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COMMUTING ANALYTIC SELF-MAPS OF THE BALL

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Under broad conditions, two analytic self-maps of the disk fixing 0 commute under composition precisely when they have the same Schroeder map, where the Schroeder map for an analytic  $\varphi : D \rightarrow D$  with  $\varphi(0) = 0$  is the unique analytic function  $\sigma$  on  $D$  solving Schroeder's equation  $\sigma \circ \varphi = \varphi'(0)\sigma$  and satisfying  $\sigma'(0) = 1$ . For analytic self-maps of the ball in  $C^N$  fixing 0 we may still seek analytic  $C^N$ -valued solutions  $\sigma$  to Schroeder's equation with  $\sigma'(0) = I$ , but considerable complications for existence and uniqueness of  $\sigma$  may ensue. Nevertheless, we show that there are reasonably general hypotheses under which it will still be the case that two analytic self-maps of the ball fixing 0 commute if and only if they share a common Schroeder map  $\sigma$  with  $\sigma'(0) = I$ .

## 1. Introduction.

If  $\varphi$  is an analytic map of the unit disk  $D$  into itself which fixes the origin and has derivative there satisfying  $0 < |\varphi'(0)| < 1$  then there exists an analytic map  $\sigma$  on  $D$  that satisfies Schroeder's functional equation

$$\sigma \circ \varphi = \varphi'(0)\sigma.$$

This "Schroeder map"  $\sigma$  is unique up to constant multiples; its existence and uniqueness was proved by Koenigs in 1884 ([3]). It is usually convenient to require that  $\sigma$  satisfy  $\sigma'(0) = 1$ . Koenigs showed that in this case  $\sigma$  can be obtained as the almost uniform limit of normalized iterates of  $\varphi$ :

$$\sigma = \lim_{n \rightarrow \infty} \frac{\varphi_n}{\varphi'(0)^n},$$

where  $\varphi_1 = \varphi$  and  $\varphi_{k+1} = \varphi \circ \varphi_k$ . When  $\varphi$  is univalent in  $D$ ,  $\sigma$  will be also, so that  $\varphi$  is conjugate to multiplication by  $\varphi'(0)$  on  $\sigma(D)$ :  $\varphi = \sigma^{-1}\varphi'(0)\sigma$ . Suppose  $\psi$  is an analytic self-map of  $D$  which commutes with  $\varphi$  under composition. Then necessarily  $\psi(0) = 0$ . Moreover,  $\varphi$  and  $\psi$  will have the same Schroeder maps, and conversely if  $\psi : D \rightarrow D$  is analytic, fixes 0 and has the same Schroeder map as  $\varphi$ , then  $\varphi \circ \psi = \psi \circ \varphi$ . These results follow from the existence and (essential) uniqueness of the Schroeder map in one variable. ([1], [4].)

If  $\varphi$  is an analytic self-map of the unit ball  $B_N$  in  $C^N$  which fixes the origin, then by a Schroeder map for  $\varphi$  we will mean an analytic map  $\sigma : B_N \rightarrow C^N$  which satisfies the functional equation

$$(1) \quad \sigma \circ \varphi = \varphi'(0)\sigma$$

where  $\varphi'(0)$  is the linear map from  $C^N$  to  $C^N$  given by the matrix whose  $ij^{th}$  entry is  $D_j\varphi_i(0)$ . By analogy to the one variable case we restrict to the case that the eigenvalues of  $\varphi'(0)$  are non-zero and of modulus strictly less than 1. In addition we exclude maps which are “unitary on a slice” of the ball; that is, maps  $\varphi$  for which there exists  $\zeta, \eta$  in  $\partial B_N$  so that  $\varphi(\lambda\zeta) = \lambda\eta$  for all  $\lambda \in D$ . We are chiefly interested in Schroeder maps  $\sigma$  which are locally univalent near 0. This is equivalent to requiring that  $\sigma'(0)$  be invertible ([5, 1.3.7 and 15.1.8]). In fact when there is a solution to Equation (1) with  $\sigma'(0)$  invertible, there will be a solution with  $\sigma'(0) = I$ . Precise conditions under which such a solution exists for a given  $\varphi$  are known ([2]; see also Theorem 1 and Corollary 2 below), but are somewhat complicated. A basic issue is whether any algebraic relationships of the form

$$\lambda_j = \lambda_1^{k_1} \lambda_2^{k_2} \cdots \lambda_N^{k_N}$$

hold between the eigenvalues  $\lambda_k$  of  $\varphi'(0)$ , where  $k_i \geq 0$  and  $\sum k_i \geq 2$ , and if any such relationships do hold, whether they in fact prevent the existence of a locally univalent Schroeder map. Such an algebraic relationship for an eigenvalue of  $\varphi$  will be called a *resonance* of  $\varphi$ .

The results which make this precise are as follows. As a convenient normalization we may assume, by a unitary change of variables, that  $\varphi'(0)$  is upper triangular.

**Theorem 1** ([2]). *Suppose  $\varphi$  is an analytic map of  $B_N$  into  $B_N$  with  $\varphi(0) = 0$  and  $A = \varphi'(0)$  an upper triangular diagonalizable matrix, with diagonal entries  $\lambda_1, \lambda_2, \dots, \lambda_N$  such that  $0 < |\lambda_j| < 1$ . Assume further that  $\varphi$  is not unitary on any slice. Suppose that  $\lambda_j = \lambda_1^{k_1} \cdots \lambda_N^{k_N}$  is the longest expression (maximal  $\sum k_i$ ) for one eigenvalue of  $A$  as a product of any number of the eigenvalues of  $A$ . Set  $m = k_1 + \cdots + k_N$  and  $M =$  the number of multi-indices for  $C^N$  of total order less than or equal to  $m$ . Let  $\mathcal{M}$  be the upper left  $M \times M$  corner of the matrix for the composition operator  $C_\varphi$  with respect to the standard (non-normalized) basis for any weighted Hardy space  $H_\beta^2(B_N)$ , ordered in the usual way. If  $\mathcal{M}$  is diagonalizable, then Schroeder’s Equation (1) has a solution  $\sigma$  with  $\sigma'(0) = I$ .*

The “standard basis” referred to in this theorem consists of the monomials  $1, z_1, z_2, \dots, z_n, z_1^2, z_1 z_2, \dots$  ordered as follows:  $z^\alpha$  precedes  $z^\gamma$  where  $\alpha = (\alpha_1, \dots, \alpha_N)$  and  $\gamma = (\gamma_1, \dots, \gamma_N)$  are multi-indices, if either  $|\alpha| < |\gamma|$  or, in the case  $|\alpha| = |\gamma|$ , if there is a  $j_0$  so that  $\alpha_j = \gamma_j$  for  $j < j_0$  and  $\alpha_{j_0} > \gamma_{j_0}$ . The matrix for the composition operator  $C_\varphi$  with respect to this basis has as

its  $j^{th}$  column the coefficients of  $\varphi^\alpha$  with respect to this basis, where  $z^\alpha$  is the  $j^{th}$  monomial in the prescribed ordering. A weighted Hardy space  $H_\beta^2(B_N)$  is a Hilbert space of analytic functions on  $B_N$  for which the monomials form a complete orthogonal set of non-zero vectors satisfying

$$\beta(|\alpha_1|) \equiv \frac{\|z^{\alpha_1}\|}{\|z^{\alpha_1}\|_2} = \frac{\|z^{\alpha_2}\|}{\|z^{\alpha_2}\|_2}$$

whenever  $|\alpha_1| = |\alpha_2|$ , where  $\|\cdot\|$  denotes the norm in  $H_\beta^2(B_N)$  and  $\|\cdot\|_2$  denotes the norm in  $L^2(\sigma_N)$ ,  $\sigma_N$  being normalized Lebesgue measure on  $B_N$ . When  $\varphi(0) = 0$  and  $\varphi$  is not unitary on any slice there exist weighted Hardy spaces on which the composition operator  $C_\varphi$  (defined by  $C_\varphi(f) = f \circ \varphi$ ) is a compact operator ([2]).

There is a converse to Theorem 1 which says, under the same hypotheses on  $\varphi$ , that if  $\varphi$  has a Schroeder map with invertible derivative at the origin, then every upper left corner of the matrix for  $C_\varphi$  is diagonalizable.

While we won't have direct need for the full strength of Theorem 1 here, the following corollary will play a crucial role in our study of commuting analytic self-maps of  $B_N$ . It gives a description of all Schroeder maps (locally univalent or not) for  $\varphi$ , based on the presence or absence of resonances for  $\varphi$ .

**Corollary 2** ([2]). *Suppose the hypotheses of Theorem 1 hold and that in addition  $A = \varphi'(0)$  is diagonal. Then all solutions to Schroeder's Equation (1) can be described as  $f = g \circ \sigma$  where  $\sigma$  is a Schroeder map with  $\sigma'(0) = I$ , as given in Theorem 1 and  $g = (g_1, g_2, \dots, g_N)$  is a mapping on  $C^N$  with polynomial coordinate functions. Moreover, if  $g_k = \sum c(\gamma)z^\gamma$ , then the coefficients  $c(\gamma)$ ,  $\gamma = (\gamma_1, \dots, \gamma_N)$  are 0 unless  $\lambda_k = \lambda_1^{\gamma_1}\lambda_2^{\gamma_2}\cdots\lambda_N^{\gamma_N}$  ( $\gamma_i \geq 0$ ), in which case  $c(\gamma)$  can be chosen arbitrarily.*

*If  $A = \varphi'(0)$  is merely diagonalizable, with  $SAS^{-1} = \text{diag}(\lambda_1, \lambda_2, \dots, \lambda_N)$ , then an arbitrary Schroeder map has the form  $S^{-1} \circ g \circ S \circ \sigma$  with  $\sigma$  and  $g$  as just described.*

Note that  $g_k$  always includes a linear term  $b_k z_k$  ( $b_k$  arbitrary), and if  $\lambda_k$  is a repeated eigenvalue of  $\varphi'(0)$  there will be other linear terms with arbitrary coefficients. The terms of  $g_k$  with order at least two correspond to the resonance relations for  $\lambda_k$ . When no resonance relations hold,  $g$  is linear. We emphasize that a resonance relation expresses an eigenvalue  $\lambda_j$  as a product  $\lambda_1^{k_1}\lambda_2^{k_2}\cdots\lambda_N^{k_N}$  where  $\sum k_i \geq 2$ ; a relation  $\lambda_j = \lambda_k$  is not a resonance relation.

The goal of this paper is to study commuting analytic self-maps of  $B_N$  and to see, by analogy with known results in one variable, to what extent it still is the case that commuting maps are those which share a locally univalent Schroeder map. Our main results (Theorems 3 and 7) will show that under natural hypotheses on  $\varphi$ , a map  $\psi$  which commutes with  $\varphi$  and

has no resonances in common with  $\varphi$  will share a locally univalent Schroeder map with  $\varphi$ . Examples will be give to show that this can fail if  $\varphi$  and  $\psi$  have resonances in common.

## 2. Non-resonant maps.

In studying the Schroeder maps for commuting self-maps  $\varphi, \psi$  of  $B_N$ , the easiest situation arises when at least one of  $\varphi, \psi$  has no resonances. This means, say, that no eigenvalue of  $\varphi'(0)$  can be written as a product of two or more of the other eigenvalues, although repeated eigenvalues are allowed.

**Theorem 3.** *Suppose  $\varphi : B_N \rightarrow B_N$  is analytic, with  $\varphi(0) = 0$ . Assume that  $A = \varphi'(0)$  is upper triangular diagonalizable with eigenvalues  $\lambda_1, \lambda_2, \dots, \lambda_N, 0 < |\lambda_j| < 1$ . Assume further that  $\varphi$  is not unitary on any slice and that no resonance relations hold for any of the  $\lambda_j$ 's. If  $\psi : B_N \rightarrow B_N$  is analytic and  $\psi \circ \varphi = \varphi \circ \psi$  then  $\varphi$  and  $\psi$  share a common Schroeder map which is locally univalent near 0.*

*Proof.* Since  $\varphi$  is not unitary on any slice and  $\varphi\psi(0) = \psi\varphi(0) = \psi(0)$  we must have  $\psi(0) = 0$ , since the fixed point set of  $\varphi$  in  $B_N$  is affine ([5, 8.2.3]). By the  $m = 1, M = N + 1$  case of Theorem 1 we know that  $\varphi$  has a Schroeder map  $\sigma_\varphi$  with  $\sigma'_\varphi(0) = I$ . Moreover,

$$(\sigma_\varphi \circ \psi) \circ \varphi = \sigma_\varphi \circ \varphi \circ \psi = \varphi'(0)(\sigma_\varphi \circ \psi)$$

so  $\sigma_\varphi \circ \psi$  is a Schroeder map for  $\varphi$ . By Corollary 2 this tells us that  $\sigma_\varphi \circ \psi = S^{-1}BS\sigma_\varphi$  where  $S$  diagonalizes  $\varphi'(0)$  and  $B$  is linear. Differentiation of this equation gives  $\sigma'_\varphi(0)\psi'(0) = S^{-1}BS\sigma'_\varphi(0)$  so that in fact  $S^{-1}BS = \psi'(0)$  and  $\sigma_\varphi$  is a Schroeder map for both  $\varphi$  and  $\psi$ , with derivative at 0 equal to  $I$ .  $\square$

It need not be the case that  $\varphi$  and  $\psi$  have the same set of Schroeder maps; see Example 1 in the next section.

As a converse to this result we have the following theorem.

**Theorem 4.** *Suppose  $\varphi, \psi$  are analytic self-maps of  $B_N$ , each fixing 0, with  $\varphi'(0)\psi'(0) = \psi'(0)\varphi'(0)$ . Suppose further that there exists an analytic  $\sigma : B_N \rightarrow C^N$  with  $\sigma'(0)$  invertible and both  $\sigma \circ \varphi = \varphi'(0)\sigma$  and  $\sigma \circ \psi = \psi'(0)\sigma$ . Then  $\varphi \circ \psi = \psi \circ \varphi$ .*

*Proof.* Since  $\sigma$  is locally univalent near 0 we may write

$$\varphi = \sigma^{-1}\varphi'(0)\sigma$$

and

$$\psi = \sigma^{-1}\psi'(0)\sigma$$

in an open neighborhood of 0. Thus near 0 we have

$$\varphi \circ \psi = \sigma^{-1}\varphi'(0)\sigma\sigma^{-1}\psi'(0)\sigma = \sigma^{-1}\varphi'(0)\psi'(0)\sigma = \sigma^{-1}\psi'(0)\varphi'(0)\sigma = \psi \circ \varphi.$$

Since  $\varphi \circ \psi = \psi \circ \varphi$  in an open neighborhood of 0 and the compositions are defined on  $B_N$  we must have  $\varphi \circ \psi = \psi \circ \varphi$  in  $B_N$ .  $\square$

The last result need not hold if the hypothesis on the commutability of the derivatives at 0 is omitted: Take  $\varphi, \psi$  to be linear maps which do not commute. They share  $\sigma(z) = z$  as a common locally univalent Schroeder map.

### 3. Resonances.

We begin with several examples which will help set the stage for Theorem 7, the main result of this section.

**Example 1.** Let  $\varphi(z_1, z_2) = (c_1 z_1, c_1^3 z_2 + c_2 z_1^2)$  where  $c_1, c_2$  are sufficiently small non-zero constants so that  $\varphi(B_2) \subset B_2$ . Note that

$$\varphi'(0) = \begin{pmatrix} c_1 & 0 \\ 0 & c_1^3 \end{pmatrix}$$

and the resonance  $\lambda_2 = \lambda_1^3$  holds for the eigenvalues  $\lambda_1 = c_1, \lambda_2 = c_1^3$  of  $\varphi'(0)$ . It is easy to check that

$$\sigma_\varphi = \left( z_1, z_2 + \frac{c_2}{c_1^3 - c_1^2} z_1^2 \right)$$

is a Schroeder map for  $\varphi$  with derivative at 0 equal to  $I$  (this example is also discussed in [2]). By Corollary 2 all Schroeder maps for  $\varphi$  are of the form  $g \circ \sigma_\varphi$  where  $g$  is a polynomial map  $(b_1 z_1, b_2 z_2 + b_3 z_1^3)$  for arbitrary constants  $b_1, b_2$  and  $b_3$ , and thus have the form

$$\left( b_1 z_1, b_2 z_2 + \frac{b_2 c_2}{c_1^3 - c_1^2} z_1^2 + b_3 z_1^3 \right).$$

Now suppose that  $\psi$  commutes with  $\varphi$ . We know from the calculations in Theorem 3 that  $\sigma_\varphi \circ \psi$  is a Schroeder map for  $\varphi$  and hence  $\sigma_\varphi \circ \psi = g \circ \sigma_\varphi$  for  $g$  as above. From this we easily determine that  $\psi$  must be of the form

$$\left( b_1 z_1, b_2 z_2 + \frac{c_2}{c_1^3 - c_1^2} (b_2 - b_1^2) z_1^2 + b_3 z_1^3 \right)$$

for some constants  $b_1, b_2, b_3$ , and moreover any choice of these constants will give a map which commutes with  $\varphi$ . If these constants are chosen sufficiently small,  $\psi(B_2) \subset B_2$ . Note that whenever  $b_3 \neq 0$  we have a commuting map which is not an iterate of  $\varphi$ , so the set of maps which commute with  $\varphi$  is considerably larger than just the natural iterates of  $\varphi$ .

If  $b_2 \neq b_1^3$  then

$$\left( z_1, z_2 + \frac{c_2}{c_1^3 - c_1^2} z_1^2 + \frac{b_3}{b_2 - b_1^3} z_1^3 \right)$$

is a common Schroeder map for  $\varphi$  and  $\psi$  with derivative at 0 equal to  $I$ . We remark that

$$\left( z_1, z_2 + \frac{c_2}{c_1^3 - c_1^2} z_1^2 + \alpha z_1^3 \right)$$

where  $\alpha \neq b_3/(b_2 - b_1^3)$  is a Schroeder map for  $\varphi$  but not for  $\psi$ , so that while  $\varphi$  and  $\psi$  have a locally univalent Schroeder map in common, their sets of Schroeder maps are not the same.

On the other hand, if  $b_2 = b_1^3, b_3 \neq 0$  and

$$\psi(z_1, z_2) = \left( b_1 z_1, b_2 z_2 + \frac{c_2}{c_1^3 - c_1^2} (b_2 - b_1^2) z_1^2 + b_3 z_1^3 \right)$$

then  $\psi$  commutes with  $\varphi$  but  $\psi$  has *no* locally univalent Schroeder map by the converse of Theorem 1. One can check that the upper left  $7 \times 7$  corner of the matrix for  $C_\psi$  has diagonal entries  $1, b_1, b_1^3, b_1^2, b_1^4, b_1^6, b_1^3$  and three non-zero off diagonal entries:  $c_2(b_1^3 - b_1^2)/(c_1^3 - c_1^2)$  in the 4-3 position,  $c_2(b_1^3 - b_1^2)b_1/(c_1^3 - c_1^2)$  in the 7-5 position, and  $b_3 \neq 0$  in the 7-3 position. This matrix is not diagonalizable. Note that the situation being considered here is that of the resonances of  $\varphi$  also being resonances of  $\psi$ , where  $\psi$  is not a natural iterate of  $\varphi$ .

We also note that this example shows that two self-maps of the ball which each commute with  $\varphi$  need not commute with each other as

$$\psi_1(z_1, z_2) = \left( b_1 z_1, b_2 z_2 + \frac{c_2}{c_1^3 - c_1^2} (b_2 - b_1^2) z_1^2 + b_3 z_1^3 \right)$$

and

$$\psi_2(z_1, z_2) = \left( b_1 z_1, b_2 z_2 + \frac{c_2}{c_1^3 - c_1^2} (b_2 - b_1^2) z_1^2 + \frac{1}{2} b_3 z_1^3 \right)$$

both commute with  $\varphi$  but fail to commute with each other if  $b_1, b_2$  and  $b_3$  are chosen to be sufficiently small non-zero values with  $b_1^3 \neq b_2$ .

In two variables only one resonance relation is possible (either  $\lambda_1 = \lambda_2^n$  or  $\lambda_2 = \lambda_1^m$ ), but as the number of dimensions increases so do the possible variety of resonance equations. The next example, describing a general situation in  $C^3$  will be instructive for formulating a general theorem.

**Example 2.** Consider any analytic mapping  $\varphi : B_3 \rightarrow B_3$ , fixing 0 and not unitary on any slice, where  $\varphi'(0)$  is diagonal, with diagonal entries  $\lambda_j$  satisfying  $1 > |\lambda_1| > |\lambda_2| > |\lambda_3| > 0$  where the resonances

$$\lambda_2 = \lambda_1^n, (n \geq 2) \quad \text{and}$$

$$\lambda_3 = \lambda_1^m \lambda_2^k = \lambda_1^{m+nk} (m, k \geq 0, m+k \geq 2, \text{ and } m < n)$$

hold. Notice that we, in fact, have  $k + 1$  different resonances for  $\lambda_3$ :

$$\lambda_3 = \lambda_1^{m+nk} = \lambda_1^{r_1} \lambda_2 = \lambda_1^{r_2} \lambda_2^2 = \cdots = \lambda_1^{r_k} \lambda_2^k$$

where

$$(2) \quad r_j + jn = m + nk$$

for  $j = 1, \dots, k$ . If  $\varphi$  satisfies the hypotheses of Corollary 2 then all Schroeder maps are of the form  $g \circ \sigma_\varphi$  where  $\sigma_\varphi$  is a Schroeder map satisfying  $\sigma'_\varphi(0) = I$  and  $g$  is a polynomial mapping with

$$g_1 = b_1 z_1, g_2 = b_2 z_2 + c_1 z_1^n$$

and

$$g_3 = b_3 z_3 + c_2 z_1^{m+nk} + c_3 z_1^{r_1} z_2 + c_4 z_1^{r_2} z_2^2 + \cdots + c_{k+2} z_1^{r_k} z_2^k$$

for arbitrary choice of the coefficients. Denote the collection of all such polynomial maps  $\mathcal{G}_\varphi$ .

Now suppose  $\psi : B_3 \rightarrow B_3$  commutes with  $\varphi$  and that no resonance of  $\varphi$  is also a resonance of  $\psi$ . We know  $\sigma_\varphi \circ \psi$  is a Schroeder map for  $\varphi$  so  $\sigma_\varphi \circ \psi = g \circ \sigma_\varphi$  for some  $g \in \mathcal{G}_\varphi$ . Taking derivatives, we see that  $\sigma'_\varphi(0)\psi'(0) = g'(0)\sigma'_\varphi(0)$  and thus  $\psi'(0) = g'(0) = \text{diag}(b_1, b_2, b_3)$ . Our hypothesis on the resonances of  $\psi$  implies that  $b_2 \neq b_1^n, b_3 \neq b_1^{m+nk}, b_3 \neq b_1^{r_1} b_2, \dots, b_3 \neq b_1^{r_k} b_2^k$ .

We claim that there exists  $\hat{g}$  in  $\mathcal{G}_\varphi$  with  $\hat{g}'(0) = I$  solving  $\hat{g} \circ g = g'(0)\hat{g}$ . Once the claim is verified we see the following holds in a neighborhood of 0:

$$\begin{aligned} (\hat{g} \circ \sigma_\varphi) \circ \psi &= (\hat{g} \circ \sigma_\varphi) \circ \sigma_\varphi^{-1} \circ g \circ \sigma_\varphi = \hat{g} \circ g \circ \sigma_\varphi \\ &= g'(0) \circ \hat{g} \circ \sigma_\varphi = \psi'(0)(\hat{g} \circ \sigma_\varphi) \end{aligned}$$

since  $\psi = \sigma_\varphi^{-1} g \sigma_\varphi$  near 0. If  $(\hat{g} \circ \sigma_\varphi) \circ \psi = \psi'(0)(\hat{g} \circ \sigma_\varphi)$  holds near 0, then it holds in  $B_3$  since  $\hat{g}$  is defined on  $C^3$ . This shows that  $\hat{g} \circ \sigma_\varphi$  is a Schroeder map for  $\psi$  with derivative at 0 equal to  $I$ ; it is also a Schroeder map for  $\varphi$  by Corollary 2.

To verify the claim we will show that coefficients  $\hat{c}_1, \hat{c}_2, \dots, \hat{c}_{k+2}$  may be determined so that  $\hat{g}$  given by

$$\hat{g}_1 = z_1, \hat{g}_2 = z_2 + \hat{c}_1 z_1^n$$

and

$$\hat{g}_3 = z_3 + \hat{c}_2 z_1^{m+nk} + \hat{c}_3 z_1^{r_1} z_2 + \hat{c}_4 z_1^{r_2} z_2^2 + \cdots + \hat{c}_{k+2} z_1^{r_k} z_2^k$$

solves  $\hat{g} \circ g = g'(0)\hat{g}$ . Notice that  $\hat{g}_1 \circ g = g_1 = b_1 z_1 = b_1 \hat{g}_1$  and that  $\hat{g}_2 \circ g = b_2 \hat{g}_2$  provided  $\hat{c}_1 = c_1/(b_2 - b_1^n)$ ; the hypothesis  $b_2 \neq b_1^n$  being used here.

Finally, we turn to

$$(3) \quad \hat{g}_3 \circ g = b_3 \hat{g}_3.$$



Using the forms of  $\hat{g}_3$  and  $g$ , we expand the left side of Equation (3) into a sum of monomials and observe that each of these monomials is a scalar multiple of a monomial which also appears in  $b_3\hat{g}_3$ , the right side of Equation (3). To see this, observe that when we expand  $g_1^{r_j} g_2^j = (b_1 z_1)^{r_j} (b_2 z_2 + c_1 z_1^n)^j$  we get terms which are scalar multiples of the monomials  $z_1^{r_j} z_2^s (z_1^n)^{j-s} = z_1^{r_j+n(j-s)} z_2^s$  ( $0 \leq s \leq j$ ). Since  $r_j + n(j-s) = r_s$ , this monomial, with some scalar coefficient, appears in  $b_3\hat{g}_3$ .

By equating in turn the coefficients of

$$z_1^{r_k} z_2^k, z_1^{r_{k-1}} z_2^{k-1}, \dots, z_1^{r_1} z_2, z_1^{m+nk},$$

we obtain equations for the unknown coefficients  $\hat{c}_{k+2}, \hat{c}_{k+1}, \dots, \hat{c}_2$ . The equation obtained from the coefficients of  $z_1^{r_k} z_2^k$  is

$$c_{k+2} + \hat{c}_{k+2} b_2^k b_1^{r_k} = b_3 \hat{c}_{k+2}$$

which may be solved for  $\hat{c}_{k+2}$  provided  $b_3 \neq b_1^{r_k} b_2^k$ ; this is guaranteed by the hypothesis on the resonances of  $\psi$ . Continuing, suppose that by comparing the coefficients of  $z_1^{r_k} z_2^k, z_1^{r_{k-1}} z_2^{k-1}, \dots, z_1^{r_{j+1}} z_2^{j+1}$  the coefficients  $\hat{c}_{k+2}, \hat{c}_{k+1}, \dots, \hat{c}_{j+3}$  have been determined. Next we compare coefficients of  $z_1^{r_j} z_2^j$  on both sides of Equation (3). None of the terms

$$\hat{c}_2 g_1^{m+nk}, \hat{c}_3 g_1^{r_1} g_2, \dots, \hat{c}_{j+1} g_1^{r_{j-1}} g_2^{j-1}$$

contribute any terms of the form  $z_1^{r_j} z_2^j$ . The expansion of  $\hat{c}_{j+2} g_1^{r_j} g_2^j$  contributes a term  $\hat{c}_{j+2} b_1^{r_j} b_2^j z_1^{r_j} z_2^j$ . The expansions of

$$\hat{c}_{j+3} g_1^{r_{j+1}} g_2^{j+1}, \dots, \hat{c}_{k+2} g_1^{r_k} g_2^k$$

contribute terms  $z_1^{r_j} z_2^j$  all of whose coefficients involve the previously determined coefficients  $\hat{c}_{j+3}, \dots, \hat{c}_{k+2}$  (and  $b_1, b_2$ ). Thus equating the coefficients of  $z_1^{r_j} z_2^j$  on both sides of Equation (3) leads to an equation of the form

$$c_{j+2} + \hat{c}_{j+2} b_1^{r_j} b_2^j + \text{known terms} = b_3 \hat{c}_{j+2}$$

where “known terms” refers to a sum involving the known values  $\hat{c}_{j+3}, \dots, \hat{c}_{k+2}$  and the  $b_i$ ’s. This equation may be solved for  $\hat{c}_{j+2}$  provided  $b_3 \neq b_1^{r_j} b_2^j$ , which is part of our hypothesis. Continuing this process we determine all of the coefficients of the second and higher order terms of  $\hat{g}_j$ . Note that the only first order term in  $\hat{g}_3 \circ g$  is  $b_3 z_3$  and this is the only first order term on the right side of Equation (3). Thus we have found a choice of coefficients so that  $\hat{g} \circ g = g'(0)\hat{g}$ , verifying the claim.

We set some notation and terminology which will be useful in the main result. We now restrict attention to the case that the eigenvalues  $\varphi'(0)$  are distinct, non-zero, and of modulus less than 1. There is no loss of generality in assuming that  $\varphi'(0)$  is upper triangular, with diagonal entries  $\lambda_1, \lambda_2, \dots, \lambda_N$  satisfying  $1 > |\lambda_1| \geq |\lambda_2| \geq \dots \geq |\lambda_N| > 0$ , since there is a

unitary map  $U$  so that  $U^*\varphi'(0)U$  is upper triangular with the eigenvalues of  $\varphi'(0)$  appearing in the prescribed order. Moreover, if  $\varphi$  and  $\psi$  commute, then so do  $U^*\varphi U$  and  $U^*\psi U$ , and  $\varphi$  and  $\psi$  have a common locally univalent Schroeder map if and only if  $U^*\varphi U$  and  $U^*\psi U$  do. This ordering on the eigenvalues of  $\varphi'(0)$  implies that  $\lambda_1$  has no resonance relations, and in general a resonance for  $\lambda_j$  is of the form

$$\lambda_j = \lambda_1^{k_1} \lambda_2^{k_2} \cdots \lambda_{j-1}^{k_{j-1}}$$

where  $k_i \geq 0$  and  $\sum k_i \geq 2$ .

For  $j \geq 2$  we say that a monomial  $cz_1^{k_1} z_2^{k_2} \cdots z_{j-1}^{k_{j-1}}$  ( $c$  any non-zero scalar) is *j-permissible* (for  $\varphi$ ) if

$$\lambda_j = \lambda_1^{k_1} \lambda_2^{k_2} \cdots \lambda_{j-1}^{k_{j-1}};$$

call the corresponding multi-index  $(k_1, k_2, \dots, k_{j-1}, 0, \dots, 0)$  *j-permissible* as well. There is a one-to-one correspondence between a resonance for  $\lambda_j$  and a *j-permissible* monomial with scalar coefficient 1 (or a *j-permissible* multi-index). For a given  $\varphi$ , let  $\Gamma_j$  denote the *j-permissible* multi-indices, so that  $(k_1, k_2, \dots, k_{j-1}, 0, \dots, 0) \in \Gamma_j$  if and only if  $\lambda_j = \lambda_1^{k_1} \lambda_2^{k_2} \cdots \lambda_{j-1}^{k_{j-1}}$  and  $\Gamma_j$  is empty if  $\lambda_j$  has no resonance relations. We order the multi-indices in  $\Gamma_j$  according to the following rule: A multi-index  $\alpha$  in  $\Gamma_j$  precedes a multi-index  $\beta$  if either the  $k_{j-1}$  entry of  $\alpha$  is greater than the  $k_{j-1}$  entry of  $\beta$ , or if the entries in the  $k_i$  through  $k_{j-1}$  positions agree for some  $i < j$ , then the  $k_{i-1}$  entry of  $\alpha$  is greater than the  $k_{i-1}$  entry of  $\beta$ . This is not the “usual” ordering on multi-indices. For example, if  $\varphi$  has resonance relations  $\lambda_2 = \lambda_1^2$ ,  $\lambda_3 = \lambda_1^3 = \lambda_1 \lambda_2$  and

$$(4) \quad \lambda_4 = \lambda_3^2 \lambda_1 = \lambda_2^2 \lambda_3 = \lambda_1^2 \lambda_2 \lambda_3 = \lambda_2^3 \lambda_1 = \lambda_1^4 \lambda_3 = \lambda_2^2 \lambda_1^3 = \lambda_1^5 \lambda_2 = \lambda_1^7$$

then the ordering on  $\Gamma_4$  is

$$(1, 0, 2, 0), (0, 2, 1, 0), (2, 1, 1, 0), (4, 0, 1, 0)$$

$$(1, 3, 0, 0), (3, 2, 0, 0), (5, 1, 0, 0), (7, 0, 0, 0).$$

Recall the notation  $\mathcal{G}_\varphi$  is used for the collection of all polynomial mappings  $g = (g_1, g_2, \dots, g_N)$  where

$$g_j(z_1, \dots, z_N) = b_j z_j + \sum_{\gamma \in \Gamma_j} c^j(\gamma) z^\gamma$$

where the coefficients  $b_j$  and  $c^j(\gamma)$  are arbitrary.

**Lemma 5.** *With  $\varphi$  as just described, suppose  $g \in \mathcal{G}_\varphi$  and  $\hat{g} \in \mathcal{G}_\varphi$  with  $\hat{g}'(0) = I$ . Then the monomials of order at least two in the expansion of  $\hat{g}_j \circ g$  are all *j-permissible*, for  $j \geq 2$ .*

*Proof.* The coordinate functions  $\hat{g}_j$  are of the form

$$\hat{g}_j = z_j + \sum \hat{c}^j(\gamma) z^\gamma$$

where the sum is over all multi-indices  $\gamma$  in  $\Gamma_j$ . Thus

$$\hat{g}_j \circ g = g_j + \sum_{\Gamma_j} \hat{c}^j(\gamma) g^\gamma$$

and it suffices to show that each monomial in the expansion of  $g^\gamma$  is  $j$ -permissible. Consider  $g^\gamma$  where  $\gamma = (\gamma_1, \gamma_2, \dots, \gamma_{j-1}, 0, \dots, 0)$ ,  $\gamma_i \geq 0$ ,  $\sum \gamma_i \geq 2$ . Computing a term of  $g^\gamma = g_1^{\gamma_1} g_2^{\gamma_2} \cdots g_{j-1}^{\gamma_{j-1}}$  involves making a choice of  $\gamma_1$  terms from  $g_1$  (necessarily each of these will be  $b_1 z_1$ ),  $\gamma_2$  terms from  $g_2$ ,  $\gamma_3$  terms from  $g_3$ , etc. Since  $\lambda_j = \lambda_1^{\gamma_1} \lambda_2^{\gamma_2} \cdots \lambda_{j-1}^{\gamma_{j-1}}$ , making these successive choices produces a  $j$ -permissible monomial.  $\square$

As an example, again suppose as above that  $\varphi$  has resonances  $\lambda_2 = \lambda_1^2$ ,  $\lambda_3 = \lambda_1^3 = \lambda_1 \lambda_2$  and  $\lambda_4$  has the resonance relations in Equation (4). In the expansion of  $\hat{g}_4 \circ g$  we obtain, for example, the terms from  $g_2^2 g_3$  since  $(0, 2, 1, 0) \in \Gamma_4$ . The monomials obtained by choosing two terms from  $g_2$  (either  $b_2 z_2$  or a multiple of  $z_1^2$ ) and one from  $g_3$  (either  $b_3 z_3$ , a multiple of  $z_1^3$  or a multiple of  $z_1 z_2$ ) are all in  $\Gamma_4$ .

Since Theorem 3 applies when  $\varphi$  has no resonances, in the next two results we consider the resonant case.

**Theorem 6.** *Let  $\varphi$  be an analytic self-map of  $B_N$  fixing 0 and not unitary on any slice with  $\varphi'(0)$  upper triangular with distinct diagonal entries  $\lambda_j$  satisfying*

$$1 > |\lambda_1| \geq |\lambda_2| \geq \cdots \geq |\lambda_N| > 0.$$

*Suppose  $\varphi$  has at least one resonance relation. Let  $g \in \mathcal{G}_\varphi$  with*

$$g'(0) = \text{diag}\{b_1, b_2, \dots, b_N\}$$

*and assume that whenever a resonance relation*

$$\lambda_j = \lambda_1^{k_1} \cdots \lambda_{j-1}^{k_{j-1}}$$

*holds then*

$$b_j \neq b_1^{k_1} \cdots b_{j-1}^{k_{j-1}}.$$

*Then there exists  $\hat{g} \in \mathcal{G}_\varphi$  with  $\hat{g}'(0) = I$  and  $\hat{g} \circ g = g'(0)\hat{g}$ .*

*Proof.* By hypotheses the coordinate functions of  $g$  are

$$g_j(z_1, \dots, z_N) = b_j z_j + \sum_{\gamma \in \Gamma_j} c^j(\gamma) z^\gamma.$$

Set

$$\hat{g}_j(z_1, \dots, z_N) = z_j + \sum_{\gamma \in \Gamma_j} \hat{c}^j(\gamma) z^\gamma$$

so that  $\hat{g}'(0) = I$ . We need only show that the coefficients  $\hat{c}^j(\gamma)$  may be determined so that

$$(5) \quad \hat{g} \circ g = g'(0) \hat{g}$$

holds. We will determine these coefficients in the order of the multi-indices in  $\Gamma_j$ . For  $\gamma = (\gamma_1, \dots, \gamma_{j-1}, 0, \dots, 0) \in \Gamma_j$ , write  $b^\gamma$  for  $b_1^{\gamma_1} b_2^{\gamma_2} \dots b_{j-1}^{\gamma_{j-1}}$ .

If  $\Gamma_j = \emptyset$ , then  $\hat{g}_j \circ g = b_j \hat{g}_j$  holds automatically. If  $\Gamma_j \neq \emptyset$ , let  $\tau_1$  be the first multi-index in  $\Gamma_j$ , and compare the coefficients of  $z^{\tau_1}$  on both sides of

$$(6) \quad \hat{g}_j \circ g = b_j \hat{g}_j$$

to obtain

$$c^j(\tau_1) + \hat{c}^j(\tau_1) b^{\tau_1} = b_j \hat{c}^j(\tau_1)$$

which can be solved for the unknown  $\hat{c}^j(\tau_1)$  since  $b_j \neq b^{\tau_1}$ .

Next compare the coefficients of  $z^{\tau_2}$  in Equation (6), where  $\tau_2$  is the second multi-index of  $\Gamma_j$ . Only for  $\gamma = \tau_1, \tau_2$  can  $g^\gamma$  contribute a  $z^{\tau_2}$  term. Thus we are led to the equation

$$c^j(\tau_2) + b^{\tau_2} \hat{c}^j(\tau_2) + \dots = b_j \hat{c}^j(\tau_2)$$

where  $\dots$  indicates terms depending only on coefficients of  $g$  and/or the just determined value  $\hat{c}^j(\tau_1)$ . This can be solved for  $\hat{c}^j(\tau_2)$  since  $b^{\tau_2} \neq b_j$ . Proceeding in this way through the multi-indices of  $\Gamma_j$  in the prescribed order we obtain equations

$$c^j(\tau_k) + b^{\tau_k} \hat{c}^j(\tau_k) + \dots = b_j \hat{c}^j(\tau_k)$$

where the omitted terms on the left are known quantities, possibly involving the coefficients  $\hat{c}^j(\tau_i)$  where  $\tau_i$  precedes  $\tau_k$ .

At this point we have determined  $\hat{c}^j(\gamma)$ ,  $\gamma \in \Gamma_j$  so that in Equation (6) the coefficients of any  $z^\tau$ ,  $\tau \in \Gamma_j$  agree on both sides. Recall that by Lemma 5 the monomials of order at least two which appear in the expansion of the left side of Equation (6) are all  $j$ -permissible, so in fact we have shown that the coefficients of  $z^\tau$  for any multi-index  $\tau$  of total order at least two on both sides of the equation agree. The only non-zero first order terms on either side of Equation (6) are  $b_j z_j$ . Hence with the determined values of  $\hat{c}^j(\gamma)$ , Equation (5) holds.  $\square$

**Theorem 7.** *Let  $\varphi : B_N \rightarrow B_N$  be analytic such that  $\varphi(0) = 0$ ,  $\varphi$  is not unitary on any slice, and  $A = \varphi'(0)$  is upper triangular with distinct diagonal entries  $\lambda_j$  satisfying  $1 > |\lambda_1| \geq |\lambda_2| \geq \dots \geq |\lambda_N| > 0$ . Assume that  $\varphi$  has resonances so  $\Gamma_j$  is non-empty for at least one  $j$ . Suppose  $\varphi$  has a Schroeder map  $\sigma_\varphi$  with  $\sigma'_\varphi(0) = I$ . If  $\varphi \circ \psi = \psi \circ \varphi$  for some analytic self-map  $\psi$  of*

$B_N$ , and the resonances of  $\varphi$  are not also resonances of  $\psi$ , then  $\varphi$  and  $\psi$  have a common Schroeder map which is locally univalent near 0.

Before giving the proof, we clarify the meaning of the hypothesis “the resonances of  $\varphi$  are not also resonances of  $\psi$ ”. Since  $\varphi$  and  $\psi$  commute, so do  $\varphi'(0)$  and  $\psi'(0)$ . Since  $\varphi'(0)$  is assumed to have distinct eigenvalues, this means that  $\varphi'(0)$  and  $\psi'(0)$  may be simultaneously diagonalized, and we may find an  $N \times N$  invertible matrix  $S$  so that

$$S\varphi'(0)S^{-1} = \text{diag}(\lambda_1, \lambda_2, \dots, \lambda_N)$$

and

$$S\psi'(0)S^{-1} = \text{diag}(\mu_1, \mu_2, \dots, \mu_N)$$

where the  $\lambda_j$ 's appear in non-increasing order, but there is no apriori ordering on the  $\mu_j$ 's. To say that the resonances of  $\varphi$  are not also resonances of  $\psi$  means that if

$$\lambda_j = \lambda_1^{k_1} \cdots \lambda_{j-1}^{k_{j-1}}$$

then

$$\mu_j \neq \mu_1^{k_1} \cdots \mu_{j-1}^{k_{j-1}}$$

(with the given ordering on the  $\mu_j$ 's).

*Proof.* If  $\varphi \circ \psi = \psi \circ \varphi$  we have already observed that  $\sigma_\varphi \circ \psi$  is a Schroeder map for  $\varphi$ . By Corollary 2, we must have

$$\sigma_\varphi \circ \psi = S^{-1} \circ g \circ S \circ \sigma_\varphi$$

where  $S$  diagonalizes  $\varphi'(0)$  and  $\psi'(0)$  as just described and  $g \in \mathcal{G}_\varphi$  so that the coordinate functions of  $g$  are

$$g_j(z_1, \dots, z_N) = b_j z_j + \sum_{\gamma \in \Gamma_j} c^j(\gamma) z^\gamma.$$

Upon differentiation of the relation  $\sigma_\varphi \circ \psi = S^{-1} g S \sigma_\varphi$  we see that  $\psi'(0) = S^{-1} g'(0) S$  so that

$$\text{diag}(\mu_1, \dots, \mu_N) = S\psi'(0)S^{-1} = g'(0) = \text{diag}(b_1, \dots, b_N)$$

and by hypothesis  $\lambda_j = \lambda_1^{k_1} \cdots \lambda_{j-1}^{k_{j-1}} \Rightarrow b_j \neq b_1^{k_1} \cdots b_{j-1}^{k_{j-1}}$ . By Theorem 6, there exists  $\hat{g} \in \mathcal{G}_\varphi$  with  $\hat{g}'(0) = I$  and  $\hat{g} \circ g = g'(0)\hat{g}$ . By Corollary 2,  $S^{-1}\hat{g}S\sigma_\varphi$  is a Schroeder map for  $\varphi$ ; its derivative at 0 is  $I$ . The following calculation shows that it is also a Schroeder map for  $\psi$ :

$$\begin{aligned} (S^{-1}\hat{g}S\sigma_\varphi)\psi &= S^{-1}\hat{g}SS^{-1}gS\sigma_\varphi &= S^{-1}\hat{g}gS\sigma_\varphi \\ &= S^{-1}g'(0)\hat{g}S\sigma_\varphi \\ &= (S^{-1}g'(0)S)S^{-1}\hat{g}S\sigma_\varphi \\ &= \psi'(0)(S^{-1}\hat{g}S\sigma_\varphi). \end{aligned}$$

□

In Example 1 we saw that Theorem 7 can fail if the resonances of  $\varphi$  are also resonances of  $\psi$ . Of course, if  $\psi$  is a natural iterate of  $\varphi$ , then  $\varphi$  and  $\psi$  will commute, have the same resonances, and have a common Schroeder map.

One application of Theorems 7 and 3 is to extract qualitative information about the maps which commute with a given map. Our next theorem is a result in this direction. It depends on the following result from [2].

**Proposition 8** ([2]). *Let  $\varphi$  be an analytic map of  $B_N$  into itself with  $\varphi(0) = 0$  and  $A = \varphi'(0)$  invertible and suppose  $\varphi$  is not unitary on any slice of  $B_N$ . If  $\sigma_\varphi$  is an analytic map of  $B_N$  into  $\mathbf{C}^N$  that solves Schroeder's functional equation  $\sigma_\varphi \circ \varphi = Af$  and  $\sigma'_\varphi(0)$  is invertible, then  $\sigma_\varphi$  is univalent on  $B_N$  if and only if  $\varphi$  is univalent on  $B_N$ .*

**Corollary 9.** *Suppose  $\varphi$  and  $\psi$  are commuting analytic self-maps of  $B_N$ , both fixing 0, not unitary on any slice, and having invertible derivative at 0. Suppose further that they satisfy the hypotheses of either Theorem 3 or Theorem 7. Then if  $\varphi$  is univalent in  $B_N$  so is  $\psi$ .*

*Proof.* There is a common locally univalent Schroeder map for  $\varphi$  and  $\psi$  which by the “if” direction of Proposition 8 is univalent in  $B_N$ . Now apply the “only if” direction of the proposition for  $\psi$  to conclude that  $\psi$  is univalent in  $B_N$ .  $\square$

The invertibility of  $\varphi'(0)$  and  $\psi'(0)$  is necessary in this corollary, as the maps  $\varphi(z_1, z_2) = (1/2z_1, 1/3z_2)$  and  $\psi(z_1, z_2) = (1/2z_1, 0)$  which share the Schroeder map  $\sigma(z_1, z_2) = (z_1, z_2)$  show.

Finally, we observe that our proof of Theorem 7 depends on the hypothesis that the eigenvalues of  $\varphi'(0)$  are distinct. This hypothesis plays a significant role in Theorem 6 as it means each coordinate function  $g_j$  has at most one non-zero linear term. We leave consideration of the repeated eigenvalue case for a later time.

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