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The subject of this paper is the study of the existence of a pendent drop. We carry out this study in full generality by exploiting the local minima of a suitable functional, chosen to represent the energy of the drop.

If we denote by $E \subset \mathbf{R}^{n+1}$ a liquid drop hanging from the fixed horizontal reference plane $\{t = 0\}$, then we can write the global energy of that configuration in the following way:

$$(0.1) \quad \mathcal{F}(E) = \int_{t < 0} |D\varphi_E| + \nu \int_{t=0} \varphi_E dH_n + \kappa \int_{t < 0} t \varphi_E(x, t) dx dt$$

Here, the first integral is the measure of that part of the boundary of E lying in the half-space $\{t < 0\}$. Physically, it corresponds to the energy due to surface tension. The second integral, proportional to the measure of the boundary of E contacting the horizontal plane, represents the energy given by the attraction between the liquid and the plane itself, while the third one corresponds to the gravitational energy.

Exact definitions of these objects will be given in the next section.

The constant ν is determined experimentally, depending on the materials in the liquid-solid-vapor interface; physically, it represents the cosine of the angle between the exterior directed normal to the liquid surface, along the intersection with the contact plane $\{t = 0\}$, and the positive (vertical) t -direction. The constant $\kappa \geq 0$ takes into account the gravitational acceleration, and is referred to as the "capillary constant" (see e.g., [12]).

We have to minimize the functional (0.1) among the sets of *finite perimeter* (in the sense of De Giorgi; see [7] or [11], where the equivalent notion of *mass* is used), having prescribed V volume. It is clear that we cannot expect a finite lower bound for (0.1) in such a class, as physical considerations may suggest. Indeed, a pendent drop is just a *local minimum* of the energy functional.

In order to prove the existence of such a local minimum, we introduce a *ground floor* (i.e., a plane $\{t = T\}$ with a suitable $T < 0$) and minimize the energy among those configurations E lying between floor and ceiling: for small gravity, we can prove that such minima do not reach the floor. We do this by observing that as gravity

decreases to zero, the corresponding minima approach—in the sense of a *good* convergence—the solution in the absence of gravity, which is part of a sphere.

We remark that when it is possible to describe the part of the boundary of E , lying below the contact plane, as the graph of a function $u \in C^2(\Omega)$, Ω open in \mathbf{R}^n , then we can write the functional (0.1) in the equivalent way

$$(0.2) \quad \mathcal{F}(u) = \int_{\Omega} \sqrt{1 + |Du|^2} + \nu |\Omega| - \frac{\kappa}{2} \int_{\Omega} u^2 dx .$$

It follows that a minimum of (0.2) satisfies the *Euler equation*

$$(0.3) \quad \operatorname{div} Tu = -\kappa u + \lambda, \quad Tu = Du \cdot (1 + |Du|^2)^{-1/2}$$

in which λ is a constant ensuing from the volume constraint.

In the physical case $n = 2$, considering only rotationally symmetric solutions (i.e., solutions symmetric about the vertical t -axis), equation (0.3) takes the simpler form:

$$(0.4) \quad \frac{1}{r} \left(\frac{ru'(r)}{\sqrt{1 + u'^2(r)}} \right)' = -\kappa u(r) + \lambda$$

with $u = u(r)$.

The behavior of a solution of (0.4), in its dependence on the initial value $u(0) = u_0 < 0$, has been studied extensively by P. Concus and R. Finn in a series of papers ([12, 2, 3, 4]; see also [28, 29]). We refer to [4] for a recent detailed exposition on this argument.

Several interesting results dealing with equation (0.4) have also been obtained from the computational point of view; see [12] and the references cited there for a general account.

The question of existence of a drop suspended from an arbitrary aperture, consisting of a simple closed curve in \mathbf{R}^3 , was studied by H. C. Wente in [34]. The method employed in that paper still involves minimization of a suitable functional, however by means of complex-variable techniques, and gives an affirmative answer for sufficiently small gravitational field; it differs completely from our method, which relies on De Giorgi's *Theory of Perimeter* and works in any dimension.

In this framework, developed in [5], [6], [7] (see also [1], [19] for a detailed treatment of the subject), several capillarity problems have been recently solved; first of all, the problem of existence of *equilibrium surfaces in a capillary tube*, whose solution was obtained by M. Emmer in [9] (see also [13], [32] for extensions of Emmer's result to capillary tube of general cross-section, and [27], [14], [15], [30], [31], [33], [17] for interior analyticity and boundary regularity of Emmer's solution). E. Gonzalez and I. Tamanini ([20], [21], [22])

studied subsequently the problem of a *sessile drop* (i.e., a liquid drop sitting on a flat plate), proving existence, regularity and convexity of the equilibrium configuration. Further contributions along this line were obtained by U. Massari and L. Pepe in [25], [26].

Definitions and notations to be used in the sequel are given in the first section. Section 2 is devoted to the *gravity-free problem*, that is to the study of the minima of (0.1) with $\kappa = 0$, while ε -gravity problems are studied in §3. Sections 4 and 5 deal with some properties of ε -solutions that are used in the last two sections to present the conclusive existence and regularity results.

For convenience of the reader, some questions related to the existence of multipliers have been quoted in the Appendix.

We wish to thank P. Concus for many helpful suggestions on the use of comparison surfaces, in §6.

When writing the manuscript we were informed by E. Giusti that he obtained an analogous existence result, without any use of comparison surfaces.

1. Notation and definitions. We denote by $z = (x, t)$, with $x \in \mathbf{R}^n$ and $t \in \mathbf{R}$ ($n \geq 2$), an arbitrary point in \mathbf{R}^{n+1} , by H_s the s -dimensional Hausdorff measure ([11]), by $BV(\Omega)$ the set of Lebesgue integrable functions $f(y)$ over the open subset Ω of \mathbf{R}^m , whose gradient, in the sense of distributions, is a vector measure with finite total variation. That is,

$$BV(\Omega) = \left\{ f \in L^1(\Omega): \int_{\Omega} |Df| < +\infty \right\}$$

where

$$\int_{\Omega} |Df| = \sup \left\{ \int_{\Omega} f(y) \operatorname{div} g(y) dy: g \in [C_0^1(\Omega)]^m, |g| \leq 1 \right\}.$$

We refer to the quantity $\int_{\Omega} |D\varphi_E|$, involving the characteristic function φ_E of a Borel set $E \subset \mathbf{R}^m$, as the *perimeter of E in Ω* ; when $\Omega = \mathbf{R}^m$, we simply write $\int |D\varphi_E|$, the *perimeter of E* .

If the boundary $\partial\Omega$ of Ω is locally Lipschitz, then ([23]) each function $f \in BV(\Omega)$ has a *trace* belonging to $L^1(\partial\Omega)$. The functional (0.1) is therefore well-defined on sets E having finite perimeter in \mathbf{R}^{n+1} ; of course, $\int_{t < 0} |D\varphi_E|$ means the perimeter of E in the half-space $\Omega = \{(x, t): t < 0\}$, while $\int_{t=0} \varphi_E dH_n$ corresponds to the trace over $\partial\Omega$.

Setting for $T < 0$

$$S_T = \{(x, t): T < t < 0\},$$

by $A \Subset S_T$ we intend that there exists $\delta > 0$ such that $A \subset S_{T+\delta}$. We call G (a set of finite perimeter in \mathbf{R}^{n+1}) a *local minimum* for the functional \mathcal{F} defined by (0.1) — a *pendent drop* — if there exists $T < 0$ such that

$$(i) \quad G \Subset S_T$$

(ii) for each F of finite perimeter, with $|F| = |G|$ and $F \triangle G \Subset S_T$, the inequality $\mathcal{F}(G) \leq \mathcal{F}(F)$ holds, where:

$$F \triangle G = (F - G) \cup (G - F) \quad \text{and} \quad |F| = H_{n+1}(F).$$

2. **Free-gravity problem.** Consider the functional

$$(2.1) \quad \mathcal{F}_0(E) = \int_{t < 0} |D\varphi_E| + \nu \int_{t=0} \varphi_E dH_n$$

in the class

$$\mathcal{E} = \left\{ E \subset \{t < 0\}: \int |D\varphi_E| < +\infty, |E| = V \right\}$$

where $V > 0$ is a fixed constant and $\nu \in \mathbf{R}$.

When $|\nu| < 1$, the isoperimetric inequality ([8], [18]) implies at once the existence of a unique minimum for \mathcal{F}_0 in \mathcal{E} . Such a minimum E_0 is the intersection of the half-space $\{t < 0\}$ with a ball, centered at the point c_0 on the t -axis and having radius R_0 ; radius and position of the center are to be determined in such a way that the measure of the intersection is V and the cosine of the contact angle is ν . That is

$$(2.2) \quad E_0 = \{(x, t) \in \mathbf{R}^{n+1}: t < 0, |x| < \rho_0(t) = (R_0^2 - (t - c_0)^2)^{1/2}\}$$

where R_0, c_0 are to be determined through the relations

$$(2.3) \quad |E_0| = V; \quad \nu = -\frac{\rho'_0(0)}{\sqrt{1 + \rho_0'^2(0)}}.$$

The *minimum height* of the solution E_0 is given by

$$(2.4) \quad Q_0 = c_0 - R_0 = \frac{-(1 + \nu)V^{1/n+1}}{\left(\omega_n \int_{-1}^{\nu} (1 - s^2)^{n/2} ds\right)^{1/n+1}}.$$

We remark moreover that every ball of volume V , lying in $\{t < 0\}$, is a solution of the problem in the case $\nu = 1$, while for $\nu = -1$ no solution can occur.

3. **ε -gravity problems.** We prove in this section that the functional

$$(3.1) \quad \mathcal{F}_\varepsilon(E) = \int_{S_T} |D\varphi_E| + \nu \int_{t=0} \varphi_E dH_n + \int_{t=T} \varphi_E dH_n + \varepsilon \int_{S_T} t\varphi_E(z) dz$$

attains its minimum in the class

$$(3.2) \quad \mathcal{E}_T = \left\{ E \subset S_T: \int |D\varphi_E| < +\infty, |E| = V \right\}$$

for every $\varepsilon > 0$.

With respect to the minimization of \mathcal{F}_ε , we may restrict ourselves to rotationally symmetric sets. Indeed, if for $E \in \mathcal{E}_T$ we define

$$\begin{aligned} \rho(t) &= \left(\omega_n^{-1} \int \varphi_E(x: t) dx \right)^{1/n} \\ E^s &= \{(x, t) \in \mathbf{R}^{n+1}, |x| < \rho(t)\} \end{aligned}$$

then we derive from Lemma 2 in [20]:

$$\mathcal{F}_\varepsilon(E^s) \leq \mathcal{F}_\varepsilon(E)$$

equality holding if and only if $E^s = E$, that is, the set E is already symmetric. We define

$$(3.3) \quad \tilde{\mathcal{E}}_T = \{E^s: E \in \mathcal{E}_T\}.$$

LEMMA 3.1. *For $|\nu| \leq 1$, the functional \mathcal{F}_ε defined by (3.1) has a finite lower bound on \mathcal{E}_T .*

Proof. Clearly we have

$$\int_{t < 0} |D\varphi_E| > \int_{t=0} \varphi_E dH_n$$

for any $E \in \mathcal{E}_T$, from which we obtain

$$\begin{aligned} 0 &\leq \frac{1-\nu}{2} \int_{t < 0} |D\varphi_E| + \frac{\nu-1}{2} \int_{t=0} \varphi_E dH_n \\ &= \left(1 - \frac{1+\nu}{2}\right) \int_{t < 0} |D\varphi_E| + \left(\nu - \frac{1+\nu}{2}\right) \int_{t=0} \varphi_E dH_n \end{aligned}$$

and then

$$(3.4) \quad \mathcal{F}_\varepsilon(E) > \frac{1+\nu}{2} \int |D\varphi_E| + \varepsilon TV$$

which concludes the proof.

LEMMA 3.2. *Let $\{E_h\}$ be a sequence of sets in $\tilde{\mathcal{E}}_T$ such that $\mathcal{F}_\varepsilon(E_h) < \text{const. } \forall h$. If $-1 < \nu \leq 1$, then a subsequence of $\{E_h\}$ converges in $L^1(\mathbf{R}^{n+1})$ to a limit set $E \in \tilde{\mathcal{E}}_T$.*

Proof. The existence of a subsequence of $\{E_h\}$, convergent in $L^1_{loc}(\mathbf{R}^{n+1})$, follows easily from (3.4) together with a known compactness result ([24]). Now, we have for almost all $t < 0$:

$$(3.5) \quad \int \varphi_{E_h}(x, t) dx \leq \int |D\varphi_{E_h}| \leq \text{const} .$$

We can thus find a ball $B_R \subset \mathbf{R}^n$, such that every set in the sequence lies in the cylinder $B_R \times (T, 0)$. Hence, convergence of the previous subsequence actually takes place in $L^1(\mathbf{R}^{n+1})$.

LEMMA 3.3. *If $|\nu| \leq 1$, then the functional \mathcal{F}_ε is lower semi-continuous with respect to $L^1(S_T)$ -convergence.*

Proof. Let $E_h \rightarrow E$ in $L^1(S_T)$, and suppose there exists $\sigma > 0$ such that the inequality

$$(3.6) \quad \mathcal{F}_\varepsilon(E) > \mathcal{F}_\varepsilon(E_h) + \sigma$$

or equivalently

$$(3.7) \quad \int_{S_T} |D\varphi_{E_h}| < \int_{S_T} |D\varphi_E| + \nu \int_{t=0} (\varphi_E - \varphi_{E_h}) dH_n \\ + \int_{t=-T} (\varphi_E - \varphi_{E_h}) dH_n + \varepsilon \int_{S_T} t(\varphi_E - \varphi_{E_h}) dz - \sigma$$

holds, for infinite indices h . Combining (3.7) with the estimate

$$\int_{\partial S_T} |f| dH_n \leq \int_{S_T(\delta)} |Df| + c(\delta) \int_{S_T} |f| dz$$

(which holds for every $f \in BV(S_T)$ and $\delta \in (0, -T/2)$, provided we define

$$S_T(\delta) = \{z \in S_T: \text{dist}(z, \partial S_T) < \delta\} ,$$

see [9]), we then obtain

$$\int_{S_T - \bar{S}_T(\delta)} |D\varphi_{E_h}| < \int_{S_T} |D\varphi_E| + \int_{S_T(\delta)} |D\varphi_E| \\ + (c(\delta) - \varepsilon T) \int_{S_T} |\varphi_E - \varphi_{E_h}| dz - \sigma$$

from which we derive a contradiction, letting $h \rightarrow +\infty$ and then $\delta \downarrow 0$, and taking into account the lower semicontinuity of perimeter functional (see [5]).

According to the previous lemmas, we can state the following existence result:

THEOREM 3.4. *For every $T < 0$, $V > 0$, $-1 < \nu \leq 1$ and for every*

$\varepsilon > 0$, the functional \mathcal{F}_ε defined by (3.1) has a minimum E_ε in the class \mathcal{E}_T , which is a rotationally symmetric set, so that we can write:

$$(3.8) \quad E_\varepsilon = \{(x, t) \in S_T : |x| < \rho_\varepsilon(t)\}$$

for a suitable function $\rho_\varepsilon \in BV((T, 0))$.

4. Some properties of ε -solutions. We study in this section the behaviour of the solutions E_ε , found in Theorem 3.4, as ε tends to zero.

From now on, we shall assume that $T = T(\nu, V)$ is a negative number, smaller than the minimum height $Q_0(\nu, V)$ of the solution E_0 in the absence of gravity (see (2.4); the data ν and V will be held fixed). Moreover, we shall write S instead of S_T .

PROPOSITION 4.1. *If $|\nu| < 1$, then*

$$(4.1) \quad E_\varepsilon \longrightarrow E_0 \text{ in } L^1(S)$$

$$(4.2) \quad \mathcal{F}_\varepsilon(E_\varepsilon) \longrightarrow \mathcal{F}_0(E_0)$$

as $\varepsilon \rightarrow 0$.

Proof. Let $A \subset S$ be a fixed set of volume V . We have for every $\varepsilon > 0$

$$\mathcal{F}_\varepsilon(E_\varepsilon) \leq \mathcal{F}_\varepsilon(A) < \text{const.},$$

so that (3.4) implies that the perimeters of the E_ε 's are uniformly bounded.

Now, the sets E_ε are all included in a fixed cylinder, as we have pointed out during the proof of Lemma 3.2. This fact, together with the compactness result already mentioned, entails the convergence in $L^1(\mathbf{R}^{n+1})$ of the E_ε 's to some set $F \in \mathcal{E}_T$. We claim that F minimizes functional \mathcal{F}_0 defined by (2.1). Indeed, by the assumption $T < Q_0$ we have $E_0 \in \mathcal{E}_T$, and moreover we have

$$(4.3) \quad \mathcal{F}_0(E_\varepsilon) \leq \mathcal{F}_0(E_0) - \varepsilon \int_S t(\varphi_{E_\varepsilon}(z) - \varphi_{E_0}(z)) dz$$

owing to the minimality of E_ε 's.

Letting $\varepsilon \rightarrow 0$ in (4.3), we obtain from the lower semicontinuity of \mathcal{F}_0 with respect to L^1 -convergence

$$\mathcal{F}_0(F) \leq \mathcal{F}_0(E_0).$$

Uniqueness of the solution E_0 allows to conclude that $F = E_0$, thus

proving (4.1).

From (4.3) we derive also

$$\mathcal{F}_0(E_0) \leq \liminf_{\varepsilon \rightarrow 0} \mathcal{F}_\varepsilon(E_\varepsilon) \leq \limsup_{\varepsilon \rightarrow 0} \mathcal{F}_\varepsilon(E_\varepsilon) \leq \mathcal{F}_0(E_0)$$

which yields (4.2).

LEMMA 4.2. *There holds for $|\nu| < 1$:*

$$(4.4) \quad \lim_{\varepsilon \rightarrow 0} \int_{t=0} \varphi_{E_\varepsilon} dH_n = \int_{t=0} \varphi_{E_0} dH_n .$$

Proof. Let $\varepsilon_h \rightarrow 0$ be such that, for the corresponding solutions $E_h \equiv E_{\varepsilon_h}$, there exists

$$\lim_{h \rightarrow +\infty} \int_{t=0} \varphi_{E_h} dH_n = a .$$

Denote by B_h , $h = 0, 1, 2, \dots$, the intersection of the half-space $\{t < 0\}$ with a suitable ball in \mathbf{R}^{n+1} , centered on the t -axis and chosen in such a way that

$$(4.5) \quad \begin{aligned} |B_h| &= V && \forall h \\ \int_{t=0} \varphi_{B_h} dH_n &= \int_{t=0} \varphi_{E_h} dH_n && \forall h \geq 1 \\ \int_{t=0} \varphi_{B_0} dH_n &= a . \end{aligned}$$

In view of the isoperimetric inequality ([8], [18]) we have, for every $h \geq 1$:

$$\int_{t < 0} |D\varphi_{E_h}| \geq \int_{t < 0} |D\varphi_{B_h}|$$

and hence

$$\mathcal{F}_{\varepsilon_h}(E_h) \geq \int_{t < 0} |D\varphi_{B_h}| + \nu \int_{t=0} \varphi_{B_h} dH_n + \varepsilon_h \int_S t\varphi_{E_h}(z) dz ;$$

letting $h \rightarrow +\infty$ and recalling (4.2) we obtain

$$\mathcal{F}_0(E_0) \geq \mathcal{F}_0(B_0) .$$

Now, E_0 is the unique minimum of \mathcal{F}_0 , so that $E_0 = B_0$ and

$$\lim_{h \rightarrow +\infty} \int_{t=0} \varphi_{E_h} dH_n = \int_{t=0} \varphi_{E_0} dH_n$$

from which (4.4) follows at once.

LEMMA 4.3. *There holds for $|\nu| < 1$*

$$(4.6) \quad \lim_{\varepsilon \rightarrow 0} \int_S |D\varphi_{E_\varepsilon}| = \int_S |D\varphi_{E_0}|$$

$$(4.7) \quad \lim_{\varepsilon \rightarrow 0} \int_{t=T} \varphi_{E_\varepsilon} dH_n = 0.$$

Proof. From (4.3) we have

$$\int_S |D\varphi_{E_\varepsilon}| + \nu \int_{t=0} \varphi_{E_\varepsilon} dH_n + \varepsilon \int_S t\varphi_{E_\varepsilon}(z) dz \leq \mathcal{F}_\varepsilon(E_\varepsilon) \leq \mathcal{F}_0(E_0)$$

and then, from (4.4)

$$\limsup_{\varepsilon \rightarrow 0} \int_S |D\varphi_{E_\varepsilon}| \leq \int_S |D\varphi_{E_0}|,$$

while

$$\int_S |D\varphi_{E_0}| \leq \liminf_{\varepsilon \rightarrow 0} \int_S |D\varphi_{E_\varepsilon}|$$

follows from (4.1) and lower semicontinuity of the perimeter functional.

REMARK 4.4. An obvious consequence of the previous results is the (pointwise a.e.) convergence of the functions defining the rotationally symmetric solutions E_ε, E_0 . We recall (see Theorem 3.4) that such functions $\rho_\varepsilon: [T, 0] \rightarrow [0, +\infty)$ are defined through the relations

$$E_\varepsilon = \{(x, t) \in S: |x| < \rho_\varepsilon(t)\}, \quad \varepsilon \geq 0.$$

Since modifications of any set of finite perimeter by sets of zero-measure do not affect its perimeter, we may assume the existence of both the one-sided limits of ρ_ε , at every point in the interval $(T, 0)$. Such limits do coincide, except for a countable set of points.

We have actually the following stronger result:

THEOREM 4.5. *The convergence $\rho_\varepsilon \rightarrow \rho_0$ is uniform on $(T, 0)$.*

Proof. We shall prove the statement in the theorem by showing that every subsequence of $\{\rho_\varepsilon\}_{\varepsilon > 0}$ admits a subsequence uniformly convergent to ρ_0 on $(T, 0)$.

Let $\varepsilon_h \rightarrow 0$; for the corresponding sequence $\rho_h = \rho_{\varepsilon_h}$ we have, in view of (4.1)

$$\lim_{h \rightarrow +\infty} \omega_n \int_T^0 |\rho_h^n(t) - \rho_0^n(t)| dt = 0.$$

Hence, a subsequence (not relabeled) of $\{\rho_h\}$ satisfies, for almost all $t \in [T, 0]$:

$$(4.8) \quad \lim_{h \rightarrow +\infty} \rho_h(t) = \rho_0(t) .$$

Assume by contradiction there exist $\sigma > 0$ and a monotone sequence $\{t_h\} \subset [T, 0]$ such that

$$(4.9) \quad \begin{cases} |\rho_h(t_h) - \rho_0(t_h)| > \sigma \\ \lim_{h \rightarrow +\infty} t_h = t_0 . \end{cases}$$

If $\{t_h\}$ is not decreasing, then choose $\bar{t} \in (T, 0)$ such that

$$(4.10) \quad \bar{t} < t_0$$

$$(4.11) \quad \lim_{h \rightarrow +\infty} \rho_h(\bar{t}) = \rho_0(\bar{t})$$

(in the opposite case, i.e., $\{t_h\}$ not increasing, choose \bar{t} s.t. (4.11) holds with $\bar{t} > t_0$).

There follows that, for sufficiently great h

$$\begin{aligned} \int_S |D\varphi_{E_h}| &= \int_{t_h < t < 0} |D\varphi_{E_h}| + \int_{\bar{t} \leq t \leq t_h} |D\varphi_{\bar{t}_h}| + \int_{T < t < \bar{t}} |D\varphi_{E_h}| \\ &\geq \int_{t_0 < t < 0} |D\varphi_{E_h}| + \omega_n |\rho_h^n(t_h) - \rho_h^n(\bar{t})| + \int_{T < t < \bar{t}} |D\varphi_{E_h}| . \end{aligned}$$

On the other hand, from (4.9), (4.11) and from the continuity of ρ_0 , we derive

$$|\rho_h(t_h) - \rho_h(\bar{t})| \geq \sigma - (|\rho_0(t_h) - \rho_0(\bar{t})| + |\rho_0(\bar{t}) - \rho_h(\bar{t})|) \geq \frac{\sigma}{2}$$

provided that \bar{t} is close to t_0 and h is great. Therefore,

$$\int_S |D\varphi_{E_h}| \geq \int_{t_0 < t < 0} |D\varphi_{E_h}| + \omega_n \left(\frac{\sigma}{2}\right)^n + \int_{T < t < \bar{t}} |D\varphi_{E_h}|$$

and letting $h \rightarrow +\infty$ we obtain

$$\liminf_{h \rightarrow \infty} \int_S |D\varphi_{E_h}| \geq \int_{t_0 < t < 0} |D\varphi_{E_0}| + \omega_n \left(\frac{\sigma}{2}\right)^n + \int_{T < t < \bar{t}} |D\varphi_{E_0}|$$

which contrasts with (4.6) as \bar{t} approaches t_0 .

5. Regularity of ε -solutions. The main result in this section is the following.

THEOREM 5.1. *Let $t_0 \in (T, 0)$. If*

$$\liminf_{t \rightarrow t_0} \rho_\varepsilon(t) > 0$$

then there exists a neighborhood U of t_0 such that $\rho_\varepsilon \in C^2(U)$.

Proof. In view of Remark 4.4 we can find a neighborhood W of t_0 , such that

$$(5.1) \quad \inf_W \rho_\varepsilon(t) = m_\varepsilon > 0.$$

Hence, there exists a ball B lying in the plane $\{x_1 = 0\}$ and centered at $(0, \dots, t_0)$, which satisfies

$$(5.2) \quad \begin{aligned} E_\varepsilon \cap (\mathbf{R}^+ \times B) &= \{(x_1, y, t) \in \mathbf{R}^{n+1}: (y, t) \in B, \\ &0 < x_1 < f_\varepsilon(y, t) = (\rho_\varepsilon^2(t) - |y|^2)^{1/2}\} \end{aligned}$$

where $y = (x_2, \dots, x_n) \in \mathbf{R}^{n-1}$ and $f_\varepsilon(y, t) \geq \psi_\varepsilon(y, t) \equiv m_\varepsilon/2$. Function f_ε belongs to $BV(B)$ and, owing to the minimum property of ρ_ε , it minimizes the functional

$$I_\varepsilon(u) = \int_B \sqrt{1 + |Du|^2} + \int_{\partial B} |u - f_\varepsilon| dH_{n-1} + \varepsilon \int_B tudydt$$

in the function class

$$H_\varepsilon = \left\{ u \in BV(B): u \geq \psi_\varepsilon, \int_B (u - \psi_\varepsilon) dydt = \int_B (f_\varepsilon - \psi_\varepsilon) dydt \right\}.$$

We then conclude that f_ε minimizes the functional

$$I_\varepsilon(u) + \lambda_\varepsilon \int_B u dydt$$

in the function class

$$K_\varepsilon = \{u \in BV(B): u \geq \psi_\varepsilon\}$$

as well, for a suitable multiplier $\lambda_\varepsilon \in \mathbf{R}$, whose existence is granted by the results in the Appendix. From known regularity results (see e.g., [10]) we obtain $f_\varepsilon \in C^{2,\alpha}(B)$, $0 \leq \alpha < 1$, thus proving the theorem.

REMARK 5.2. As a consequence of the theorem we can state, that if $\liminf_{t \rightarrow t_0} \rho_\varepsilon(t) > 0 \forall t_0 \in (a, b) \subset (T, 0)$, then $\rho_\varepsilon \in C^2(a, b)$ and there exists a real number A_ε such that

$$(5.3) \quad \left(\frac{\rho'_\varepsilon}{\sqrt{1 + \rho_\varepsilon'^2}} \right)' - \frac{n-1}{\rho_\varepsilon \sqrt{1 + \rho_\varepsilon'^2}} = \varepsilon t + A_\varepsilon \quad \forall t \in (a, b).$$

Now, let $[a, b]$ be a sub-interval of $(Q_0, 0)$ (we recall that Q_0 denotes

the height of the apex of the drop below the contact plane, in the absence of gravity; see (2.4)). In view of Theorem 4.5 we conclude, that for small ε each function ρ_ε is greater than a positive constant, independent of ε , in the whole interval (a, b) . This fact allows the proof of the following result, concerning the convergence of the A_ε 's defined in the preceding remark.

PROPOSITION 5.3. *There holds*

$$(5.4) \quad \lim_{\varepsilon \rightarrow 0} A_\varepsilon = -\frac{n}{R_0} = A_0$$

with R_0 defined by (2.2), (2.3).

Proof. As in the proof of Theorem 5.1, we can find a ball $B \subset \{x_1 = 0\}$ (with radius independent of ε) and a function ψ positive on B , such that

$$(5.5) \quad \begin{cases} B \cap \{y = 0\} \subset (a, b) \\ f_\varepsilon \geq 2\psi \quad \text{on } B \end{cases}$$

and f_ε minimizes the functional

$$(5.6) \quad \int_B \sqrt{1 + |Du|^2} + \int_{\partial B} |u - f_\varepsilon| dH_{n-1} + \varepsilon \int_B tudydt + \lambda_\varepsilon \int_B udydt$$

in the class $\{u \in BV(B): u \geq \psi\}$.

About the multipliers λ_ε we know the estimate (see (A6), (A7) in the Appendix):

$$C_2(B) \leq \lambda_\varepsilon \leq \frac{C_1(\psi, B) + \int_{\partial B} |f_\varepsilon| dH_{n-1}}{V_\varepsilon} + c(B)$$

where

$$V_\varepsilon = \int_B (f_\varepsilon - \psi) dydt .$$

We may then assume $\lambda_\varepsilon \rightarrow \lambda_0$, from which we conclude that the function f_0 , limit of the f_ε 's, minimizes functional (5.6) with $\varepsilon = 0$. From (5.3) we have therefore:

$$\lambda_0 = -\frac{n}{R_0}, \quad \lambda_\varepsilon = A_\varepsilon .$$

6. **The existence theorem.** So far we obtained a minimum E_ε of the functional \mathcal{F}_ε defined by (3.1), for every $\varepsilon > 0$. In order to

get the existence of a pendent drop it suffices to show, that when ε is small enough, there exists $t_\varepsilon \in (T, 0)$ such that $\liminf_{t \rightarrow t_\varepsilon} \rho_\varepsilon(t) = 0$.

Indeed, in this situation the set

$$G_\varepsilon = E_\varepsilon \cap \{t > t_\varepsilon\}$$

clearly yields the minimum value for \mathcal{F}_ε , when compared with the subsets of the strip S having finite perimeter and volume equal to $|G_\varepsilon|$. In other words, following the definition given in §1, we can assert that G_ε is a *local minimum* of functional (0.1), with $\kappa = \varepsilon$.

To this end, we compare the boundaries of solutions E_ε with surfaces of constant mean curvature (see [4]).

We begin by noting that if $\rho'_\varepsilon(\bar{t}) > 0$, then there exists, in a suitable neighborhood U of the point $\bar{r} = \rho_\varepsilon(\bar{t})$, the inverse function $u_\varepsilon(r)$ of ρ_ε ; moreover, from (5.3) we deduce

$$(6.1) \quad \frac{1}{r^{n-1}} \left(\frac{r^{n-1} u'_\varepsilon}{\sqrt{1 + u_\varepsilon'^2}} \right)' = -\varepsilon u_\varepsilon - A_\varepsilon \quad (r \in U).$$

On the contrary, if $\rho'_\varepsilon(\bar{t}) < 0$, then the following equation holds instead of (6.1)

$$(6.2) \quad \frac{1}{r^{n-1}} \left(\frac{r^{n-1} u'_\varepsilon}{\sqrt{1 + u_\varepsilon'^2}} \right)' = \varepsilon u_\varepsilon + A_\varepsilon \quad (r \in U).$$

In either case, denoting by $\psi_{u_\varepsilon}(r)$ the angle between the tangent line to u_ε and the r -axis, measured counterclockwise from the r -axis itself, we can combine (6.1) and (6.2) in the single equation

$$(6.3) \quad \frac{1}{r^{n-1}} (r^{n-1} \sin \psi_{u_\varepsilon}(r))' = -\varepsilon u_\varepsilon(r) - A_\varepsilon.$$

We also need the following results (see [4]):

LEMMA 6.1. *Let $u(r), v(r)$ be functions defined over $0 \leq a \leq r \leq b$, s.t.*

$$(r^{n-1} \sin \psi_v(r))' \geq (r^{n-1} \sin \psi_u(r))'$$

If $\sin \psi_v(a) \geq \sin \psi_u(a)$, then $\sin \psi_v(b) \geq \sin \psi_u(b)$, and equality holds if and only if $u = v + \text{const}$ in $[a, b]$.

LEMMA 6.2. *For every $H > 0$ and for every a such that*

$$(6.4) \quad 0 < a < \frac{n-1}{nH}$$

there exists an unique b with

$$(6.5) \quad \frac{n-1}{nH} < b < \frac{1}{H}$$

and an increasing function $v: [a, b] \rightarrow \mathbf{R}$ for which

$$(6.6) \quad \frac{1}{r^{n-1}}(r^{n-1} \sin \psi_v(r))' = nH \quad (r \in (a, b))$$

$$(6.7) \quad \sin \psi_v(a) = \sin \psi_v(b) = 1.$$

Moreover, there holds

$$(6.8) \quad v(b) - v(a) \leq \frac{c(n)}{H}.$$

The proof of Lemma 6.1 is quite obvious, so we outline only that of Lemma 6.2. We point out that the surface obtained by rotating the graph of function $v(r)$ in the lemma (which is part of an *ellipse roulade* when $n = 2$, see [4]) about the vertical axis, has constant mean curvature H .

Integrating (6.6) yields

$$(6.9) \quad v'(r) = \frac{g(r)}{\sqrt{1-g^2(r)}}$$

where $g(r) = Hr + Br^{1-n}$.

For the constant of integration B we have, in view of (6.7)

$$B = a^{n-1} - Ha^n = b^{n-1} - Hb^n > 0$$

and letting $c = (n-1)/nH$ we get from (6.9)

$$(6.10) \quad v(b) - v(a) \leq \frac{1}{\sqrt{2}} \left(\int_a^c \frac{dr}{\sqrt{1-g(r)}} + \int_c^b \frac{dr}{\sqrt{1-g(r)}} \right).$$

Defining

$$\alpha(r) = (r-a)(1-g(r))^{-1}$$

$$\beta(r) = (b-r)(1-g(r))^{-1}$$

and noting that $\alpha'(r) > 0 > \beta'(r)$, we have from (6.10)

$$v(b) - v(a) \leq \sqrt{2} (\sqrt{\alpha(c)}\sqrt{c-a} + \sqrt{\beta(c)}\sqrt{b-c}).$$

Now, functions

$$\sqrt{\alpha(c)(c-a)} = (c-a)(1-g(c))^{-1/2} \quad 0 < a < c$$

$$\sqrt{\beta(c)(b-c)} = (b-c)(1-g(c))^{-1/2} \quad c < b < \frac{1}{H}$$

are decreasing and we can conclude

$$v(b) - v(a) \leq \frac{c(n)}{H}.$$

We are now in position to prove the following

THEOREM 6.3. *If $T < Q_0 - 3c(n)R_0$ and $0 < \varepsilon < \varepsilon_0(T)$, then there exists $t_\varepsilon \in (T, 0)$ such that*

$$\liminf_{t \rightarrow t_\varepsilon} \rho_\varepsilon(t) = 0 .$$

Proof. Assume by contradiction the existence of a sequence $\{\varepsilon_h\}$ converging to zero, such that

$$\liminf_{s \rightarrow t} \rho_h(s) > 0$$

for every $t \in (T, 0)$ and every h , with $\rho_h \equiv \rho_{\varepsilon_h}$.

In view of Remark 5.2, for every h we have $\rho_h \in C^2(T, 0)$; moreover ρ_h satisfies the Euler equation

$$(6.11) \quad \left(\frac{\rho'_h}{\sqrt{1 + \rho_h'^2}} \right)' - \frac{n - 1}{\rho_h \sqrt{1 + \rho_h'^2}} = \varepsilon_h t + A_h$$

in which

$$(6.12) \quad \lim_{h \uparrow +\infty} A_h = -\frac{n}{R_0}$$

(see (5.4)). Now, let $T_1 = T + (Q_0 - T)/3$ and $T_2 = T + (2(Q_0 - T))/3$, and denote by $t_h \in (T_1, T_2)$ a point such that $\rho'_h(t_h) \neq 0$. For the sake of definiteness, assume $\rho'_h(t_h) > 0$, since in the opposite case one can proceed analogously.

Let $a_h = \rho_h(t_h)$ and $H_h = (-\varepsilon_h T - A_h)/n$; there follows from Theorem 4.5 and (6.12) that

$$(6.13) \quad \lim_{h \rightarrow +\infty} a_h = 0 ; \quad \lim_{h \rightarrow +\infty} H_h = \frac{1}{R_0} ,$$

so that, in view of Lemma 6.2, we can find (for sufficiently great h) $b_h \in ((n - 1)/nH_h, 1/H_h)$ and $v_h : [a_h, b_h] \rightarrow \mathbf{R}$ satisfying (6.6), (6.7) together with $v_h(a_h) = t_h$.

An application of Lemma 6.1 to u_h, v_h (u_h denotes the inverse function of ρ_h) allows the conclusion that u_h is defined over the whole interval $[a_h, b_h]$, and verifies there

$$u_h(r) \leq v_h(r) \quad \forall r \in [a_h, b_h] .$$

Finally, as a consequence of the choice of T , we have from (6.8)

$$v_h(b_h) \leq \frac{c(n)}{H_h} + T_2 < Q_0 + c(n) \left(\frac{1}{H_h} - R_0 \right)$$

and therefore, for great h

$$v_h(b_h) < Q_0 .$$

Now, in the interval (T, Q_0) the sequence $\{\rho_h\}$ tends uniformly to zero (Theorem 4.5), so contradicting the fact that

$$\rho_h(u_h(b_h)) = b_h > \frac{n-1}{nH_h} \longrightarrow \frac{n-1}{n}R_0 > 0 .$$

From the minimality of E_ε there follows that the set of points $t_\varepsilon \in (T, 0)$ such that $\liminf_{t \rightarrow t_\varepsilon} \rho_\varepsilon(t) = 0$ is a closed interval in $(T, 0)$ (of course, it may consist of a single point). Denoting by Q_ε the maximum of such interval, we conclude that ρ_ε is positive and regular in $(Q_\varepsilon, 0)$; Q_ε thus represents the minimum height of the pendent drop

$$(6.14) \quad G_\varepsilon = \{(x, t) : t \in (Q_\varepsilon, 0), |x| < \rho_\varepsilon(t)\} .$$

7. **Regularity at the vertex.** At this point we have proved the existence of a local minimum of (0.1) for sufficiently small κ ; nevertheless, we cannot say anything about the effective volume of such solution, nor about the effective smallness of κ .

It is however clear, that by a homothetic transformation of the co-ordinate system we can expand (or contract) our solution so that it becomes a local minimum of (0.1) (for a different κ , of course) among sets of prescribed volume V ; or, so that it becomes a local minimum of (0.1) for a prescribed value of $\kappa > 0$, in the class of sets having its own volume (which remains unspecified).

Thus, we can obtain by this method pendent drops of fixed volume in a weak gravitational field, or pendent drops of small volume in a prescribed gravitational field.

The following considerations are devoted to the study of the behavior of our solution near its minimum height Q_ε .

Let $\tau \in (Q_\varepsilon, 0)$, and denote by $B(\tau)$ the ball centered on the t -axis, passing through the boundary point $(\rho_\varepsilon(\tau), 0, \dots, 0, \tau) \in \mathbf{R}^{n+1}$ and satisfying (see (6.14)):

$$(7.1) \quad H_{n+1}[G_\varepsilon - B(\tau)] \cap \{Q_\varepsilon < t < \tau\} = H_{n+1}[B(\tau) \cap \{t < Q_\varepsilon\}]$$

$R(\tau)$ will denote the radius of $B(\tau)$.

LEMMA 7.1. *If*

$$(7.2) \quad \lim_{t \downarrow Q_\varepsilon} \rho_\varepsilon(t) = L > 0$$

then there exists $\tau \in (Q_\varepsilon, 0)$ such that

$$(7.3) \quad B(\tau) \cap \{Q_\varepsilon < t < \tau\} \subset G_\varepsilon \cap \{Q_\varepsilon < t < \tau\} .$$

Proof. It follows from (7.2) that

$$(7.4) \quad \lim_{\tau \downarrow Q_\varepsilon} R(\tau) = +\infty .$$

If the assertion in the lemma were false, then we could find, in view of (7.4) and the mean value theorem, a sequence $\{t_j\}$ converging to Q_ε from above, such that

$$(7.5) \quad \lim_{j \rightarrow +\infty} \rho'_\varepsilon(t_j) = +\infty .$$

Taking into account (7.2), (7.5), we derive from (5.3)

$$\frac{\rho''_\varepsilon(t_j)}{\sqrt{(1 + \rho'^2_\varepsilon(t_j))^3}} = \frac{n-1}{\rho_\varepsilon(t_j)\sqrt{1 + \rho'^2_\varepsilon(t_j)}} + \varepsilon t_j + A_\varepsilon < 0$$

when $j > j_0$, since, as we shall see in the following remark, there holds $\varepsilon Q_\varepsilon + A_\varepsilon < 0$. Then we can easily conclude that $\rho''_\varepsilon(t) < 0 \forall t \in (Q_\varepsilon, t_{j_0})$. Hence, we can describe the boundary of G_ε , lying in a neighborhood of the point $(L, 0, \dots, 0, Q_\varepsilon) \in \mathbf{R}^{n+1}$, as the graph of a suitable function $t = u(x)$ which, in view of the results in §5, would be analytic over its domain of definition; but this leads to a contradiction, since the (not-identically constant) function u should be constant ($\equiv Q_\varepsilon$) over an open set.

COROLLARY 7.2. *There holds*

$$(7.6) \quad \lim_{t \downarrow Q_\varepsilon} \rho_\varepsilon(t) = 0 .$$

Proof. If not (see Remark 4.4), we get from the lemma the existence of a ball $B(\tau)$ satisfying (7.3). Defining

$$\tilde{G}_\varepsilon = \begin{cases} G_\varepsilon & \text{for } t \geq \tau \\ B(\tau) & \text{for } t \leq \tau \end{cases}$$

there follows from the isoperimetric inequality

$$H_{n+1}(\tilde{G}_\varepsilon) = H_{n+1}(G_\varepsilon); \quad \mathcal{F}_\varepsilon(\tilde{G}_\varepsilon) < \mathcal{F}_\varepsilon(G_\varepsilon) ,$$

a contradiction.

REMARK 7.3. The inequality $\varepsilon Q_\varepsilon + A_\varepsilon < 0$ we used in the proof of Lemma 7.1 really holds for a local minimum G_ε . Assuming the contrary, we derive from (5.3): $\rho''_\varepsilon(t) > 0$ for every $t \in (Q_\varepsilon, 0)$, and then either ∂G_ε is a (regular) graph in the t -direction, or there exists τ close to Q_ε s.t. the corresponding ball $B(\tau)$ satisfies (7.3). In either case, as we saw just before, we are led to a contradiction.

COROLLARY 7.4. *There exists $\delta > 0$ such that $\partial G_\varepsilon \cap B_\delta(0, \dots, 0, Q_\varepsilon)$ is representable as a graph of an analytic function $t = u(x)$.*

Proof. Assume the contrary; then it would be possible to find a sequence $\{t_j\}$ of local maximum points for ρ_ε , decreasing to Q_ε . On the other hand, we have from (5.3)

$$\rho_\varepsilon''(t_j) \geq \frac{n-1}{\rho_\varepsilon(t_j)} + \varepsilon Q_\varepsilon + A_\varepsilon$$

which is positive for j great enough.

The assertion about analiticity follows from the methods of §5.

The preceding results, together with that of §5, can be summarized in the following

THEOREM 7.5. *If G is a local minimum of the functional \mathcal{F} defined by (0.1), then $\partial G \cap \{t < 0\}$ is an analytic n -dimensional manifold.*

A. Appendix: Existence of multipliers. For convenience of the reader, we quote the proof of the following result, which we used in the proof of Theorem 5.1.

THEOREM A.1. *If f minimizes functional*

$$I(u) = \int_B \sqrt{1 + |Du|^2} + \int_{\partial B} |u - f| dH_{n-1} + \int_B tudydt$$

in the class

$$H = \left\{ u \in BV(B) : u \geq \psi, \int_B (u - \psi) dydt = \int_B (f - \psi) dydt = V > 0 \right\}$$

(ψ denotes a Lipschitz function over B), then there exists $\lambda \in \mathbf{R}$ such that f minimizes $I(u) + \lambda \int_B u dydt$ in the class $K = \{u \in BV(B) : u \geq \psi\}$.

The method of proof appears in [16], and involves various steps.

LEMMA A.2. *For any $\eta > 0$ and $\lambda \in \mathbf{R}$, the functional*

$$I_{\eta,\lambda}(u) = I(u) + \frac{\eta}{2} \int_B u^2 dydt + \lambda \int_B u dydt$$

attains its unique minimum $u_{\eta,\lambda} \in C^{0,1}(B)$ in the class K .

Proof. From the inequality $2ab \leq \sigma a^2 + b^2/\sigma (\sigma > 0)$ we derive

$$I_{\eta,\lambda}(u) \geq \int_B \sqrt{1 + |Du|^2} + \int_{\partial B} |u - f| dH_{n-1} - \frac{1}{\eta} \left(\int_B t^2 dydt + \lambda^2 |B| \right)$$

that is, a lower bound for the functional and the necessary compactness property.

Lower semicontinuity and strict convexity of $I_{\eta,\lambda}$ allow the conclusion about existence and uniqueness of the minimum in K . Its regularity can be obtained from a wellknown gradient estimate.

LEMMA A.3. *The function*

$$h(\lambda) = \int_B (u_{\eta,\lambda} - \psi) dydt$$

is continuous over \mathbf{R} and satisfies

$$(A.1) \quad \lim_{\lambda \rightarrow +\infty} h(\lambda) = 0$$

$$(A.2) \quad \lim_{\lambda \rightarrow -\infty} h(\lambda) = +\infty .$$

Proof. It is easy to show, that as λ_j tends to λ_0 , the sequence of corresponding minima $\{u_{\eta,\lambda_j}\}$ tends in $L^1(B)$ to a function v , which minimizes I_{η,λ_0} . Uniqueness implies therefore $v = u_{\eta,\lambda_0}$.

On the other hand, letting $c = \max_B \{-t\}$, we have

$$(A.3) \quad \begin{aligned} (\lambda - c) \int_B (u_{\eta,\lambda} - \psi) dydt &\leq I_{\eta,\lambda}(u_{\eta,\lambda}) - \int_B t u_{\eta,\lambda} dydt \\ &- c \int_B (u_{\eta,\lambda} - \psi) dydt - \lambda \int_B \psi dydt \\ &\leq \int_B \sqrt{1 + |D\psi|^2} + \int_{\partial B} |\psi - f| dH_{n-1} + \frac{\eta}{2} \int_B \psi^2 dydt \\ &- \int_B (t + c)(u_{\eta,\lambda} - \psi) dydt \leq c_1(\psi, B) + \int_{\partial B} |f| dH_{n-1} \end{aligned}$$

and this yields (A.1).

As far as (A.2) is concerned, define for $\delta \geq 0$ and $\varphi \in H^{1,1}(B)$, $\varphi \geq 0$:

$$\begin{aligned} u_\delta &= u_{\eta,\lambda} + \delta\varphi \\ \alpha(\delta) &= I_{\eta,\lambda}(u_\delta) - \int_{\partial B} |u_\delta - f| dH_{n-1} \end{aligned}$$

so that

$$\alpha(0) \leq \alpha(\delta) + \delta \int_{\partial B} \varphi dH_{n-1} .$$

Hence, the function $\alpha(\delta) + \delta \int_{\partial B} \varphi dH_{n-1}$ ($\delta \geq 0$) attains its minimum at $\delta = 0$, from which we obtain

$$\begin{aligned} &\int_B \frac{Du_{\eta,\lambda} \cdot D\varphi}{\sqrt{1 + |Du_{\eta,\lambda}|^2}} dydt + \eta \int_B u_{\eta,\lambda} \varphi dydt \\ &+ \lambda \int_B \varphi dydt + \int_{\partial B} \varphi dH_{n-1} + \int_B t\varphi dydt \geq 0 . \end{aligned}$$

In particular, choosing $\varphi \equiv 1$ we have

$$\eta \int_B u_{\eta, \lambda} dydt + \lambda |B| \geq -H_{n-1}(\partial B) - \int_B t dydt$$

that is

$$\begin{aligned} (A.4) \quad & \eta \int_B (u_{\eta, \lambda} - \psi) dydt + \lambda |B| \\ & \geq -H_{n-1}(\partial B) - \int_B t dydt - \eta \int_B \psi dydt \geq c_2(\psi, B) \end{aligned}$$

and hence (A.2).

Proof of Theorem A.1. In view of Lemma A.3, for every $\eta > 0$ there exists λ_η s.t. $\int_B (u_{\eta, \lambda_\eta} - \psi) dydt = V$. Moreover, from (A.3), (A.4) we have:

$$(A.5) \quad \frac{c_2 - \eta V}{|B|} \leq \lambda_\eta \leq \frac{c_1 + \int_{\partial B} |f| dH_{n-1}}{V} + c$$

hence we can assume $\lambda_\eta \rightarrow \lambda$ as $\eta \rightarrow 0$. There follows that u_{η, λ_η} converges in $L^1(B)$ to a function u_0 minimizing $I_{0, \lambda}$ in the class K .

From the relation

$$\int_B u_0 dydt = \int_B f dydt$$

we derive $I_{0, \lambda}(u_0) = I_{0, \lambda}(f)$, which concludes the proof of the theorem.

REMARK A.4. From (A.5) we easily derive the following estimates concerning multipliers λ :

$$(A.6) \quad \lambda \geq C_2(B) = \frac{-H_{n-1}(\partial B) - \int_B t dydt}{|B|}$$

$$\begin{aligned} (A.7) \quad \lambda & \leq \frac{C_1 + \int_{\partial B} |f| dH_{n-1}}{V} + c \\ & = \frac{\int_B \sqrt{1 + |D\psi|^2} + \int_{\partial B} |\psi| dH_{n-1} + \int_{\partial B} |f| dH_{n-1}}{V} + \max_B \{-t\}. \end{aligned}$$

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