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We show that the power 3/2 appearing in the estimate of analytic capacity is best possible.

1. Introduction. For a compact set E in the complex plane \mathbb{C} , $H^{\infty}(E^c)$ denotes the Banach space of bounded analytic functions in $E^c = \mathbb{C} \cup \{\infty\} - E$ with supremum norm $\|\cdot\|_{H^{\infty}}$. The analytic capacity of E is defined by

$$\gamma(E) = \sup\{|f'(\infty)|; \, \|f\|_{H^{\infty}} \le 1, \, f \in H^{\infty}(E^c)\},\$$

where $f'(\infty) = \lim_{z \to \infty} z(f(z) - f(\infty))$, i.e., $f'(\infty)$ is the (1/z)-coefficient of the Taylor expansion of f(z) at infinity. It is easily seen that $\gamma(E) \leq |E|$, where |E| is the (generalized) length of E; if E is a subset of the real line \mathbb{R} , then |E| equals its 1-dimension Lebesgue measure (cf. Garnett [4, Chap. III]). Vitushkin [12] constructed an example Q_{∞} such that $\gamma(Q_{\infty}) = 0$ and $|Q_{\infty}| > 0$ (cf. [4, p. 87]). Denjoy [3] showed that $\gamma(E) > 0$ if E is a subset of a rectifiable curve such that |E| > 0. But his proof has a serious gap, and his theorem was, for a while, called the Denjoy conjecture. As is easily seen, we may assume that E is a subset of a rectifiable graph. Let prE denote the projection of E to \mathbb{R} . Since pr is a contraction [6, p. 377], it is natural to try the lower estimate of $\gamma(E)$ by $\gamma(\operatorname{pr} E)$. Pommerenke [11] showed that $\gamma(\operatorname{pr} E) = |\operatorname{pr} E|/4$. Hence this approach is equivalent to comparing $\gamma(E)$ with |pr E|. To do this, the study of the Cauchy-Hilbert transform on C^1 graphs is necessary (Davie [2]). In 1977, Calderón [1] succeeded in proving its boundedness, and, using his theorem, Marshall [8] finally settled the Denjoy conjecture in the affirmative. After Marshall's theorem, we are concerned with studying further relations between $\gamma(E)$ and $|\operatorname{pr} E|$. Using an estimate of the Cauchy-Hilbert transform on Lipschitz graphs [10, p. 53], the author [9] showed that

$$\gamma(E) \ge C_0 |\operatorname{pr} E|^{3/2}$$

if E is a subset of a rectifiable graph Γ satisfying $|\Gamma| = 1$, where C_0 is an absolute constant. The main purpose of this paper is to show

that the power 3/2 is best possible. Our method gives a new approach to the computation of analytic capacity, and suggests that analytic capacity is related to the theory of fractals (Mandelbrot [7]).

For an integer $p \ge 2$, we put

$$B_p(x) = \frac{1}{2p} \{1 - (-1)^k\} \qquad \left(\frac{k}{p} \le x < \frac{k+1}{p}, \ 0 \le k \le p-1\right).$$

For an *n*-tuple (p_1, \ldots, p_n) of integers larger than or equal to 2, we put

$$A(x; p_1, ..., p_n) = \sum_{j=1}^n B_{p_1 \cdots p_j}(x)$$

A set $\Gamma \subset \mathbb{C}$ is called a crank of degree *n* if it is expressed in the form

$$\Gamma = \Gamma(p_1, \ldots, p_n) = \{x + iA(x; p_1, \ldots, p_n); 0 \le x < 1\}$$

for some *n*-tuple (p_1, \ldots, p_n) of integers larger than or equal to 2. (The class of cranks in this paper is smaller than a class defined in [10, Chap. III].) We shall show

THEOREM. For any $n \ge 1$, there exists a crank Γ_n of degree n such that

$$\frac{1}{C_1}\frac{1}{\sqrt{n}} \le \gamma(\Gamma_n) \le C_1 \frac{1}{\sqrt{n}} \,,$$

where C_1 is an absolute constant.

Once this theorem is established, we can deduce the exactness of the power 3/2 as follows. Adding some segments (perpendicular to the x-axis) to Γ_n , we obtain an arc connecting 0 and 1. Then the length of this arc is less than or equal to n+1. Hence we can define a rectifiable graph Γ'_n so that $|\Gamma'_n| \leq 3n$, $|\operatorname{pr} E'_n| \geq 1/2$, where $E'_n = \Gamma_n \cap \Gamma'_n$. Then $\gamma(E'_n) \leq \gamma(\Gamma_n) \leq C_1/\sqrt{n}$. Contracting E'_n , Γ'_n , we define E''_n , Γ''_n so that $|\Gamma''_n| = 1$. Then

$$\begin{split} \gamma(E_n'') &= \gamma(E_n')/|\Gamma_n'| \le \sqrt{3}C_1 |\Gamma_n'|^{-3/2} \\ &\le 2^{3/2}\sqrt{3}C_1 \{|\operatorname{pr} E_n'|/|\Gamma_n'|\}^{3/2} = 2^{3/2}\sqrt{3}C_1 |\operatorname{pr} E_n''|^{3/2}, \end{split}$$

which shows that the power 3/2 cannot be replaced by any number less than 3/2.

To prove our theorem, it is necessary to investigate cranks carefully. In $\S2$, we shall give a formula ((1) in Proposition 1) to compute analytic capacity. Proposition 2 is a generalization of Garnett's example [4, p. 87], and will be used to prove our theorem. Using the THE POWER 3/2

method in the proof of the formula, we shall, in $\S3$, give the proof of our theorem. In the last section, we shall give a new proof of Pommerenke's theorem [11] as another application of Proposition 1; our method shows how to construct the extremal functions.

2. A formula for the computation of $\gamma(\cdot)$. Let $L^2(\Gamma)$ denote the L^2 space of functions on a finite union Γ of smooth arcs with respect to the length element |dz|. The norm is denoted by $\|\cdot\|_{L^2(\Gamma)}$. The Cauchy-Hilbert transform \mathscr{H}_{Γ} from $L^2(\Gamma)$ to itself is defined by

$$\mathscr{H}_{\Gamma}f(z) = \frac{1}{\pi} \operatorname{p.v.} \int_{\Gamma} \frac{f(w)}{w-z} |dw|,$$

where p.v. is the principal value. This is a bounded operator and the norm is denoted by $\|\mathscr{H}_{\Gamma}\|_{L^{2}(\Gamma), L^{2}(\Gamma)}$. An operator $\overline{\mathscr{H}}_{\Gamma}$ is defined by $\overline{\mathscr{H}}_{\Gamma}f = \overline{\mathscr{H}_{\Gamma}\overline{f}}$, and \mathscr{I}_{Γ} is the identity operator. We show

PROPOSITION 1. Let Γ be a finite union of smooth arcs. Then, for any $0 < \varepsilon < 1/||\mathscr{H}_{\Gamma}||_{L^{2}(\Gamma), L^{2}(\Gamma)}$,

(1)
$$\gamma(\Gamma) = \frac{1}{\pi} \left\{ |\Gamma| + \sum_{m=0}^{\infty} (1 - \varepsilon^2)^m \sum_{l=0}^{\infty} \varepsilon^{2l+2} \frac{(l+1)(l+2)\cdots(l+m)}{m!} d_{2l+2}(\mathscr{H}_{\Gamma}) \right\},$$

where

$$d_{2l}(\mathscr{H}_{\Gamma}) = \int_{\Gamma} (\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma})^{l} 1 |dz| \qquad (l \ge 0, \ (\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma})^{0} = \mathscr{I}_{\Gamma})$$

and $(l+1)\cdots(l+m)/m! = 1$ if m = 0. (First $\sum_{l=0}^{\infty}$ is taken, and next $\sum_{m=0}^{\infty}$ is taken.) If $\|\mathscr{I}_{\Gamma} + \overline{\mathscr{H}}_{\Gamma}\mathscr{H}_{\Gamma}\|_{L^{2}(\Gamma), L^{2}(\Gamma)} < 2$, then

(2)
$$\gamma(\Gamma) = \frac{1}{\pi} \left\{ |\Gamma| + \sum_{m=0}^{\infty} 2^{-m-1} \sum_{l=0}^{\infty} \binom{m}{l} d_{2l+2}(\mathscr{H}_{\Gamma}) \right\},$$

where $\binom{0}{0} = 1$. If $\lim_{l\to\infty} d_{2l}(\mathscr{H}_{\Gamma}) = 0$, then

(3)
$$\gamma(\Gamma) = \frac{1}{\pi} \sum_{l=0}^{\infty} d_{2l}(\mathscr{H}_{\Gamma}).$$

This is a version of Garabedian's theorem [4, p. 22] to \mathscr{H}_{Γ} . Equality (3) is applicable to give a new proof of Pommerenke's theorem. (See §4.) Notice that $\mathscr{I}_{\mathbb{R}} + \widetilde{\mathscr{H}}_{\mathbb{R}} \mathscr{H}_{\mathbb{R}} = 0$, where $\mathscr{H}_{\mathbb{R}}$ is the Hilbert transform on \mathbb{R} . Hence (2) is applicable to compact sets Γ on a Lipschitz graph which is a small perturbation of \mathbb{R} . For any M > 0, there exists a

crank Γ' such that $d_2(\mathscr{H}_{\Gamma'}) \ge M$ [10, p. 84]. Then Cauchy-Schwarz' inequality yields that

$$d_{2^{l}}(\mathscr{H}_{\Gamma'}) \ge d_{2}(\mathscr{H}_{\Gamma'})^{2^{l-1}} \ge M^{2^{l-1}} \qquad (l \ge 1).$$

Hence (1) is necessary in this case.

Proof of Proposition 1. Let

(4)
$$\gamma^*(\mathscr{H}_{\Gamma}) = \inf\{\|1 + \mathscr{H}_{\Gamma}h\|_{L^2(\Gamma)}^2 + \|h\|_{L^2(\Gamma)}^2; h \in L^2(\Gamma)\}.$$

We begin by showing that

(5)
$$\gamma(\Gamma) = \frac{1}{\pi} \gamma^* (\mathscr{H}_{\Gamma}).$$

For a compact set E bounded by a finite number of smooth Jordan curves, we have

(6)
$$\gamma(E) = \frac{1}{2\pi} \inf \left\{ \int_{\partial E} |g(z)|^2 |dz|; g(\infty) = 1, g \text{ is analytic in } E^c \right\}$$

[4, p. 22]. Hence a standard argument yields that (6) holds with E replaced by Γ ; in this case, the boundary $\partial\Gamma$ has two sides. We define a smooth curve \mathscr{L} tending to infinity so that $\Gamma \subset \mathscr{L}$ and that $\mathscr{L} = \mathbb{R}$ outside a large disk. Then \mathscr{L} divides \mathbb{C} into two domains Ω_{\pm} . For an analytic function g(z) in Γ^c such that $g(\infty) = 1$ and $\int_{\partial\Gamma} |g(z)|^2 |dz| < \infty$, we can write

$$g(z) = 1 + \frac{1}{\pi} \int_{\Gamma} \frac{h(w)}{w - z} \, dw \,,$$

where the orientation of dw is chosen so that Ω_+ lies to the left. Let $g_{\pm}(z)$ be the nontangential limits of g at $z \in \Gamma$ with respect to Ω_{\pm} , respectively. Then

$$g_{+}(z) = 1 + \frac{1}{\pi} \operatorname{p.v.} \int_{\Gamma} \frac{h(w)}{w - z} dw + ih(z)$$

= 1 + $\mathscr{H}_{\Gamma}(h\psi)(z) + ih(z)$ $(z \in \Gamma)$,

where $\psi(z) = dz/|dz|$. Analogously,

$$g_{-}(z) = 1 + \mathscr{H}_{\Gamma}(h\psi)(z) - ih(z) \qquad (z \in \Gamma).$$

Thus

$$\begin{split} \int_{\partial\Gamma} |g(z)|^2 |dz| &= \|g_+\|_{L^2(\Gamma)}^2 + \|g_-\|_{L^2(\Gamma)}^2 \\ &= \|1 + \mathscr{H}_{\Gamma}(h\psi) + ih\|_{L^2(\Gamma)}^2 + \|1 + \mathscr{H}_{\Gamma}(h\psi) - ih\|_{L^2(\Gamma)}^2 \\ &= 2\{\|1 + \mathscr{H}_{\Gamma}(h\psi)\|_{L^2(\Gamma)}^2 + \|h\|_{L^2(\Gamma)}^2\} \\ &= 2\{\|1 + \mathscr{H}_{\Gamma}(h\psi)\|_{L^2(\Gamma)}^2 + \|h\psi\|_{L^2(\Gamma)}^2\} \end{split}$$

because $|\psi(z)| = 1$ $(z \in \Gamma)$. This shows that the quantity in the right-hand side of (6) $(E = \Gamma)$ equals $\frac{1}{\pi}\gamma^*(\mathscr{H}_{\Gamma})$, i.e., (5) holds.

We next compute $\gamma^*(\mathscr{H}_{\Gamma})$. Fatou's lemma shows that there exists $h_{\Gamma} \in L^2(\Gamma)$ which attains the infimum in (4). A variational method yields that $(1 + \mathscr{H}_{\Gamma}h_{\Gamma}, \mathscr{H}_{\Gamma}h) + (h_{\Gamma}, h) = 0$ for all $h \in L^2(\Gamma)$, where (\cdot, \cdot) is the (complex) inner product with respect to |dz|. Since the adjoint operator of \mathscr{H}_{Γ} is $-\widetilde{\mathscr{H}}_{\Gamma}$, this shows that

(7)
$$(\mathscr{F}_{\Gamma} - \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}) h_{\Gamma} = \overline{\mathscr{H}}_{\Gamma} \mathbf{1}.$$

Suppose that $h'_{\Gamma} \in L^2(\Gamma)$ also attains the infimum in (4). Then h'_{Γ} satisfies (7), and hence

$$0 = ((\mathscr{I}_{\Gamma} - \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma})(h_{\Gamma} - h'_{\Gamma}), h_{\Gamma} - h'_{\Gamma})$$

= $\|h_{\Gamma} - h'_{\Gamma}\|^2_{L^2(\Gamma)} + \|\mathscr{H}_{\Gamma}(h_{\Gamma} - h'_{\Gamma})\|^2_{L^2(\Gamma)}$

This shows that $h'_{\Gamma} = h_{\Gamma}$. Thus h_{Γ} is uniquely determined. By (7), we have

(8)
$$\gamma^{*}(\mathscr{H}_{\Gamma}) = \|1 + \mathscr{H}_{\Gamma}h_{\Gamma}\|_{L^{2}(\Gamma)}^{2} + \|h_{\Gamma}\|_{L^{2}(\Gamma)}^{2}$$
$$= (1 + \mathscr{H}_{\Gamma}h_{\Gamma}, 1) + ((\mathscr{H}_{\Gamma} - \overline{\mathscr{H}}_{\Gamma}\mathscr{H}_{\Gamma})h_{\Gamma} - \overline{\mathscr{H}}_{\Gamma}1, h_{\Gamma})$$
$$= \int_{\Gamma} \{1 + \mathscr{H}_{\Gamma}h_{\Gamma}\}|dz|.$$

Let

$$T_{\Gamma} = (\mathscr{I}_{\Gamma} - \varepsilon^2 \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma})^{-1}.$$

Then we can write

$$T_{\Gamma} = \sum_{l=0}^{\infty} \varepsilon^{2l} (\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma})^{l}$$

because $0 < \varepsilon < 1/||\mathscr{H}_{\Gamma}||_{L^{2}(\Gamma), L^{2}(\Gamma)}$. We have, for any $h \in L^{2}(\Gamma)$,

$$\begin{split} \|T_{\Gamma}h\|_{L^{2}(\Gamma)}^{2} &\leq \|T_{\Gamma}h\|_{L^{2}(\Gamma)}^{2} + \varepsilon^{2}\|\mathscr{H}_{\Gamma}T_{\Gamma}h\|_{L^{2}(\Gamma)}^{2} \\ &= ((\mathscr{I}_{\Gamma} - \varepsilon^{2}\widetilde{\mathscr{H}}_{\Gamma}\mathscr{H}_{\Gamma})T_{\Gamma}h, \ T_{\Gamma}h) = (h, \ T_{\Gamma}h) \leq \|h\|_{L^{2}(\Gamma)}\|T_{\Gamma}h\|_{L^{2}(\Gamma)}, \end{split}$$

which shows that $||T_{\Gamma}||_{L^{2}(\Gamma), L^{2}(\Gamma)} \leq 1$. Equality (7) can be rewritten as

(9)
$$(\mathscr{I}_{\Gamma} - \varepsilon^2 \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma}) h_{\Gamma} = (1 - \varepsilon^2) h_{\Gamma} + \varepsilon^2 \overline{\mathscr{H}}_{\Gamma} 1.$$

Observing this equality, we inductively define $(h_m)_{m=0}^{\infty}$ by $h_0 = 0$,

$$h_m = T_{\Gamma}\{(1-\varepsilon^2)h_{m-1} + \varepsilon^2 \overline{\mathscr{H}}_{\Gamma}1\} \qquad (m \ge 1).$$

Then

$$\|h_{m+1} - h_m\|_{L^2(\Gamma)} = (1 - \varepsilon^2) \|T_{\Gamma}(h_m - h_{m-1})\|_{L^2(\Gamma)}$$

$$\leq (1 - \varepsilon^2) \|h_m - h_{m-1}\|_{L^2(\Gamma)}.$$

Hence $\lim_{m\to\infty} h_m$ exists and satisfies (9), i.e., (7). Thus $h_{\Gamma} = \lim_{m\to\infty} h_m$. Since

$$h_{m+1} - h_m = (1 - \varepsilon^2) T_{\Gamma} (h_m - h_{m-1}) = \dots = (1 - \varepsilon^2)^m T_{\Gamma}^m h_1$$

= $\varepsilon^2 (1 - \varepsilon^2)^m T_{\Gamma}^{m+1} \overline{\mathscr{H}}_{\Gamma} 1$,

we have

$$h_{\Gamma} = \sum_{m=0}^{\infty} (h_{m+1} - h_m) = \varepsilon^2 \sum_{m=0}^{\infty} (1 - \varepsilon^2)^m T_{\Gamma}^{m+1} \overline{\mathscr{H}}_{\Gamma} 1.$$

Consequently, (8) yields that

$$\begin{split} \gamma^*(\mathscr{H}_{\Gamma}) &= |\Gamma| + \sum_{m=0}^{\infty} (1 - \varepsilon^2)^m \int_{\Gamma} \varepsilon^2 \mathscr{H}_{\Gamma} T_{\Gamma}^{m+1} \overline{\mathscr{H}}_{\Gamma} 1 |dz| \\ &= |\Gamma| + \sum_{m=0}^{\infty} (1 - \varepsilon^2)^m \int_{\Gamma} \varepsilon^2 \mathscr{H}_{\Gamma} \left\{ \sum_{l=0}^{\infty} \varepsilon^{2l} (\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma})^l \right\}^{m+1} \overline{\mathscr{H}}_{\Gamma} 1 |dz| \\ &= |\Gamma| + \sum_{m=0}^{\infty} (1 - \varepsilon^2)^m \sum_{l=0}^{\infty} \varepsilon^{2l+2} \frac{(l+1)\cdots(l+m)}{m!} \int_{\Gamma} (\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma})^{l+1} 1 |dz| \\ &= |\Gamma| + \sum_{m=0}^{\infty} (1 - \varepsilon^2)^m \sum_{l=0}^{\infty} \varepsilon^{2l+2} \frac{(l+1)\cdots(l+m)}{m!} d_{2l+2} (\mathscr{H}_{\Gamma}). \end{split}$$

Using (5), we obtain (1).

We can write

$$\mathcal{I}_{\Gamma} - \overline{\mathcal{H}}_{\Gamma} \mathcal{H}_{\Gamma} = 2\{\mathcal{I}_{\Gamma} - \frac{1}{2}(\mathcal{I}_{\Gamma} + \overline{\mathcal{H}}_{\Gamma} \mathcal{H}_{\Gamma})\}.$$

Hence, if $\|\mathscr{F}_{\Gamma} + \overline{\mathscr{H}}_{\Gamma}\mathscr{H}_{\Gamma}\|_{L^{2}(\Gamma), L^{2}(\Gamma)} < 2$, then

$$h_{\Gamma} = \sum_{m=0}^{\infty} 2^{-m} (\mathscr{I}_{\Gamma} + \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma})^m (\frac{1}{2} \overline{\mathscr{H}}_{\Gamma}) 1.$$

Thus (5) and (8) yield (2).

Equality (7) shows that $\mathscr{H}_{\Gamma}h_{\Gamma} = \mathscr{H}_{\Gamma}\overline{\mathscr{H}}_{\Gamma}1 + \mathscr{H}_{\Gamma}\overline{\mathscr{H}}_{\Gamma}h_{\Gamma}$, and hence, by (8),

$$\gamma^*(\mathscr{H}_{\Gamma}) = \int_{\Gamma} \{1 + \mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma} 1 + \mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma} h_{\Gamma} \} |dz|.$$

Repeating this argument, we have

$$\begin{split} \gamma^*(\mathscr{H}_{\Gamma}) &= \int_{\Gamma} \left\{ \sum_{l=0}^{L} (\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma})^l 1 + \mathscr{H}_{\Gamma} (\overline{\mathscr{H}}_{\Gamma} \mathscr{H}_{\Gamma})^L h_{\Gamma} \right\} |dz| \\ &= \sum_{l=0}^{L} d_{2l} (\mathscr{H}_{\Gamma}) - \int_{\Gamma} \{ (\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma})^L \mathscr{H}_{\Gamma} \} 1(z) h_{\Gamma}(z) |dz|. \end{split}$$

If $\lim_{L\to\infty} d_{2L}(\mathscr{H}_{\Gamma}) = 0$, then

$$\begin{split} \lim_{L \to \infty} \left| \int_{\Gamma} \{ (\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma})^{L} \mathscr{H}_{\Gamma} \} \mathbf{1}(z) h_{\Gamma}(z) |dz| \right| \\ &\leq \lim_{L \to \infty} \| (\mathscr{H}_{\Gamma} \overline{\mathscr{H}}_{\Gamma})^{L} \mathscr{H}_{\Gamma} \mathbf{1} \|_{L^{2}(\Gamma)} \| h_{\Gamma} \|_{L^{2}(\Gamma)} \\ &= \lim_{L \to \infty} d_{4L+2} (\mathscr{H}_{\Gamma})^{1/2} \| h_{\Gamma} \|_{L^{2}(\Gamma)} = 0. \end{split}$$

Hence (5) gives (3). This completes the proof of our proposition.

We now give a remark. There exists an analytic function $g_{\Gamma}(z)$ in Γ^c such that $g_{\Gamma}(\infty) = 1$ and $\gamma(\Gamma) = (1/2\pi) \int_{\partial \Gamma} |g_{\Gamma}(z)| |dz|$ [4, p. 19]. This is called the Garabedian function of Γ . Equality (5) shows that

$$g_{\Gamma}(z) = \left\{ 1 + \frac{1}{\pi} \int_{\Gamma} \frac{h_{\Gamma}(w)}{w - z} |dw| \right\}^2.$$

There exists $f_{\Gamma} \in H^{\infty}(\Gamma^c)$ such that $||f_{\Gamma}||_{H^{\infty}} = 1$ and $f'_{\Gamma}(\infty) = \gamma(\Gamma)$ [4, p. 18]. This is called the Ahlfors function of Γ . We have

(10)
$$f_{\Gamma}(z) = \frac{-\frac{1}{\pi} \left\{ \int_{\Gamma} \frac{|dw|}{w-z} + \int_{\Gamma} \frac{\overline{\mathscr{H}}_{\Gamma} \overline{h}_{\Gamma}(w)}{w-z} |dw| \right\}}{\left\{ 1 + \frac{1}{\pi} \int_{\Gamma} \frac{h_{\Gamma}(w)}{w-z} |dw| \right\}}$$

To see this, let f(z) denote the function in the right-hand side. Since $g_{\Gamma}(z)$ does not take 0 in Γ^c , f(z) is analytic in Γ^c [4, p. 21]. We have $f'(\infty) = \frac{1}{\pi} \gamma^*(\mathscr{H}_{\Gamma}) = \gamma(\Gamma)$ and

$$f_{\pm}(z) = -\frac{\mathscr{H}_{\Gamma}\mathbf{1}(z) \pm i\tilde{\psi}(z) + \mathscr{H}_{\Gamma}\overline{\mathscr{H}}_{\Gamma}\overline{h}_{\Gamma}(z) \pm i\overline{\mathscr{H}}_{\Gamma}\overline{h}_{\Gamma}(z)\tilde{\psi}(z)}{1 + \mathscr{H}_{\Gamma}h_{\Gamma}(z) \pm ih_{\Gamma}(z)\tilde{\psi}(z)},$$

where $\tilde{\psi}(z) = |dz|/dz$ and $f_{\pm}(z)$ are the nontangential limits of f at $z \in \Gamma$ with respect to Ω_{\pm} , respectively. Equality (7) shows that

$$\begin{split} \mathcal{H}_{\Gamma} 1 &+ i\tilde{\psi} + \mathcal{H}_{\Gamma} \overline{\mathcal{H}}_{\Gamma} \overline{h}_{\Gamma} + i(\overline{\mathcal{H}}_{\Gamma} \overline{h}_{\Gamma}) \tilde{\psi} \\ &= \mathcal{H}_{\Gamma} 1 + i\tilde{\psi} + (\overline{h}_{\Gamma} - \mathcal{H}_{\Gamma} 1) + i(\overline{\mathcal{H}}_{\Gamma} \overline{h}_{\Gamma}) \tilde{\psi} \\ &= i\tilde{\psi} + i(\overline{\mathcal{H}}_{\Gamma} \overline{h}_{\Gamma}) \tilde{\psi} + \overline{h}_{\Gamma} = i\tilde{\psi} \{ \overline{1 + \mathcal{H}_{\Gamma} h_{\Gamma} + ih_{\Gamma} \tilde{\psi}} \} \,, \end{split}$$

which yields that $|f_+(z)| = 1$ on Γ . Analogously, $|f_-(z)| = 1$ on Γ . Thus $||f||_{H^{\infty}} = 1$. This shows that $f = f_{\Gamma}$.

For the proof of our theorem, we note

PROPOSITION 2. Let $0 < \delta_0 < 1$ and let $(q_n)_{n=1}^{\infty}$ be a sequence of integers larger than or equal to 2 such that

$$\sum_{n=j}^{\infty} (q_j \cdots q_n)^{-1} \le \delta_0 \qquad (j \ge 1).$$

Then

$$\lim_{n\to\infty}\sup\gamma(\Gamma(p_1,\ldots,p_n))=0,$$

where the supremum is taken over all *n*-tuples (p_1, \dots, p_n) satisfying $p_j \ge q_j$ $(1 \le j \le n)$.

This is a generalization of Garnett's example [4, p. 87], and used later. Notice that $\sum_{n=1}^{\infty} 2^{-n} = 1$. A sequence $(\Gamma(\mathbf{2}_n))_{n=1}^{\infty}$ ($\mathbf{2}_n$ is the *n*tuple of 2) topologically converges to a segment $\{x + ix; 0 \le x < 1\}$, and these cranks behave like cranks of degree 1 with respect to this segment. Hence we have $\limsup_{n\to\infty} \gamma(\Gamma(\mathbf{2}_n)) > 0$. This shows that our proposition is sharp in a sense. Since a minor change of the argument in [10, p. 81] yields the required equality, we omit the proof (cf. Jones [5]).

3. Proof of Theorem. In this section, we give the proof of our theorem. Let L^q denote the L^q space of functions on [0, 1) with respect to the 1-dimension Lebesgue measure $|\cdot|$ $(1 \le q < \infty)$. For a kernel K = K(x, y) on $[0, 1) \times [0, 1)$, we simply write by the same notation K an operator defined by this kernel, and write by \overline{K} an operator defined by $\overline{K(x, y)}$; $||K||_{L^q, L^{q'}}$ denotes the norm of K as an operator from L^q to $L^{q'}$. The identity operator is denoted by Id. A kernel K is anti-symmetric if K(x, y) = -K(y, x) $(x \ne y)$. A kernel K is of type 0 if

$$\sup_{x,y\in[0,1)}\left\{|K(x,y)|+\left|\frac{\partial}{\partial x}K(x,y)\right|+\left|\frac{\partial}{\partial y}K(x,y)\right|\right\}<\infty.$$

A kernel K is of type 1 if $||K||_{L^4, L^4} < \infty$ and if there exists a sequence $(K_j)_{j=1}^{\infty}$ of kernels of type 0 such that

$$\lim_{j\to\infty} \|K_j - K\|_{L^4, L^2} = 0, \quad \sup_{j\geq 1} \|K_j\|_{L^4, L^4} < \infty.$$

Kernels used in this section are bounded as operators from L^q to itself for all $1 < q < \infty$. Let

$$\gamma^*(K) = \inf\{\|1 + Kh\|_{L^2}^2 + \|h\|_{L^2}^2; h \in L^2\},\$$

$$d_{2l}(K) = \int_0^1 (K\overline{K})^l 1 \, dx \qquad (l \ge 0, \ (K\overline{K})^0 = \mathrm{Id}).$$

Recall the function $A(x; p_1, \ldots, p_n)$ in the introduction. Let

$$H(x, y) = \mathscr{H}_{\mathbb{R}}(x, y) = \frac{1}{\pi} \frac{1}{y - x}$$

$$H[p_1, \ldots, p_n](x, y) = \frac{1}{\pi} \frac{1}{(y - x) + i(A(y; p_1, \ldots, p_n) - A(x; p_1, \ldots, p_n))},$$

 $\Delta[p_1, \dots, p_n] = H[p_1, \dots, p_n] - H[p_1, \dots, p_{n-1}] \qquad (n \ge 1),$ where $H[p_1, \dots, p_{n-1}] = H$ if n = 1. Then

$$H[p_1,\ldots,p_n]=H+\sum_{j=1}^n\Delta[p_1,\ldots,p_j].$$

Since all components/segments of $\Gamma(p_1, \dots, p_n)$ are parallel to the x-axis, we can identify $\mathscr{H}_{\Gamma(p_1,\dots,p_n)}$, $L^2(\Gamma(p_1,\dots,p_n))$ with $H[p_1,\dots,p_n]$, L^2 , respectively. We have $||H[p_1,\dots,p_n]||_{L^2,L^2} \leq C_2\sqrt{n}$ for some absolute constant C_2 [10, p. 84]. Hence Proposition 1 shows that

(11)
$$\gamma(\Gamma(p_1, \dots, p_n)) = \frac{1}{\pi} \gamma^* (H[p_1, \dots, p_n])$$

$$= \frac{1}{\pi} \left\{ 1 + \sum_{m=0}^{\infty} (1 - \varepsilon_n^2)^m \times \sum_{l=0}^{\infty} \varepsilon_n^{2l+2} \frac{(l+1)\cdots(l+m)}{m!} d_{2l+2} (H[p_1, \dots, p_n]) \right\},$$

where $\varepsilon_n = (2C_2\sqrt{n})^{-1}$. We shall inductively estimate

$$\lim_{p_1\to\infty}\cdots\lim_{p_n\to\infty}\gamma^*(H[p_1,\ldots,p_n]),$$

where $\lim_{p_n\to\infty}$ is taken first and $\lim_{p_1\to\infty}$ is taken last. For $E \subset \mathbb{R}$, χ_E denotes its characteristic function, and, for $x \in \mathbb{R}$, $\iota(x)$ denotes its integral part. Here are some lemmas necessary for the estimate.

LEMMA 3. For two kernels K and K',

$$\gamma^*(K+K') \leq 2(1+\|K'\|_{L^2,L^2}^2)\gamma^*(K).$$

Proof. We have, for any $h \in L^2$,

 $\|1 + (K + K')h\|_{L^2}^2 + \|h\|_{L^2}^2 \le 2(1 + \|K'\|_{L^2, L^2}^2) \{\|1 + Kh\|_{L^2}^2 + \|h\|_{L^2}^2\},$ which yields the required inequality.

LEMMA 4. Let K be an anti-symmetric kernel such that

$$\lim_{l\to\infty}d_{2l}(K)=0.$$

Then

(12)
$$\gamma^*(K) = \sum_{l=0}^{\infty} d_{2l}(K).$$

Since this is a version of (3) to K, we omit the proof.

LEMMA 5. For an anti-symmetric kernel K, $0 < \varepsilon_0 \le (3 \|K\|_{L^2, L^2})^{-1}$ and $w \in U = \{\zeta \in \mathbb{C}; |\zeta| < 2, |\arg \zeta| < \pi/4\}$,

(13)
$$1 + \sum_{m=0}^{\infty} (1 - \varepsilon_0^2)^m \sum_{l=0}^{\infty} w^{2l+2} \varepsilon_0^{2l+2} \frac{(l+1)\cdots(l+m)}{m!} d_{2l+2}(K)$$
$$(= \gamma^*(w; K), \ say)$$

exists and $\gamma^*(w; K)$ is analytic in U.

Proof. Let

$$T(w; K) = (\operatorname{Id} - w^2 \varepsilon_0^2 \overline{K} K)^{-1}.$$

Then

$$w^{2}\varepsilon_{0}^{2}\int_{0}^{1} KT(w;K)^{m+1}\overline{K} \, dx$$

= $\sum_{l=0}^{\infty} w^{2l+2}\varepsilon_{0}^{2l+2} \frac{(l+1)\cdots(l+m)}{m!} d_{2l+2}(K)$

because $2\varepsilon_0 ||K||_{L^2, L^2} < 1$. Evidently, this is analytic in U. Since K is anti-symmetric and $\operatorname{Re} w^2 > 0$ ($w \in U$, $\operatorname{Re} w^2$ is the real part of w^2), we have, in the same manner as in the proof of (1),

(14)
$$||T(w; K)||_{L^2, L^2} \le 1 \quad (w \in U).$$

Thus the convergence of $\sum_{m=0}^{\infty}$ in (13) is uniform in U, which shows that $\gamma^*(w; K)$ exists and is analytic in U.

LEMMA 6. For any $l \ge 0$,

$$\lim_{p \to \infty} d_{2l}(\Delta[p]) \quad (= d_{2l}(\Delta[\infty]), \ say)$$

exists and

(15)
$$d_2(\Delta[\infty]) \le -\frac{1}{25\pi^2}.$$

Proof. We put

$$R(s, t) = \frac{1}{\pi} \left\{ \frac{1}{t-s+1+i} - \frac{1}{t-s+1} \right\} - \frac{2}{\pi} \sum_{m=1}^{\infty} \left\{ \frac{t-s+1+i}{4m^2 - (t-s+1+i)^2} - \frac{t-s+1}{4m^2 - (t-s+1)^2} \right\}$$

and show that

(16)
$$\lim_{p \to \infty} d_{2l}(\Delta[p]) = \frac{1}{2} \int_0^1 \{R^{2l} 1 + \overline{R}^{2l} 1\} \, ds.$$

Let

$$W_p = \bigcup_{m, \text{ odd}} \left[\frac{m}{p}, \frac{m+1}{p} \right), \quad W'_p = \bigcup_{m, \text{ even}} \left[\frac{m}{p}, \frac{m+1}{p} \right),$$
$$X_p = \bigcup_{m=l(\log p)}^{p-l(\log p)-1} \left[\frac{m}{p}, \frac{m+1}{p} \right),$$
$$s_x = px - l(px) \quad (0 \le x < 1, \ p \ge 2).$$

Notice that $|[0, 1) - X_p| \le 2 \ \iota(\log p)/p$ and $||\Delta[p]||_{L^4, L^4} \le 10$. Since $\Delta[p](x, y) = 0 \ (x, y \in W_p; x, y \in W'_p)$, we have

$$\begin{split} d_{2l}(\Delta[p]) &= \int_0^1 (\Delta[p]\overline{\Delta[p]})^{l-1} \Delta[p] \{ \chi_{W_p} \overline{\Delta[p]} \chi_{W'_p} + \chi_{W'_p} \overline{\Delta[p]} \chi_{W_p} \} \, dx \\ &= \int_0^1 (\Delta[p]\overline{\Delta[p]})^{l-1} \Delta[p] \{ \chi_{W_p \cap X_p} \overline{\Delta[p]} \chi_{W'_p} + \chi_{W'_p \cap X_p} \overline{\Delta[p]} \chi_{W_p} \} \, dx \\ &+ O\left(\left(\frac{\log p}{p} \right)^{1/4} \right). \end{split}$$

We now study $\overline{\Delta[p]}\chi_{W'_p}(x)$ $(x \in W_p \cap X_p)$. Without loss of generality, we may assume that p is even. Since $x \in W_p \cap X_p$, $\iota(px)$ is even and

 $\iota(\log p) \le \iota(px) \le p - \iota(\log p) - 1$. We may assume that $\iota(\log p) \le \iota(px) \le p/2$. We have

$$\begin{split} \overline{\Delta[p]}\chi_{W_p'}(x) &= \frac{1}{\pi} \int_{W_p'} \left\{ \frac{1}{y - x - i/p} - \frac{1}{y - x} \right\} dy \\ &= \frac{1}{\pi} \int_{W_p} \left\{ \frac{1}{y - x + 1/p - i/p} - \frac{1}{y - x + 1/p} \right\} dy \\ &= \frac{1}{\pi} \sum_{m=0}^{(p/2)-1} \int_0^{1/p} \left\{ \frac{1}{(2m/p + y) - (i(px)/p + x - i(px)/p) + 1/p - i/p} \right. \\ &\qquad - \frac{1}{(2m/p + y) - (i(px)/p + x - i(px)/p) + 1/p} \right\} dy \\ &= \frac{1}{\pi} \sum_{m=0}^{i(px)} + \frac{1}{\pi} \sum_{m=i(px)+1}^{(p/2)-1} = L_1 + L_2 \,, \end{split}$$

$$\begin{split} L_1 &= \frac{1}{\pi} \int_0^1 \left\{ \frac{1}{t - s_x + 1 - i} - \frac{1}{t - s_x + 1} \right\} dt + \frac{1}{\pi} \sum_{0 \le m \le i(px), \ m \ne i(px)/2} \\ &= \frac{1}{\pi} \int_0^1 \left\{ \frac{1}{t - s_x + 1 - i} - \frac{1}{t - s_x + 1} \right\} dt \\ &- \frac{2}{\pi} \sum_{m=1}^{i(px)/2} \int_0^1 \left\{ \frac{t - s_x + 1 - i}{4m^2 - (t - s_x + 1 - i)^2} \right. \\ &- \frac{t - s_x + 1}{4m^2 - (t - s_x + 1)^2} \right\} dt \\ &= \overline{R} 1(s_x) + O\left(\frac{1}{\log p}\right), \\ L_2 &= -\frac{i}{\pi} \sum_{m=i(px)+1}^{(p/2)-1} \int_0^1 \frac{1}{(2m - i(px)) + (t - s_x + 1 - i)} \\ &\times \frac{dt}{(2m - i(px)) + (t - s_x + 1)} = O\left(\frac{1}{\log p}\right), \end{split}$$

which shows that $\overline{\Delta[p]}\chi_{W'_p}(x) = \overline{R}1(s_x) + O(1/\log p)$ $(x \in W_p \cap X_p)$.

In the same manner, $\overline{\Delta[p]}\chi_{W_p}(x) = R1(s_x) + O(1/\log p)$ $(x \in W'_p \cap X_p)$. Thus

$$\begin{split} d_{2l}(\Delta[p]) &= \int_0^1 (\Delta[p]\overline{\Delta[p]})^{l-1} \Delta[p] \{ \chi_{W_p \cap X_p} \overline{R} 1(s_{\cdot}) + \chi_{W'_p \cap X_p} R 1(s_{\cdot}) \} \, dx \\ &+ O\left(\frac{1}{\log p}\right) \\ &= \int_0^1 (\Delta[p]\overline{\Delta[p]})^{l-1} \Delta[p] \{ \chi_{W_p} \overline{R} 1(s_{\cdot}) + \chi_{W'_p} R 1(s_{\cdot}) \} \, dx \\ &+ O\left(\frac{1}{\log p}\right) \\ &= \int_0^1 (\Delta[p]\overline{\Delta[p]})^{l-1} \{ \chi_{W_p \cap X_p} \Delta[p] (\chi_{W'_p} R 1(s_{\cdot})) \\ &+ \chi_{W'_p \cap X_p} \Delta[p] (\chi_{W_p} \overline{R} 1(s_{\cdot})) \} \, dx \\ &+ O\left(\frac{1}{\log p}\right). \end{split}$$

Since $R1(s_x)$ is a periodic function with period 1/p, we have, in the same manner as above,

$$\Delta[p](\chi_{W'_p}R1(\underline{s}))(x) = R^2 1(s_x) + O\left(\frac{1}{\log p}\right) \qquad (x \in W_p \cap X_p),$$

$$\Delta[p](\chi_{W_p}\overline{R}1(\underline{s}))(x) = \overline{R}^2 1(s_x) + O\left(\frac{1}{\log p}\right) \qquad (x \in W'_p \cap X_p).$$

Repeating this argument, we have

$$d_{2l}(\Delta[p]) = \int_0^1 \{\chi_{W_p}(x) R^{2l} 1(s_x) + \chi_{W'_p}(x) \overline{R}^{2l} 1(s_x)\} dx + O\left(\frac{1}{\log p}\right)$$
$$= \frac{1}{2} \int_0^1 \{R^{2l} 1 + \overline{R}^{2l} 1\} ds + O\left(\frac{1}{\log p}\right),$$

which gives (16).

We have

$$R(s, t) = \frac{1}{\pi} \sum_{m=-\infty}^{\infty} \left\{ \frac{1}{2m+1+t-s+i} - \frac{1}{2m+1+t-s} \right\}$$
$$= -\frac{1}{\pi} \sum_{m=-\infty}^{\infty} \frac{1}{(2m+1+t-s)\{(2m+1+t-s)^2+1\}}$$
$$-\frac{i}{\pi} \sum_{m=-\infty}^{\infty} \frac{1}{(2m+1+t-s)^2+1}$$
$$= -R'(s, t) - iR''(s, t), \text{ say.}$$

Then R' is anti-symmetric and R'' is symmetric, i.e., R''(s, t) = R''(t, s). Thus

$$d_{2}(\Delta[\infty]) = \operatorname{Re} \int_{0}^{1} R^{2} 1 \, ds$$

= $\operatorname{Re} \int_{0}^{1} (-R'1 + iR''1)(R'1 + iR''1) \, ds$
= $-\int_{0}^{1} \{(R'1)^{2} + (R''1)^{2}\} \, ds$
 $\leq -\frac{1}{\pi^{2}} \int_{0}^{1} \left\{ \int_{0}^{1} \frac{dt}{(1+t-s)^{2}+1} \right\}^{2} \, ds \leq -\frac{1}{25\pi^{2}}$

Thus (15) holds.

LEMMA 7. Let K be an anti-symmetric kernel of type 1, and let $(g_p)_{p=2}^{\infty}$, $(h_p)_{p=2}^{\infty}$ be two sequences in L^4 such that $||g_p||_{L^4} \leq 1$, $||h_p||_{L^4} \leq 1$. Then, for any $l \geq 0$,

(17)
$$\lim_{p \to \infty} \left\{ \int_0^1 K g_p \cdot (\Delta[p] \overline{\Delta[p]})^l \overline{K} h_p \, dx - d_{2l} (\Delta[\infty]) \int_0^1 K g_p \cdot \overline{K} h_p \, dx \right\} = 0,$$

(18)
$$\lim_{p\to\infty}\int_0^1 Kg_p\cdot\overline{\Delta[p]}(\Delta[p]\overline{\Delta[p]})^l Kh_p\,dx=0.$$

Equalities (17) and (18) hold with Kg_p replaced by 1.

Proof. First we assume that K is of type 0. Let

$$\Delta'[p](x, y) = \Delta[p](x, y)\chi_{[0, N/p)}(|y - x|) \qquad (p \ge 2),$$

where $N = \iota(\log p)$. Then $\|\Delta[p] - \Delta'[p]\|_{L^2, L^2} = O(1/\log p)$ (cf. Lemma 6), and hence

$$\int_0^1 K g_p \cdot (\Delta[p]\overline{\Delta[p]})^l \overline{K} h_p \, dx$$

= $\int_0^1 K g_p \cdot (\Delta'[p]\overline{\Delta'[p]})^l \overline{K} h_p \, dx + O\left(\frac{1}{\log p}\right).$

Notice that

$$(\Delta'[p]\overline{\Delta'[p]})^l(x, y) = 0 \qquad (|y - x| > 2lN/p),$$

and that $(\Delta'[p]\overline{\Delta'[p]})^l 1$ is a periodic function on [2lN/p, 1-(2lN/p)) with period 2/p. Let

$$\eta_p^{(m)} = \frac{p}{2} \int_{2m/p}^{2(m+1)/p} (\Delta'[p] \overline{\Delta'[p]})^l \, 1 \, dx - d_{2l}(\Delta[\infty])$$
$$(0 \le m \le \iota(p/2) - 1).$$

Then $\eta_p^{(m)} = \eta_p^{(lN)}$ $(lN \le m \le \iota(p/2) - lN - 1)$. Lemma 6 shows that

$$\limsup_{p \to \infty} |\eta_p^{(lN)}| = \limsup_{p \to \infty} |d_{2l}(\Delta'[p]) - d_{2l}(\Delta[\infty])|$$
$$= \limsup_{p \to \infty} |d_{2l}(\Delta[p]) - d_{2l}(\Delta[\infty])| = 0.$$

Since K is of type 0, we have

$$\sup \frac{|Kh(y) - Kh(x)|}{|y - x|} \leq \sup_{s, t \in [0, 1]} \left| \frac{\partial}{\partial s} K(s, t) \right| < \infty,$$

where the supremum in the left-hand side is taken over all $x, y \in [0, 1)$ and all $h \in L^4$ satisfying $||h||_{L^4} \leq 1$. Thus

$$\begin{split} &\int_{0}^{1} Kg_{p} \cdot (\Delta[p]\overline{\Delta[p]})^{l}\overline{K}h_{p} dx \\ &= \int_{2lN/p}^{1-(2lN/p)} Kg_{p}(x) \\ &\times \int_{|y-x| \leq 2lN/p} (\Delta'[p]\overline{\Delta'[p]})^{l}(x, y)\overline{K}h_{p}(y) dy dx + o(1) \\ &= \int_{2lN/p}^{1-(2lN/p)} Kg_{p} \cdot \overline{K}h_{p} \cdot (\Delta'[p]\overline{\Delta'[p]})^{l} 1 dx + o(1) \\ &= \frac{2}{p} \sum_{m=lN}^{i(p/2)-lN-1} Kg_{p} \left(\frac{2m}{p}\right) \overline{K}h_{p} \left(\frac{2m}{p}\right) \{\eta_{p}^{(m)} + d_{2l}(\Delta[\infty])\} + o(1) \\ &= d_{2l}(\Delta[\infty]) \int_{0}^{1} Kg_{p} \cdot \overline{K}h_{p} dx + O(\eta_{p}^{(lN)}) + o(1) \\ &= d_{2l}(\Delta[\infty]) \int_{0}^{1} Kg_{p} \cdot \overline{K}h_{p} dx + o(1) \,, \end{split}$$

which shows that (17) holds. Let K be of type 1. Then there exists a

sequence $(K_j)_{i=1}^{\infty}$ of kernels of type 0 such that

$$\|K - K_j\|_{L^4, L^2} \le \frac{1}{j}$$
 $(j \ge 1), \quad \sup_{j \ge 1} \|K_j\|_{L^4, L^4} < \infty.$

Then

$$\left|\int_0^1 Kg_p \cdot (\Delta[p]\overline{\Delta[p]})^l \overline{K}h_p \, dx - \int_0^1 K_j g_p \cdot (\Delta[p]\overline{\Delta[p]})^l \overline{K}_j h_p \, dx\right| \leq C_3/j \,,$$

$$\left|\int_0^1 Kg_p \cdot \overline{K}h_p \, dx - \int_0^1 K_j g_p \cdot \overline{K}_j h_p \, dx\right| \le C_3/j \qquad (j \ge 1)$$

for some constant C_3 independent of p and j. Since (17) holds for all K_j $(j \ge 1)$, this shows that (17) holds.

Since $\overline{\Delta[p]}(\Delta[p]\overline{\Delta[p]})^l$ is anti-symmetric, we have

$$\int_0^1 \overline{\Delta[p]} (\Delta[p] \overline{\Delta[p]})^l 1 \, dx = 0.$$

Hence, in the same manner as above, we obtain (18). Analogously, we can replace Kg_p by 1.

LEMMA 8. Let K be an anti-symmetric kernel of type 1. Then, for any $l \ge 0$,

$$\lim_{p \to \infty} d_{2l}(\Delta[p] + K) \quad (= d_{2l}(\Delta[\infty] + K), \ say)$$

exists, and we can write

(19)
$$d_{2l}(\Delta[\infty] + K) = \sum_{k=0}^{l} c_{2k}^{(2l)} d_{2l-2k}(K)$$

so that $c_{2k}^{(m)}$ $(0 \le k \le \iota(m/2), m \ge 0)$ satisfy

(20)
$$c_0^{(m)} = 1$$
, $c_{2k}^{(2k)} = d_{2k}(\Delta[\infty])$ $(m \ge 0, k \ge 0)$,

(21)
$$c_{2k}^{(m)} = \sum_{j=0}^{k} c_{2k-2j}^{(m-2j-1)} d_{2j}(\Delta[\infty])$$

 $\left(0 \le k \le \iota\left(\frac{m-1}{2}\right), \ m \ge 0\right).$

Proof. We say that a 2*l*-tuple $(\tau_1, \ldots, \tau_{2l}), \tau_j = \pm 1$ is negligible if there exist two integers j_0, j'_0 $(1 \le j_0 < j'_0 \le 2l)$ such that $j'_0 - j_0 - 1$ is odd, $\tau_j = -1$ $(j_0 \le j \le j'_0)$ and $\tau_{j_0-1} = \tau_{j'_0+1} = 1$. (We put $\tau_0 = \tau_{2l+1} = 1$. Hence $\tau_{j_0-1} = 1$ if $j_0 = 1$, and $\tau_{j'_0+1} = 1$ if $j'_0 = 2l$.) Let $\tau(\Delta[p]) = -1$ $(p \ge 2), \tau(K) = 1$. Lemmas 6 and 7 show that $d_{2l}(\Delta[\infty] + K)$ exists and

$$d_{2l}(\Delta[\infty] + K) = \lim_{p \to \infty} \sum_{(K_1, \dots, K_{2l}), K_j = \Delta[p], K} \int_0^1 K_1 \overline{K}_2 \cdots K_{2l-1} \overline{K}_{2l} \mathbf{1} \, dx$$
$$= \lim_{p \to \infty} \sum_{(p)} \int_0^1 K_1 \overline{K}_2 \cdots K_{2l-1} \overline{K}_{2l} \mathbf{1} \, dx \,,$$

where $\sum_{(p)}$ is the summation over all 2*l*-tuples (K_1, \ldots, K_{2l}) , $K_j = \Delta[p]$, K such that $(\tau(K_1), \ldots, \tau(K_{2l}))$ is not negligible. If $(\tau(K_1), \ldots, \tau(K_{2l}))$ is not negligible, then K appears even times in (K_1, \ldots, K_{2l}) . We can choose $j_1 < j_2 < \cdots < j_{2\nu}$ so that $K_{j_{\mu}} = K$ $(1 \le \mu \le 2\nu)$, $K_j = \Delta[p]$ $(j \notin \{j_{\mu}\}_{\mu=1}^{2\nu})$. Then j_1-1 , $j_{\mu+1}-j_{\mu}-1$ $(1 \le \mu \le 2\nu-1)$, $2l - j_{2\nu}$ are even. Notice that

$$d_{2j}(K) = \int_0^1 (\overline{K}K)^j 1 \, dx \quad (j \ge 0).$$

Thus we can write

$$d_{2l}(\Delta[\infty] + K) = \sum_{k=0}^{l} c_{2k}^{(2l)} d_{2l-2k}(K).$$

Let κ_0 be an operator defined by $h \in L^2 \to (\int_0^1 h \, dx) \chi_{[0,1)}$. We put $Y_{p,-1}(t) = 1$,

$$Y_{p,m}(t) = \begin{cases} \int_0^1 K_{p,t}^{m/2} 1 \, dx & (m \text{ is even}), \\ \int_0^1 (\kappa_0 + t \overline{\Delta[p]}) K_{p,t}^{(m-1)/2} 1 \, dx & (m \text{ is odd}), \end{cases}$$

where $K_{p,t} = (\kappa_0 + t\Delta[p])(\kappa_0 + t\overline{\Delta[p]})$. Then $Y_{\infty,m}(t) = \lim_{p \to \infty} Y_{p,m}(t)$ exists, and $c_{2k}^{(2l)}$ equals the t^{2k} -coefficient of $Y_{\infty,2l}(t)$. Evidently, (20) holds. Since $\int_0^1 \Delta[p] (\overline{\Delta[p]} \Delta[p])^j 1 \, dx = 0 \quad (j \ge 0)$, we have inductively

$$\begin{split} Y_{p,2l}(t) &= Y_{p,2l-1}(t) + t \int_0^1 \Delta[p](\kappa_0 + t\overline{\Delta[p]}) K_{p,t}^{l-1} \mathbf{1} \, dx \\ &= Y_{p,2l-1}(t) + t^2 \int_0^1 \Delta[p] \overline{\Delta[p]} K_{p,t}^{l-1} \mathbf{1} \, dx \\ &= Y_{p,2l-1}(t) + t^2 d_2(\Delta[p]) Y_{p,2l-3}(t) \\ &+ t^3 \int_0^1 \Delta[p] \overline{\Delta[p]} \Delta[p] (\kappa_0 + t\overline{\Delta[p]}) K_{p,t}^{l-2} \mathbf{1} \, dx \\ &= Y_{p,2l-1}(t) + t^2 d_2(\Delta[p]) Y_{p,2l-3}(t) \\ &+ t^4 \int_0^1 (\Delta[p] \overline{\Delta[p]})^2 K_{p,t}^{l-2} \mathbf{1} \, dx \\ &= \cdots = \sum_{j=0}^l t^{2j} d_{2j}(\Delta[p]) Y_{p,2l-2j-1}(t). \end{split}$$

Letting p tend to infinity, we have

$$Y_{\infty,2l}(t) = \sum_{j=0}^{l} t^{2j} d_{2j}(\Delta[\infty]) Y_{\infty,2l-2j-1}(t).$$

In the same manner,

$$Y_{\infty,2l+1}(t) = \sum_{j=0}^{l} t^{2j} d_{2j}(\Delta[\infty]) Y_{\infty,2l-2j}(t).$$

Thus

$$Y_{\infty,m}(t) = \sum_{j=0}^{l(m/2)} t^{2j} d_{2j}(\Delta[\infty]) Y_{\infty,m-2j-1}(t).$$

Comparing the t^{2k} -coefficients of both sides, we obtain (21).

LEMMA 9. Let K be an anti-symmetric kernel of type 1. Then, for any $0 < \delta \leq 1$,

$$\lim_{p \to \infty} \gamma^* (\delta \Delta[p] + \delta K) \quad (= \gamma^* (\delta \Delta[\infty] + \delta K), \ say)$$

exists; we write $\gamma^*(\delta\Delta[\infty])$ if K = 0. Moreover,

(22)
$$\gamma^*(\delta\Delta[\infty] + \delta K) = \gamma^*(\delta\Delta[\infty])\gamma^*(\gamma^*(\delta\Delta[\infty])\delta K).$$

Proof. First we show that $\gamma^*(\delta\Delta[\infty] + \delta K)$ and $\gamma^*(\delta\Delta[\infty])$ exist. Define $\gamma^*(w; \Delta[p] + K)$, $T(w; \Delta[p] + K)$ ($w \in U$) for $\varepsilon_0 =$

 $(12 + 3 \|K\|_{L^2, L^2})^{-1}$ in the same manner as in Lemma 5; we have $\varepsilon_0 \le (3 \|\Delta[p] + K\|_{L^2, L^2})^{-1}$ because $\|\Delta[p]\|_{L^2, L^2} \le 4$. Lemma 8 shows that

$$\lim_{p \to \infty} w^2 \varepsilon_0^2 \int_0^1 (\Delta[p] + K) T(w; \Delta[p] + K)^{m+1} (\overline{\Delta[p] + K})^1 dx$$
$$= \sum_{l=0}^\infty w^{2l+2} \varepsilon_0^{2l+2} \frac{(l+1)\cdots(l+m)}{m!} d_{2l+2} (\Delta[\infty] + K) \qquad (m \ge 0).$$

Since (14) holds with K replaced by any $\Delta[p] + K$ $(p \ge 2)$, (13) exists with K replaced by $\Delta[\infty] + K$, i.e.,

(23)
$$\lim_{p \to \infty} \gamma^*(w; \Delta[p] + K) \quad (= \gamma^*(w; \Delta[\infty] + K), \text{ say})$$

exists. Since

$$\gamma^*(\delta\,;\,\Delta[p]+K)=\gamma^*(\delta\Delta[p]+\delta K)\qquad (p\geq 2)\,,$$

 $\gamma^*(\delta\Delta[\infty] + \delta K) \quad (= \gamma^*(\delta; \Delta[\infty] + K))$ exists. Putting K = 0, we see that $\gamma^*(\delta\Delta[\infty])$ exists.

Next we show that $\gamma^*(w; \Delta[\infty] + K)$ and $\gamma^*(\gamma^*(w; \Delta[\infty])w; K)$ are analytic in a domain containing (0, 1]. The convergence of (23) is uniform in U. By Lemma 5, $\gamma^*(w; \Delta[p] + K)$ is analytic in U, and hence $\gamma^*(w; \Delta[\infty] + K)$ is analytic in U. The definition of $\gamma^*(\cdot)$ immediately shows that

$$\gamma^*(\operatorname{Re} w; \Delta[p]) = \gamma^*(\operatorname{Re} w\Delta[p]) \le 1 \qquad (w \in U).$$

Letting p tend to infinity, we have $\gamma^*(\operatorname{Re} w; \Delta[\infty]) \leq 1 \quad (w \in U)$. Since $\gamma^*(w; \Delta[\infty])$ is analytic in U, there exists $0 < \eta < \pi/8$ such that

$$|\gamma^*(w; \Delta[\infty])| \le \frac{4}{3}, \qquad |\arg \gamma^*(w; \Delta[\infty])| \le \frac{\pi}{8}$$

in $U_{\eta} = \{w \in \mathbb{C}; |w| < 4/3, |\arg w| < \eta\}$. Then $\gamma^*(w; \Delta[\infty])w \in U$ $(w \in U_{\eta})$. Thus, by Lemma 5, $\gamma^*(\gamma^*(w; \Delta[\infty])w; K)$ is analytic in U_{η} .

By the theorem of identity, it is sufficient to show that (22) holds for $0 < \delta < (8 + 2 \|K\|_{L^2, L^2})^{-1}$. Since

$$\lim_{l\to\infty} d_{2l}(\delta\Delta[p]) = \lim_{l\to\infty} d_{2l}(\delta\Delta[p] + \delta K) = 0,$$

(12) holds for $\delta\Delta[p]$, $\delta\Delta[p] + \delta K$ $(p \ge 2)$. Letting p tend to infinity, we have

$$\gamma^*(\delta\Delta[\infty]) = \sum_{l=0}^{\infty} d_{2l}(\delta\Delta[\infty]) = \sum_{l=0}^{\infty} \delta^{2l} d_{2l}(\Delta[\infty]),$$

$$\gamma^*(\delta\Delta[\infty] + \delta K) = \sum_{l=0}^{\infty} \delta^{2l} d_{2l}(\Delta[\infty] + K).$$

Let

$$\mu_m = \sum_{k=0}^{\infty} \delta^{2k} c_{2k}^{(m+2k)} \qquad (m \ge 0) \,,$$

where $c_{2k}^{(m)}$ $(0 \le k \le \iota(m/2), m \ge 0)$ are numbers in Lemma 8. Then

$$\mu_0 = \sum_{k=0}^{\infty} \delta^{2k} c_{2k}^{(2k)} = \gamma^* (\delta \Delta[\infty]) ,$$

by (20). Equality (21) yields that

$$\begin{split} \mu_m &= \sum_{k=0}^{\infty} \delta^{2k} \sum_{j=0}^k c_{2k-2j}^{(m+2k-2j-1)} d_{2j}(\Delta[\infty]) \\ &= \sum_{j=0}^{\infty} \delta^{2j} d_{2j}(\Delta[\infty]) \sum_{k=j}^{\infty} \delta^{2(k-j)} c_{2(k-j)}^{(m-1+2(k-j))} \\ &= \mu_{m-1} \mu_0 \qquad (m \ge 1) \,, \end{split}$$

which gives

$$\mu_m = \mu_0^{m+1} = \gamma^* (\delta \Delta[\infty])^{m+1} \qquad (m \ge 1).$$

Thus, by (21),

$$\begin{split} \gamma^*(\delta\Delta[\infty] + \delta K) &= \sum_{l=0}^{\infty} \delta^{2l} d_{2l}(\Delta[\infty] + K) \\ &= \sum_{l=0}^{\infty} \delta^{2l} \sum_{k=0}^{l} c_{2k}^{(2l)} d_{2l-2k}(K) = \sum_{k=0}^{\infty} \sum_{l=k}^{\infty} \delta^{2l} c_{2k}^{(2l)} d_{2l-2k}(K) \\ &= \sum_{k=0}^{\infty} \sum_{j=0}^{\infty} \delta^{2j+2k} c_{2k}^{(2j+2k)} d_{2j}(K) = \sum_{j=0}^{\infty} \mu_{2j} \delta^{2j} d_{2j}(K) \\ &= \sum_{j=0}^{\infty} \gamma^* (\delta\Delta[\infty])^{2j+1} \delta^{2j} d_{2j}(K) = \gamma^* (\delta\Delta[\infty]) \gamma^* (\gamma^* (\delta\Delta[\infty]) \delta K). \end{split}$$

LEMMA 10. We inductively define a sequence $(\gamma_n^*)_{n=1}^{\infty}$ of positive numbers by

$$\gamma_1^* = \gamma^*(\Delta[\infty]), \quad \gamma_n^* = \gamma_{n-1}^* \gamma^*(\gamma_{n-1}^* \Delta[\infty]) \qquad (n \ge 2).$$

Then

$$\lim_{p_1\to\infty}\lim_{p_2\to\infty}\cdots\lim_{p_n\to\infty}\gamma^*(H[p_1,\ldots,p_n]-H)=\gamma_n^* \qquad (n\geq 1)\,,$$

where $\lim_{p_n\to\infty}$ is taken first and $\lim_{p_n\to\infty}$ is taken last.

Proof. We define a sequence $(\lambda_n)_{n=1}^{\infty}$ of positive numbers by $\lambda_1 = \gamma^*(\Delta[\infty])$, $\lambda_n = \gamma^*(\lambda_1 \dots \lambda_{n-1}\Delta[\infty])$ $(n \ge 2)$. Then $\gamma_n^* = \lambda_1 \dots \lambda_n$ $(n \ge 1)$. Fixing an (n-1)-tuple (p_1, \dots, p_{n-1}) $(n \ge 2)$ of integers larger than or equal to 3, we study

$$\lim_{p_n\to\infty}\gamma^*(H[p_1,\ldots,p_n]-H)$$

=
$$\lim_{p_n\to\infty}\gamma^*(\Delta[p_1,\ldots,p_n]+(H[p_1,\ldots,p_{n-1}]-H)).$$

Put $I_0 = [0, 1/(p_1 \cdots p_{n-1})), I_j = (I_0 + j/(p_1 \cdots p_{n-1})) \quad (0 \le j \le (p_1 \dots p_{n-1}) - 1).$ Then

$$(H[p_1, \ldots, p_{n-1}] - H)(x, y) = 0$$
 $(x, y \in I_j),$

$$\begin{aligned} |(H[p_1, \dots, p_{n-1}] - H)(x, y)| \\ &\leq \frac{2}{\pi} p_1 \cdots p_{n-1} + \frac{1}{\pi} \frac{1}{|y - x|} \qquad (x \in I_j, y \in I_k, j \neq k), \end{aligned}$$

which shows that $H[p_1, \ldots, p_{n-1}] - H$ is of type 1. Let

$$\Delta'[p_1, \dots, p_n](x, y) = \Delta[p_1, \dots, p_n](x, y)\chi_{[0, \frac{N'}{p_1 \dots p_n})}(|y - x|),$$
$$\Delta'\left[\prod_{j=1}^n p_j\right](x, y) = \Delta\left[\prod_{j=1}^n p_j\right](x, y)\chi_{[0, \frac{N'}{p_1 \dots p_n})}(|y - x|)$$

 $(N' = \iota(\log(p_1 \dots p_n)))$. Then

$$\lim_{p_n \to \infty} \|\Delta[p_1, \dots, p_n] - \Delta'[p_1, \dots, p_n]\|_{L^4, L^4} = 0,$$
$$\lim_{p_n \to \infty} \left\|\Delta\left[\prod_{j=1}^n p_j\right] - \Delta'\left[\prod_{j=1}^n p_j\right]\right\|_{L^4, L^4} = 0$$

(cf. Lemmas 6 and 7). Since

$$\Delta'[p_1, ..., p_n](x, y) = \Delta' \left[\prod_{j=1}^n p_j\right](x, y)$$

(x, y \in I_j, 0 \le j \le (p_1 \dots p_{n-1}) - 1),

we have

$$\lim_{p_n \to \infty} \left\| \Delta[p_1, \dots, p_n] - \Delta\left[\prod_{j=1}^n p_j\right] \right\|_{L^4, L^2}$$
$$= \lim_{p_n \to \infty} \left\| \Delta'[p_1, \dots, p_n] - \Delta'\left[\prod_{j=1}^n p_j\right] \right\|_{L^4, L^2} = 0,$$

and hence, in the same manner as in the proof of the existence of (23),

$$\lim_{p_n \to \infty} \gamma^* (\Delta[p_1, \dots, p_n] + (H[p_1, \dots, p_{n-1}] - H))$$
$$= \lim_{p_n \to \infty} \gamma^* \left(\Delta \left[\prod_{j=1}^n p_j \right] + (H[p_1, \dots, p_{n-1}] - H) \right).$$

Using (22) with $\delta = 1$, $K = H[p_1, \dots, p_{n-1}] - H$, we have

$$\lim_{p_n \to \infty} \gamma^* (H[p_1, \dots, p_n] - H) = \gamma^* (\Delta[\infty] + (H[p_1, \dots, p_{n-1}] - H))$$
$$= \lambda_1 \gamma^* (\lambda_1 (H[p_1, \dots, p_{n-1}] - H)).$$

In the same manner, using (22) with $\delta = \lambda_1$, $K = H[p_1, \ldots, p_{n-2}] - H$, we have

$$\lim_{p_{n-1}\to\infty}\lim_{p_n\to\infty}\gamma^*(H[p_1,\ldots,p_n]-H)$$

= $\lambda_1\gamma^*(\lambda_1\Delta[\infty])\gamma^*(\gamma^*(\lambda_1\Delta[\infty])\lambda_1(H[p_1,\ldots,p_{n-2}]-H))$
= $\lambda_1\lambda_2\gamma^*(\lambda_1\lambda_2(H[p_1,\ldots,p_{n-2}]-H)).$

Repeating this argument,

$$\lim_{p_1 \to \infty} \cdots \lim_{p_n \to \infty} \gamma^* (H[p_1, \ldots, p_n] - H)$$

= $\lambda_1 \cdots \lambda_{n-1} \lim_{p_1 \to \infty} \gamma^* (\lambda_1 \cdots \lambda_{n-1} \Delta[p_1]) = \lambda_1 \ldots \lambda_n = \gamma_n^*.$

This completes the proof of our lemma.

We now give the proof of our theorem. By Proposition 2, there exists a positive integer n_0 such that

(24)
$$\sup \gamma(\Gamma(p_1,\ldots,p_n)) \leq 10^{-5} \qquad (n \geq n_0),$$

where the supremum is taken over all *n*-tuples (p_1, \ldots, p_n) of integers larger than or equal to 3. By Lemma 10, we can inductively choose a sequence $(p_n^0)_{n=1}^{\infty}$ of integers larger than or equal to 3 so that

$$\frac{1}{2}\gamma_n^* \leq \gamma^*(H[p_1^0,\ldots,p_n^0]-H) \leq 2\gamma_n^* \qquad (n\geq 1),$$

where $(\gamma_n^*)_{n=1}^{\infty}$ is the sequence in Lemma 10. We show that $\Gamma_n = \Gamma(p_1^0, \ldots, p_n^0)$ $(n \ge 1)$ are required cranks. We may assume that $n \ge n_0$. Lemma 3 shows that

$$\frac{1}{4}\gamma^*(H[p_1^0, \dots, p_n^0] - H) \le \gamma^*(H[p_1^0, \dots, p_n^0]) \\ \le 4\gamma^*(H[p_1^0, \dots, p_n^0] - H),$$

and hence

$$\frac{1}{8}\gamma_n^* \le \gamma^*(H[p_1^0, \ldots, p_n^0]) \le 8\gamma_n^*$$

Thus, by (11),

(25)
$$\frac{1}{8\pi}\gamma_n^* \le \gamma(\Gamma_n) \le \frac{8}{\pi}\gamma_n^*.$$

Using (24) and (25), we have $\gamma_n^* \leq 8\pi \cdot 10^{-5}$. Recall (15), and notice that $d_{2l}(\Delta[\infty]) \leq 4^l$ $(l \geq 1)$. Since $\lim_{l\to\infty} d_{2l}(\gamma_n^*\Delta[p]) = 0$, (12) holds for $\gamma_n^*\Delta[p]$. Letting p tend to infinity, we have

$$\begin{aligned} \gamma_{n+1}^{*} &= \gamma_{n}^{*} \gamma^{*} (\gamma_{n}^{*} \Delta[\infty]) = \gamma_{n}^{*} \sum_{l=0}^{\infty} d_{2l} (\gamma_{n}^{*} \Delta[\infty]) \\ &= \gamma_{n}^{*} \sum_{l=0}^{\infty} \gamma_{n}^{*^{2l}} d_{2l} (\Delta[\infty]) \le \gamma_{n}^{*} - \frac{1}{25\pi^{2}} \gamma_{n}^{*^{3}} + \sum_{l=2}^{\infty} 4^{l} \gamma_{n}^{*^{2l+1}} \\ &\le \gamma_{n}^{*} - 10^{-3} \gamma_{n}^{*^{3}}, \\ \gamma_{n+1}^{*} &\ge \gamma_{n}^{*} - \sum_{l=1}^{\infty} 4^{l} \gamma_{n}^{*^{2l+1}} \ge \gamma_{n}^{*} - 10 \gamma_{n}^{*^{3}}, \quad \text{i.e.}, \\ \gamma_{n}^{*} - 10 \gamma_{n}^{*^{3}} \le \gamma_{n+1}^{*} \le \gamma_{n}^{*} - 10^{-3} \gamma_{n}^{*^{3}}. \end{aligned}$$

Since this holds for all $n \ge n_0$, a simple induction yields that

$$\frac{1}{C_4}\frac{1}{\sqrt{n}} \le \gamma_n^* \le C_4 \frac{1}{\sqrt{n}} \qquad (n \ge n_0)$$

for some absolute constant C_4 . Using (25) again,

$$\frac{1}{8\pi C_4} \frac{1}{\sqrt{n}} \leq \gamma(\Gamma_n) \leq \frac{8}{\pi} C_4 \frac{1}{\sqrt{n}} \qquad (n \geq n_0).$$

This completes the proof of our theorem.

REMARK 11. It is not known whether $\gamma(\cdot)$ is semi-additive [4, p. 11]. For $0 < \eta \le 1$, we define $B_p^{\eta}(x)$ replacing 1/2p by $\eta/2p$ in the definition of $B_p(x)$. Then cranks $\Gamma^{\eta}(p_1, \ldots, p_n)$ of degree *n* are

analogously defined. We see that there exists a crank Γ_n^{η} of degree n such that $\gamma(\Gamma_n^{\eta}) \leq C_{\eta}/\sqrt{n}$, where C_{η} is a constant depending only on η . Adding some segments (perpendicular to the x-axis) to Γ_n^{η} , we obtain an arc $\tilde{\Gamma}_n^{\eta}$ connecting 0 and 1. Then the diameter of $\tilde{\Gamma}_n^{\eta}$ is larger than or equal to 1. Since $\tilde{\Gamma}_n^{\eta}$ is connected, $\gamma(\tilde{\Gamma}_n^{\eta}) \geq 1/4$ [4, p. 9]. Hence, from the point of view of the above semi-additive problem, it seems interesting to compute $\gamma(\tilde{\Gamma}_n^{\eta} - \Gamma_n^{\eta})$.

4. Another application of Proposition 1. In this section, we show another application of our method. Let E be a compact set on \mathbb{R} . Pommerenke [11] showed that

(26)
$$\gamma(E) = |E|/4,$$

(27)
$$f_E(z) = \left\{ 1 - \exp\left(\frac{1}{2}\int_E \frac{dt}{t-z}\right) \right\} / \left\{ 1 + \exp\left(\frac{1}{2}\int_E \frac{dt}{t-z}\right) \right\}.$$

We deduce (26), (27) from (3), (10); our method explains a quarter and (27). Let $L^2(\mathbb{R})$ denote the L^2 space of functions on \mathbb{R} , and let M_E denote the multiplier: $h \in L^2(\mathbb{R}) \to \chi h \in L^2(\mathbb{R})$, where $\chi = \chi_E$. We inductively define a sequence $(H_E^{(m)})_{m=0}^{\infty}$ of operators from $L^2(\mathbb{R})$ to itself by $H_E^{(0)} = M_E$, $H_E^{(m)} = HM_EH_E^{(m-1)}$ $(m \ge 1)$. Notice that

$$\gamma(E) = \frac{1}{\pi} \gamma^* (M_E H M_E),$$

$$d_{2l}(M_E H M_E) = \int_E H_E^{(2l)} \chi \, dx \qquad (l \ge 0, \, \chi = \chi_E).$$

We also remark that

(28)
$$H(g \cdot Hh) + H(Hg \cdot h) = Hg \cdot Hh - gh \qquad (g, h \in L^{2}(\mathbb{R})).$$

We first show that, for any $m \ge 1$,

(29)
$$\chi H\chi \cdot H_E^{(m)}\chi = (m+1)\chi H_E^{(m+1)}\chi + m\chi H_E^{(m-1)}\chi.$$

Equality (28) shows that $2H(\chi H\chi) = (H\chi)^2 - \chi\chi$, which gives $\chi H\chi \cdot H_E^{(1)}\chi = 2\chi H_E^{(2)}\chi + \chi H_E^{(0)}\chi$. Suppose that (29) holds for *m*. Using (28) with $g = \chi$, $h = \chi H_E^{(m)}\chi$, we have

$$\begin{split} \chi H \chi \cdot H_E^{(m+1)} \chi &= \chi H \chi \cdot H(\chi H_E^{(m)} \chi) \\ &= \chi H\{\chi H(\chi H_E^{(m)} \chi) + H \chi \cdot \chi H_E^{(m)} \chi\} + \chi \{\chi \cdot \chi H_E^{(m)} \chi\} \\ &= \chi H_E^{(m+2)} \chi + \chi H\{(m+1)\chi H_E^{(m+1)} \chi + m \chi H_E^{(m-1)} \chi\} + \chi H_E^{(m)} \chi \\ &= (m+2)\chi H_E^{(m+2)} \chi + (m+1)\chi H_E^{(m)} \chi, \quad \text{i.e.} \,, \end{split}$$

(29) holds for m + 1. Thus (29) holds for all $m \ge 1$.

We next show that

(30)
$$\int_E H_E^{(2l)} \chi \, dx = \frac{(-1)^l}{2l+1} |E| \qquad (l \ge 0).$$

We put $a_{2l} = \int_E H_E^{(2l)} \chi \, dx$ $(l \ge 0)$. Evidently, $a_0 = |E|$. Suppose that $a_{2l-2} = \{(-1)^{l-1}/(2l-1)\}|E|$. Equality (29) (m = 2l-1) shows that

$$\int_{E} H\chi \cdot H_{E}^{(2l-1)}\chi \, dx = 2l \int_{E} H_{E}^{(2l)}\chi \, dx + (2l-1) \int_{E} H_{E}^{(2l-2)}\chi \, dx$$
$$= 2la_{2l} + (2l-1)a_{2l-2}.$$

Since the adjoint operator of H equals -H, we have

$$\int_E H\chi \cdot H_E^{(2l-1)}\chi \, dx = -\int_E H\{\chi H_E^{(2l-1)}\chi\} \, dx = -a_{2l}.$$

Thus $-a_{2l} = 2la_{2l} + (2l-1)a_{2l-2}$, which yields that

$$a_{2l} = -\frac{2l-1}{2l+1}a_{2l-2} = \frac{(-1)^l}{2l+1}|E|.$$

Now the deduction of (26) is immediate. By (30),

$$\lim_{l\to\infty} d_{2l}(M_E H M_E) = \lim_{l\to\infty} \int_E H_E^{(2l)} \chi \, dx = 0.$$

Hence we can apply (3). Leibniz's formula and (30) yield that

$$\begin{split} \gamma(E) &= \frac{1}{\pi} \gamma^* (M_E H M_E) = \frac{1}{\pi} \sum_{l=0}^{\infty} d_{2l} (M_E H M_E) \\ &= \frac{1}{\pi} \sum_{l=0}^{\infty} \frac{(-1)^l}{2l+1} |E| = \frac{1}{4} |E|. \end{split}$$

Last, we deduce (27) from (10). Equality (10) gives that

$$f_E(z) = -\frac{1}{\pi} \left\{ \int_E \frac{ds}{s-z} + \int_E \frac{H_E^{(1)} h_E(s)}{s-z} \, ds \right\} / \left\{ 1 + \frac{1}{\pi} \int_E \frac{h_E(s)}{s-z} \, ds \right\} ,$$

where $h_E(s)$ is the function which attains $\gamma^*(M_E H M_E)$. We show that this equals the function in the right-hand side of (27). Let

$$u_0(z) = 1, \qquad u_m(z) = \frac{1}{\pi} \int_E \frac{H_E^{(m-1)}\chi(s)}{s-z} \, ds,$$
$$v_m(z) = \frac{1}{\pi} \int_E \frac{H\chi(s)H_E^{(m-2)}\chi(s)}{s-z} \, ds \qquad (m \ge 1),$$

where $H_E^{(-1)}\chi = \chi$. Let

$$P_t(z) = \sum_{m=0}^{\infty} t^m u_m(z) \qquad (t \in \mathbb{C}, |t| < 1).$$

We begin by showing that

(31)
$$(1+t^2)\frac{\partial}{\partial t}P_t(z) = u_1(z)P_t(z) \quad (0 < t < 1).$$

Let $m \ge 1$. We have, on \mathbb{R} ,

$$\begin{split} \lim_{\eta \downarrow 0} \{ u_{m+1}(\cdot + i\eta) + v_{m+1}(\cdot + i\eta) \} \\ &= H(\chi H_E^{(m)}\chi) + i\chi H_E^{(m)}\chi \\ &+ H\{\chi H\chi \cdot H_E^{(m-1)}\chi\} + i\chi H\chi \cdot H_E^{(m-1)}\chi \\ &= H\{\chi H(\chi H_E^{(m-1)}\chi) + H\chi \cdot \chi H_E^{(m-1)}\chi\} \\ &+ i\{\chi H(\chi H_E^{(m-1)}\chi) + H\chi \cdot \chi H_E^{(m-1)}\chi\} \,, \end{split}$$

$$\begin{split} \lim_{\eta \downarrow 0} u_1(\cdot + i\eta) u_m(\cdot + i\eta) \\ &= \{H\chi + i\chi\}\{H(\chi H_E^{(m-1)}\chi) + i\chi H_E^{(m-1)}\chi\} \\ &= H\chi \cdot H(\chi H_E^{(m-1)}\chi) - \chi \cdot \chi H_E^{(m-1)}\chi \\ &+ i\{\chi H(\chi H_E^{(m-1)}\chi) + H\chi \cdot \chi H_E^{(m-1)}\chi\}. \end{split}$$

Hence (28) $(g = \chi, h = \chi H_E^{(m-1)}\chi)$ shows that

$$\lim_{\eta \downarrow 0} \{ u_{m+1}(\cdot + i\eta) + v_{m+1}(\cdot + i\eta) - u_1(\cdot + i\eta)u_m(\cdot + i\eta) \} = 0$$

on \mathbb{R} . In particular, this holds on $\mathbb{R} - E$. Hence, by the theorem of identity, $u_{m+1}(z) + v_{m+1}(z) - u_1(z)u_m(z) = 0$. Equality (29) shows that $v_{m+1}(z) = mu_{m+1}(z) + (m-1)u_{m-1}(z)$. Thus

$$(m+1)u_{m+1}(z) + (m-1)u_{m-1}(z) - u_1(z)u_m(z) = 0 \qquad (m \ge 1),$$

which yields that

$$\sum_{m=0}^{\infty} mt^m u_m(z) + t^2 \sum_{m=0}^{\infty} mt^m u_m(z) = tu_1(z) \sum_{m=0}^{\infty} t^m u_m(z), \quad \text{i.e.},$$
$$t \frac{\partial}{\partial t} P_t(z) + t^3 \frac{\partial}{\partial t} P_t(z) = tu_1(z) P_t(z).$$

This is the required equality (31).

We can choose $x_0 \in \mathbb{R} - E$, $\eta > 0$ so that $P_t(x) > 0$, $u_1(x) > 0$ for all $x \in (x_0 - \eta, x_0 + \eta)$, 0 < t < 1. Equality (31) shows that

$$\frac{1}{1+t^2}u_1(x) = \frac{\partial}{\partial t}P_t(x) / P_t(x) \qquad (x \in (x_0 - \eta, x_0 + \eta), \ 0 < t < 1),$$

which gives that

$$P_t(x) = \exp\left\{\int_0^t \frac{ds}{1+s^2} u_1(x)\right\} \qquad (x \in (x_0 - \eta, x_0 + \eta), \ 0 < t < 1)$$

because $P_0 = 1$. By the theorem of identity,

$$P_t(z) = \exp\left\{\int_0^t \frac{ds}{1+s^2} u_1(z)\right\} \qquad (0 < t < 1).$$

Since $P_t(z)$ and $\exp\{(\int_0^t (ds/(1+s^2))u_1(z))\}$ are analytic in the unit disk as functions of t, this equality holds for -1 < t < 0 also. Thus

$$\begin{split} 1 &+ \frac{1}{\pi} \int_{E} \frac{h_{E}(s)}{s-z} \, ds = 1 + \frac{1}{\pi} \int_{E} \frac{1}{s-z} \sum_{l=1}^{\infty} H_{E}^{(2l-1)} \chi(s) \, ds \\ &= \lim_{t \uparrow 1} \sum_{l=0}^{\infty} t^{2l} u_{2l}(z) = \frac{1}{2} \lim_{t \uparrow 1} \{ P_{-t}(z) + P_{t}(z) \} \\ &= \frac{1}{2} \left\{ \exp\left(-\frac{\pi}{4} u_{1}(z)\right) + \exp\left(\frac{\pi}{4} u_{1}(z)\right) \right\} , \\ &- \frac{1}{\pi} \int_{E} \frac{ds}{s-z} - \frac{1}{\pi} \int_{E} \frac{H_{E}^{(1)} h_{E}(s)}{s-z} \, ds = \frac{1}{2} \lim_{t \uparrow 1} \{ P_{-t}(z) - P_{t}(z) \} \\ &= \frac{1}{2} \left\{ \exp\left(-\frac{\pi}{4} u_{1}(z)\right) - \exp\left(\frac{\pi}{4} u_{1}(z)\right) \right\} , \end{split}$$

which gives (27).

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