, or (ii) $E \in \mathcal{F}$ and $P(E) > 0$. Using the fact that the real line is separable and the countable additivity of the probability P , it follows that there exist a real number x_0 and $\rho_0 > 0$ such that $E_1 = \{\omega \in E: \text{there is a } C^2\text{-solution to the heat equation in } S(x_0, -1; \rho_0) \text{ which agrees with } u_\omega \text{ where they are both defined}\}$ satisfies either (i) $E_1 \notin \mathcal{F}$, or (ii) $E_1 \in \mathcal{F}$ and $P(E_1) > 0$. For $i = 1, 2, \dots$ define

$$E_{2,i} = \left\{ \omega \in \Omega': \left| \frac{\partial^m u_\omega}{\partial x^m}(x, t) \right| \leq i^m m^m \text{ for } (x, t) \text{ in } S\left(x_0, -1; \frac{\rho_0}{2}\right), \right. \\ \left. |t| < 1, \text{ and } m = i, i+1, \dots \right\}$$

and let $E_2 = \bigcup_{i=1}^{\infty} E_{2,i}$. E_2 is in the tail σ -field generated by the independent X_n 's. From (2.3) it follows that $E_1 \subseteq E_2$. By Kolmogorov's zero-one law $P(E_2) = 1$. Let A and B be as in (3.1). Take i_0 sufficiently large that $P(E_{2,i_0}) > A$ and let n_0 correspond to E_{2,i_0} as in (3.1). Now let $m \geq \max\{n_0, i_0\}$ and let (x, t) be in $S(x_0, -1; \rho_0/2)$ with $|t| < 1$. Then by (3.1) for $k = 1, 2, \dots$

$$\sum_{n=m}^{m+k} \left[\frac{n!}{(n-m)!} a_n v_{n-m}(x, t) \right]^2 \\ \leq B \int_{E_{2,i_0}} \left[\sum_{n=m}^{m+k} \frac{n!}{(n-m)!} a_n v_{n-m}(x, t) X_n(\omega) \right]^2 dP(\omega).$$

Making use of the independence and symmetry of the random variables and of condition (i) we see that the integrand of the last integral is Cauchy in the variable k in $L^1(\Omega)$ and thus in $L^1(E_{2,i_0})$. Hence

$$\sum_{n=m}^{\infty} \left[\frac{n!}{(n-m)!} a_n v_{n-m}(x, t) \right]^2 \\ \leq B \int_{E_{2,i_0}} \left[\sum_{n=m}^{\infty} \frac{n!}{(n-m)!} a_n v_{n-m}(x, t) X_n(\omega) \right]^2 dP(\omega) \\ = B \int_{E_{2,i_0}} \left| \frac{\partial^m u_\omega}{\partial x^m}(x, t) \right|^2 dP(\omega) \leq B i_0^{2m} m^{2m}$$

with the last inequality following from the definition of E_{2,i_0} . We conclude that for every $m \geq \max\{n_0, i_0\}$, every $n \geq m$, and every (x, t) in $S(x_0, -1; \rho_0/2)$ with $|t| < 1$; we have

$$(3.2) \quad \frac{n!}{(n-m)!} |a_n| |v_{n-m}(x, t)| \leq B^{1/2} i_0^m m^m.$$

It follows from Theorem 3.1 of [2] that there exist numbers A and l_0 such that for $n \geq l_0$

$$\sup_{|x-x_0| < \rho_0/2} |v_n(x, -1)| \geq A[2n/e]^{n/2}.$$

Thus from (3.2) we have for $n > m + l_0 > m \geq \max\{n_0, i_0\}$

$$|a_n| \frac{n!}{(n-m)!} A[2(n-m)/e]^{(n-m)/2} \leq B^{1/2} i_0^m m^m.$$

Employing Stirling's theorem we see there is a number C such that for $n > m + l_0 > m \geq \max\{n_0, i_0\}$

$$(3.3) \quad |a_n| (2n/e)^{n/2} \leq \left[\frac{Cm}{\sqrt{n-m}} \right]^m \cdot ((n-m)/n)^{(n+1)/2}.$$

Let r be a number which is strictly greater than the classification number of $\{a_n\}$ and strictly less than $1/2$. Let m be related to n by $m = [4n^r] + 1$ where the brackets denote the greatest integer function. Then from (3.3), for sufficiently large n ,

$$(3.4) \quad |a_n| (2n/e)^{n/2} \leq (1 - 4/n^{1-r})^{(n^{1-r/4}) \cdot 2 \cdot n^r}.$$

For large enough n , $(1 - 4/n^{1-r})^{(n^{1-r/4}) \cdot 2} \leq 1/e$ and thus from (3.4) we have for such n , $|a_n| (2n/e)^{n/2} \leq 1/e^{n^r}$. Hence $\{a_n\} \in E_r$ which contradicts the fact that r is strictly greater than the classification number of $\{a_n\}$ and concludes the proof for the line $t = -1$.

For the last part of the proof we find it convenient to introduce the probability space $(R^\omega, \mathcal{A}', \mu')$ which we now describe.

$$R^\omega = \prod_{n=0}^{\infty} R_n$$

where each R_n is the set of real numbers. Let \mathcal{A}_0 be the field of all subsets of R^ω of the form $B \times (\prod_{n=n_0+1}^{\infty} R_n)$ where n_0 is a positive integer and B is a Borel set in $\prod_{n=0}^{n_0} R_n$. Let \mathcal{A} be the σ -field generated by \mathcal{A}_0 . Let μ be the probability on (R^ω, \mathcal{A}) which is induced by the X_n 's. Then $(R^\omega, \mathcal{A}', \mu')$ is the completion of $(R^\omega, \mathcal{A}, \mu)$.

Let $\{\eta_i\}_{i=0}^{\infty}$ be a sequence of ± 1 's. Define $T: R^\omega \rightarrow R^\omega$ by

$$T((\xi_0, \xi_1, \dots)) = (\eta_0 \xi_0, \eta_1 \xi_1, \dots).$$

Notice that

$$\begin{aligned} \mu \left(\prod_{n=0}^{n_0} (a_n, b_n] \times \prod_{n=n_0+1}^{\infty} R_n \right) &= \prod_{n=0}^{n_0} P[X_n \in (a_n, b_n]] \\ &= \prod_{n=0}^{n_0} P[X_n \in \eta_n(a_n, b_n)] = \mu \left(T \left(\prod_{n=0}^{n_0} (a_n, b_n] \times \prod_{n=n_0+1}^{\infty} R_n \right) \right) \end{aligned}$$

where we have used both the independence and symmetry of the X_n 's. From this it follows that for $A \in \mathcal{N}'$, $\mu'(A) = \mu'(T(A))$. We will make use of this fact twice in the remainder of this proof.

To finish the proof it suffices to show that for a.e. $p \in R^\omega$ the line $t = 1$ is part of the natural boundary for

$$u_p(x, t) = \sum_{n=0}^{\infty} \pi_n(p) a_n v_n(x, t)$$

where the π_n 's are the projection random variables. Suppose this is false. From the first paragraph of the present proof we know $R^{\omega'} = \{p \in R^\omega: \sum_{n=0}^{\infty} \pi_n(p) a_n v_n(x, t) \text{ converges in } |t| < 1\}$ has μ' -measure 1. Now let $F' = \{p \in R^{\omega'}: t = 1 \text{ is not part of the natural boundary for } u_p\}$. Then either (i) $F' \in \mathcal{N}'$, or (ii) $F' \in \mathcal{N}'$ and $\mu'(F') > 0$. It follows that there exist numbers a, b, ρ with $a < b$ and $\rho > 0$ such that $F_1 = \{p \in R^{\omega'}: \text{there is a function } v_p(x, t) \text{ which is continuous on } a \leq x \leq b, 0 \leq t \leq 1 + \rho; \text{ is a } C^2\text{-solution to the heat equation for } a < x < b, 0 < t < 1 + \rho; \text{ and agrees with } u_p(x, t) \text{ in } a \leq x \leq b, 0 \leq t < 1\}$ satisfies either (i) $F_1 \in \mathcal{N}'$, or (ii) $F_1 \in \mathcal{N}'$ and $\mu'(F_1) > 0$. But $F_1 = \{p \in R^{\omega'}: \lim_{t \uparrow 1} u_p(a, t) \text{ and } \lim_{t \uparrow 1} u_p(b, t) \text{ both exist}\}$. F_1 is in the tail σ -field generated by the independent π_n 's. From the zero-one law, $\mu'(F_1) = 1$.

Either $a \neq 0$ or $b \neq 0$ and for definiteness we assume $a \neq 0$. Then $F_2 = \{p \in R^{\omega'}: \lim_{t \uparrow 1} u_p(a, t) \text{ exists}\}$ has $\mu'(F_2) = 1$. Let $T: R^\omega \rightarrow R^\omega$ be defined by $T((\xi_0, \xi_1, \dots)) = (\xi_0, -\xi_1, \xi_2, -\xi_3, \dots)$. By our earlier comments concerning such mappings we have $\mu'(F_2 \cap T(F_2)) = 1$. For $p \in R^{\omega'}$ and $|t| < 1$ one checks that $u_{T(p)}(-a, t) = u_p(a, t)$. Hence for $p \in F_2 \cap T(F_2)$, $\lim_{t \uparrow 1} u_p(-a, t)$ and $\lim_{t \uparrow 1} u_p(a, t)$ both exist. Thus for $p \in F_2 \cap T(F_2)$ there is a function $w_p(x, t)$ which is continuous in $|x| \leq a, 0 \leq t \leq 2$; is a C^2 -solution to the heat equation in $|x| < a, 0 < t < 2$; and agrees with u_p in $|x| \leq a, 0 \leq t < 1$. For $p \in F_2 \cap T(F_2)$ and $0 \leq t \leq 2$ let $\phi_p(t) = w_p(0, t)$ and $\psi_p(t) = (\partial w_p / \partial x)(0, t)$. Then, employing (2.3), we see that ϕ_p and ψ_p are in class $C\{(2n)!\}$ on $[0, 2]$ (a function f is in class $C\{(2n)!\}$ on an interval I if f is in class C^∞ on I and there exist constants β and B such that for every t in I , $|f^{(n)}(t)| \leq \beta B^n (2n)!, n = 0, 1, \dots$).

Now let $T': R^\omega \rightarrow R^\omega$ be defined by

$$T'((\xi_0, \xi_1, \dots)) = (\xi_0, \xi_1, -\xi_2, -\xi_3, \xi_4, \xi_5, -\xi_6, -\xi_7, \dots).$$

Then for $p \in R^{\omega'}$ and $|t| < 1$, $u_p(0, t) = u_{T'(p)}(0, -t)$ and

$$\frac{\partial u_p}{\partial x}(0, t) = \frac{\partial u_{T'(p)}}{\partial x}(0, -t).$$

For p in the almost sure set $T'(F_2 \cap T(F_2))$ we have $T'(p) \in F_2 \cap T(F_2)$ and we define ϕ'_p and ψ'_p on $[-2, 0]$ by $\phi'_p(t) = \phi_{T'(p)}(-t)$ and

$$\psi'_p(t) = \psi_{T'(p)}(-t)$$

thereby obtaining class $C\{(2n)!\}$ extensions of $u_p(0, t)$ and $(\partial u_p / \partial x)(0, t)$ on $[-1, 0]$. Thus for $p \in T'(F_2 \cap T(F_2))$

$$u'_p(x, t) = \sum_{n=0}^{\infty} \frac{\phi'_p{}^{(n)}(t)x^{2n}}{(2n)!} + \sum_{n=0}^{\infty} \frac{\psi'_p{}^{(n)}(t)x^{2n+1}}{(2n+1)!}$$

provides a solution to the heat equation which is a C^2 -extension of u_p into some rectangle $|x| < r, -2 < t < 0$ which contradicts the first part of the proof.

4. THEOREM 2. *Let $\{X_n\}$ be a sequence of independent random variables over a probability space (Ω, \mathcal{F}, P) which satisfies (i) and (ii) of Theorem 1. Assume $\{a_n\}$ satisfies $\limsup |a_n|^{2/n}(2n/e) = 1$ and has its classification number in $(1/2, 1]$. Then for almost every ω in Ω the following holds: $|t| < 1$ is the strip of convergence of $\sum_{n=0}^{\infty} X_n(\omega)a_nv_n(x, t)$ which for every $\mathcal{L} > 0$ can be extended as a C^2 -solution of the heat equation into $\{|t| < 1\} \cup \{|x| < \mathcal{L}\}$.*

Proof. We will first show for almost every ω in Ω that $|t| < 1$ is the strip of convergence of $\sum_{n=0}^{\infty} X_n(\omega)a_nv_n(x, t)$. By (2.2) we must show that almost surely $\limsup |X_n(\omega)a_n|^{2/n}(2n/e) = 1$. The argument given in the first part of the proof of Theorem 1 shows that almost surely the last limit superior does not exceed 1. Let $\{n_j\}$ be a strictly increasing sequence of positive integers such that

$$\lim_{j \rightarrow \infty} |a_{n_j}|^{2/n_j}(2n_j/e) = 1.$$

Then $\limsup |X_n(\omega)a_n|^{2/n}(2n/e) \geq \limsup_{j \rightarrow \infty} |X_{n_j}(\omega)a_{n_j}|^{2/n_j}(2n_j/e) \geq \limsup_{j \rightarrow \infty} |X_{n_j}(\omega)|^{2/n_j}$ which by the zero-one law is equal to some number c almost surely. Suppose $c < 1$. Then $X_{n_j} \rightarrow 0$ almost surely. By (ii) for $A > 0$ and $j = 0, 1, \dots$

$$N \leq \int_{\{|X_{n_j}| \leq A\}} |X_{n_j}(\omega)| dP(\omega) + A^{-1} \int_{\{|X_{n_j}| > A\}} |X_{n_j}(\omega)|^2 dP(\omega).$$

By the Lebesgue dominated convergence theorem the next to the last integral tends to 0 as j tends to ∞ . From (i) the last term is uniformly bounded by $A^{-1}M$. Thus for every $A > 0, N \leq A^{-1}M$ which is a contradiction. We conclude that $c \geq 1$. Thus almost surely

$$\limsup |X_n(\omega)a_n|^{2/n}(2n/e) \geq 1$$

which concludes the proof that almost surely this limit superior is 1.

In order to establish Theorem 2 for the line $t = 1$ we first construct a function which is C^∞ on the closed strip $|t| \leq 1$ and has a heat polynomial expansion in $|t| < 1$. Let r be a number which is strictly greater than $1/2$ and strictly less than the classification num-

ber of $\{a_n\}$. For $n = 0, 1, \dots$ define $\alpha_n = (2n)e^{-n^r}$. Define f on $[-1, 1]$ by $f(t) = \sum_{k=0}^{\infty} \alpha_k t^k$. We will show this definition makes sense and obtain some bounds on the derivatives of f .

Let n be a nonnegative integer. Differentiating $\sum_{k=0}^{\infty} \alpha_k t^k$ term by term n times yields $\sum_{k=n}^{\infty} k!/(k-n)! \alpha_k t^{k-n}$. For $|t| \leq 1$ the k^{th} term of this series is dominated by $2 k^{n+1} e^{-k^r}$. One checks that

$$g_n(x) = x^{n+1} e^{-x^r}$$

is increasing on $(0, (n+1/r)^{1/r})$ and decreasing on $((n+1/r)^{1/r}, \infty)$. Hence

$$\sum_{k=n}^{\infty} k^{n+1} e^{-k^r} \leq \int_n^{\infty} g_n(x) dx + g_n\left(\left(\frac{n+1}{r}\right)^{1/r}\right) \leq 3\Gamma((n+2)/r)/r.$$

We conclude that f is a C^∞ -function with $|f^{(n)}(t)| \leq 6\Gamma((n+2)/r)/r$ for $n = 0, 1, \dots$ and $|t| \leq 1$.

Now define

$$(4.1) \quad u(x, t) = \sum_{n=0}^{\infty} \frac{f^{(n)}(t)x^{2n}}{(2n)!} + \sum_{n=0}^{\infty} \frac{f^{(n+1)}(t)x^{2n+1}}{(2n+1)!}.$$

Because of the bounds obtained in the preceding paragraph it can be shown that the series of (4.1) can be differentiated term by term and that $u(x, t)$ is a C^∞ -solution to the heat equation in the closed strip $|t| \leq 1$. Since both $u(0, t)$ and $\partial u/\partial x(0, t)$, as functions of t on $(-1, 1)$, are given by their Maclaurin expansions, u has a heat polynomial expansion in $|t| < 1$ (see [5]). Thus

$$(4.2) \quad \begin{aligned} u(x, t) &= \sum_{n=0}^{\infty} b_n v_n(x, t), \\ b_{2n} &= f^{(n)}(0)/(2n)!, \\ b_{2n+1} &= f^{(n+1)}(0)/(2n+1)!. \end{aligned}$$

According to the first paragraph of the proof of Theorem 1, $\bigcup_{k=1}^{\infty} \bigcap_{n=k}^{\infty} [|X_n| \leq nM^{1/2}]$ has probability 1. Let ω be in this almost sure set. Let k_0 be a positive integer such that for $n \geq k_0$, $|X_n(\omega)| \leq nM^{1/2}$. Since r is less than the classification number of $\{a_n\}$, there is a number K such that $|a_n|(2n/e)^{n/2} \leq Ke^{-n^r}$, $n = 1, 2, \dots$. Using Stirling's theorem we have for $2n \geq k_0$

$$b_{2n}(4n/e)^n \geq |X_{2n}(\omega)a_{2n}| (4n/e)^n (1/2)^{3/2}/KM^{1/2}.$$

Similarly for $2n+1 \geq k_0$

$$b_{2n+1}(2(2n+1)/e)^{(2n+1)/2} \geq |X_{2n+1}(\omega)a_{2n+1}| (2(n+1)/e)^{(2n+1)/2} e^{-1/2}/KM^{1/2}.$$

Letting $K' = K(Me)^{1/2}$ we have

$$|X_n(\omega)a_n| \leq K'b_n \text{ for } n \geq k_0.$$

Let $\mathcal{L} > 0$. Then for $0 < t < 1$ we have

$$\begin{aligned} \left| \frac{\partial}{\partial t} \sum_{n=k_0}^{\infty} X_n(\omega)a_nv_n(\pm \mathcal{L}, t) \right| &= K' \sum_{n=k_0}^{\infty} b_nn(n-1) |v_{n-2}(\pm \mathcal{L}, t)| \\ &\leq K' \sum_{n=k_0}^{\infty} b_nn(n-1)v_{n-2}(\mathcal{L}, t) \leq K' \frac{\partial u}{\partial t}(\mathcal{L}, 1) < \infty. \end{aligned}$$

Thus $\lim_{t \uparrow 1} \sum_{n=k_0}^{\infty} X_n(\omega)a_nv_n(\pm \mathcal{L}, t)$ both exist as is easily seen from the mean value theorem and the Cauchy criterion. Hence we can obtain an extension of $\sum_{n=0}^{\infty} X_n(\omega)a_nv_n(x, t)$ into

$$\{(x, t): |t| < 1\} \cup \{(x, t): |x| < \mathcal{L}, 0 < t\}$$

which is a C^2 -solution of the heat equation. (Notice at this point that we can also obtain an extension which is a bounded C^2 -solution in $\{(x, t): |x| < \mathcal{L}, 0 \leq t\}$.) Since ω was from the almost sure set

$$\bigcap_{k=1}^{\infty} \bigcap_{n=k}^{\infty} [|X_n| \leq nM^{1/2}],$$

this establishes the result for the line $t = 1$.

We now turn to the line $t = -1$. Define $\{Y_n\}_{n=0}^{\infty}$ on Ω by $Y_{2n} = (-1)^n X_{2n}$ and $Y_{2n+1} = (-1)^n X_{2n+1}$. Then, applying the first part of the proof, there is a set F in \mathcal{F} with $P(F) = 1$ such that for ω in F and $\mathcal{L} > 0$ the solution $v_{\omega}(x, t) = \sum_{n=0}^{\infty} Y_n(\omega)a_nv_n(x, t)$ can be extended into $\{|t| < 1\} \cup \{|x| < \mathcal{L} \text{ and } 0 < t\}$ so as to be a bounded C^2 -solution of the heat equation in $\{(x, t): |x| < \mathcal{L} \text{ and } 0 < t\}$. One easily checks that for ω in F ,

$$\sum_{n=0}^{\infty} X_n(\omega)a_nv_n(0, t) = \sum_{n=0}^{\infty} Y_n(\omega)a_nv_n(0, -t)$$

and $\sum_{n=1}^{\infty} X_n(\omega)a_nv_{n-1}(0, t) = \sum_{n=1}^{\infty} Y_n(\omega)a_nv_{n-1}(0, -t)$. Using these facts and (2.3) we see that for ω in F and $\mathcal{L} > 0$ the functions $\phi(t) = \sum_{n=0}^{\infty} X_n(\omega)a_nv_n(0, t)$ and $\psi(t) = \sum_{n=1}^{\infty} X_n(\omega)a_nv_{n-1}(0, t)$ on $(-1, 1)$ possess sufficiently well behaved extensions ϕ' and ψ' to $(-\infty, 1)$ that

$$\sum_{n=0}^{\infty} \frac{\phi'^{(n)}(t)x^{2n}}{(2n)!} + \sum_{n=0}^{\infty} \frac{\psi'^{(n)}(t)x^{2n+1}}{(2n+1)!}$$

is an extension of $\sum_{n=0}^{\infty} X_n(\omega)a_nv_n(x, t)$ in $|t| < 1$ to

$$\{(x, t): |t| < 1\} \cup \{(x, t): |x| < \mathcal{L} \text{ and } -\infty < t < 1\}.$$

5. Examples. The first example will show that our two theorems are best possible with respect to the allowable values of the classification number.

EXAMPLE 1. We will take $[0, 1]$ with Lebesgue measure as the probability space and the sequence of Rademacher functions, $\{\phi_n\}_{n=0}^\infty$, for the random variables.

For $k = 0, 1, \dots$ define $\alpha_k = e^{-\sqrt{k}}$. Then, as in the proof of Theorem 2, defining f on $[-1, 1]$ by $f(t) = \sum_{k=0}^\infty \alpha_k t^k$ yields a C^∞ -function whose n^{th} derivative on $[-1, 1]$ is bounded in absolute value by $6F(2(2n+1))$. In the strip $|t| < 1$ define $u(x, t) = \sum_{n=0}^\infty (f^{(n)}(t)x^{2n})/(2n)!$. To see that this definition makes sense and that term by term partial differentiation is permitted, we note that for every closed interval $I \subseteq (-1, 1)$, f is in class $C\{n!\}$ on I . Because of the bounds on the derivatives of f we see from the defining series for u that u may be extended as a C^∞ -solution of the heat equation to

$$\{|t| < 1\} \cup \{(x, 1): |x| < 1\}.$$

Since $u(0, t)$ and $\partial u/\partial x(0, t)$ are both given by their Maclaurin expansions in $|t| < 1$, u possesses a heat polynomial expansion in the strip $|t| < 1$ (see [5]). Thus for $|t| < 1$, $u(x, t) = \sum_{n=0}^\infty a_n v_n(x, t)$; $a_{2n} = (e^{-\sqrt{n}}n!)/(2n)!$, $a_{2n+1} = 0$. One checks that $\limsup |a_n|^{2/n}(2n/e) = 1$. Also it is easily seen that $\lim |a_{2n}|(4n/e)^n e^{\sqrt{2n}} = \infty$ which implies $\{a_n\} \notin E_{1/2}$ and thus the classification number of $\{a_n\}$ is in $[0, 1/2]$. As in the proof of Theorem 2, $\lim_{t \uparrow 1} u_\omega(\pm 1/2, t)$ both exist for every ω in $[0, 1]$. Thus for every $\omega \in [0, 1]$ the line $t = 1$ is not part of the natural boundary for $u_\omega(x, t)$. Using Theorem 1, we conclude that the classification number of $\{a_n\}$ is $1/2$ and that in Theorem 1 we cannot replace $[0, 1/2)$ by $[0, 1/2]$ as the allowable range for the classification number.

We will next show that the conclusion of Theorem 2 does not hold for $\sum_{n=0}^\infty \phi_n(\omega)a_n v_n(x, t)$. Assume there is a set A in $[0, 1]$ with $m(A) = 1$ such that for each ω in A no interval of the line $t = 1$ is part of the natural boundary for $u_\omega(x, t)$. Thus for ω in A , $g_\omega(x) = \lim_{t \uparrow 1} u_\omega(x, t)$ is well defined and is the restriction of an entire function to the real axis (this last assertion can be seen by employing (2.3)). Thus for ω in A , $\limsup (|g_\omega^{(n)}(0)|/n!)^{1/n} = 0$. For ω in A , $|g_\omega^{(2n+1)}(0)| = 0$ and $|g_\omega^{(2n)}(0)| = |\sum_{k=2n}^\infty \phi_k(\omega)a_k(k!/(k-2n)!)v_{k-2n}(0, 1)| = |\sum_{k=n}^\infty \phi_{2k}(\omega)(k!/(k-n)!)e^{-\sqrt{k}}|$. Thus for ω in A ,

$$\limsup \left[\frac{\left| \sum_{k=n}^\infty \phi_{2k}(\omega) \frac{k!}{(k-n)!} e^{-\sqrt{k}} \right|}{(2n)!} \right]^{1/n} = 0.$$

Let $\delta > 0$. For $m = 0, 1, \dots$ let

$$\begin{aligned} F_m &= \left\{ \omega \in A : \left(\left| \sum_{k=n}^\infty \phi_{2k}(\omega) \frac{k!}{(k-n)!} e^{-\sqrt{k}} \right| / (2n)! \right)^{1/n} \right. \\ &\quad \left. \leq \delta \text{ for } n = m, m+1, \dots \right\} \end{aligned}$$

and note $F_m \uparrow A$. Let A and B be two numbers associated with the sequence $\{\phi_{2m}\}_{m=0}^\infty$ as in (3.1). Let m_0 be sufficiently large that $m(F_{m_0}) > A$. Let n_0 be an integer larger than m_0 with n_0 corresponding to F_{m_0} as in (3.1). Thus for $n \geq n_0$ and $k \geq 1$

$$(5.1) \quad \sum_{j=n}^{n+k} \left[\frac{j!}{(j-n)!} e^{-\sqrt{j}} \right]^2 \leq B \int_{F_{m_0}} \left(\sum_{j=n}^{n+k} \phi_{2j}(\omega) \frac{j!}{(j-n)!} e^{-\sqrt{j}} \right)^2 d\mu(\omega).$$

As in the proof of Theorem 1, letting k tend to ∞ yields (5.1) with $n+k$ replaced by ∞ . Using the definition of F_{m_0} , we have

$$\sum_{j=n}^{\infty} \left[\frac{j!}{(j-n)!} e^{-\sqrt{j}} \right]^2 \leq B((2n)! \delta^n)^2,$$

for $n \geq n_0$. From this we conclude that

$$\limsup \left[\frac{\left[\sum_{k=n}^{\infty} \left(\frac{k!}{(k-n)!} e^{-\sqrt{k}} \right)^2 \right]^{1/2}}{(2n)!} \right]^{1/n} = 0.$$

On the other hand, letting L denote this last limit superior, we have

$$L \geq \limsup \left[\frac{\left[\sum_{k=n}^{\infty} (k-n)^{2n} \exp(-2\sqrt{k-n}) \exp(-(2\sqrt{k} - 2\sqrt{k-n})) \right]^{1/2}}{(2n)!} \right]^{1/n}.$$

But $\exp(-(2\sqrt{k} - 2\sqrt{k-n})) \geq e^{-2\sqrt{n}}$ for $k \geq n$ and $\lim (e^{-\sqrt{n}})^{1/n} = 1$. Hence $L \geq \limsup ((\sum_{k=0}^{\infty} k^{2n} e^{-2\sqrt{k}})^{1/2} / (2n)!)^{1/n}$. Define h_n on $(0, \infty)$ by $h_n(x) = x^{2n} e^{-2\sqrt{x}}$. One checks that h_n is increasing on $(0, (2n)^2)$ and decreasing on $((2n)^2, \infty)$. Thus $\sum_{k=0}^{\infty} k^{2n} e^{-2\sqrt{k}} \geq \int_0^{\infty} h_n(x) dx - h_n((2n)^2) = (\Gamma(4n+2) - 2(4n)^{4n} e^{-4n}) / (2 \cdot 4^{2n})$. Thus

$$L \geq \frac{1}{4} \limsup \left[\left(\frac{\Gamma(4n+2)}{(4n)!} - \frac{2(4n)^{4n} e^{-4n}}{(4n)!} \right) ((4n)! / ((2n)!)^2) \right]^{1/2n} > 0.$$

This is a contradiction. Hence in Theorem 2 we cannot replace $(1/2, 1]$ by $[1/2, 1]$ as the allowable range for the classification number.

The next example shows that in Theorem 1 we cannot omit the symmetry of the random variables.

EXAMPLE 2. Let $k(x, t) = e^{-x^2/4t} / (4\pi t)^{1/2}$ for $t > 0$ and define

$$u(x, t) = k(x, t + 1)$$

in the strip $|t| < 1$. Then [2, Th. 4.2, p. 227]

$$u(x, t) = (4\pi)^{-1/2} \sum_{n=0}^{\infty} \frac{(-1)^n}{n! 4^n} v_{2n}(x, t).$$

Let $\{a_n\}_{n=0}^{\infty}$ be defined by $a_{2n} = (-1)^n/n! 4^n$ and $a_{2n+1} = 0$. One easily checks that $\limsup |a_n|^{2/n} (2n/e) = 1$ and that the classification number of $\{a_n\} = 0$. Let $X_n = 1, n = 0, 1, \dots$ on some complete probability space. Then for every ω, u_ω can be continued above the line $t = 1$.

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Received July 23, 1968. This paper is based on part of the author's doctoral dissertation at the University of California at Riverside which was written under the direction of Victor L. Shapiro. Research sponsored by the Air Force Office of Scientific Research, Office of Aerospace Research, United States Air Force, under AFOSR Grant No. AF-AFOSR 694-66.

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