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LOGARITHMIC SOBOLEV INEQUALITIES
ALONG RICCI-HARMONIC MAP FLOW**

ABIMBOLA ABOLARINWA

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This paper introduces a new family of entropy functionals which is proved to be monotonically nondecreasing along the Ricci-harmonic map heat flow. Some of the consequences of the monotonicity are combined to derive gradient estimates and Harnack inequalities for all positive solutions to the associated conjugate heat equation. We relate the entropy monotonicity and the ultracontractivity property of the heat semigroup, and as a result we obtain the equivalence of logarithmic Sobolev inequalities, conjugate heat kernel upper bounds and uniform Sobolev inequalities under Ricci-harmonic map heat flow.

1. Introduction

Let (M, g) and (N, ξ) be compact Riemannian manifolds (without boundary) of dimensions m and n respectively. Let a smooth map $u : M \rightarrow N$ be a critical point of the Dirichlet energy integral $E(u) = \int_M |\nabla u|^2 d\mu_g$, where N is isometrically embedded in \mathbb{R}^d , $d \geq n$, by the Nash embedding theorem. The configuration $(g(x, t), u(x, t))$, $t \in [0, T)$, of a one-parameter family of Riemannian metrics $g(x, t)$ and a family of smooth maps $u(x, t)$ is defined to be a Ricci-harmonic map flow if it satisfies the coupled system of nonlinear parabolic equations denoted by $(RH)_\alpha$

$$(1-1) \quad \begin{cases} \frac{\partial}{\partial t} g(x, t) = -2 \operatorname{Rc}(x, t) + 2\alpha \nabla u(x, t) \otimes \nabla u(x, t), \\ \frac{\partial}{\partial t} u(x, t) = \tau_g u(x, t), \end{cases}$$

where $\operatorname{Rc}(x, t)$ is the Ricci curvature tensor for the metric g , $\alpha(t) \equiv \alpha > 0$ is a time-dependent coupling constant and $\tau_g u$ is the intrinsic Laplacian of u which denotes the tension field of the map u . The system (1-1) was first studied by B. List [2008] in a special case, $N \subseteq \mathbb{R}$ and $\alpha = 2$, where the flow was modified by the Lie derivative of g with respect to a gradient vector field to give a gradient flow

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of an energy functional whose stationary points are solutions to the static Einstein vacuum equations arising in general relativity. This has since been generalised by R. Müller [2012] to the general case $N \hookrightarrow \mathbb{R}^d$, for sufficiently large d . Precisely, the system couples together the Ricci flow of Hamilton [1982] and the heat flow for harmonic maps of Eells and Sampson [1964]. The system $(RH)_\alpha$ is closer to the former in behaviours, such as in existence and singularities, though may be less singular than both. Hence, the analysis of the flow is usually done along the line of Ricci flow and for this, Perelman’s works [2002; 2003b; 2003a] on Ricci flow are very applicable to the theory and applications of the Ricci-harmonic map flow.

In this paper we study the behaviour of all positive solutions to the associated conjugate heat equation along the Ricci-harmonic map flow. Let $h, H : M \times [0, T] \rightarrow (0, \infty)$ satisfy

$$\left(\frac{\partial}{\partial t} - \Delta_g\right)h = 0 \quad \text{and} \quad \left(-\frac{\partial}{\partial t} - \Delta_g + R - \alpha|\nabla u|_g^2\right)H = 0,$$

with

$$\int_0^T \int_M \left(\frac{\partial}{\partial t} - \Delta_g\right)hH \, d\mu_g \, dt = \int_0^T \int_M h\left(-\frac{\partial}{\partial t} - \Delta_g + R - \alpha|\nabla u|_g^2\right)H \, d\mu_g \, dt,$$

where Δ_g is the usual Laplace–Beltrami operator and $\square^* := -\partial/\partial t - \Delta_g + R - \alpha|\nabla u|_g^2$ is the standard conjugate to the heat operator $\square := \partial/\partial t - \Delta_g$. We say h and H are respectively solutions to the heat equation and conjugate heat equation. The main idea here is to solve the Ricci-harmonic map flow forward in time and solve the conjugate heat equation backward in time. Fixing the coordinate (y, s) , $H = H(x, t; y, s)$ will be called the conjugate heat kernel (the positive minimal solution) if it tends to a δ -function as $t \rightarrow T$.

Our main results in the first part of this paper are Perelman’s differential Harnack estimates for $f \in C^\infty(M \times [0, T])$ satisfying $H(x, \tau; y, s) = (4\pi\tau)^{-m/2}e^{-f(x,\tau)}$, $\tau = T - t$,

$$(1-2) \quad -\frac{d}{dt}f(\gamma(\tau), \tau) \leq \frac{1}{2}\left(|\gamma'(\tau)|^2 + S_g(\gamma(\tau), \tau) - \frac{m}{2\tau}\right),$$

and Li–Yau Harnack estimates for all positive solutions to the conjugate heat equation

$$(1-3) \quad \frac{H(x_2, t_2)}{H(x_1, t_1)} \leq \left(\frac{\tau_1}{\tau_2}\right)^{m/4} \exp\left(\frac{1}{2} \int_{t_1}^{t_2} (|\gamma'(t)|^2 + S_g(\gamma(t), t)) \, dt\right),$$

where $S_g = R_g - \alpha|\nabla u|_g^2$. (The proofs of (1-2) and (1-3) are delayed until Section 4). Both results stated above are consequences of a monotonicity formula for a new entropy functional $\mathcal{W}_{\alpha,\epsilon}$ introduced in Section 3, where we obtain the Harnack

inequality for $0 < \epsilon^2 \leq 4\pi$,

$$(1-4) \quad \frac{\epsilon^2 \tau}{4\pi} (2\Delta f - |\nabla f|^2 + S_g) + f - \frac{m\epsilon^2}{4\pi} \leq 0$$

for all time t and prove that

$$(1-5) \quad \lim_{\tau \rightarrow 0} \int_M \frac{\epsilon^2 \tau}{4\pi} (2\Delta f - |\nabla f|^2 + S_g) + f - \frac{m\epsilon^2}{4\pi} hH d\mu_g \rightarrow 0,$$

with the condition that $\epsilon^2 \rightarrow 4\pi$ as $t \rightarrow T$. Monotonicity formulas are generally useful in controlling solutions of evolution equations. This entropy is also intimately related to the logarithmic Sobolev inequality of Gross [1975]. Perelman used this property to obtain upper bounds for the fundamental solution to the adjoint heat equation via his reduced length. This leads to the proof of the noncollapsing theorem on Riemannian manifolds and, consequently, to the completion of R. Hamilton's program on the Poincaré conjecture. See [Perelman 2002; 2003b; 2003a; Cao et al. 2003]. Among several examples, Perelman's entropy and the gradient estimates of Li and Yau [1986] are important ones that show close relations between entropy monotonicity and the gradient estimate for the heat equation (forward or backward in time). Lei Ni [2004] has also considered a case for the linear heat equation on a static manifold with nonnegative Ricci curvature. We notice that coupling a heat-type equation with geometric flow began with [Hamilton 1993] and it has since become a very active research area and has led to numerous physical and geometric applications; for examples, see [Băileşteanu et al. 2010; Băileşteanu and Tran 2013; Cao and Zhang 2011; Kuang and Zhang 2008; List 2008; Müller 2012; Ni 2006; Zhang 2007] and the references therein.

Another important application of Perelman's \mathcal{W} -entropy monotonicity is in the derivation of uniform Sobolev inequalities by Q. Zhang [2007]; see also [Hsu 2008; Ye 2007]. In the second part of this paper, we relate the entropy monotonicity and the ultracontractivity property of the heat semigroup, and as a result we establish the equivalence of logarithmic Sobolev inequalities, conjugate heat kernel upper bounds and uniform Sobolev inequalities under Ricci-harmonic map heat flow. Precisely, let A_0 and B_0 be finite positive constants depending only on m , g_0 , the lower bound for the Ricci curvature and the injectivity radius of M . For any $v \in W^{1,2}(M, g_0)$ such that

$$(1-6) \quad \|v\|_{2m/(m-2)} \leq A_0 \|\nabla v\|_2 + B_0 \|v\|_2,$$

where $m \geq 3$ and $\|\cdot\|_q = \left(\int_M |\cdot|^q d\mu_g\right)^{1/q}$, $1 \leq p < \infty$, we have the following result.

Theorem. *Let M be a compact Riemannian manifold of dimension $m \geq 3$. Let the solution to the $(\text{RH})_\alpha$ -flow exist for all times $t \in [0, T)$. Assume the Sobolev*

embedding (1-6) holds; then for finite positive constants A and B depending on m, A_0, B_0 , the lower bound for R_{g_0} and T ,

$$(1-7) \quad \left(\int_M v^{2m/(m-2)} d\mu_g \right)^{(m-2)/2} \leq A \int_M (|\nabla v|^2 + \frac{1}{4} S_g v^2) d\mu_g + B \int_A v^2 d\mu_g$$

and

$$(1-8) \quad \int_M v^2 \ln v^2 d\mu_{g(t)} \leq \sigma^2 \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g(t)} - \frac{m}{2} \ln \sigma^2 + (t + \sigma^2)\beta_1 + \frac{m}{2} \ln \frac{mA}{2e},$$

hold for each $t \in [0, T)$ and $v \in W^{1,2}(M)$ if $\lambda_\alpha = \inf_{\|v\|_2=1} \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0}$; that is, $\lambda_{\alpha 0}$ is the first eigenvalue of the operator $-\Delta + \frac{1}{4} S_g$.

Finally, for some constant C depending on m, t, T, A_0, B_0 and $\sup S_g(\cdot, 0)$, the estimate

$$(1-9) \quad H(x, T; y) \leq CT^{-m/2}$$

for the positive solution to the conjugate heat equation associated to $(RH)_\alpha$ holds.

The three results in the above theorem are essentially equivalent, and their proofs occupy Sections 5 – 7. The approach to the proof here is *Sobolev inequality (1-7) \implies log-Sobolev inequality (1-8) \implies heat kernel upper bound (1-9) \implies Sobolev inequality (1-7)*. Indeed, any of them can be derived from the other. The results of the above form [Hsu 2008; Ye 2007; Zhang 2007] yield a long time κ -noncollapsing estimate which generalises Perelman’s short time result [2002] along the Ricci flow.

We recall that the nonnegativity of the scalar curvature R_g is preserved along Ricci flow [Chow and Knopf 2004], so the nonnegativity of S_g is also preserved as long as $(RH)_\alpha$ exists. Indeed, S_g evolves by a reaction-diffusion equation which helps to visualise its behaviour up to singular time (we discuss this in the next section). The condition $S_g = R_g - \alpha|\nabla u|_g^2 \geq 0$ at the starting time $t = 0$ must now be considered. The assumption is not necessary for the derivation of (1-7) since additional geometric data are not usually required to derive a Sobolev inequality from either a log-Sobolev inequality or the heat kernel bound. The assumption is required for the condition that a certain eigenvalue λ_α for the initial metric is positive, which is required to pass to (1-8). The class of manifold (M, g_0) with $\lambda_{\alpha 0} > 0$ is a very large one and significant from a geometric point of view. Moreover, if $\lambda_{\alpha 0} > 0$ for $S_g(0) > 0$ (i.e., $R_g(0) > \alpha(0)|\nabla u(0)|^2$) then A, B are independent of time and $B = 0$. Corollary 7.5 below presents corresponding versions of (1-7) and (1-8) in this case.

In the next section we discuss necessary background on Perelman–Müller entropy monotonicity formulas for $(RH)_\alpha$.

2. Background on entropy formulas for $(\text{RH})_\alpha$ -flow

Let (M, g) be a compact Riemannian manifold. For the metric g , any smooth functions $u \in C^\infty(M, N)$, $u(M) \subseteq N \hookrightarrow \mathbb{R}^n$, $f \in C^\infty(M)$ and constant $\alpha > 0$, Perelman and Müller's energy functional [Müller 2012] is defined on the triple (g, u, f) by

$$(2-1) \quad \mathcal{F}_\alpha(g, u, f) := \int_M (R_g + |\nabla f|_g^2 - \alpha |\nabla u|_g^2) e^{-f} d\mu_g,$$

which can also be written in two other ways,

$$\begin{aligned} \mathcal{F}_\alpha(g, u, f) &= \int_M (S_g + \Delta_g f) e^{-f} d\mu_g \\ &= \int_M (2\Delta_g f - |\nabla f|_g^2 + S_g) e^{-f} d\mu_g, \end{aligned}$$

since $\int_M \Delta(e^{-f}) = 0 = \int_M (-\Delta f + |\nabla f|_g^2) e^{-f} d\mu_g$. For any diffeomorphism $\phi : M \rightarrow M$, we have $\mathcal{F}_\alpha(\phi^*g, \phi^*u, \phi^*f) = \mathcal{F}_\alpha(g, u, f)$. If (g, u) is a solution to the system (1-1), Müller [2012] proved that the \mathcal{F}_α -functional is nondecreasing under the flow and showed that the system is equivalent (after pulling back with a diffeomorphism generated by a vector field) to the gradient flow system for the energy functional \mathcal{F}_α , locally written as,

$$(2-2) \quad \begin{cases} \frac{\partial}{\partial t} g_{ij} = -2R_{ij} + 2\alpha \nabla_i u \otimes \nabla_j u + 2\nabla_i \nabla_j f, \\ \frac{\partial}{\partial t} u = \tau_g u - \langle \nabla u, \nabla f \rangle, \\ \frac{\partial}{\partial t} f = -R + \alpha |\nabla u|^2 - \Delta f. \end{cases}$$

More precisely,

$$(2-3) \quad \begin{aligned} \frac{d}{dt} \mathcal{F}_\alpha(g, u, f) \\ = 2 \int_M (|\text{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f|^2 + \alpha |\tau_g u - \langle \nabla u, \nabla f \rangle|^2) e^{-f} d\mu_g \geq 0. \end{aligned}$$

An application of this is that \mathcal{F}_α is constant if and only if (g, u) is a steady gradient soliton.

Define

$$\lambda_\alpha(g) = \inf \left\{ \mathcal{F}_\alpha(g, u, f) : f \in C^\infty(M), \int_M e^{-f} d\mu_g = 1 \right\}.$$

Then $\lambda_\alpha(g)$ is the first eigenvalue of the operator $-4\Delta + S_g$, where the nondecreasing property of \mathcal{F}_α implies $\lambda_\alpha(g)$ is nondecreasing and we have, by setting $v = e^{-f/2}$,

the corresponding normalised eigenvector,

$$-4\Delta v + S_g v = \lambda_\alpha(g) v.$$

Hence

$$\lambda_\alpha(g, u) = \inf \left\{ \int_M (4|\nabla v|^2 + S_g v^2) d\mu_g : \int_M v^2 d\mu_g = 1 \right\}.$$

Similar to the case of Hamilton’s Ricci flow, all geometric quantities associated with the source manifold evolve along $(RH)_\alpha$. For instance, we consider those quantities that are directly relevant at the present; the metric inverse, volume element, Laplace–Beltrami operator and S_g evolve as follows:

$$\begin{aligned} \frac{\partial}{\partial t} g^{ij} &= 2S^{ij}, & \frac{\partial}{\partial t} \Delta_g &= 2S^{ij} \nabla_i \nabla_j - 2\alpha \tau_g u \langle \nabla u, \nabla \cdot \rangle, \\ \frac{\partial}{\partial t} d\mu_g &= -S_g d\mu_g, & \frac{\partial}{\partial t} S_g &= \Delta S_g + 2|S_{ij}|^2 + 2\alpha |\tau_g u|_g^2, \end{aligned}$$

where $S_{ij} = Rc - \alpha \nabla u \otimes \nabla u$ and $g^{ij} S_{ij} = S_g$. The nonnegativity of the curvature operator and S_g are preserved during the flow; for example, the evolution of $S_g = R_g - \alpha |\nabla u|^2$ is governed by the differential inequality

$$\frac{\partial}{\partial t} S_g \geq \Delta S_g + \frac{2}{m} S_g^2,$$

since $\alpha \geq 0$ and $|S_{ij}|^2 \geq (1/m) S_g^2$. Suppose $S_{g_0} \geq \rho$. We can use the maximum principle by comparing the solution of the above inequality with that of the ODE

$$(2-4) \quad \begin{cases} \frac{d\psi(t)}{dt} = \frac{2}{m} (\psi(t))^2, \\ \psi(0) = \rho, \end{cases}$$

solving to

$$\psi(t) = \frac{1}{\frac{1}{\rho} - \frac{2}{m} t}.$$

Therefore,

$$(2-5) \quad S_{g(t)} \geq \psi(t) = \frac{1}{\frac{1}{\rho} - \frac{2}{m} t}$$

for all $t \geq 0$ as long as the flow exists. We remark that (2-5) implies

$$S_{g(t)\min} \geq \frac{S_{g(0)\min}}{1 - (2t/m) S_{g(0)\min}}.$$

Clearly, $S_{g(0)\min} > 0$ implies $S_{g(t)\min} \rightarrow \infty$ in finite time $T_\epsilon \leq m/(2S_{g(0)\min}) < \infty$. This also implies that $R_{g(t)\min} \rightarrow \infty$ as $t \rightarrow T_\epsilon$, and thus $g(t)$ becomes singular in finite time $T_{\text{singular}} \leq T_\epsilon < \infty$.

Recall the Perelman–Müller \mathcal{W}_α -entropy functional also introduced in [Müller 2012] as

$$(2-6) \quad \mathcal{W}_\alpha(g, u, f, \tau) := \int_M (\tau(S_g + |\nabla f|_g^2) + f - m) \frac{e^{-f}}{(4\pi\tau)^{m/2}} d\mu_g,$$

where $\tau \in \mathbb{R}$ is a real number, $f \in C^\infty(M \times [0, T])$, $\alpha > 0$ is a constant and $u \in C^\infty(M, N)$ is a harmonic map between the m -dimensional manifold M and the n -dimensional manifold N , which by the Nash embedding theorem is isometrically embedded in \mathbb{R}^d for sufficiently large d . The above entropy functional is analogous to Perelman’s \mathcal{W} -entropy for shrinkers [2002] under the Ricci flow. \mathcal{W}_α is equally used for shrinkers under Ricci-harmonic map flow as can be traced back to List [2008]. As pointed out in [Perelman 2002], such an entropy is invariant and monotone. In fact, given a constant $\lambda > 0$ and a diffeomorphism ϕ of M , under simultaneous scaling of g and τ , we have

$$\mathcal{W}_\alpha(\lambda g, u, f, \lambda\tau) = \mathcal{W}_\alpha(g, u, f, \tau),$$

and under the pullback of g , u and f , we have

$$\mathcal{W}_\alpha(\phi^*g, \phi^*u, \phi^*f, \tau) = \mathcal{W}_\alpha(g, u, f, \tau).$$

More importantly, we have the following monotonicity formula.

Proposition 2.1 [List 2008; Müller 2012]. *Let $(g(t), u(t), f(t), \tau(t))$, $t \in [0, T)$ be a solution of the system*

$$(2-7) \quad \begin{cases} \frac{\partial}{\partial t} g = -2 \operatorname{Rc} + 2\alpha \nabla u \otimes \nabla u, \\ \frac{\partial}{\partial t} u = \tau_g u, \\ \left(-\frac{\partial}{\partial t} - \Delta_g + R - \alpha |\nabla u|_g^2 \right) \frac{e^{-f}}{(4\pi\tau)^{m/2}} = 0, \\ \frac{\partial}{\partial t} \tau = -1. \end{cases}$$

Then the \mathcal{W}_α -entropy is nondecreasing with

$$(2-8) \quad \begin{aligned} \frac{d}{dt} \mathcal{W}_\alpha(g, u, f, \tau) &= 2\tau \int_M \left| \operatorname{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f - \frac{1}{2\tau} g \right|^2 \frac{e^{-f}}{(4\pi\tau)^{m/2}} d\mu_g \\ &\quad + 2\tau \int_M \alpha |\tau_g u - \langle \nabla u, \nabla f \rangle|^2 \frac{e^{-f}}{(4\pi\tau)^{m/2}} d\mu_g. \end{aligned}$$

Notice that the third equation in the above system is equivalent to the following backward heat equation

$$(2-9) \quad \frac{\partial f}{\partial t} = -\Delta_g f + |\nabla f|_g^2 - R + \alpha |\nabla u|_g^2 + \frac{m}{2\tau},$$

A monotonicity formula of the above type is used to rule out any periodic solution to the $(\text{RH})_\alpha$ -flow other than those that are striking and Einstein [List 2008; Müller 2012; Perelman 2002].

Similar to $\lambda_\alpha(g, u)$ above, define the minimizing problem

$$\mu_\alpha(g, u, \tau) := \inf \left\{ \mathcal{W}_\alpha(g, u, f, \tau) : f \in C^\infty(M), \int_M (4\pi\tau)^{-m/2} e^{-f} d\mu_g = 1 \right\},$$

replacing f by $v = e^{-f/2}$. We have an equivalent minimizing integral

$$\mathcal{W}_\alpha(g, u, v, \tau) = \int_M (\tau(4|\nabla v|^2 + S_g v^2) - v^2 \ln v^2 - m v^2)(4\pi\tau)^{-m/2} d\mu_g$$

for functions $v \in H^1(M)$ with $\int_M v^2(4\pi\tau)^{-m/2} d\mu_g = 1$. Then v satisfies the Euler–Lagrange equation, and it follows that $\mu_\alpha(g, u, \tau)$ is achieved by a minimizer f_τ satisfying

$$\tau(2\Delta f_\tau - |\nabla f_\tau|^2 + S_g) + f_\tau - n = \mu(g, \tau).$$

By the result of Perelman, it is well understood that for any metric g on a compact manifold M and $\tau > 0$, we have $\mu(g, u, \tau) > -\infty$ and it approaches zero as $\tau \rightarrow 0$.

3. A new entropy monotonicity formula

In this section we introduce a new family of dual entropy formulas, which are dual in the sense that they generalise Ni’s entropy formula [2004] for the forward heat equation on the one hand and generalise Perelman and Müller’s \mathcal{W}_α -entropy on the other hand. A similar family of entropy functionals was constructed by Kuang and Zhang [2008]. The monotonicity property discussed here is very crucial to the derivation of our results in the rest of this paper.

Definition 3.1. Let $f : M \times [0, T] \rightarrow \mathbb{R}$ be smoothly defined with normalisation condition

$$\int_M \frac{e^{-f}}{(4\pi\tau)^{m/2}} d\mu_g = 1.$$

We define a generalised family of entropy by

(3-1)

$$\mathcal{W}_{\alpha,\epsilon}(g, u, f, \tau) = \int_M \left(\frac{\epsilon^2\tau}{4\pi}(S_g + |\nabla f|_g^2) + f - \frac{m\epsilon^2}{4\pi} + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2} \right) \frac{e^{-f}}{(4\pi\tau)^{m/2}} d\mu_g,$$

where $\tau(t) = T - t > 0$, $0 < \epsilon^2 \leq 4\pi$ and $S_g = S_g(x, t) = (R_g - \alpha|\nabla u|_g^2)(x, t)$.

Let $H = H(x, t)$ be a positive solution to the conjugate heat equation on a complete compact manifold with metric $g = g(x, t)$, evolving by the $(\text{RH})_\alpha$. Let

$H = (4\pi\tau)^{-m/2}e^{-f}$ and $\int_M H d\mu_g = 1$. Then

$$(3-2) \quad \left(-\frac{\partial}{\partial t} - \Delta_g + S_g\right)H = 0.$$

Theorem 3.2. *Suppose that $(g(t), u(t)), t \in [0, T)$, solves $(\text{RH})_\alpha$ with $\alpha(t) \equiv \alpha > 0$ and τ is a backward time with $\partial\tau/\partial t = -1$. Suppose that $H : M \times [0, T) \rightarrow (0, \infty)$ solves the conjugate heat equation $(-\partial/\partial t - \Delta_g + S_g)H = 0$. The entropy functional $\mathcal{W}_{\alpha,\epsilon}$ is nondecreasing by the formula*

$$(3-3) \quad \begin{aligned} &\frac{d}{dt}\mathcal{W}_{\alpha,\epsilon}(g, u, f, \tau) \\ &\geq \frac{\epsilon^2\tau}{2\pi} \int_M \left(\left| \text{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f - \frac{1}{2\tau} g \right|^2 + \alpha |\tau_g u - \langle \nabla u, \nabla f \rangle|^2 \right) H d\mu_g \end{aligned}$$

for $0 < \epsilon^2 \leq 4\pi$.

Remark 3.3. We remark that if $\epsilon^2 = 4\pi$, we recover Perelman and Müller’s \mathcal{W}_α -entropy.

Scaling and diffeomorphism invariance of $\mathcal{W}_{\alpha,\epsilon}$. Before we prove the monotonicity formula (3-3), we shall first establish the invariance of our new entropy with respect to dilation and diffeomorphism.

Lemma 3.4 [Chow and Knopf 2004, Lemma 6.57]. *If g and h are two Riemannian metrics on an n -dimensional Riemannian manifold and they are related by the time-scale factor λ (i.e., $g = \lambda h$), then the various geometric quantities scale as follows:*

$$\begin{aligned} g^{ij} &= \frac{1}{\lambda} h^{ij}, & \Gamma_{ij(g)}^k &= \Gamma_{ij(h)}^k, \\ R^l_{ijk}(g) &= R^l_{ijk}(h), & R_{ijkl}(g) &= \phi R_{ijkl}(h), \\ R_{ij}(g) &= R_{ij}(h), & R(g) &= \frac{1}{\lambda} R(h), & d\mu_{(g)} &= \lambda^{n/2} d\mu_{(h)}. \end{aligned}$$

Lemma 3.5. *Let $\lambda > 0$ be any constant and $\phi : M \rightarrow M$ be a one-parameter family of diffeomorphisms. Then*

$$\begin{aligned} \mathcal{W}_{\alpha,\epsilon}(\lambda g, u, f, \lambda\tau) &= \mathcal{W}_{\alpha,\epsilon}(g, u, f, \tau), \\ \mathcal{W}_{\alpha,\epsilon}(\phi^*g, \phi^*u, \phi^*f, \tau) &= \mathcal{W}_{\alpha,\epsilon}(g, u, f, \tau). \end{aligned}$$

Proof. By a straightforward computation, we have

$$\begin{aligned} \mathcal{W}_{\alpha,\epsilon}(\lambda g, u, f, \lambda\tau) &= \int_M \left(\frac{\epsilon^2\lambda\tau}{4\pi} (R(\lambda g) - \alpha(\lambda g)^{ij} \nabla_i u \otimes \nabla_j u + (\lambda g)^{ij} \nabla_i f \nabla_j f) \right. \\ &\quad \left. + f - \frac{m\epsilon^2}{4\pi} + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2} \right) \frac{e^{-f}}{(4\pi\lambda\tau)^{m/2}} \sqrt{\det(\lambda g)} dx \end{aligned}$$

$$\begin{aligned}
 &= \int_M \left(\frac{\epsilon^2 \lambda \tau}{4\pi} (\lambda^{-1} R(g) - \alpha \lambda^{-1} g^{ij} \nabla_i u \otimes \nabla_j u + \lambda^{-1} g^{ij} \nabla_i f \nabla_j f) \right. \\
 &\quad \left. + f - \frac{m\epsilon^2}{4\pi} + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2} \right) \frac{e^{-f}}{\lambda^{m/2} (4\pi \tau)^{m/2}} \sqrt{\lambda^m \det(g)} dx \\
 &= \int_M \left(\frac{\epsilon^2 \tau}{4\pi} (R_g - \alpha |\nabla u|_g^2 + |\nabla f|_g^2) \right. \\
 &\quad \left. + f - \frac{m\epsilon^2}{4\pi} + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2} \right) \frac{e^{-f}}{(4\pi \tau)^{m/2}} d\mu_g \\
 &= \mathcal{W}_{\alpha, \epsilon}(g, u, f, \tau).
 \end{aligned}$$

The invariance under diffeomorphisms is trivial since $(RH)_\alpha$ -flow is equivalent to the flow modified by the time-dependent diffeomorphism ϕ generated by the gradient of f , where ϕ^*g is the pulled-back metric and $\phi^*f = f \circ \phi$. For the harmonic map u , the invariance holds if we combine the following facts: ϕ is a C^∞ -diffeomorphism and $u \in C^\infty(M, N)$ is a harmonic map with respect to (M, g) ; then $\phi^*u = u \circ \phi \in C^\infty(M, N)$ is a harmonic map with respect to (M, ϕ^*g) with the identity

$$\int_M |\nabla u|_g^2 d\mu_g = \int_M |\nabla(u \circ \phi)|_{\phi^*g}^2 d\mu_{\phi^*g}.$$

Then, all the geometric quantities are invariant under $(RH)_\alpha$ -flow and the diffeomorphism invariance of $\mathcal{W}_{\alpha, \epsilon}$ follows. □

Proof of Theorem 3.2 (the monotonicity formula for $\mathcal{W}_{\alpha, \epsilon}$).

Proof. The entropy functional can be rewritten as

$$\begin{aligned}
 &\mathcal{W}_{\alpha, \epsilon}(g, u, f, \tau) \\
 &= \frac{\epsilon^2}{4\pi} \int_M (\tau(S_g + |\nabla f|^2) + f - m)H d\mu_g + \left(1 - \frac{\epsilon^2}{4\pi}\right) \int_M fH d\mu_H + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2}.
 \end{aligned}$$

By direct computation we obtain the evolution equation

$$(3-4) \quad \frac{d}{dt} \mathcal{W}_{\alpha, \epsilon}(g, u, f, \tau) = \frac{\epsilon^2}{4\pi} \frac{\partial}{\partial t} \left(\int_M V d\mu_g \right) + \left(1 - \frac{\epsilon^2}{4\pi}\right) \frac{\partial}{\partial t} \left(\int_M fH d\mu_g \right),$$

where

$$(3-5) \quad V := (\tau(2\Delta_g f + S_g - |\nabla f|^2) + f - m)H$$

since $\int_M (\Delta_g f - |\nabla f|_g^2) e^{-f} d\mu_g = 0$ on a closed manifold M . We make two claims here, which we shall prove in the next two propositions, namely,

$$(3-6) \quad \frac{\partial}{\partial t} \left(\int_M V d\mu_g \right) = \int_M -\square^* V d\mu_g = \frac{d}{dt} \mathcal{W}_{\alpha}(g, u, f, \tau)$$

and

$$(3-7) \quad \frac{\partial}{\partial t} \left(\int_M f H d\mu_g \right) = -\mathcal{F}_\alpha(g, u, f) + \frac{m}{2\tau} \geq 0.$$

With the above two claims, we arrive at

$$(3-8) \quad \frac{d}{dt} \mathcal{W}_{\alpha, \epsilon}(g, u, f, \tau) = \frac{\epsilon^2}{4\pi} \frac{d}{dt} \mathcal{W}_\alpha(g, u, f, \tau) + \frac{m}{2\tau} - \mathcal{F}_\alpha(g, u, f),$$

which proves the monotonicity formula (3-3) for $0 < \epsilon^2 \leq 4\pi$. \square

Proposition 3.6. *With the assumptions of Theorem 3.2, the quantity*

$$V := (\tau(2\Delta_g f + S_g - |\nabla f|^2) + f - m)H$$

satisfies

$$(3-9) \quad \square^* V = -2\tau \left(\left| \text{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f - \frac{1}{2\tau} g \right|^2 + \alpha |\tau_g u - \langle \nabla u, \nabla f \rangle|^2 \right) H$$

and

$$(3-10) \quad \frac{d}{dt} \mathcal{W}_\alpha(g, u, f, \tau) = - \int_M \square^* V d\mu_g.$$

Moreover if H tends to a δ -function as $t \rightarrow T$, then $V \leq 0$ for all $t < T$ with $H(x, \tau)$ replaced with $H(x, \tau; y, \sigma)$, the fundamental solution.

Proof. Let $P = \tau(2\Delta f - |\nabla f|^2 + S_g) + f - n$, and $\partial_t \tau = -1$ since $\tau = T - t$. Thus,

$$\begin{aligned} \square^* V &= (-\partial_t - \Delta + S_g)(PH) \\ &= -(\partial_t + \Delta)PH - 2\langle \nabla P, \nabla H \rangle \end{aligned}$$

and

$$H^{-1} \square^* V = -(\partial_t + \Delta)P + 2\langle \nabla P, \nabla f \rangle$$

since $f = -\ln H - (m/2) \ln(4\pi\tau)$ implies that $\nabla f = -H^{-1} \nabla H$. Let us compute $(\partial_t + \Delta)P$ as follows:

$$(3-11) \quad \frac{\partial P}{\partial t} = -(2\Delta f - |\nabla f|^2 + S_g) + \tau \frac{\partial}{\partial t} (2\Delta f - |\nabla f|^2) + \tau \frac{\partial}{\partial t} S_g + \frac{\partial}{\partial t} f.$$

Note that

$$\begin{aligned} \partial_t f &= -\Delta_g f - S_g + |\nabla f|_g^2 + \frac{m}{2\tau}, \\ \partial_t S_g &= \Delta S_g + 2|S_{ij}|_g^2 + 2\alpha |\tau_g u|_g^2. \end{aligned}$$

Then a straightforward computations yields

$$(3-12) \quad 2 \frac{\partial}{\partial t} (\Delta f) = 4S_{ij} \nabla_i \nabla_j f - 4\alpha \tau_g u \langle \nabla u, \nabla f \rangle + 2\Delta(-\Delta f + |\nabla f|^2 - S_g),$$

$$(3-13) \quad \frac{\partial}{\partial t} |\nabla f|^2 = \frac{\partial}{\partial t} (g^{ij} \nabla_i f \nabla_j f) = 2S_{ij} \nabla_i f \nabla_j f + 2 \left\langle \nabla f, \nabla \frac{\partial}{\partial t} f \right\rangle.$$

Combining (3-11)–(3-13) with the identity $\Delta P = \tau(2\Delta(\Delta f) - \Delta|\nabla f|^2 + \Delta S_g) + \Delta f$, we have

$$\begin{aligned} \left(\frac{\partial}{\partial t} + \Delta\right)P &= -2\Delta f - 2|\nabla f|^2 - 2S_g + \frac{m}{2\tau} \\ &\quad + \tau(4S_{ij}\nabla_i\nabla_j f - 4\alpha\tau_g u \langle \nabla u, \nabla f \rangle + \Delta|\nabla f|^2 + 2\langle \nabla f, \nabla \Delta f \rangle \\ &\quad - 2S_{ij}\nabla_i f \nabla_j f + 2|S_{ij}|^2 + 2\alpha|\tau_g u|^2 - 2\langle \nabla f, \nabla|\nabla f|^2 \rangle + 2\langle \nabla f, \nabla S_g \rangle). \end{aligned}$$

Similarly,

$$\begin{aligned} 2\langle \nabla P, \nabla f \rangle &= 2\langle \nabla(\tau(2\Delta f - |\nabla f|^2 + S - g) + f), \nabla f \rangle \\ &= 2\tau(2\langle \nabla \Delta f, \nabla f \rangle - \langle \nabla|\nabla f|^2, \nabla f \rangle + \langle \nabla S_g, \nabla f \rangle) + 2|\nabla f|^2. \end{aligned}$$

Therefore,

$$\begin{aligned} -\left(\frac{\partial}{\partial t} + \Delta\right)P + 2\langle \nabla P, \nabla f \rangle &= \left(2\Delta f + 2S_g - \frac{m}{2\tau}\right) - \tau(4S_{ij}\nabla_i\nabla_j f - 4\alpha\tau_g u \langle \nabla u, \nabla f \rangle + \Delta|\nabla f|^2 \\ &\quad - 2\langle \nabla f, \nabla \Delta f \rangle - 2S_{ij}\nabla_i f \nabla_j f + 2|S_{ij}|^2 + 2\alpha|\tau_g u|^2) \\ &= \left(2\Delta f + 2S_g - \frac{m}{2\tau}\right) - \tau(4S_{ij}\nabla_i\nabla_j f + 2|\nabla \nabla f|^2 + 2|S_{ij}|^2) \\ &\quad - 2\tau\alpha(|\tau_g u + \langle \nabla u, \nabla f \rangle|^2 - 2\tau_g u \langle \nabla u, \nabla f \rangle) \\ &= -2\tau\left(2S_{ij}\nabla_i\nabla_j f + |\nabla \nabla f|^2 + |S_{ij}|^2 - \frac{1}{\tau}\left(\Delta f + R - \frac{m}{4\tau}\right)\right) \\ &\quad - 2\tau\alpha(|\tau_g u - \langle \nabla u, \nabla f \rangle|^2) \\ &= -2\tau\left((S_{ij} + \nabla_i\nabla_j f)^2 - \frac{1}{\tau}\left(\Delta f + R - \frac{m}{4\tau}\right)\right) - 2\tau\alpha(|\tau_g u - \langle \nabla u, \nabla f \rangle|^2) \\ &= -2\tau\left|S_{ij} + \nabla_i\nabla_j f - \frac{1}{2\tau}g_{ij}\right|^2 - 2\tau\alpha(|\tau_g u - \langle \nabla u, \nabla f \rangle|^2), \end{aligned}$$

where we have used the following calculation by Bochner’s identity:

$$\begin{aligned} \Delta|\nabla f|^2 - 2\langle \nabla f, \nabla \Delta f \rangle - 2S_{ij}\nabla_i f \nabla_j f &= 2|\nabla \nabla f|^2 + 2(R_{ij} - S_{ij})\nabla_i f \nabla_j f \\ &= 2|\nabla \nabla f|^2 + 2\alpha\langle \nabla u, \nabla f \rangle^2. \end{aligned}$$

Hence,

$$H^{-1}\square^*V = -2\tau\left|S_{ij} + \nabla_i\nabla_j f - \frac{1}{2\tau}g_{ij}\right|^2 - 2\tau\alpha(|\tau_g u - \langle \nabla u, \nabla f \rangle|^2)$$

and

$$\square^*V = -2\tau\left|S_{ij} + \nabla_i\nabla_j f - \frac{1}{2\tau}g_{ij}\right|^2 H - 2\tau\alpha(|\tau_g u - \langle \nabla u, \nabla f \rangle|^2)H.$$

The consequence of which is a localised version of Perelman’s \mathcal{W} -entropy monotonicity formula. Thus,

$$\begin{aligned}
 \frac{d\mathcal{W}}{dt} &= \frac{\partial}{\partial t} \int_M V d\mu_g = \int_M (\partial_t V - RV + \alpha |\nabla u|_g^2 V) d\mu \\
 &= \int_M (-\square^* V - \Delta_g V) d\mu_g = \int_M -\square^* V d\mu_g \\
 &= 2(T-t) \int_M \left(|S_{ij} + \nabla_i \nabla_j f - \frac{1}{2(T-t)} g_{ij}|^2 \right. \\
 &\quad \left. + \alpha (|\tau_g u - \langle \nabla u, \nabla f \rangle|^2) \right) \frac{e^{-f}}{(4\pi\tau)^{-m/2}} d\mu_g. \quad \square
 \end{aligned}$$

Proposition 3.7. *With the assumptions of Theorem 3.2, we have*

$$(3-14) \quad \frac{\partial}{\partial t} \left(\int_M f H d\mu_g \right) \geq 0.$$

Proof. By direct computation,

$$\begin{aligned}
 \frac{\partial}{\partial t} \left(\int_M f H d\mu \right) &= \int_M \left(\frac{\partial}{\partial t} f H + f \frac{\partial}{\partial t} H - S_g f H \right) d\mu_g \\
 &= \int_M \left(-\Delta_g f + |\nabla f|_g^2 - S_g + \frac{m}{2\tau} \right) H d\mu_g \\
 &\quad + \int_M f (-\Delta_g H + S_g H) d\mu - \int_M S_g f H d\mu_g \\
 &= \int_M (-2\Delta_g f + |\nabla f|_g^2) H d\mu + \int_M \left(\frac{m}{2\tau} - S_g \right) H d\mu_g,
 \end{aligned}$$

where we used integration by parts on $-\int_M \Delta_g f H = -\int_M f \Delta_g H$. Rearranging the above, we have

$$\begin{aligned}
 \frac{\partial}{\partial t} \left(\int_M f H d\mu_g \right) &= \int_M (-S_g - 2\Delta_g f + |\nabla f|_g^2) H d\mu_g + \frac{m}{2\tau} \int_M H d\mu_g \\
 &= - \int_M (S_g + |\nabla f|_g^2) H d\mu_g + \frac{m}{2\tau} \\
 &= -\mathcal{F}_\alpha + \frac{m}{2\tau},
 \end{aligned}$$

where $\mathcal{F}_\alpha = \int_M (S_g + |\nabla f|_g^2) H d\mu_g$ is the Perelman energy functional introduced in [Müller 2012], which we discussed in Section 2. Next is to show that

$$(3-15) \quad \frac{\partial}{\partial t} \left(\int_M f u d\mu \right) = -\mathcal{F}_\alpha + \frac{m}{2\tau} \geq 0.$$

Recall the evolution of \mathcal{F}_α :

$$(3-16) \quad \frac{d}{dt} \mathcal{F}_\alpha(g, f) = 2 \int_M (|\text{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f|^2 + \alpha |\tau_g u - \langle \nabla u, \nabla f \rangle|^2) H d\mu_g.$$

Straightforward analysis, using an elementary inequality and the Cauchy–Schwarz inequality, gives

$$(3-17) \quad |\text{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f|^2 \geq \frac{1}{m} (R_g - \alpha |\nabla u|_g^2 + \Delta_g f)^2$$

so that

$$\int (S_g + \Delta_g f) H d\mu_g \leq \left(\int (S_g + \Delta_g f)^2 H d\mu_g \right)^{1/2} \left(\int H d\mu_g \right)^{1/2},$$

which implies

$$\left(\int_M (S_g + \Delta_g f) H d\mu \right)^2 \leq \int_M (S_g + \Delta_g f)^2 H d\mu_g.$$

Hence by (3-16) and (3-17), we obtain

$$(3-18) \quad \frac{d}{dt} \mathcal{F}_\alpha \geq \frac{2}{m} \int_M (S_g + \Delta_g f)^2 H d\mu_g + \int_M 2\alpha |\tau_g u - \langle \nabla u, \nabla f \rangle|^2 H d\mu_g.$$

We can then solve

$$\frac{d}{dt} \mathcal{F}_\alpha \geq \frac{2}{m} \mathcal{F}_\alpha^2, \quad \mathcal{F}_\alpha \geq 0.$$

This implies

$$\begin{aligned} \frac{d\mathcal{F}_\alpha}{\mathcal{F}_\alpha^2} \geq \frac{2}{m} dt &\implies -\frac{1}{\mathcal{F}_\alpha} \Big|_t^T \geq \frac{2}{m} (T-t) \implies \frac{1}{\mathcal{F}_\alpha(t)} - \frac{1}{\mathcal{F}_\alpha(T)} \geq \frac{2}{m} \tau \\ &\implies \frac{1}{\mathcal{F}_\alpha(t)} \geq \frac{2}{m} \tau + \frac{1}{\mathcal{F}_\alpha(T)}. \end{aligned}$$

From here we can conclude as follows:

(i) Suppose $\mathcal{F}_\alpha(T) > 0$. Then

$$\frac{1}{\mathcal{F}(t)} \geq \frac{2\tau}{m}; \quad \text{i.e., } \mathcal{F}_\alpha(t) \leq \frac{m}{2\tau}.$$

(ii) Suppose $\mathcal{F}_\alpha(T) \leq 0$. Then $\mathcal{F}_\alpha(t) \leq 0 \leq m/(2\tau)$ for all $t \in [0, T)$, since we know that $d\mathcal{F}_\alpha/dt \geq 0$.

Hence

$$\mathcal{F}_\alpha(t) \leq \frac{m}{2\tau} \quad \text{for } t \in [0, T),$$

which proves the claim (3-15). □

4. Differential Harnack estimates

In this section we obtain Perelman's differential Harnack-type estimate which holds for the fundamental solution and, of course, all positive solutions to the conjugate heat equation coupled to the Ricci-harmonic map flow. There is an improvement over some known results as there is no explicit restriction on the curvature and no

recourse to Perelman's reduced distance. In what follows, we want to show that the local entropy satisfies a pointwise differential inequality for the positive minimal solution. Define a differential Harnack quantity

$$P_\epsilon := \frac{\epsilon^2 \tau}{4\pi} (2\Delta f - |\nabla f|^2 + S_g) + f + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2} - \frac{m\epsilon^2}{4\pi}.$$

Theorem 4.1. *Let M be a closed manifold with bounded Ricci curvature and $H(x, y, t) = H = (4\pi t)^{-n/2} e^{-f}$ satisfy $\square^* H = 0$, where H tends to a δ -function as $t \rightarrow T$ and satisfies $\int_M H d\mu_g = 1$. Then for all $t < T$ and $\epsilon^2 \rightarrow 4\pi$ as $t \rightarrow T$, we have*

$$(4-1) \quad \frac{\epsilon^2 \tau}{4\pi} (2\Delta f - |\nabla f|^2 + S_g) + f - \frac{m\epsilon^2}{4\pi} \leq 0.$$

Proof. Let h be any compactly supported smooth function for all $t_0 > 0$. Suppose $h(\cdot, t)$ is a positive solution to the ordinary heat equation $(\partial_t - \Delta)h = 0$ (this is Perelman's argument in [2002, Corollary 9.3]). Then, it is clear that

$$\frac{\partial}{\partial t} \int_M H h dV = 0$$

and we have by direct calculation that

$$\begin{aligned} \frac{\partial}{\partial t} \int_M h P_\epsilon H d\mu_g &= \int_M (\partial_t h (P_\epsilon H) + h \partial_t (P_\epsilon H) - S_g P_\epsilon H) d\mu_g \\ &= \int_M ((\partial_t - \Delta)h (P_\epsilon H) + h(\partial_t + \Delta - S_g) P_\epsilon H) d\mu_g \\ &= - \int_M h \square^* (P_\epsilon H) d\mu_g \\ &= - \frac{\epsilon^2}{4\pi} \int_M h \square^* V_\epsilon d\mu_g \geq 0. \end{aligned}$$

The inequality is due to [Proposition 3.6](#). We are left to show that for the everywhere positive function $h(\cdot, t)$, the limit of $\int_M h V_\epsilon d\mu_g$ is nonpositive as $t \rightarrow T$. We assume the claim a priori, i.e, $\lim_{t \rightarrow T} \int_M h V_\epsilon d\mu_g = 0$, with

$$V_\epsilon = \left(\tau(2\Delta f - |\nabla f|^2 + S_g) + \frac{4\pi}{\epsilon^2} f - m \right) H,$$

and conclude the result. □

For completeness, we devote the next effort to justifying the claim

$$(4-2) \quad \lim_{t \rightarrow T} \int_M h V_\epsilon d\mu_g \leq 0 \iff \lim_{t \rightarrow T} \int_M h P_\epsilon H d\mu_g \leq \frac{m}{2} \lim_{t \rightarrow T} \left(\ln \frac{4\pi}{\epsilon^2} \int_M h H d\mu_g \right).$$

Our argument follows from [\[Ni 2006; Perelman 2002\]](#) and can be compared with the recent preprint [\[Băileşteanu and Tran 2013, Proposition 4.2\]](#) (see also [\[Chow](#)

et al. 2008, Section 16.4]), where we know that $\lim_{t \rightarrow T} \int_M Vh \, d\mu_g \leq 0$ (where V is as defined in Proposition 3.6). To see this clearly, we write

$$P_\epsilon H = \frac{\epsilon^2}{4\pi} V_\epsilon + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2} H,$$

which implies

$$(4-3) \quad \lim_{t \rightarrow T} \int_M h P_\epsilon H \, d\mu_g = \frac{\epsilon^2}{4\pi} \lim_{t \rightarrow T} \int_M V_\epsilon h \, d\mu_g + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2} \lim_{t \rightarrow T} \int_M h H \, d\mu_g.$$

If H tends to a Dirac δ -function, say at a point $p \in M$, for $\tau \rightarrow T$, then f satisfies $f(x, \tau) \rightarrow d^2(p, x)/4\tau$. This is in relation to the l -length of Perelman. This yields

$$(4-4) \quad \lim_{\tau \rightarrow 0} \int_M f h H \, d\mu_g \leq \limsup_{\tau \rightarrow 0} \int_M \frac{d^2(p, x)}{4\tau} h H \, d\mu_g = \frac{m}{2} h(p, T).$$

Meanwhile, by the strong maximum principle, we have $h(x, T) > 0$ and

$$\lim_{\tau \rightarrow 0} \int_M h H \, d\mu_g = h(x, T).$$

Hence by a scaling argument, we assume that $h(x, T) = 1$. Rewriting P_ϵ and using integration by parts, we have

$$\begin{aligned} \int_M P_\epsilon h H \, d\mu_g &= \int_M \frac{\epsilon^2 \tau}{4\pi} \left(|\nabla f|^2 + S_g \, d\mu_g - \frac{m}{2\tau} \right) h H \, d\mu_g - \frac{\epsilon^2 \tau}{2\pi} \int_M \langle \nabla f, \nabla h \rangle H \, dV \\ &\quad + \int_M f H h \, d\mu_g + \frac{m}{2} \left(\ln \frac{4\pi}{\epsilon^2} - \frac{\epsilon^2}{4\pi} \right) \int_M H h \, d\mu_g. \end{aligned}$$

We should also note that since $h(\cdot, t_0)$ is compactly supported and by the strong maximum principle, we have that $h(\cdot, t_0)$, $|\nabla h(\cdot, t_0)|$ and $|\Delta h(\cdot, t_0)|$ are bounded on M . This implies that there exists a bounded solution $h(\cdot, t_0)$. Now we claim that the first three terms on the right-hand side of the last equation vanish as $\tau \rightarrow 0$. We can see this, for instance, in the following argument: By integration by parts and the fact that $\nabla H = -H \nabla f$, we have

$$(4-5) \quad -\tau \int_M \langle \nabla f, \nabla h \rangle H \, d\mu_g = \tau \int_M \langle \nabla H, \nabla h \rangle \, d\mu_g = -\tau \int_M H \Delta h \, d\mu_g$$

is bounded since $|\Delta h|$ is bounded as stated earlier. Thus, the second term in right-hand side of the preceding equation is bounded and goes to zero as $\tau \rightarrow 0$, so the same is true for first terms (which follows from gradient estimates [Chow et al. 2008, Lemma 16.47]). Thus the analysis is reduced to showing that

$$(4-6) \quad \lim_{\tau \rightarrow 0} \int V_\epsilon h \, d\mu_g < C(m) \leq 0.$$

By the monotonicity formula for $W_{\alpha,\epsilon}$, we have

$$\frac{\partial}{\partial t} \int_M P_\epsilon h H d\mu_g = \frac{\epsilon^2 \tau}{4\pi} \frac{\partial}{\partial t} \int_M V_\epsilon h d\mu_g \geq 0.$$

By the mean value theorem, there exists a sequence $\tau_k \rightarrow 0$ such that

$$\lim_{\tau_k \rightarrow 0} \tau_k \int_M \left(\left| \text{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f - \frac{n}{2\tau_k} g \right|^2 + \alpha |\tau_k u - \langle \nabla u, \nabla f \rangle|^2 \right) H h d\mu_g = 0.$$

Applying the Cauchy–Schwarz and Hölder inequalities, we have

$$\left| \text{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f - \frac{1}{2\tau_k} g \right|^2 \geq \frac{1}{m} \left(R_g - \alpha |\nabla u|_g^2 + \Delta_g f - \frac{n}{2\tau_k} g \right)^2$$

and

$$\begin{aligned} & \int_M \tau_k \left(S_g + \Delta f - \frac{n}{2\tau_k} \right) H h d\mu_g \\ & \leq \left(\tau_k^2 \int_M \left(S_g + \Delta f - \frac{n}{2\tau_k} \right)^2 H h d\mu_g \right)^{1/2} \left(\int_M H h d\mu_g \right)^{1/2} \\ & \leq \sqrt{m} \left(\tau_k^2 \int_M \left| \text{Rc} - \alpha \nabla u \otimes \nabla u + \nabla \nabla f - \frac{1}{2\tau_k} g \right|^2 H h d\mu_{g(\tau_k)} \right)^{1/2} \left(\int_M H h d\mu_g \right)^{1/2} \rightarrow 0 \end{aligned}$$

as $\tau_k \rightarrow 0$, since $\alpha |\tau_k u - \langle \nabla u, \nabla f \rangle|^2 \geq 0$ and $\lim_{\tau_k \rightarrow 0} \int_M H h d\mu_{g(\tau_k)}$ is finite.

Then we have

$$\begin{aligned} & \lim_{t \rightarrow T} \int V_\epsilon h d\mu_g \\ & = \lim_{t \rightarrow T} \int_M \left(\frac{\epsilon^2}{4\pi} \tau_k (2\Delta f - |\nabla f|^2 + S_g) + \frac{4\pi}{\epsilon^2 \tau} f - m \right) H h d\mu_g \\ & = \lim_{t \rightarrow T} \int_M \left(\frac{\epsilon^2 \tau_k}{4\pi} \left(\Delta f + S_g - \frac{n}{2\tau_k} \right) \right) H h d\mu_{g(\tau_k)} \\ & \quad + \lim_{t \rightarrow T} \int_M \left(\frac{\epsilon^2 \tau_k}{4\pi} (\Delta f - |\nabla f|^2) \right) H h d\mu_g + \lim_{t \rightarrow T} \int_M \left(f - \frac{m\epsilon^2}{8\pi} \right) H h d\mu_g \\ & = \lim_{t \rightarrow T} \int_M \left(f - \frac{m\epsilon^2}{8\pi} \right) H h d\mu_g, \end{aligned}$$

where we have used the identity

$$\int_M (\Delta f - |\nabla f|^2) H d\mu = - \int_M \Delta H d\mu = 0$$

for any positive solution H and the fact that each quantity in (4-5) is bounded to obtain $\lim_{t \rightarrow T} \tau_k \int_M (\Delta f - |\nabla f|^2) H h d\mu_g = 0$.

By (4-4) and the asymptotic behaviour of the heat kernel, i.e, $f \approx d^2/(4\tau)$ as $\tau \rightarrow 0$, we have (see [Ni 2006, Theorem 2.1])

$$H(x, y, \tau) \sim (4\pi\tau)^{-m/2} \exp\left(\frac{d^2(x, y)}{4\tau}\right) \sum_{j=0}^{\infty} u_j(x, y, \tau)\tau^j := w_k(x, y, \tau)$$

as $\tau \rightarrow 0$, where $d^2(x, y)$ is the distance function and $w_k(x, y, t)$ satisfies

$$w_k(x, y, \tau) = O\left(\tau^{k+1-m/2} \exp\left(\frac{\delta d^2(x, y)}{4\tau}\right)\right)$$

uniformly for all $x, y \in M$ and δ is just a number depending only on the geometry of (M, g) . The function can be chosen such that $u_0(x, y, 0) = 1$. Though, the above asymptotic result does not require any curvature assumption, a result due to Cheeger and Yau [1981] states that on a manifold with bounded Ricci curvature (which is our case), the heat kernel satisfies

$$H(x, y, \tau) \geq (4\pi\tau)^{-m/2} \exp\left(\frac{d^2(x, y)}{4\tau}\right),$$

which implies

$$f(x, \tau) \leq \frac{d^2(x, y)}{4\tau}.$$

Therefore,

$$\begin{aligned} \lim_{\tau \rightarrow 0} \int_M \left(f - \frac{m\epsilon^2}{8\pi}\right) h H d\mu_g &\leq \limsup_{\tau \rightarrow 0} \int_M \left(\frac{d^2(x, y)}{4\tau} - \frac{m\epsilon^2}{8\pi}\right) h(y, t) H(x, y, \tau) d\mu_g \\ &= \limsup_{\tau \rightarrow 0} \int_M \left(\frac{d^2(x, y)}{4\tau} - \frac{m\epsilon^2}{8\pi}\right) \frac{e^{-d^2(x, y)/4\tau}}{(4\pi\tau)^{m/2}} h(y, t) d\mu_g. \end{aligned}$$

It is easy to see that for any $\delta > 0$, the integration of the above integrand in the domain $d(x, y) \geq \delta$ converges to zero. Therefore,

$$\begin{aligned} (4-7) \quad \lim_{t \rightarrow 0} \int_M \left(f - \frac{m\epsilon^2}{8\pi}\right) h H d\mu_g \\ \leq \lim_{t \rightarrow 0} \int_{d(x, y) \leq \delta} \left(\frac{d^2(x, y)}{4t} - \frac{m\epsilon^2}{8\pi}\right) \frac{e^{-d^2(x, y)/(4t)}}{(4\pi t)^{m/2}} h(y, t) d\mu_g. \end{aligned}$$

Whenever δ is chosen sufficiently small, $d(x, y)$ is asymptotically sufficiently close to the Euclidean distance. Then by a standard approximation using local coordinates, we have

$$\begin{aligned} (4-8) \quad \lim_{t \rightarrow 0} \int_M \left(\frac{d^2(x, y)}{4t} - \frac{m\epsilon^2}{8\pi}\right) h H d\mu_g \\ = \lim_{t \rightarrow 0} \int_{\mathbb{R}^m} \left(\frac{|x - y|^2}{4t} - \frac{m\epsilon^2}{8\pi}\right) \frac{e^{-|x - y|^2/(4t)}}{(4\pi t)^{m/2}} h_p(y) dy, \end{aligned}$$

where h_p is the pullback of $h(\cdot, 0)$ from the region $d(x, y) \leq \delta$ to the Euclidean space.

Splitting the last integrand as in [Kuang and Zhang 2008], we are left with

$$\begin{aligned} \lim_{t \rightarrow 0} \int_M \left(f - \frac{m\epsilon^2}{8\pi} \right) h H d\mu_g &\leq h_p(x) \lim_{t \rightarrow 0} \int_{\mathbb{R}^m} \left(\frac{|x-y|^2}{4\tau} - \frac{m\epsilon^2}{8\pi} \right) \frac{e^{-|x-y|^2/(4\tau)}}{(4\pi\tau)^{m/2}} dy \\ &= h_p(\cdot) \lim_{\tau \rightarrow 0} \int_{\mathbb{R}^m} \left(\frac{|y|^2}{4\tau} \frac{e^{-|y|^2/(4\tau)}}{(4\pi\tau)^{m/2}} \right) dy - \frac{m\epsilon^2}{8\pi} h_p(\cdot). \end{aligned}$$

Lastly, we have that the right-hand side evaluates to a constant $C(m) \leq 0$ by using the standard Gauss integral

$$\int_{\mathbb{R}^m} \left(\frac{|y|^2}{4\tau} \frac{e^{-|y|^2/(4\tau)}}{(4\pi\tau)^{n/2}} \right) dy = \frac{m}{2}$$

and the condition $\epsilon \rightarrow 2\sqrt{\pi}$ as $\tau \rightarrow 0$. The claim then follows.

Finally in this section we prove Perelman's differential Harnack estimates for f as an application of [Theorem 4.1](#) and the monotonicity of $\mathcal{W}_{\epsilon, \alpha}$. A corollary to this gives estimates of Li–Yau type for all positive solutions $H(x, \tau)$.

Proposition 4.2. *Let the assumptions of [Theorem 4.1](#) hold. Then for any smooth curve $\gamma(\tau)$ in M , we have the estimate*

$$(4-9) \quad -\frac{d}{dt} f(\gamma(\tau), \tau) \leq \frac{1}{2} \left(|\gamma'(\tau)|^2 + S_g(\gamma(\tau), \tau) - \frac{m}{2\tau} \right).$$

After the usual integration of (4-9) along the path $\gamma(\tau)$ and exponentiation, we have the following result.

Corollary 4.3. *With the notation and assumptions of [Theorem 4.1](#), the following Li–Yau Harnack estimate holds:*

$$(4-10) \quad \frac{H(x_2, t_2)}{H(x_1, t_1)} \leq \left(\frac{T - t_1}{T - t_2} \right)^{m/4} \exp \left(\frac{1}{2} \int_{t_1}^{t_2} (|\gamma'(t)|^2 + S_g(\gamma(t), t)) dt \right).$$

Proof of [Proposition 4.2](#). Precisely from (4-1), we have

$$f \leq \frac{m\epsilon^2}{4\pi} - \frac{\epsilon^2\tau}{4\pi} (2\Delta f - |\nabla f|^2 + S_g) \leq \frac{m\epsilon^2}{8\pi}$$

since $\Delta f + S_g - m/(2\tau) \geq 0$ by the monotonicity formula (3-3). Now multiplying (4-1) through by $2\pi/(\epsilon^2\tau)$, we have

$$\Delta f - \frac{1}{2} |\nabla f|^2 + \frac{1}{2} S_g + \frac{2\pi}{\epsilon^2 t} f - \frac{m}{2\tau} \leq 0.$$

Using $\Delta f = -\partial_t f + |\nabla f|^2 - S_g + m/(2\tau)$ from (2-9), we obtain

$$(4-11) \quad -\partial_t f + \frac{1}{2} |\nabla f|^2 \leq \frac{1}{2} S_g - \frac{2\pi}{\epsilon^2 t} f.$$

By Young’s inequality, we have

$$\begin{aligned}
 -\frac{d}{dt} f(\gamma(\tau), \tau) &= -\partial_t f(\gamma(\tau), \tau) - \langle \nabla f(\gamma(\tau), \tau), \gamma'(\tau) \rangle \\
 &\leq -\partial_t f + \frac{1}{2} |\nabla f|^2 + \frac{1}{2} |\gamma'(\tau)|^2 \\
 &= \frac{1}{2} |\gamma'(\tau)|^2 + \frac{1}{2} S_g(\gamma(\tau), \tau) - \frac{2\pi}{\epsilon^2 \tau} f(\gamma(\tau)\tau)
 \end{aligned}$$

on the path $\gamma(\tau)$. The result follows by using the fact that $f \leq m\epsilon^2/(8\pi)$. □

5. Log-Sobolev inequalities along $(RH)_\alpha$ -flow

By the results of Aubin [1976] and Hebey [1996] for complete manifolds whose Ricci curvature is bounded from below and injectivity radius is positive and bounded from above, we can assume the Sobolev embedding on the initial metric, since $(M, g(0))$ is a compact Riemannian manifold. Let $A_0, B_0 < \infty$ be positive constants such that for all $v \in W^{1,2}(M, g_0)$,

$$(5-1) \quad \|v\|_{2m/(m-2)} \leq A_0 \|\nabla v\|_2 + B_0 \|v\|_2,$$

where A_0 and B_0 depend only on m, g_0 , the lower bound for the Ricci curvature and the injectivity radius. We can then write (5-1) as

$$(5-2) \quad \left(\int_M v^{2m/(m-2)} d\mu_{g_0} \right)^{(m-2)/m} \leq A \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0} + B \int_M v^2 d\mu_{g_0},$$

where

$$A = \frac{1}{4} A_0 \quad \text{and} \quad B = \frac{1}{4} A_0 \sup S_g^-(\cdot, 0) + B_0$$

since $S_g(x, 0) + \sup S_g^-(\cdot, 0) = S_g^+(x, 0) - S_g^-(x, 0)$. We will assume that (5-2) holds uniformly for $g(t), t > 0$, and different A and B in order to prove the logarithmic Sobolev inequalities.

The usual way of deriving logarithmic Sobolev inequalities follows from a careful application of Hölder’s and Jensen’s inequalities since $\log v$ is a concave function, in which case

$$\int v^2 \ln v^{q-2} d\mu \leq \ln \int v^q d\mu$$

with the assumption that $\int v^2 d\mu = 1$. Then

$$\int v^2 \ln v d\mu \leq \frac{q}{q-2} \ln \left(\int v^q d\mu \right)^{1/q}.$$

Taking $q = 2m/(m-2)$, we have

$$\int v^2 \ln v d\mu \leq \frac{m}{2} \ln \left(\int v^{2m/(m-2)} d\mu \right)^{(m-2)/2m},$$

and by multiplying both sides by 2 we obtain the following result.

Lemma 5.1. *For any $v \in W^{1,2}(M, g_0)$ with $\|v\|_2 = 1$,*

$$(5-3) \quad \int_M v^2 \ln v^2 d\mu_{g_0} \leq \frac{m}{2} \ln \left(A \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0} + B \right).$$

See [Hsu 2008; Ye 2007; Zhang 2007] for similar proofs. Inequalities of the form (5-3) are usually estimated further by the application of an elementary inequality of the form $\ln y \leq \theta y - \ln \theta - 1$, where $\theta, y \geq 0$. Precisely, taking

$$y = A \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0} + B$$

in (5-3) gives us

$$(5-4) \quad \begin{aligned} \int_M v^2 \ln v^2 d\mu_{g_0} &\leq \frac{m\theta}{2} \left(A \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0} + B \right) - \frac{m}{2} (1 + \ln \alpha) \\ &= \frac{m\theta A}{2} \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0} + \frac{m\theta B}{2} - \frac{m}{2} - \frac{m}{2} \ln \alpha. \end{aligned}$$

We will now modify the arguments in both [Ye 2007] and [Zhang 2007] to prove the following result which says the monotonicity of the $\mathcal{W}_{\alpha, \epsilon}$ -entropy implies a logarithmic Sobolev inequality (not with the best constant). Here we assume the flow exists for all time.

Theorem 5.2. *Let (M, g) be a compact Riemannian manifold of dimension $m \geq 3$ and the metric $g(t)$ evolved by the $(\text{RH})_\alpha$ -flow. Assume that an L^2 -Sobolev embedding (5-2) holds true with respect to the initial metric $g(0) = g_0$. Then, we have*

$$(5-5) \quad \int_M v^2 \ln v^2 d\mu_{g(t)} \leq \int_M \sigma^2 (4|\nabla v|^2 + S_g v^2) d\mu_{g(t)} - \frac{m}{2} \ln \sigma^2 + (t + \sigma^2) \beta_1 + \frac{m}{2} \ln \frac{mA}{2e},$$

where $\sigma > 0$, $\beta_1 = 4A_0^{-1} B_0 + \sup S_g^-(\cdot, 0)$ and

$$\lambda_{\alpha 0} = \inf_{\|v\|_2=1} \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0};$$

that is, $\lambda_{\alpha 0}$ is the first eigenvalue of the operator $-4\Delta + S_g$.

Moreover, if $\lambda_{\alpha 0}$ is strictly positive for $S_g(\cdot, 0) > 0$ (i.e., $R(\cdot, 0) > \alpha(0)|\nabla u(0)|^2$), then

$$(5-6) \quad \int_M v^2 \ln v^2 d\mu_{g(t)} \leq \int_M \sigma^2 (4|\nabla v|^2 + S_g v^2) d\mu_{g(t)} - \frac{m}{2} \ln \sigma^2 + (t + \sigma^2) \beta_2 + \frac{m}{2} \ln \frac{mA}{2e}$$

holds with $B_0 = 0$, i.e., $\beta_2 = \sup S_g^-(\cdot, 0)$.

We first discuss some vital issues that will help put the proof of the above theorem in perspective. Now take an L^2 -solution $H = H(x, t)$ of the conjugate heat equation

$$(5-7) \quad \partial_t H = -\Delta H + S_g H$$

to be $H = (4\pi\tau)^{m/2} e^{-f}$. Relating the entropy $\mathcal{W}_{\alpha, \epsilon}$ with the idea of logarithmic Sobolev inequalities, we consider a function

$$(5-8) \quad v = \sqrt{H} = \frac{e^{-f/2}}{(4\pi\tau)^{m/4}}$$

such that $\int_M v^2 d\mu = 1$. We also notice that (5-8) implies $f = -\ln v^2 - (m/2) \ln \tau - (m/2) \ln(4\pi)$; hence the entropy (3-1) is rewritten as

$$(5-9) \quad \mathcal{W}_\epsilon(g, v, \tau) = \frac{\epsilon^2}{4\pi} \int_M (\tau(4|\nabla v|^2 + S_g v^2) - v^2 \ln v^2) d\mu - \frac{\epsilon^2}{4\pi} \frac{m}{2} \ln \tau \\ - \frac{\epsilon^2}{4\pi} \frac{m}{2} \ln(4\pi) + \left(1 - \frac{\epsilon^2}{4\pi}\right) \int_M f v^2 d\mu - \frac{m\epsilon^2}{4\pi} + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2}.$$

Define

$$(5-10) \quad \mathcal{W}_\epsilon^*(g, v, \tau) = \frac{\epsilon^2}{4\pi} \int_M (\tau(4|\nabla v|^2 + S_g v^2) - v^2 \ln v^2) d\mu$$

and

$$(5-11) \quad \mu_\epsilon^*(g, v, \tau) = \inf \left\{ \mathcal{W}_\epsilon^*(g, v, \tau) : \int_M v^2 d\mu = 1 \right\}.$$

Set $T^* = t^* + \sigma^2$ and $\tau(t) = T^* - t$ for $0 \leq t \leq t^*$, $\sigma > 0$. Then

$$\frac{d}{dt} \mathcal{W}_\epsilon(g, v, \tau) \\ = \frac{d}{dt} \mathcal{W}_\epsilon^*(g, v, \tau) - \frac{m\epsilon^2}{8\pi} \frac{d}{dt} \ln \tau + \left(1 - \frac{\epsilon^2}{4\pi}\right) \frac{\partial}{\partial t} \int_M f v^2 d\mu + \frac{m}{2} \ln \frac{4\pi}{\epsilon^2} \geq 0,$$

where the last inequality is due to the monotonicity of $\mathcal{W}_\epsilon(g, f, \tau)$, the proof of which also reveals that

$$\frac{\partial}{\partial t} \int_M f v^2 d\mu = -\mathcal{F}_\alpha + \frac{m}{2\tau},$$

where $\mathcal{F}_\alpha = \int_M (S_g + |\nabla f|^2) v^2 d\mu$ is Perelman and Müller's energy functional. Let $\lambda_{\alpha 0}$ be the first eigenvalue of the operator $-4\Delta + S_g$. Then, we know that $\lambda_{\alpha 0} = \inf_{\|u\|_2=1} \mathcal{F}_\alpha$. Therefore we arrive at

$$\frac{d}{dt} \mathcal{W}_\epsilon^* \geq \frac{n\epsilon^2}{8\pi} \frac{d}{dt} \ln \tau + \left(1 - \frac{\epsilon^2}{4\pi}\right) \lambda_{\alpha 0}.$$

To continue this argument, we should note that either (5-7) or (5-8) implies that the function $f = f(t)$ satisfies the following backward heat equation

$$(5-12) \quad \frac{\partial f}{\partial t} = -\Delta f + |\nabla f|^2 - S_g + \frac{m}{2\tau},$$

with $v = v(x, t)$ satisfying

$$(5-13) \quad \frac{\partial v}{\partial t} = -\Delta v + \frac{|\nabla v|^2}{v} + \frac{S_g}{2}v$$

on $[0, t^*]$ with a given terminal value at $t + t^*$ with $g = g(t^*)$.

Let v_0 be a minimizer of the entropy $\mathcal{W}_\epsilon(g, f, \tau_0)$ for all v with $\int_M v_0^2 d\mu_{g(t_0)} = 1$. We can then solve heat equation (5-12) backward in time with initial data $f(t_0) = f_0$ and v_0 chosen at $t = t_0$. Let u_j be the value of the conjugate heat equation (5-13) at $t = t_j$. We can define functions $f_j, j = 1, 2$, by

$$u_j = \frac{e^{-f_j/2}}{(4\pi\tau_j)^{n/4}}, \quad j = 1, 2.$$

Then by the monotonicity of $\mathcal{W}_{\alpha\epsilon}(g, f, \tau)$ -entropy, using Perelman's approach we have

$$\begin{aligned} \mu_\epsilon(g(t_1), \tau(t_1)) &= \inf_{\|v_0\|_{g(t_1)}=1} \mathcal{W}_\epsilon(g(t_1), f_0, \tau_1) \leq \mathcal{W}_\epsilon(g(t_1), f_1, \tau_1) \\ &\leq \mathcal{W}_\epsilon(g(t_2), f_2, \tau_2) = \inf_{\|v_0\|_{g(t_2)}=1} \mathcal{W}_\epsilon(g(t_2), f, \tau_2) = \mu_\epsilon(g(t_2), \tau(t_2)). \end{aligned}$$

It follows from the above that

$$\mu_\epsilon^*(g(t_1), \tau(t_1)) \leq \mu_\epsilon^*(g(t_2), \tau(t_2)) + \frac{n\epsilon^2}{8\pi} \ln \frac{\tau_1}{\tau_2}$$

for any $t_1 < t_2$, where $\tau_j = \tau(t_j), j = 1, 2$. Choosing $t_1 = 0$ and $t_2 = t^*$, we then obtain

$$(5-14) \quad \mu_\epsilon^*(g(0), t^* + \sigma^2) \leq \mu_\epsilon^*(g(t^*), \sigma^2) + \frac{n\epsilon^2}{8\pi} \ln \frac{t^* + \sigma^2}{\sigma}.$$

Since $0 < t^* < T$ is arbitrary, we can write (5-14) as

$$(5-15) \quad \mu_\epsilon^*(g(t), \sigma^2) \geq \mu_\epsilon^*(g(0), t + \sigma^2) + \frac{n\epsilon^2}{8\pi} \ln \frac{\sigma^2}{t + \sigma^2}$$

for all $t \in [0, T)$.¹ We now state the proof.

Proof. We now apply (5-4) with $g = g_0$ to estimate $\mu_\epsilon^*(g(0), t + \sigma^2)$. For any function $v \in W^{1,2}(M, g)$ with $\|v\|_2 = 1$ and using

$$\frac{m\theta A}{2} = t + \sigma^2 \quad \Rightarrow \quad \theta = \frac{8(t + \sigma^2)}{nA_0},$$

¹ The case $t = 0$ is optimal, as equality is attained.

the inequality in (5-4) becomes

$$\begin{aligned} \int_M v^2 \ln v^2 d\mu_{g_0} &\leq (t + \sigma^2) \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0} + \frac{m}{2} \frac{8(t + \sigma^2)B}{mA_0} \\ &\quad - \frac{m}{2} \ln \frac{8(t + \sigma^2)}{nA_0} - \frac{m}{2} \\ &= (t + \sigma^2) \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0} + 4(t + \sigma^2)BA_0^{-1} \\ &\quad - \frac{m}{2} \ln(t + \sigma^2) + \frac{m}{2}(\ln A_0 + \ln m - 3 \ln 2 - 1). \end{aligned}$$

Choosing $\epsilon^2 \leq 4\pi$ as before, it then follows that

$$(5-16) \quad \mu_\epsilon^*(g(0), t + \sigma^2) \geq \frac{m\epsilon^2}{4\pi} \left(\frac{1}{2} \ln(t + \sigma^2) - \frac{4}{m}(t + \sigma)BA_0^{-1} - \frac{1}{2}(\ln A_0 + \ln m - 3 \ln 2 - 1) \right).$$

Combining (5-15) and (5-16), we obtain

$$(5-17) \quad \mu_\epsilon^*(g(t), \sigma^2) \geq \frac{m\epsilon^2}{8\pi} \ln \sigma^2 - \frac{m\epsilon^2}{\pi}(t + \sigma^2)BA_0^{-1} - \frac{m\epsilon^2}{8\pi}(\ln A_0 + \ln m - 3 \ln 2 - 1),$$

which implies

$$\begin{aligned} \frac{\epsilon^2}{4\pi} \int_M (\sigma^2(4|\nabla v|^2 + S_g v^2) - v^2 \ln v^2) d\mu \\ \geq \frac{m\epsilon^2}{8\pi} \ln \sigma^2 - \frac{m\epsilon^2}{\pi}(t + \sigma^2)BA_0^{-1} - \frac{m\epsilon^2}{8\pi} \ln \frac{mA_0}{8e}. \end{aligned}$$

Therefore,

$$(5-18) \quad \int_M v^2 \ln v^2 d\mu \leq \int_M \sigma^2(4|\nabla v|^2 + S_g v^2) d\mu - \frac{m}{2} \ln \sigma^2 + 4(t + \sigma^2)BA_0^{-1} - \frac{m}{2} \ln \frac{nA_0}{8e}.$$

Choosing $\beta_1 = 4BA_0^{-1} = 4A_0^{-1}(B_0 + A \sup S_g^-(x, 0))$ and $A = A_0/4$, we obtain the result. We can also derive (5-6) in a similar manner. \square

6. Heat kernel bound via log-Sobolev inequalities

We apply the logarithmic Sobolev inequality obtained in the last section to derive an upper bound for the conjugate heat kernel along the Ricci flow, demonstrating that there is a lot of geometric information embedded in such inequalities. The basic ideas, due to Davies and Simon [1984], relate Nelson's hypercontractivity (see [Gross 1975]) to ultracontractivity (see also [Davies 1989]). These ideas always yield estimates with sharp constants. We modify the argument in [Zhang 2007] (see also [Lieb and Loss 1997; Zhang 2011]) to prove our result.

Theorem 6.1. *Suppose there exists a solution to the $(RH)_\alpha$ -flow with $m \geq 2$ and let $H(x, t; y)$ be the fundamental solution to the conjugate heat equation*

$$(6-1) \quad (-\partial_t - \Delta + S_g(x, \tau))w(x, \tau) = 0.$$

Then, for some nonnegative finite constant C depending on n, t, T, A_0, B_0 and $\sup S_g^-(\cdot, 0)$, the estimate

$$(6-2) \quad H(x, T; y) \leq CT^{-m/2}$$

holds, where $\partial_t \tau = -1$ and A_0, B_0 are as defined in the last section.

Without loss of generality, we may assume $w = w(x, t)$ to be a nonnegative solution of the conjugate heat equation (6-1) on the interval $[0, T]$, where $\partial_t \tau = -1$. Let $T > 0$ and $r(\tau) : [0, T] \rightarrow [1, \infty]$ be a continuously differentiable increasing function such that $r(0) = \infty$ and $r(T) = 1$. The function $r(\tau) = T/\tau$ gives a perfect example as we shall see below.

The idea here follows from the fact that if

$$w(x, t) = \int H(x, t; y)w_0(y) d\mu(y)$$

solves the heat equation, where $H(x, t; y)$ is the heat kernel, then

$$\sup_{w \neq 0} \frac{\|w(\cdot, t)\|_\infty}{\|w(\cdot, 0)\|_1} = \sup_{x, y} H(x, t; y).$$

We may obtain an estimation of the time derivative for the logarithms of the quantity

$$\|w\|_{r(t)} = \left(\int_M |w|^{r(t)} d\mu_{g(t)} \right)^{1/r(t)}$$

as follows:

$$\int_0^T \frac{\partial}{\partial t} \ln \|w\|_{r(t)} dt = \ln \frac{\|w(\cdot, t)\|_\infty}{\|w(\cdot, 0)\|_1}.$$

Proof. By routine computation,

$$\begin{aligned} \partial_t \|w\|_{r(t)} &= \partial_t \left(\int_M |w|^{r(t)} d\mu_{g(\tau)} \right)^{1/r(\tau)} \\ &= -\frac{\dot{r}(\tau)}{r^2(\tau)} \|w\|_{r(\tau)} \ln \|w\|_{r(\tau)}^{r(\tau)} + \frac{\|w\|_{r(\tau)}^{1-r(\tau)}}{r(\tau)} \left(\dot{r}(\tau) \int_M w^{r(\tau)} \ln w d\mu_{g(\tau)} \right. \\ &\quad \left. + r(\tau) \int_M (w^{r(\tau)-1} (-\Delta w + S_g w) + w^{r(\tau)} (-S_g)) d\mu_{g(\tau)} \right). \end{aligned}$$

Multiplying both sides by $r^2(\tau) \|w\|_{r(\tau)}^{r(\tau)}$, we have

$$\begin{aligned} r^2(\tau) \|w\|_{r(\tau)}^{r(\tau)} \partial_t \|w\|_{r(\tau)} &= -\dot{r}(\tau) \|w\|_{r(\tau)}^{r(\tau)+1} \ln \|w\|_{r(\tau)}^{r(\tau)} + r(\tau) \dot{r}(\tau) \|w\|_{r(\tau)} \int_M w^{r(\tau)} \ln u \, d\mu_g(\tau) \\ &\quad + r^2(\tau) \|w\|_{r(\tau)} \int_M w^{r(\tau)-1} (-\Delta w) \, d\mu_g(\tau) + r^2(\tau) \|w\|_{r(\tau)} \\ &\quad \times \int_M w^{r(\tau)-1} (S_g w) \, d\mu_g(\tau) - r(\tau) \|w\|_{r(\tau)} \int_M w^{r(\tau)} S_g \, d\mu_g(\tau). \end{aligned}$$

By the application of integration by parts, we have

$$\begin{aligned} r^2(\tau) \|w\|_{r(\tau)} \int_M w^{r(\tau)-1} (-\Delta w) \, d\mu_g(\tau) &= r^2(\tau) \|u\|_{r(\tau)} \int_M \nabla (u^{r(\tau)-1}) \nabla w \, d\mu_g(\tau) \\ &= r^2(\tau) (r(\tau) - 1) \|w\|_{r(\tau)} \int_M w^{r(\tau)-2} |\nabla w|^2 \, d\mu_g(\tau). \end{aligned}$$

Hence,

$$\begin{aligned} r^2(\tau) \|w\|_{r(\tau)}^{r(\tau)} \partial_t \|w\|_{r(\tau)} &= -\dot{r}(\tau) \|w\|_{r(\tau)}^{r(\tau)+1} \ln \|w\|_{r(\tau)}^{r(\tau)} + r(\tau) \dot{r}(\tau) \|w\|_{r(\tau)} \int_M w^{r(\tau)} \ln w \, d\mu_g(\tau) \\ &\quad + r^2(\tau) (r(\tau) - 1) \|w\|_{r(\tau)} \int_M w^{r(\tau)-2} |\nabla w|^2 \, d\mu_g(\tau) \\ &\quad + r(\tau) (r(\tau) - 1) \|w\|_{r(\tau)} \int_M S_g w^{r(\tau)} \, d\mu_g(\tau). \end{aligned}$$

Further dividing both sides by $\|w\|_{r(\tau)}$, we obtain

$$\begin{aligned} (6-3) \quad r^2(\tau) \|w\|_{r(\tau)}^{r(\tau)} \partial_t (\ln \|w\|_{r(\tau)}) &= -\dot{r}(\tau) \|u\|_{r(\tau)}^{r(\tau)} \ln \|w\|_{r(\tau)}^{r(\tau)} + r(\tau) \dot{r}(\tau) \int_M w^{r(\tau)} \ln w \, d\mu_g(\tau) \\ &\quad + r^2(\tau) (r(\tau) - 1) \int_M w^{r(\tau)-2} |\nabla w|^2 \, d\mu_g(\tau) \\ &\quad + r(\tau) (r(\tau) - 1) \int_M S_g w^{r(\tau)} \, d\mu_g(\tau). \end{aligned}$$

Using

$$v = \frac{w^{r(\tau)/2}}{\|w^{r(\tau)/2}\|_2} \implies v^2 = \frac{w^{r(\tau)}}{\|w\|_{r(\tau)}^{r(\tau)}},$$

we have

$$|\nabla v|^2 = \frac{r^2(\tau)}{4\|w\|_{r(\tau)}^{r(\tau)}} w^{r(\tau)-2} |\nabla w|^2$$

and

$$\ln v^2 = \ln w^{r(\tau)} - \ln \|w\|_{r(\tau)}^{r(\tau)}.$$

Therefore,

$$\begin{aligned} \dot{r}(\tau) \int_M v^2 \ln v^2 d\mu_{g(\tau)} &= \dot{r}(\tau) \int_M \frac{w^{r(\tau)}}{\|w\|_{r(\tau)}^{r(\tau)}} (\ln w^{r(\tau)} - \ln \|w\|_{r(\tau)}^{r(\tau)}) d\mu_{g(\tau)} \\ &= \frac{\dot{r}(\tau) r(\tau)}{\|w\|_{r(\tau)}^{r(\tau)}} \int_M w^{r(\tau)} \ln w^{r(\tau)} d\mu_{g(\tau)} - \dot{r} \ln \|w\|_{r(\tau)}^{r(\tau)}. \end{aligned}$$

Plugging these into (6-3), we arrive at

$$\begin{aligned} r^2(\tau) \partial_t (\ln \|w\|_{r(\tau)}) &= \dot{r}(\tau) \int_M v^2 \ln v^2 d\mu_{g(\tau)} + 4(r(\tau) - 1) \int_M |\nabla v|^2 d\mu_{g(\tau)} \\ &\quad + r(\tau)(r(\tau) - 1) \int_M Rv^2 d\mu_{g(\tau)} \\ &= \dot{r}(\tau) \int_M v^2 \ln v^2 d\mu_{g(\tau)} + (r(\tau) - 1) \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g(\tau)} \\ &\quad + (r(\tau) - 1)^2 \int_M S_g v^2 d\mu_{g(\tau)}. \end{aligned}$$

Using the choice $r(\tau) = T/\tau$, we have $\dot{r}(\tau) = -T/\tau^2$ and $r(\tau) - 1 = (T - \tau)/\tau$ so that we write the last equality as

$$\begin{aligned} r^2(\tau) \partial_t (\ln \|w\|_{r(\tau)}) &= -\frac{T}{\tau^2} \int_M v^2 \ln v^2 d\mu_{g(\tau)} + \frac{T - \tau}{\tau} \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g(\tau)} \\ &\quad + \left(\frac{T - \tau}{\tau}\right)^2 \int_M S_g v^2 d\mu_{g(\tau)} \\ &= \frac{T}{\tau^2} \left(\frac{\tau(T - \tau)}{T} \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g(\tau)} - \int_M v^2 \ln v^2 d\mu_{g(\tau)} \right) \\ &\quad + \left(\frac{T - \tau}{\tau}\right)^2 \int_M S_g v^2 d\mu_{g(\tau)}. \end{aligned}$$

From the log-Sobolev inequality (5-5) point of view, we may choose

$$\sigma^2 = \frac{4\tau(T - \tau)}{T} \leq \frac{T}{4},$$

and we get

$$(6-4) \quad r^2(\tau) \partial_t (\ln \|w\|_{r(t)}) \geq \frac{T}{\tau^2} \left(\frac{m}{2} \ln \sigma^2 - \frac{m}{2} \ln \frac{mA}{2e} - (t_0 + \sigma^2) \beta_1 \right) + \left(\frac{T-\tau}{\tau} \right)^2 \int_M S_g v^2 d\mu_{g(\tau)}$$

and

$$(6-5) \quad \partial_t (\ln \|w\|_{r(t)}) \geq \frac{1}{T} \left(\frac{m}{2} \ln \frac{4\pi\tau(T-\tau)}{T} - \frac{m}{2} \ln \frac{m\pi A}{2e} - (t_0 + \sigma^2) \beta_1 - T \sup S_g^-(\cdot, 0) \right).$$

Notice that (since $\sigma^2 \leq T/4$)

$$(t + \sigma^2) \beta_1 + T \sup R^-(\cdot, 0) = 4(t_0 + \sigma^2) (A_0^{-1} B_0 + \frac{1}{4} \sup S_g^-(\cdot, 0)) + T \sup S_g^-(\cdot, 0) \leq (4t_0 + T) A_0^{-1} B_0 + \frac{1}{4} (4t_0 + 5T) \sup S_g^-(\cdot, 0).$$

Denoting D by

$$D \equiv \frac{m}{2} \ln \frac{m\pi A}{2e} + (4t_0 + T) A_0^{-1} B_0,$$

substituting into (6-5) and integrating the result from 0 to T , we have

$$\begin{aligned} \ln \frac{\|w(\cdot, T)\|_{r(T)}}{\|w(\cdot, T)\|_{r(0)}} &\geq \frac{m}{2T} \int_0^T \ln \frac{4\pi\tau(T-\tau)}{T} dt - D - \frac{1}{4} (4t_0 + 5T) \sup R^-(\cdot, 0) \\ &= \frac{m}{2} \ln(4\pi) - \frac{n}{2} \ln T - n + n \ln T - D - \frac{1}{4} (4t_0 + 5T) \sup S_g^-(\cdot, 0) \\ &= \frac{m}{2} \ln(4\pi T) - m - D - (4t_0 + 5T) \sup S_g^-(\cdot, 0). \end{aligned}$$

This then yields

$$\ln \frac{\|w(\cdot, T)\|_1}{\|w(\cdot, T)\|_\infty} \geq \frac{m}{2} \ln(4\pi T) - m - D - \frac{1}{4} (4t_0 + 5T) \sup S_g^-(\cdot, 0),$$

which implies

$$\|w(\cdot, T)\|_\infty \leq \|w(\cdot, T)\|_1 \frac{\exp(\frac{1}{4}(4t_0 + 5T) \sup S_g^-(\cdot, 0) + D + m)}{(4\pi T)^{m/2}}.$$

Because

$$w(x, T) = \int_M H(x, T; y) w(y, 0) d\mu(y)_{g(\tau)},$$

where $H(x, T; y)$ is the conjugate heat kernel,

$$H(x, T; y) \leq \frac{\exp(mD)}{(4\pi T)^{m/2}} \exp(\frac{1}{4}(4t_0 + 5T) \sup S_g^-(\cdot, 0)).$$

This ends the proof of the estimate (6-2). □

7. The Sobolev inequality along $(RH)_\alpha$ -flow

In this section we show that global bounds on the heat kernel to the conjugate heat equation imply a uniform Sobolev inequality under Ricci-harmonic map flow. This type of proof is standard as contained in [Davies 1989, Chapter 2]. The same procedures have been adapted in [Zhang 2007] for Kähler–Ricci flow; see also [Ye 2007; Hsu 2008]. For completeness we give the summary of the approach.

For any $t \in [0, T)$, we define the operator

$$(7-1) \quad L := -\Delta_g + \frac{S_g + \sup_M S_g^-}{4}.$$

Since $R_g(\cdot, \tau) \geq -\sup_M R_g(\cdot, \tau)$, we know that $\Phi = \frac{1}{4}(S_g + \sup_M S_g^-) \geq 0$, $\Phi \in L^\infty(M)$. Then $L \geq 0$ and is essentially a self-adjoint operator on $L^2(M)$ with the associated quadratic form

$$(7-2) \quad \mathcal{Q}(v) = \int_M (|\nabla v|^2 + \Phi v^2) d\mu_g \quad \forall v \in W^{1,2}(M).$$

By the heat kernel convolution property, we have

$$(7-3) \quad e^{-tL}w_0 = \int_M H(x, t; y)w_0(y) d\mu_g(y),$$

where e^{-tL} is a self-adjoint positivity preserving semigroup for all $t \geq 0$. It is also a contraction on $L^\infty(M)$ and $L^1(M)$ for all $t \geq 0$. Then

$$(7-4) \quad \|e^{-tL}w_0\|_\infty \leq C_0 t^{-m/2} \|w_0\|_1.$$

The next step is to apply a theorem in [Davies 1989], which we state below as a lemma.

Lemma 7.1. *If $m \geq 2$, then a bound of the form*

$$(7-5) \quad \|e^{-tL}w_0\|_\infty \leq C_1 t^{-m/4} \|w_0\|_2$$

for all $t > 0$ and all $w_0 \in L^2(M)$ is equivalent to a bound of the form

$$(7-6) \quad \|w_0\|_{2m/(m-2)}^2 \leq C_2 \mathcal{Q}(w_0) \quad \forall w_0 \in W^{1,2}(M).$$

By Lemma 7.1 we can prove that

$$(7-7) \quad \left(\int_M v^{2m/(m-2)} d\mu_g \right)^{(m-2)/2} \leq A_0 \int_M (|\nabla v|^2 + \frac{1}{4}(S_g + \sup_M S_g^-)v^2) d\mu_g$$

using an estimate of the form (1-9). The only thing remaining for us to show is that estimates (7-4) and (7-5) are equivalent. We do this via the following lemma and the Hölder inequality.

Lemma 7.2. *Suppose $m \geq 2$ and $T < \infty$. Let $C_1 > 0$ be the same as C_1 in (7-5). Then we have*

$$(7-8) \quad \|e^{-tL}w_0\|_2 \leq C_1 t^{-m/4} \|w_0\|_1 \quad \forall w_0 \in L^1(M).$$

Now write $e^{-tL}w_0 = e^{-1/2tL}e^{-1/2tL}w_0$ and by assuming (7-5), we have

$$\|e^{-tL}w_0\|_\infty \leq C_1 t^{-m/4} \|e^{-1/2tL}w_0\|_2 \leq C_1^2 t^{-m/2} \|w_0\|_1.$$

Similarly, combining the fact that e^{-tL} is a contraction on $L^\infty(M)$ with bound (7-4) gives us (7-5). Indeed,

$$\begin{aligned} \|e^{-tL}w_0\|_\infty &= \left| \int_M H(x, t; y)w_0(y) d\mu_g(y) \right| \\ &\leq \left(\int_M H^{q'}(x, t; y)\mu_g(y) \right)^{1/q'} \left(\int_M w_0^q \mu_g(y) \right)^{1/q} \leq C t^{-m/2q} \|w_0\|_q, \end{aligned}$$

for all $w_0 \in L^q(M)$ with $1/q = 1 - 1/q'$ and $\int_M H(x, t; y) d\mu_g \leq 1$. Here we take q to satisfy $1 \leq q < m$ for obvious reason. (Though, by the Riez–Thorin interpolation theorem, the above holds for any $1 \leq q < \infty$ since e^{-tL} is a contraction on $L^1(M)$ and $L^\infty(M)$.)

The main result of this section is as follows.

Theorem 7.3. *With the conditions of the theorem in the introduction, we claim that estimate (1-8) implies the uniform Sobolev inequality (1-7).*

Proof. Based on the previous argument and a modification of the calculation in [Zhang 2007], we define the operator $\tilde{L} = L + 1$, which also has all the properties of L , ($\tilde{L} \geq 0$ and generates a symmetric Markov semigroup). Then for any positive constant c depending on m, T , a lower bound for R_{g_0} and an upper bound for A_0 such that for all $t \in [0, T)$ and $v \in \text{Dom}(\tilde{L}) \subseteq W^{1,q}(M)$,

$$(7-9) \quad \|\tilde{L}^{-1/2}w\|_{mq/(m-q)} \leq c \|w\|_q \quad \forall w \in W_0^{1,2}(M)$$

holds for $m \geq 3$. Since $\tilde{L}^{-1/2}$ is of weak type (p, q) , $p = mq/(m - q)$ for any $1 < q < m$. A simple analysis and the Marcinkiewicz interpolation theorem tell us that $\tilde{L}^{-1/2}$ is a bounded operator from L^q to L^p and that (7-9) holds.

Define $v(x, t) = \tilde{L}^{-1/2}w(x, t)$, which implies $w(x, t) = \tilde{L}^{1/2}v(x, t)$. Taking $q = 2$, we have

$$\|w\|_2^2 = \int_M \tilde{L}^{1/2}v \tilde{L}^{1/2}v d\mu_g = \int_M (\tilde{L}v)v d\mu_g = \int_M ((L + 1)v)v d\mu_g.$$

Combining with (7-9) and (7-6), we obtain the Sobolev inequality

$$(7-10) \quad \|v\|_{2m/(m-2)}^2 \leq c C_2 \left(\mathcal{Q}(v) + \int_M v^2 \mu_g \right),$$

whereby (1-7) follows with $A = c C_2$ and $B = \frac{1}{4}c C_2(\sup_M S_g + 1)$. □

Remark 7.4. Fixing t_0 during $(\text{RH})_\alpha$ -flow, it is clear that $\tilde{H} = e^{-1}H$ is the heat kernel of \tilde{L} and that

$$\int_M \tilde{H}(x, t; y) d\mu_g(y) \leq \int_M H(x, t; y) d\mu_g(y) \leq 1.$$

By the upper bound for H , we are sure that \tilde{H} obeys the global upper bound

$$\tilde{H}(x, t; y) d\mu_g(y) \leq \tilde{C}t^{-m/2}, \quad t > 0,$$

where \tilde{C} depends on m, A_0, B_0, t_0 and T . Similarly,

$$\|e^{-t\tilde{L}}w\|_\infty = \|e^{-t}e^{-t\tilde{L}}\|_\infty \leq e^{-t}Ct^{-m/2}\|w\|_1 = \tilde{C}t^{-m/2}\|w\|_1.$$

As a corollary, suppose

$$\lambda_{\alpha 0} = \inf_{\|v\|_2=1} \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0} > 0.$$

It can be proved by following [Zhang 2007] that Sobolev inequality (5-2) holds with $B = B(t = 0) = 0$ on a compact manifold (M, g_0) ; i.e.,

$$(7-11) \quad \left(\int_M v^{2m/(m-2)} d\mu_{g_0} \right)^{(m-2)/m} \leq \tilde{A}_0 \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g_0},$$

where \tilde{A} depends only on m, g_0 and $\lambda_{\alpha 0}$. Therefore, we have the following result.

Corollary 7.5. *Let (M, g) be a compact Riemannian manifold of dimension $m \geq 3$ and the metric $g(t)$ evolved by the $(\text{RH})_\alpha$ -flow. Assume that L^2 -Sobolev embedding (7-11) holds true with respect to the initial metric $g(0) = g_0$. Then, there exists a positive constant \tilde{A} depending on \tilde{A}_0 such that for all $v \in W^{1,2}(M, g(t)), t \in [0, T)$,*

$$(7-12) \quad \left(\int_M v^{2m/(m-2)} d\mu_{g(t)} \right)^{(m-2)/m} \leq \tilde{A} \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g(t)},$$

and

$$(7-13) \quad \int_M v^2 \ln v^2 d\mu_{g(t)} \leq \sigma^2 \int_M (4|\nabla v|^2 + S_g v^2) d\mu_{g(t)} - \frac{m}{2} \ln \sigma^2 + \frac{m}{2} \ln \frac{mA}{2e},$$

where $\sigma > 0$.

Remark 7.6. The smallest eigenvalue is an important quantity that gives a better understanding of the geometric nature of the underlying manifold. For instance, consider the operator semigroup e^{-tL} generated by $L := -\Delta + \Phi$, with $\Phi \in L^\infty(M, g)$. By spectral decomposition, we write a positive solution on M as

$$U = e^{-tL}u = \sum_{j=1}^{\infty} e^{-\lambda_j t} \psi_j \langle u, \psi_j \rangle_{L^2(M)}$$

for $u \in L^2(M)$ satisfying the Cauchy problem

$$\begin{aligned} \frac{\partial}{\partial t}(e^{-tL}u) &= -Le^{-tL}u, \\ U|_{t=0} &= u, \end{aligned}$$

and the eigenvalue problem $L\psi = \lambda\psi$, where $\{\psi_j\}_{j=1}^\infty$ forms a complete set of L^2 -orthonormal eigenfunctions of L and the corresponding eigenvalues can be arranged in a nondecreasing order $\lambda_1 \leq \lambda_2 \leq \dots$, with $\lambda_j \rightarrow \infty$. An interested reader will find the books [Davies 1989] and [Schoen and Yau 1994] useful in this respect.

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ABIMBOLA ABOLARINWA
DEPARTMENT OF MATHEMATICS
UNIVERSITY OF SUSSEX
BRIGHTON
BN1 9QH
UNITED KINGDOM
a.abolarinwa@sussex.ac.uk

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Department of Mathematics
University of California
Los Angeles, CA 90095-1555
popa@math.ucla.edu

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Department of Mathematics
University of California
Riverside, CA 92521-0135
chari@math.ucr.edu

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Department of Mathematics
University of California
Los Angeles, CA 90095-1555
liu@math.ucla.edu

Jie Qing
Department of Mathematics
University of California
Santa Cruz, CA 95064
qing@cats.ucsc.edu

Daryl Cooper
Department of Mathematics
University of California
Santa Barbara, CA 93106-3080
cooper@math.ucsb.edu

Jiang-Hua Lu
Department of Mathematics
The University of Hong Kong
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