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Fröhlich’s polaron Hamiltonian describes an electron coupled to the quantized phonon field of an ionic crystal. We show that in the strong coupling limit the dynamics of the polaron are approximated by an effective nonlinear partial differential equation due to Landau and Pekar, in which the phonon field is treated as a classical field.

1. Introduction and main result

1A. Setting of the problem. In this paper we are interested in the dynamics of a strongly coupled polaron. A polaron is a model of an electron in an ionic lattice interacting with its surrounding polarization field. Fröhlich [1937] proposed a quantum-mechanical Hamiltonian, given in (1-1) below, in order to describe the dynamics of a polaron. In this model the phonon field is treated as a quantum field. The Fröhlich Hamiltonian depends on a single parameter $\alpha > 0$ which describes the strength of the coupling between the electron and the phonon field. Landau and Pekar [1948] proposed a system of nonlinear PDEs, see (1-8), (1-9) below, to describe the dynamics of a polaron and used this in their famous computation of the effective polaron mass (see [Spohn 1987] for an alternative approach). They treat the phonons as a classical field. The derivation of their equations is phenomenological and they do not comment on the relation between their equations and the dynamics generated by Fröhlich’s Hamiltonian. Our purpose in this paper is to establish a connection between the two dynamics and to rigorously derive the Landau–Pekar equations from the Fröhlich dynamics in the strong coupling limit $\alpha \rightarrow \infty$ for a natural class of initial conditions and on certain time scales.

In order to describe this result in detail, we recall that the Fröhlich Hamiltonian acts in $\mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}$, where $\mathcal{L}^2(\mathbb{R}^3)$ corresponds to the electron and $\mathcal{F} = \mathcal{F}(\mathcal{L}^2(\mathbb{R}^3))$, the bosonic Fock space over $\mathcal{L}^2(\mathbb{R}^3)$, corresponds to the phonon field. The Hamiltonian is given by

$$p^2 + \sqrt{\alpha} \int_{\mathbb{R}^3} [e^{-ik \cdot x} a_k + e^{ik \cdot x} a_k^*] \frac{dk}{|k|} + \int_{\mathbb{R}^3} a_k^* a_k dk, \tag{1-1}$$

where $p := -i \nabla_x$ and x are momentum and position of the electron and a_k and a_k^* are annihilation and creation operators in \mathcal{F} satisfying the commutation relations

$$[a_k, a_{k'}^*] = \delta(k - k'), \quad [a_k, a_{k'}] = 0, \quad \text{and} \quad [a_k^*, a_{k'}^*] = 0. \tag{1-2}$$

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As mentioned before, the scalar $\alpha > 0$ describes the strength of the coupling between the electron and the phonon field and will be large in our study.

To facilitate later discussions we rescale the variables, as in [Frank and Schlein 2014],

$$x \mapsto \alpha^{-1}x, \quad k \mapsto \alpha k, \tag{1-3}$$

and find that the Hamiltonian in (1-1) is unitarily equivalent to $\alpha^2 \tilde{H}_\alpha^F$, where the new Hamiltonian \tilde{H}_α^F , acting again in $\mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}$, is defined as

$$\tilde{H}_\alpha^F := p^2 + \int_{\mathbb{R}^3} [e^{-ik \cdot x} b_k + e^{ik \cdot x} b_k^*] \frac{dk}{|k|} + \int_{\mathbb{R}^3} b_k^* b_k dk. \tag{1-4}$$

The new annihilation and creation operators $b_k := \alpha^{1/2} a_{\alpha k}$ and $b_k^* := \alpha^{1/2} a_{\alpha k}^*$ satisfy the commutation relations

$$[b_k, b_{k'}^*] = \alpha^{-2} \delta(k - k'), \quad [b_k, b_{k'}] = 0, \quad \text{and} \quad [b_k^*, b_{k'}^*] = 0. \tag{1-5}$$

We emphasize the α -dependence in (1-5).

We will discuss the dynamics generated by \tilde{H}_α^F for initial conditions of the product form

$$\psi_0 \otimes W(\alpha^2 \varphi_0) \Omega. \tag{1-6}$$

Here, Ω denotes the vacuum in \mathcal{F} and $W(f)$ denotes the Weyl operator,

$$W(f) := e^{b^*(f) - b(f)}, \tag{1-7}$$

so that $W(\alpha^2 \varphi) \Omega$ is a coherent state. This particular choice of initial conditions is motivated by Pekar’s approximation [1946; 1951] to the ground state energy, which uses exactly states of this form. Pekar’s approximation was made mathematically rigorous by Donsker and Varadhan [1983] (see [Lieb and Thomas 1997] for an alternative approach).

Clearly, the time-evolved state $e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega$ with $t \neq 0$ will in general no longer have an exact product structure. However, we will see that for large α (and t of order one, or even larger) it can be approximated, in a certain sense, by a state of the product form $\psi_t \otimes W(\alpha^2 \varphi_t) \Omega$, where ψ_t and φ_t solve the Landau–Pekar equations

$$i \partial_t \psi_t(x) = \left[-\Delta + \int_{\mathbb{R}^3} [e^{-ik \cdot x} \varphi_t(k) + e^{ik \cdot x} \overline{\varphi_t(k)}] \frac{dk}{|k|} \right] \psi_t(x), \tag{1-8}$$

$$i \alpha^2 \partial_t \varphi_t(k) = \varphi_t(k) + |k|^{-1} \int_{\mathbb{R}^3} |\psi_t(x)|^2 e^{ik \cdot x} dx \tag{1-9}$$

with initial data ψ_0 and φ_0 . Using standard methods one can show that for any $\psi_0 \in \mathcal{H}^1(\mathbb{R}^3)$, $\varphi_0 \in \mathcal{L}^2(\mathbb{R}^3)$ and $\alpha > 0$, the system (1-8), (1-9) has a global solution (ψ_t, φ_t) , which satisfies

$$\|\psi_t\|_{\mathcal{L}^2(\mathbb{R}^3)} = \|\psi_0\|_{\mathcal{L}^2(\mathbb{R}^3)} \quad \text{and} \quad \mathcal{E}(\psi_t, \varphi_t) = \mathcal{E}(\psi_0, \varphi_0) \quad \text{for all } t \in \mathbb{R}$$

with the energy

$$\mathcal{E}(\psi, \varphi) := \int_{\mathbb{R}^3} |\nabla \psi|^2 dx + \int_{\mathbb{R}^3} |\psi(x)|^2 \int_{\mathbb{R}^3} (e^{-ik \cdot x} \varphi(k) + e^{ik \cdot x} \overline{\varphi(k)}) \frac{dk}{|k|} dx + \int_{\mathbb{R}^3} |\varphi(k)|^2 dk. \tag{1-10}$$

We refer to [Lemma 2.1](#) and [Proposition 2.2](#) for more details about the solution (ψ_t, φ_t) . In the original work of Landau and Pekar the equations are given in a different, but equivalent form, and we explain this connection in [Section 1D](#).

1B. Main result. In order to prove our main result we need the following regularity and decay assumptions on the initial data. We denote by $\mathcal{H}^m(\mathbb{R}^3)$ the Sobolev space of order m and by

$$\mathcal{L}_{(m)}^2(\mathbb{R}^3) := \mathcal{L}^2(\mathbb{R}^3, (1+k^2)^m dk) \tag{1-11}$$

the weighted \mathcal{L}^2 space with norm

$$\|\varphi\|_{\mathcal{L}_{(m)}^2} = \left(\int_{\mathbb{R}^3} (1+k^2)^m |\varphi(k)|^2 dk \right)^{\frac{1}{2}}.$$

Our main result will be valid under the following:

Assumption 1.1. We assume $\psi_0 \in \mathcal{H}^4(\mathbb{R}^3)$ and $\varphi_0 \in \mathcal{L}_{(3)}^2(\mathbb{R}^3)$ with $\|\psi_0\|_{\mathcal{L}^2(\mathbb{R}^3)} = 1$.

A first version of our main result concerns the approximation of the reduced density matrices of $e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega$ in the trace norm.

Theorem 1.2. Assume ψ_0 and φ_0 satisfy [Assumption 1.1](#) and let (ψ_t, φ_t) be the solution of (1-8), (1-9) with initial condition (ψ_0, φ_0) . Define

$$\begin{aligned} \gamma_t^{\text{particle}} &:= \text{Tr}_{\mathcal{F}} |e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega\rangle \langle e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega|, \\ \gamma_t^{\text{field}} &:= \text{Tr}_{\mathcal{L}^2(\mathbb{R}^3)} |e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega\rangle \langle e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega|. \end{aligned}$$

Then, for all $\alpha \geq 1$ and all $t \in [-\alpha, \alpha]$,

$$\begin{aligned} \text{Tr}_{\mathcal{L}^2(\mathbb{R}^3)} |\gamma_t^{\text{particle}} - |\psi_t\rangle \langle \psi_t|| &\leq C \alpha^{-2} (1+t^2), \\ \text{Tr}_{\mathcal{F}} |\gamma_t^{\text{field}} - |W(\alpha^2 \varphi_t) \Omega\rangle \langle W(\alpha^2 \varphi_t) \Omega|| &\leq C \alpha^{-2} (1+t^2). \end{aligned}$$

Note that $\gamma_t^{\text{particle}}$, γ_t^{field} , $|\psi_t\rangle \langle \psi_t|$ and $|W(\alpha^2 \varphi_t) \Omega\rangle \langle W(\alpha^2 \varphi_t) \Omega|$ all have trace norm equal to one (in fact, they are nonnegative operators with trace one) and therefore [Theorem 1.2](#) gives a nontrivial approximation up to times $t = o(\alpha)$. Already the approximation up to times of order one is significant since this is the time scale on which ψ_t changes. It is a bonus that the same approximation is in fact valid for much longer times.

We emphasize that the Landau–Pekar approximation to the Fröhlich dynamics depends on α (through (1-9)). As we will explain in [Section 1C](#), without allowing for an α -dependence one cannot approximate $\gamma_t^{\text{particle}}$ with accuracy α^{-2} for times of order one.

We next present a more precise result which comes at the expense of a more complicated formulation. We approximate the state $e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega$ itself in $\mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}$, and not only its reduced density matrices. However, it turns out that up to the desired order α^{-2} this is *not* possible in terms of simple product states. Instead, we need to include an explicit nonproduct state of order α^{-1} which takes correlations between the particle and the field into account. The key observation is that this term

satisfies an almost orthogonality condition, so that it does not contribute to the reduced density matrices to order α^{-1} . For the statement we need the real scalar function ω defined as

$$\omega(t) := \alpha^2 \operatorname{Im}(\varphi_t, \partial_t \varphi_t) + \|\varphi_t\|_{\mathcal{L}^2(\mathbb{R}^3)}^2. \tag{1-12}$$

It will follow from [Lemma 2.1](#) below that this function is uniformly bounded in $t \in \mathbb{R}$.

The following is our main result.

Theorem 1.3. *Assume ψ_0 and φ_0 satisfy [Assumption 1.1](#) and let (ψ_t, φ_t) be the solution of [\(1-8\)](#), [\(1-9\)](#) with initial condition (ψ_0, φ_0) . Then there is a decomposition*

$$e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega = e^{-i \int_0^t \omega(s) ds} \psi_t \otimes W(\alpha^2 \varphi_t) \Omega + R(t) \tag{1-13}$$

and a constant $C > 0$ such that for all $\alpha \geq 1$ and all $t \in [-\alpha, \alpha]$,

$$\|\langle \Omega, W^*(\alpha^2 \varphi_t) R(t) \rangle_{\mathcal{F}}\|_{\mathcal{L}^2(\mathbb{R}^3)} \leq C \alpha^{-2} |t| (1 + |t|), \tag{1-14}$$

$$\|\langle \psi_t, W^*(\alpha^2 \varphi_t) R(t) \rangle_{\mathcal{L}^2(\mathbb{R}^3)}\|_{\mathcal{F}} \leq C \alpha^{-2} |t| (1 + |t|) \tag{1-15}$$

and

$$\|R(t)\|_{\mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}} \leq C \alpha^{-1} (1 + |t|). \tag{1-16}$$

More precisely, [\(1-13\)](#) holds with $R(t) = R_1(t) + R_2(t)$ and with the bounds

$$\|\langle \Omega, W^*(\alpha^2 \varphi_t) R_1(t) \rangle_{\mathcal{F}}\|_{\mathcal{L}^2(\mathbb{R}^3)} \leq C \alpha^{-2} t^2, \tag{1-17}$$

$$\|\langle \psi_t, W^*(\alpha^2 \varphi_t) R_1(t) \rangle_{\mathcal{L}^2(\mathbb{R}^3)}\|_{\mathcal{F}} \leq C \alpha^{-2} t^2 \tag{1-18}$$

and

$$\|R_2(t)\|_{\mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}} \leq C \alpha^{-2} |t| (1 + |t|), \quad \|R_1(t)\|_{\mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}} \leq C \alpha^{-1} (1 + |t|). \tag{1-19}$$

Similarly as before, we note that for $t = o(\alpha)$ the term $R(t)$ is of lower order than the main term $e^{-i \int_0^t \omega(s) ds} \psi_t \otimes W(\alpha^2 \varphi_t) \Omega$, which has constant norm equal to one.

The message of [Theorem 1.3](#) is that, while $R(t)$ is in general not of order α^{-2} (for times of order one), it can be split into a piece which is, namely $R_2(t)$, and a piece which satisfies almost orthogonality conditions, so that it does not contribute to the reduced particle or field density matrices at order α^{-1} either. The term $R_1(t)$ is given explicitly in [\(2-16\)](#) below.

[Theorem 1.3](#) implies [Theorem 1.2](#) by a simple abstract argument, which we explain in [Appendix D](#). In the following we concentrate on proving [Theorem 1.3](#).

In [Section 1C](#) we compare [Theorem 1.3](#) with a similar approximation in [\[Frank and Schlein 2014\]](#) where φ_t is independent of t . In [Lemma 1.4](#) we show that this simpler approximation does not yield the same accuracy in terms of powers of α^{-1} as [Theorem 1.3](#). In this sense [Theorem 1.3](#) derives the Landau–Pekar dynamics from the Fröhlich dynamics and answers an open question in [\[Frank and Schlein 2014\]](#).

While it is necessary to take the time dependence of φ_t into account, this dependence is still weak for times of order α as considered in our theorems. The field φ_t changes by order one only on times of order α^2 , and it would be desirable to extend [Theorems 1.2](#) and [1.3](#) to this time scale, at least for a certain class of initial conditions. This remains an open problem.

The key point in [Theorem 1.3](#) and novel aspect of this work are the almost orthogonality relations (1-14) and (1-15). As we will see in [Section 1C](#), they will be crucial for deriving [Theorem 1.2](#). Inequality (1-16) is not sufficient for this purpose. Let us discuss the motivation behind the almost orthogonality relations in more detail. We introduce the function

$$\tilde{\psi}_t := e^{-i \int_0^t \omega(s) ds} \psi_t \quad (1-20)$$

and consider the problem of approximating $e^{-i \tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega$ by a function of the form $\tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega$. (We do *not* assume at this point that $\tilde{\psi}_t$ and φ_t satisfy an equation.) Since $W(\alpha^2 \varphi_t)$ is unitary, this is the same as the problem of choosing $\tilde{\psi}_t$ and φ_t so as to minimize the norm of the vector

$$W^*(\alpha^2 \varphi_t) e^{-i \tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega - \tilde{\psi}_t \otimes \Omega. \quad (1-21)$$

Clearly, for given ψ_0, φ_0 and φ_t , the optimal choice for $\tilde{\psi}_t$ is

$$\tilde{\psi}_t = \langle \Omega, W^*(\alpha^2 \varphi_t) e^{-i \tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega \rangle_{\mathcal{F}}. \quad (1-22)$$

In order to determine φ_t we only solve the simpler problem of minimizing the norm of the projection of (1-21) onto the subspace $\text{span}\{\tilde{\psi}_t\} \otimes \mathcal{F}$. This norm could be made zero if we could achieve

$$\Omega = \langle \tilde{\psi}_t, W^*(\alpha^2 \varphi_t) e^{-i \tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega \rangle_{\mathcal{L}^2}. \quad (1-23)$$

While it may not be possible to have exact equalities in (1-22) and (1-23), we will see that the Landau–Pekar equations yield almost equalities. In fact, the almost orthogonality relations (1-14) and (1-15) in our main theorem state exactly that:

$$\tilde{\psi}_t - \langle \Omega, W^*(\alpha^2 \varphi_t) e^{-i \tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega \rangle_{\mathcal{F}} = O_{\mathcal{L}^2}(\alpha^{-2} |t| (1 + |t|)), \quad (1-24)$$

$$\Omega - \langle \tilde{\psi}_t, W^*(\alpha^2 \varphi_t) e^{-i \tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega \rangle_{\mathcal{L}^2} = O_{\mathcal{F}}(\alpha^{-2} |t| (1 + |t|)). \quad (1-25)$$

1C. Comparison with earlier results. The problem of approximating the Fröhlich dynamics of a polaron was studied before in [\[Frank and Schlein 2014\]](#). There a different and simpler effective equation is proposed in which only the particles move and the phonon field remains constant. In this subsection we show that [Theorem 1.2](#) is *not* valid for these effective dynamics from [\[Frank and Schlein 2014\]](#), in the sense that the reduced phonon density matrix cannot be approximated to within an error α^{-2} for times of order one. The fact that our [Theorem 1.2](#) achieves an approximation at this accuracy is because the phonon motion is taken into account in the Landau–Pekar equations. Technically this is reflected in the orthogonality conditions (1-14) and (1-15).

To be more specific we recall that in [\[Frank and Schlein 2014\]](#) it was shown that

$$\|e^{-i \tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega - e^{-i \|\varphi_0\|_2^2 t} \zeta_t \otimes W(\alpha^2 \varphi_0) \Omega\|_{\mathcal{L}^2 \otimes_{\mathcal{F}}} \leq C \alpha^{-1} (e^{C|t|} - 1)^{\frac{1}{2}}, \quad (1-26)$$

where ζ_t denotes the solution of the linear equation

$$i \partial_t \zeta_t(x) = \left[-\Delta + \int_{\mathbb{R}^3} [e^{-ikx} \varphi_0(k) + e^{ik \cdot x} \overline{\varphi_0(k)}] \frac{dk}{|k|} \right] \zeta_t(x)$$

with initial condition ψ_0 . We stress again that in this approximation, φ_0 does not evolve in time. An anonymous referee, to whom we are most grateful, has explained to us that the method of [Frank and Schlein 2014] actually leads to the bound

$$\|e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega - e^{-i\|\varphi_0\|_2^2 t} \zeta_t \otimes W(\alpha^2 \varphi_0) \Omega\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq C(e^{C|t|/\alpha} - 1)^{\frac{1}{2}}, \tag{1-27}$$

which provides an approximation even up to times of order $o(\alpha)$. With his/her kind permission we reproduce the argument in Appendix E.

As an aside we note that we recover a similar bound as a simple consequence of Theorem 1.3. (In fact, our new bound is better by a power of $\alpha^{-\frac{1}{2}}$ for times $t \gg 1$.) Namely, (1-16) says that

$$\|e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega - e^{-i \int_0^t \omega(s) ds} \psi_t \otimes W(\alpha^2 \varphi_t) \Omega\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq C\alpha^{-1}(1 + |t|). \tag{1-28}$$

(In [Frank and Schlein 2014] weaker regularity and decay assumptions are imposed on ψ_0 and φ_0 , but we emphasize that (1-16) is also valid under weaker assumptions than those in Assumption 1.1. In fact, the latter assumption is needed to bound $R_2(t)$, whereas for (1-16) one can avoid the use of Duhamel’s principle in Proposition 2.3.)

For the reduced density matrices, inequalities (1-26) and (1-27) give, using (D-1) and possibly changing the value of C ,

$$\begin{aligned} \text{Tr}_{\mathcal{L}^2} |\gamma_t^{\text{particle}} - |\zeta_t\rangle\langle\zeta_t|| &\leq C \min\{\alpha^{-1}(e^{C|t|} - 1)^{\frac{1}{2}}, (e^{C|t|/\alpha} - 1)^{\frac{1}{2}}\}, \\ \text{Tr}_{\mathcal{L}^2} |\gamma_t^{\text{field}} - |W(\alpha^2 \varphi_0) \Omega\rangle\langle W(\alpha^2 \varphi_0) \Omega|| &\leq C \min\{\alpha^{-1}(e^{C|t|} - 1)^{\frac{1}{2}}, (e^{C|t|/\alpha} - 1)^{\frac{1}{2}}\}. \end{aligned}$$

These bounds behave like α^{-1} for times of order one.

The next result shows that in this approximation of γ_t^{field} by a time-independent φ_0 the order α^{-1} (for times of order one) cannot be improved in general.

Lemma 1.4. *In addition to Assumption 1.1 suppose that $\varphi_0 \not\equiv -\sigma\psi_0$ in the notation (2-2). Then there are $\varepsilon > 0$, $C > 0$ and $c > 0$ such that for all $|t| \in [C\alpha^{-1}, \varepsilon]$ and all $\alpha \geq C/\varepsilon$,*

$$\text{Tr}_{\mathcal{F}} |\gamma_t^{\text{field}} - |W(\alpha^2 \varphi_0) \Omega\rangle\langle W(\alpha^2 \varphi_0) \Omega|| \geq c\alpha^{-1}|t|.$$

This lemma should be contrasted with Theorem 1.2, which says that the time-dependent approximation $|W(\alpha^2 \varphi_t) \Omega\rangle\langle W(\alpha^2 \varphi_t) \Omega|$ is correct to order α^{-2} (for times of order one). This argument shows the importance of the orthogonality conditions (1-14) and (1-15). Indeed, if we would only use (1-16), we would arrive at (1-28) and this would again only give an approximation to order α^{-1} (for times of order one).

Since Theorem 1.2 is a consequence of Theorem 1.3 and since we showed that one cannot replace φ_t by φ_0 in Theorem 1.2, the same applies also to Theorem 1.3.

Let us consider our problem from a wider perspective. We have a composite quantum system $\mathcal{H}_1 \otimes \mathcal{H}_2$ and a Hamiltonian which couples the two subsystems. Each system has an effective “Planck constant” and the characteristic feature of the problem is that the Planck constant of one system goes to zero, whereas that of the other system remains fixed. Thus, one of the systems becomes classical, whereas the other one

remains quantum-mechanical, and Ginibre, Nironi and Velo [Ginibre et al. 2006] used the term “partially classical limit” in a closely related context. (For us, the “Planck constant” of the phonons is α^{-2} , as can be seen from the commutation relations, whereas that of the electron is of order one.) A prime example of such a problem is the Born–Oppenheimer approximation, where the inverse square root of the nuclear mass plays the role of the small Planck constant.

Here, however, we consider the case where $\mathcal{H}_1 \otimes \mathcal{H}_2$ has infinitely many degrees of freedom. As is well known, our Hamiltonian is the Wick quantization of an energy functional on an infinite-dimensional phase space and the notion of “Planck constant” has a well-defined meaning through the commutation relations of the fields. (We emphasize that in our problem we can imagine that we have also a field Ψ for the electrons, but that we only consider the sector of a single electron.)

Although there is an enormous literature concerning the classical limit, starting with Hepp’s work [1974], and although we believe that the question of a partially classical limit is a very natural one which appears in many models, we are only aware of the single work [Ginibre et al. 2006] prior to [Frank and Schlein 2014] on this question, and it studies fluctuation dynamics. Closer to our focus here are the works [Falconi 2013; Ammari and Falconi 2014] about the Nelson model with a cut-off where, however, a classical limit on *both* systems is taken. On the level of results, one obtains equations similar to the Landau–Pekar equations (without the factor α^2 in (1-9)), but the proofs are completely different, as [Ammari and Falconi 2014] relies on the Wigner measure approach from [Ammari and Nier 2008; 2009].

The polaron model, in contrast to the Nelson model, does not require a cut-off, although this is not obvious since the operator $\int e^{ik \cdot x} b_k |k|^{-1} dk$ and its adjoint are not bounded relative to the number operator. Lieb and Yamazaki [1958] devised a method to deal with this problem in the stationary case, but it is not clear to us how to apply their argument in a dynamical setting and we consider our solution of this problem as a technical novelty in this paper. Our methods apply equally well to a partially classical limit in the cut-off Nelson model and, in fact, the proofs in that case would be considerably shorter.

1D. An equivalent form of the Landau–Pekar equations. Often the Landau–Pekar equations are stated in the form

$$i \partial_t \psi_t = (-\Delta + |x|^{-1} * P_t) \psi_t, \tag{1-29}$$

$$\alpha^4 \partial_t^2 P_t = -P_t - (2\pi)^2 |\psi_t|^2 \tag{1-30}$$

for a real-valued polarization field P_t ; see, e.g., [Landau and Pekar 1948; Devreese and Alexandrov 2009]. Let us show that this pair of equations is equivalent to the pair of equations that we discussed so far. In fact, assume that ψ_t and φ_t solve (1-8) and (1-9) and define

$$P_t(x) := (2\pi)^{-1} \operatorname{Re} \int_{\mathbb{R}^3} |k| \varphi_t(k) e^{-ik \cdot x} dk,$$

as well as the auxiliary function

$$Q_t(x) := (2\pi)^{-1} \operatorname{Im} \int_{\mathbb{R}^3} |k| \varphi_t(k) e^{-ik \cdot x} dk.$$

If we multiply (1-9) by $|k|$ and integrate with respect to $e^{-ik \cdot x}$, we obtain

$$i\alpha^2 \partial_t (P_t + iQ_t) = P_t + iQ_t + (2\pi)^2 |\psi_t|^2.$$

Since P_t and Q_t are real, this equation is equivalent to the pair of equations

$$\alpha^2 \partial_t P_t = Q_t, \quad \alpha^2 \partial_t Q_t = -P_t - (2\pi)^2 |\psi_t|^2.$$

Here we can eliminate Q_t by differentiating the first equation and arrive at (1-30).

Moreover, the inversion formula

$$\varphi_t(k) = (2\pi)^{-2} |k|^{-1} \int_{\mathbb{R}^3} (P_t + iQ_t) e^{ik \cdot x} dx$$

implies

$$\int_{\mathbb{R}^3} (e^{-ik \cdot x} \varphi_t(k) + e^{ik \cdot x} \overline{\varphi_t(k)}) \frac{dk}{|k|} = |x|^{-1} * P_t,$$

which yields (1-29).

2. Outline of the proof

2A. Well-posedness of the Landau–Pekar equations. We begin by discussing the well-posedness of the equations for ψ_t and φ_t in (1-8) and (1-9). We use the following abbreviations for the coupling terms in these equations,

$$V_\varphi(x) := \int_{\mathbb{R}^3} [e^{-ik \cdot x} \varphi(k) + e^{ik \cdot x} \overline{\varphi(k)}] \frac{dk}{|k|} \tag{2-1}$$

and

$$\sigma_\psi(k) := |k|^{-1} \int_{\mathbb{R}^3} |\psi(x)|^2 e^{ik \cdot x} dx. \tag{2-2}$$

The following lemma, which is proved in Appendix C, states global well-posedness in the energy space $\mathcal{H}^1(\mathbb{R}^3) \times \mathcal{L}^2(\mathbb{R}^3)$.

Lemma 2.1. *For any $(\psi_0, \varphi_0) \in \mathcal{H}^1(\mathbb{R}^3) \times \mathcal{L}^2(\mathbb{R}^3)$ there is a unique global solution (ψ_t, φ_t) of (1-8), (1-9). One has the conservation laws*

$$\|\psi_t\|_{\mathcal{L}^2} = \|\psi_0\|_{\mathcal{L}^2} \quad \text{and} \quad \mathcal{E}(\psi_t, \varphi_t) = \mathcal{E}(\psi_0, \varphi_0) \quad \text{for all } t \in \mathbb{R}.$$

Moreover, for all $\alpha > 0$ and all $t \in \mathbb{R}$,

$$\|\psi_t\|_{\mathcal{H}^1} \lesssim 1, \quad \|\varphi_t\|_{\mathcal{L}^2} \lesssim 1 \tag{2-3}$$

and

$$\|\partial_t \varphi_t\|_{\mathcal{L}^2} \lesssim \alpha^{-2}, \quad \|\varphi_t - \varphi_s\|_{\mathcal{L}^2} \lesssim \alpha^{-2} |t - s|, \quad \|\sigma_{\psi_t}\|_{\mathcal{L}^2} \lesssim 1. \tag{2-4}$$

In the proof of our main result we need to go beyond the energy space $\mathcal{H}^1(\mathbb{R}^3) \times \mathcal{L}^2(\mathbb{R}^3)$. The following proposition states that if the initial conditions have more regularity and decay then, at least for a certain

(long) time interval, we have bounds on the solution in the corresponding spaces. We will also need some bounds on the auxiliary functions $g_{s,t} : \mathbb{R}^3 \rightarrow \mathbb{C}$ defined by

$$g_{s,t}(x) := \int_{\mathbb{R}^3} [\overline{\varphi_t(k)} - \overline{\varphi_s(k)}] e^{ik \cdot x} \frac{dk}{|k|} \tag{2-5}$$

and $g_s : \mathbb{R}^3 \rightarrow \mathbb{C}$ defined by

$$g_s(x) := -\partial_s g_{s,t}(x) = \int_{\mathbb{R}^3} e^{ik \cdot x} \overline{\partial_s \varphi_s(k)} \frac{dk}{|k|}. \tag{2-6}$$

The following proposition will also be proved in [Appendix C](#).

Proposition 2.2. *Let $\tau > 0$. If (ψ_0, φ_0) satisfies [Assumption 1.1](#), then for all $\alpha > 0$ and for all $t, s \in [-\tau\alpha^2, \tau\alpha^2]$ we have*

$$\|\psi_t\|_{\mathcal{H}^4} \lesssim_\tau 1, \quad \|\varphi_t\|_{\mathcal{L}^2_{(3)}} \lesssim_\tau 1. \tag{2-7}$$

Moreover,

$$\|\partial_t \psi_t\|_{\mathcal{H}^2} \lesssim_\tau 1, \quad \|\partial_t \sigma \psi_t\|_{\mathcal{L}^2} \lesssim_\tau 1 \tag{2-8}$$

and

$$\|g_{s,t}\|_\infty \lesssim_\tau \alpha^{-2} |t - s|, \quad \|g_s\|_\infty \lesssim \alpha^{-2}. \tag{2-9}$$

2B. Decomposition of the solution. We now decompose the solution $e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega$ as claimed in [Theorem 1.3](#). In order to state this, we need to introduce some notations.

It will be convenient to work with the function $\tilde{\psi}_t$ from (1-20). Clearly, the bounds from [Lemma 2.1](#) and [Proposition 2.2](#) hold for $\tilde{\psi}_t$ as well. (For the bounds on $\partial_t \tilde{\psi}_t$ we use the fact that $|\omega(t)| \lesssim 1$ by [Lemma 2.1](#).) Moreover, we note that $\tilde{\psi}_t$ and φ_t satisfy the modified equations

$$i \partial_t \tilde{\psi}_t(x) = \left[-\Delta + \int_{\mathbb{R}^3} [e^{-ik \cdot x} \varphi_t(k) + e^{ik \cdot x} \bar{\varphi}_t(k)] \frac{dk}{|k|} + \omega(t) \right] \tilde{\psi}_t(x), \tag{2-10}$$

$$i \alpha^2 \partial_t \varphi_t(k) = \varphi_t(k) + |k|^{-1} \int_{\mathbb{R}^3} |\tilde{\psi}_t(x)|^2 e^{ik \cdot x} dx. \tag{2-11}$$

Next, we define for $\psi \in \mathcal{L}^2(\mathbb{R}^3)$ with $\|\psi\| = 1$ the orthogonal projections in $\mathcal{L}^2(\mathbb{R}^3)$,

$$P_\psi := |\psi\rangle\langle\psi|, \quad P_\psi^\perp := 1 - P_\psi = 1 - |\psi\rangle\langle\psi|.$$

The effective Schrödinger operator H_φ in $\mathcal{L}^2(\mathbb{R}^3)$ is defined by

$$H_\varphi := -\Delta + V_\varphi + \int_{\mathbb{R}^3} |\varphi(k)|^2 dk \tag{2-12}$$

with V_φ from (2-1). Moreover, let us introduce the operator

$$\tilde{H}_\varphi := W^*(\alpha^2 \varphi) \tilde{H}_\alpha^F W(\alpha^2 \varphi) \tag{2-13}$$

in $\mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}$. Using the commutation relations (see [Lemma A.1](#)) we find that

$$\tilde{H}_\varphi = H_\varphi + \int_{\mathbb{R}^3} [e^{ik \cdot x} b_k^* + e^{-ik \cdot x} b_k] \frac{dk}{|k|} + \int_{\mathbb{R}^3} [\varphi(k) b_k^* + \bar{\varphi}(k) b_k] dk + \int_{\mathbb{R}^3} b_k^* b_k dk. \tag{2-14}$$

Finally, we introduce the vector

$$F_{t,s} := P_{\psi_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \tag{2-15}$$

and define

$$D_0 := \int_0^t e^{iH_{\varphi_t} s} F_{t,s} ds$$

and

$$\begin{aligned} D_1 &:= \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} (e^{i\tilde{H}_{\varphi_t}(s+s_1)} e^{ik \cdot x} b_k^* e^{-iH_{\varphi_t} s_1} F_{t,s}) \frac{dk}{|k|} ds_1 ds, \\ D_2 &:= \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} (e^{i\tilde{H}_{\varphi_t}(s+s_1)} e^{-ik \cdot x} b_k e^{-iH_{\varphi_t} s_1} F_{t,s}) \frac{dk}{|k|} ds_1 ds, \\ D_3 &:= \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} (e^{i\tilde{H}_{\varphi_t}(s+s_1)} \varphi_t(k) b_k^* e^{-iH_{\varphi_t} s_1} F_{t,s}) dk ds_1 ds, \\ D_4 &:= \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} (e^{i\tilde{H}_{\varphi_t}(s+s_1)} \overline{\varphi_t(k)} b_k e^{-iH_{\varphi_t} s_1} F_{t,s}) dk ds_1 ds, \\ D_5 &:= \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} (e^{i\tilde{H}_{\varphi_t}(s+s_1)} b_k^* b_k e^{-iH_{\varphi_t} s_1} F_{t,s}) dk ds_1 ds. \end{aligned}$$

While these definitions might seem formal, we will show in [Theorem 2.5](#) that each of D_0, \dots, D_5 belongs to $\mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}$.

With these notations, the promised representation formula for the solution looks as follows.

Proposition 2.3. *Assume that $(\tilde{\psi}_t, \varphi_t)$ satisfy (2-10), (2-11) with initial conditions (ψ_0, φ_0) where $\|\psi_0\|^2 = 1$. Then for any $t \in \mathbb{R}$ one has the decomposition*

$$e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega = \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + R_1(t) + R_2(t)$$

with

$$R_1(t) := -i W(\alpha^2 \varphi_t) e^{-iH_{\varphi_t} t} D_0$$

and

$$R_2(t) := -W(\alpha^2 \varphi_t) e^{-i\tilde{H}_{\varphi_t} t} (D_1 + D_2 + D_3 + D_4 + D_5).$$

Clearly, in terms of the original function ψ_t , the term R_1 is explicitly given by

$$R_1(t) = -i W(\alpha^2 \varphi_t) \int_0^t \left[e^{-iH_{\varphi_t}(t-s) - i \int_0^s \omega(s_1) ds_1} P_{\psi_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* \psi_s \otimes \Omega) \frac{dk}{|k|} \right] ds. \tag{2-16}$$

The proof of [Proposition 2.3](#) makes use of equations (2-10), (2-11) for $(\tilde{\psi}_t, \varphi_t)$ as well as the Duhamel formula. We single out the use of the equations in the following lemma.

Lemma 2.4. *Assume that $(\tilde{\psi}_t, \varphi_t)$ satisfy (2-10), (2-11) with initial conditions (ψ_0, φ_0) where $\|\psi_0\|^2 = 1$. Then for any $t \in \mathbb{R}$ one has*

$$e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega = \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega - i \int_0^t e^{-i\tilde{H}_\alpha^F (t-s)} W(\alpha^2 \varphi_t) F_{t,s} ds. \tag{2-17}$$

Proof of Lemma 2.4. Applying the operator $e^{i\tilde{H}_\alpha^F t}$ to both sides of (2-17) we see that we need to prove

$$\psi_0 \otimes W(\alpha^2 \varphi_0) \Omega = e^{i\tilde{H}_\alpha^F t} \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega - i \int_0^t e^{i\tilde{H}_\alpha^F s} W(\alpha^2 \varphi_t) F_{t,s} ds.$$

This is clearly true at $t = 0$ and therefore we only need to show that the time derivatives of both sides coincide for all t ; that is, in view of definition (2-15) of $F_{t,s}$,

$$0 = e^{i\tilde{H}_\alpha^F t} \left[i \tilde{H}_\alpha^F \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \partial_t \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \tilde{\psi}_t \otimes \partial_t W(\alpha^2 \varphi_t) \Omega - i W(\alpha^2 \varphi_t) P_{\tilde{\psi}_t}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_t \otimes \Omega) \frac{dk}{|k|} \right].$$

This is, of course, the same as

$$i \tilde{H}_\alpha^F \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \partial_t \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \tilde{\psi}_t \otimes \partial_t W(\alpha^2 \varphi_t) \Omega = i W(\alpha^2 \varphi_t) P_{\tilde{\psi}_t}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_t \otimes \Omega) \frac{dk}{|k|}, \quad (2-18)$$

which is what we are going to show now.

We begin by rewriting the first term on the left side. Using (2-13) and (2-14) we obtain

$$i \tilde{H}_\alpha^F \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega = i H_{\varphi_t} \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \tilde{\psi}_t W(\alpha^2 \varphi_t) \left[i b^*(\varphi_t) + i \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* \frac{dk}{|k|} \right] \Omega.$$

In order to rewrite the third term on the left side of (2-18) we use the formula for $\partial_t W(\alpha^2 \varphi_t)$ from (A-4) below and find

$$\tilde{\psi}_t \otimes \partial_t W(\alpha^2 \varphi_t) \Omega = i \alpha^2 (\text{Im}(\varphi_t, \partial_t \varphi_t)) \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \alpha^2 \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) b^*(\partial_t \varphi_t) \Omega.$$

Thus, recalling the definition of ω in (1-12), we have shown that

$$i \tilde{H}_\alpha^F \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \partial_t \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \tilde{\psi}_t \otimes \partial_t W(\alpha^2 \varphi_t) \Omega = [\partial_t + i(-\Delta + V_{\varphi_t} + \omega(t))] \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega \quad (2-19)$$

$$+ W(\alpha^2 \varphi_t) \left[\alpha^2 b^*(\partial_t \varphi_t) + i b^*(\varphi_t) + i \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* \frac{dk}{|k|} \right] (\tilde{\psi}_t \otimes \Omega). \quad (2-20)$$

At this point in the proof we use the equations for $\tilde{\psi}_t$ and φ_t . It follows from (2-10) that line (2-19) vanishes identically. For line (2-20) we use (2-11) to obtain

$$\begin{aligned} i \tilde{H}_\alpha^F \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \partial_t \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega + \tilde{\psi}_t \otimes \partial_t W(\alpha^2 \varphi_t) \Omega &= i W(\alpha^2 \varphi_t) \left[\int_{\mathbb{R}^3} \left(- \int_{\mathbb{R}^3} |\tilde{\psi}_t(y)|^2 e^{ik \cdot y} dy + e^{ik \cdot x} \right) b_k^* \frac{dk}{|k|} \right] (\tilde{\psi}_t \otimes \Omega) \\ &= i W(\alpha^2 \varphi_t) P_{\tilde{\psi}_t}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_t \otimes \Omega) \frac{dk}{|k|}. \end{aligned} \quad (2-21)$$

Here we used the fact that $\|\tilde{\psi}_t\| = \|\psi_0\| = 1$ by assumption and [Lemma 2.1](#), and therefore

$$P_{\tilde{\psi}_t}^\perp = 1 - |\tilde{\psi}_t\rangle\langle\tilde{\psi}_t|.$$

[Equation \(2-21\)](#) proves [\(2-18\)](#) and completes the proof. □

Having proved [Lemma 2.4](#) we turn to the proof of [Proposition 2.3](#).

Proof of Proposition 2.3. It follows from [Lemma 2.4](#) and [\(2-13\)](#) that

$$e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega = \tilde{\psi}_t \otimes W(\alpha^2 \varphi_t) \Omega - i W(\alpha^2 \varphi_t) \int_0^t e^{-i\tilde{H}_{\varphi_t}(t-s)} F_{t,s} ds.$$

In the time integral on the right side we use Duhamel’s principle and [\(2-14\)](#),

$$e^{-i\tilde{H}_{\varphi_t}(t-s)} = e^{-iH_{\varphi_t}(t-s)} - i \int_0^{t-s} e^{-i\tilde{H}_{\varphi_t}(t-s-s_1)} \left(\int_{\mathbb{R}^3} [e^{ik \cdot x} b_k^* + e^{-ik \cdot x} b_k] \frac{dk}{|k|} + \int_{\mathbb{R}^3} b_k^* b_k dk + \int_{\mathbb{R}^3} [\varphi_t(k) b_k^* + \bar{\varphi}_t(k) b_k] dk \right) e^{-iH_{\varphi_t} s_1} ds_1.$$

[Proposition 2.3](#) now follows easily from the definition of D_0, \dots, D_5 . □

2C. Reduction of the proof of the main result. In the remainder of this paper we will prove the following.

Theorem 2.5. *Assume that ψ_0 and φ_0 satisfy [Assumption 1.1](#), let $(\tilde{\psi}_t, \varphi_t)$ be the solution of [\(2-10\)](#), [\(2-11\)](#) with initial condition (ψ_0, φ_0) and let D_0, \dots, D_5 be as in [Proposition 2.3](#). Then there is a constant $C > 0$ such that for all $\alpha \geq 1$ and $t \in [0, \alpha^2]$,*

$$\|D_0\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq C\alpha^{-1}(1+t), \tag{2-22}$$

$$\|D_1\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq C\alpha^{-2}t(1+t), \tag{2-23}$$

$$\|D_2\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq \alpha^{-2}t(1+t)(1+\alpha^{-1}t), \tag{2-24}$$

$$\|D_3\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq C\alpha^{-2}t(1+t)(1+\alpha^{-1}t), \tag{2-25}$$

$$\|D_4\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq C\alpha^{-2}t^2(1+\alpha^{-1}t), \tag{2-26}$$

$$\|D_5\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq C\alpha^{-3}t(1+t)(1+\alpha^{-2}t^2), \tag{2-27}$$

$$\|\langle \Omega, e^{-iH_{\varphi_t} t} D_0 \rangle_{\mathcal{F}}\|_{\mathcal{L}^2(\mathbb{R}^3)} \leq C\alpha^{-2}t^2, \tag{2-28}$$

$$\|\langle \tilde{\psi}_t, e^{-iH_{\varphi_t} t} D_0 \rangle_{\mathcal{L}^2(\mathbb{R}^3)}\|_{\mathcal{F}} \leq C\alpha^{-2}t^2(1+\alpha^{-2}t^2). \tag{2-29}$$

This theorem (and its analogue for $t \in [-\alpha^2, 0]$), together with the decomposition from [Proposition 2.3](#) and the fact that the operators $W(\alpha^2 \varphi_t)$, $e^{-iH_{\varphi_t} t}$ and $e^{-i\tilde{H}_{\varphi_t} t}$ are unitary, implies [Theorem 1.3](#). In fact, [\(2-22\)](#) implies the second bound in [\(1-19\)](#), [\(2-23\)](#)–[\(2-27\)](#) imply the first bound in [\(1-19\)](#), [\(2-28\)](#) implies [\(1-17\)](#) and [\(2-29\)](#) implies [\(1-18\)](#).

We emphasize that [Theorem 2.5](#) is valid up to times α^2 . (In fact, since the proof only relies on [Proposition 2.2](#), it is valid up to times $\tau\alpha^2$ for an arbitrary $\tau > 0$ with C depending on τ .) Consequently, the bounds in [Theorem 1.3](#) are also valid up to times α^2 . However, since the evolved state and the main

term in the approximation both have norm one, the bounds are only meaningful for times up to $\varepsilon\alpha$ for some small $\varepsilon > 0$.

The basic intuition behind the bounds on D_k , $k = 0, \dots, 5$, is that each annihilation or creation operator is of order α^{-1} and therefore D_0 , which contains only one creation operator, is of order α^{-1} , D_1, D_2, D_3, D_4 , which contain two creation or annihilation operators, are of order α^{-2} and D_5 , which contains three creation or annihilation operators, is of order α^{-3} . We illustrate this intuition in more detail in Section 2E with the simplest possible terms.

While this basic principle is true, it is oversimplifying the situation considerably as it does not take the slow-decaying terms $|k|^{-1}$ into account. The operator $\int e^{ik \cdot x} b_k^* |k|^{-1} dk$ and its adjoint are *not* bounded relative to the number operator $\int b_k^* b_k dx$. In fact, the treatment of these operators is the major difficulty that we have to overcome here.

At this point we have reduced the proof of Theorem 1.3 to the proof of Theorem 2.5, and the remainder of the paper is concerned with this. We bound D_0 in Section 3, D_1 in Section 4 and D_2 in Section 5. The terms D_3, D_4 and D_5 , which are easier to bound than D_1 and D_2 , are briefly discussed in Section 6. Finally, the bounds (2-28) and (2-29) will be proved in Subsections 7A and 7B, respectively.

2D. A further decomposition. Using the fact that $P_{\tilde{\psi}_t}^\perp = 1 - |\tilde{\psi}_t\rangle\langle\tilde{\psi}_t|$ (see the proof of Lemma 2.4), we have the decomposition

$$F_{t,s} = F_{t,s}^{(1)} - F_{t,s}^{(2)},$$

where

$$F_{t,s}^{(1)} := \int_{\mathbb{R}^3} (e^{ik \cdot x} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|}$$

and, with the notation σ_ψ from (2-2),

$$F_{t,s}^{(2)} := \tilde{\psi}_s \otimes W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b^*(\sigma_{\tilde{\psi}_s}) \Omega.$$

Correspondingly, we define

$$D_i = D_{i1} - D_{i2} \quad \text{for } i = 0, 1, 2, 3, 4, 5.$$

In general, the terms D_{i2} are easier to deal with than the terms D_{i1} . The reason for this is that $e^{ik \cdot x} |k|^{-1} \notin L^2(\mathbb{R}^3)$, whereas $\sigma_{\tilde{\psi}_t} \in L^2(\mathbb{R}^3)$ by Lemma 2.1, so the operator $\int e^{ik \cdot x} b_k^* |k|^{-1} dk$ in $F_{t,s}^{(1)}$ is harder to control than the operator $b^*(\sigma_{\tilde{\psi}_s})$ in $F_{t,s}^{(2)}$.

For $k = 1, \dots, 5$, both operators D_{i1} and D_{i2} involve an operator b_k^*, b_k or $b_k^* b_k$ to the left of $F_{t,s}^{(1)}$ or $F_{t,s}^{(2)}$, which in turn involves an operator $W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s)$. We now have the decomposition

$$D_{ij} = D_{ij1} + D_{ij2} \quad \text{for } i = 1, 2, 3, 4, 5 \text{ and } j = 1, 2,$$

where D_{ij1} denotes the expression with b_k, b_k^* or $b_k^* b_k$ commuted through the operator $W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s)$ and D_{ij2} denotes the expression coming from the commutator. To be explicit, we display some exemplary

cases,

$$D_{111} = \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} e^{ik \cdot x} e^{-iH_{\varphi_t} s_1} e^{ik' \cdot x} \\ \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|} ds_1 ds, \quad (2-30)$$

$$D_{121} = \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} e^{ik \cdot x} e^{-iH_{\varphi_t} s_1} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* b^*(\sigma_{\tilde{\psi}_s}) \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds_1 ds, \quad (2-31)$$

$$D_{211} = \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} e^{ik \cdot x} e^{-iH_{\varphi_t} s_1} e^{ik' \cdot x} \\ \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|} ds_1 ds, \quad (2-32)$$

$$D_{221} = \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} e^{ik \cdot x} e^{-iH_{\varphi_t} s_1} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k b^*(\sigma_{\tilde{\psi}_s}) \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds_1 ds. \quad (2-33)$$

The commutator terms can be computed with the help of [Corollary A.2](#). Recalling the definition of the function $g_{s,t}$ in (2-5), we have, for instance,

$$D_{112} = - \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} g_{s,t} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) e^{-iH_{\varphi_t} s_1} e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds_1 ds, \quad (2-34)$$

$$D_{122} = - \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} g_{s,t} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) e^{-iH_{\varphi_t} s_1} b^*(\sigma_{\tilde{\psi}_s}) \tilde{\psi}_s \otimes \Omega ds_1 ds, \quad (2-35)$$

$$D_{212} = - \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} \overline{g_{s,t}} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) e^{-iH_{\varphi_t} s_1} e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds_1 ds, \quad (2-36)$$

$$D_{222} = - \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} \overline{g_{s,t}} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) e^{-iH_{\varphi_t} s_1} b^*(\sigma_{\tilde{\psi}_s}) \tilde{\psi}_s \otimes \Omega ds_1 ds. \quad (2-37)$$

2E. Some warm-up bounds. In order to prepare for the rather technical sections that follow, we will first focus on the terms that do not include a term of the form $|k|^{-1}$, that is, on the terms D_{02} , D_{32} , D_{42} and D_{52} . We hope that this explains the underlying mechanism of our proof and the intuition that each annihilation or creation operator is of size α^{-1} .

Bound on D_{02} . We recall that

$$D_{02} = \int_0^t (e^{iH_{\varphi_t} s} \tilde{\psi}_s) \otimes (W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b^*(\sigma_{\tilde{\psi}_s}) \Omega) ds$$

and, therefore, by [Lemma 2.1](#),

$$\|D_{02}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq \int_0^t \|\tilde{\psi}_s\|_2 \|b^*(\sigma_{\tilde{\psi}_s}) \Omega\|_{\mathcal{F}} ds = \alpha^{-1} \int_0^t \|\sigma_{\tilde{\psi}_s}\|_2 ds \lesssim \alpha^{-1} t. \quad (2-38)$$

Bound on D_{32} . We have

$$D_{321} = \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \otimes (W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b^*(\varphi_t) b^*(\sigma_{\tilde{\psi}_s}) \Omega) ds_1 ds$$

and, according to [Corollary A.2](#),

$$D_{322} = -\int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \otimes ((\varphi_t - \varphi_s, \varphi_t) W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b^*(\sigma_{\tilde{\psi}_s}) \Omega) ds_1 ds.$$

By the bounds from [Lemma 2.1](#) we have

$$\|b^*(\varphi_t) b^*(\sigma_{\tilde{\psi}_s}) \Omega\|_{\mathcal{F}} = \alpha^{-2} (\|\varphi_t\|_2^2 \|\sigma_{\tilde{\psi}_s}\|_2^2 + |(\varphi_t, \sigma_{\tilde{\psi}_s})|^2)^{\frac{1}{2}} \lesssim \alpha^{-2},$$

and therefore, using also the conservation of the \mathcal{L}^2 -norm of $\tilde{\psi}_s$,

$$\|D_{321}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t^2.$$

On the other hand, the bounds from [Lemma 2.1](#) imply

$$\|b^*(\sigma_{\tilde{\psi}_s}) \Omega\|_{\mathcal{F}} = \alpha^{-1} \|\sigma_{\tilde{\psi}_s}\|_2 \lesssim \alpha^{-1}, \quad |(\varphi_t - \varphi_s, \varphi_t)| \lesssim \alpha^{-2} |t - s|,$$

and therefore, using again the conservation of the \mathcal{L}^2 -norm of $\tilde{\psi}_s$,

$$\|D_{322}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3} t^3.$$

Thus, we have shown that

$$\|D_{32}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t^2 (1 + \alpha^{-1} t). \quad (2-39)$$

Bound on D_{42} . We have

$$D_{421} = \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \otimes (W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b(\varphi_t) b^*(\sigma_{\tilde{\psi}_s}) \Omega) ds_1 ds$$

and, according to [Corollary A.2](#),

$$D_{422} = -\int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \otimes ((\varphi_t, \varphi_t - \varphi_s) W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b^*(\sigma_{\tilde{\psi}_s}) \Omega) ds_1 ds.$$

We commute once again and obtain

$$D_{421} = \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \otimes (\alpha^{-2} (\varphi_t, \sigma_{\tilde{\psi}_s}) W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \Omega) ds_1 ds.$$

According to [Lemma 2.1](#) we have $|(\varphi_t, \sigma_{\tilde{\psi}_s})| \lesssim 1$. This and computations similar to those in the bound of D_{32} yield

$$\|D_{421}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t^2, \quad \|D_{422}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3} t^3.$$

Thus, we have shown that

$$\|D_{42}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t^2 (1 + \alpha^{-1} t). \quad (2-40)$$

Bound on D_{52} . To simplify the notation, let us introduce

$$\mathcal{N} := \int_{\mathbb{R}^3} b_k^* b_k dk. \tag{2-41}$$

We have

$$D_{521} = \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \otimes (W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \mathcal{N} b^*(\sigma_{\tilde{\psi}_s}) \Omega) ds_1 ds.$$

Moreover, by [Corollary A.2](#),

$$\begin{aligned} [\mathcal{N}, W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s)] &= -W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (b(\varphi_t) - b(\varphi_s)) \\ &\quad - W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (b^*(\varphi_t) - b^*(\varphi_s)) + W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \|\varphi_t - \varphi_s\|_2^2, \end{aligned}$$

so

$$\begin{aligned} D_{522} &= \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \\ &\quad \otimes \left(W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (-b(\varphi_t - \varphi_s) - b^*(\varphi_t - \varphi_s) + \|\varphi_t - \varphi_s\|_2^2) b^*(\sigma_{\tilde{\psi}_s}) \Omega \right) ds_1 ds. \end{aligned}$$

We use $\mathcal{N} b^*(\sigma_{\tilde{\psi}_s}) = b^*(\sigma_{\tilde{\psi}_s}) \mathcal{N} + \alpha^{-2} b^*(\sigma_{\tilde{\psi}_s})$ and obtain

$$D_{521} = \alpha^{-2} \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \otimes (W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b^*(\sigma_{\tilde{\psi}_s}) \Omega) ds_1 ds.$$

Therefore, much as before,

$$\|D_{521}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3} t^2.$$

For D_{522} we commute again to get

$$\begin{aligned} D_{522} &= \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s) \otimes \left(W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \right. \\ &\quad \left. \times (-\alpha^{-2} (\varphi_t - \varphi_s, \sigma_{\tilde{\psi}_s}) \Omega - b^*(\varphi_t - \varphi_s) b^*(\sigma_{\tilde{\psi}_s}) \Omega + \|\varphi_t - \varphi_s\|_2^2 b^*(\sigma_{\tilde{\psi}_s}) \Omega) \right) ds_1 ds. \end{aligned}$$

For the second term on the right side we compute

$$\|b^*(\varphi_t - \varphi_s) b^*(\sigma_{\tilde{\psi}_s}) \Omega\|_{\mathcal{F}} = \alpha^{-2} (\|\varphi_t - \varphi_s\|_2^2 \|\sigma_{\tilde{\psi}_s}\|_2^2 + |(\varphi_t - \varphi_s, \sigma_{\tilde{\psi}_s})|^2)^{\frac{1}{2}}.$$

Using the bounds from [Lemma 2.1](#) for $\|\varphi_t - \varphi_s\|_2$ we obtain that

$$\|D_{522}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-4} t^3 (1 + \alpha^{-1} t).$$

Thus, we have shown that

$$\|D_{52}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3} t^2 (1 + \alpha^{-2} t^2). \tag{2-42}$$

3. Bound on D_0

We have already controlled D_{02} in [\(2-38\)](#), so it remains to consider D_{01} .

Bound on D_{01} . We recall that

$$D_{01} = \int_0^t e^{iH_{\varphi_t}s} \int_{\mathbb{R}^3} (e^{ik \cdot x} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} ds.$$

The main difficulty here, which we will encounter in various forms throughout this paper, is the unboundedness of the operator $\int e^{ik \cdot x} b_k^* |k|^{-1} dk$ (for any fixed $x \in \mathbb{R}^3$), since $e^{ik \cdot x} |k|^{-1} \notin \mathcal{L}^2(\mathbb{R}^3)$.

To overcome this difficulty we make use of the oscillatory behavior of $e^{ik \cdot x}$ via the formula

$$e^{ik \cdot x} = \frac{1 - ik \cdot \nabla_x}{1 + |k|^2} e^{ik \cdot x} \quad (3-1)$$

and aim at integrating by parts with respect to x . However, this integration by parts creates a new difficulty: the resulting operator ∇_x is unbounded and has to be controlled.

To overcome this new difficulty, it will be desirable to have an operator $(-\Delta + 1)^{-1}$ somewhere in the expression of D_{01} so that we can use it to control ∇_x , since obviously $\nabla_x (-\Delta + 1)^{-1}$ is bounded. It is equivalent and technically more convenient to work with $(H_{\varphi_t} + M)^{-1}$, where $M > 0$ is a large constant (independent of α and t), instead of $(-\Delta + 1)^{-1}$. In order to create this term we first integrate by parts in s and make use of the identity

$$e^{iH_{\varphi_t}s} = -i(H_{\varphi_t} + M)^{-1} e^{-iMs} \partial_s [e^{i(H_{\varphi_t} + M)s}]. \quad (3-2)$$

We obtain, using the fact that H_{φ_t} commutes with $W(\alpha^2 \varphi_s)$,

$$\begin{aligned} D_{01} &= -i e^{iH_{\varphi_t}t} (H_{\varphi_t} + M)^{-1} \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* \tilde{\psi}_t \otimes \Omega \frac{dk}{|k|} \\ &\quad + i W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_0) (H_{\varphi_t} + M)^{-1} \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* \tilde{\psi}_0 \otimes \Omega \frac{dk}{|k|} \\ &\quad + M \int_0^t e^{iH_{\varphi_t}s} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (H_{\varphi_t} + M)^{-1} \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds \\ &\quad + i \int_0^t e^{iH_{\varphi_t}s} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (H_{\varphi_t} + M)^{-1} \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* \partial_s \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds \\ &\quad + i \int_0^t e^{iH_{\varphi_t}s} W^*(\alpha^2 \varphi_t) (\partial_s W(\alpha^2 \varphi_s)) (H_{\varphi_t} + M)^{-1} \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds \\ &= D_{011} + D_{012} + D_{013} + D_{014} + D_{015}, \end{aligned}$$

where the terms D_{01k} are defined in a natural way. We will prove the following lemma.

Lemma 3.1. For $u \in \mathcal{H}^1(\mathbb{R}^3)$ and $f \in \mathcal{L}^2(\mathbb{R}^3)$,

$$\left\| (-\Delta + 1)^{-\frac{1}{2}} \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* u \otimes \Omega \frac{dk}{|k|} \right\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-1} \|u\|_{\mathcal{H}^1}$$

and

$$\left\| (-\Delta + 1)^{-\frac{1}{2}} \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^*(f) b_k^* u \otimes \Omega \frac{dk}{|k|} \right\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} \|u\|_{\mathcal{H}^1} \|f\|_2.$$

We defer the proof of this lemma to the end of this section and first show how to use it to control D_{01} . By [Corollary B.2](#) and [Lemma 2.1](#), we can choose M large enough so that $(H_{\varphi_t} + M)^{-\frac{1}{2}}(-\Delta + 1)^{\frac{1}{2}}$ is bounded uniformly in $t \in \mathbb{R}$. Moreover, by [Proposition 2.2](#), $\tilde{\psi}_t$ and $\partial_t \tilde{\psi}_t$ belong to $H^1(\mathbb{R}^3)$ and have uniformly bounded norms for $t \in [0, \alpha^2]$; see also the remark at the beginning of [Section 2B](#) concerning the bounds on $\partial_t \tilde{\psi}_t$. These facts, together with the unitarity of $e^{iH_{\varphi_t} s}$, $W^*(\alpha^2 \varphi_t)$ and $W(\alpha^2 \varphi_s)$, imply that

$$\|D_{011}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-1}, \quad \|D_{012}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-1}$$

and

$$\|D_{013}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-1}t, \quad \|D_{014}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-1}t.$$

In order to deal with the term D_{015} we make use of [\(A-4\)](#) and find

$$\begin{aligned} D_{015} &= -\int_0^t (\text{Im}(\varphi_s, \alpha^2 \partial_s \varphi_s)) e^{iH_{\varphi_t} s} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (H_{\varphi_t} + M)^{-1} \int_{\mathbb{R}^3} e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds \\ &\quad + i \int_0^t e^{iH_{\varphi_t} s} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (H_{\varphi_t} + M)^{-1} \int_{\mathbb{R}^3} e^{ik \cdot x} b^*(\alpha^2 \partial_s \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds \\ &\quad - i \int_0^t e^{iH_{\varphi_t} s} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (H_{\varphi_t} + M)^{-1} \int_{\mathbb{R}^3} e^{ik \cdot x} b(\alpha^2 \partial_s \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|} ds \\ &= D_{0151} + D_{0152} + D_{0153}. \end{aligned}$$

From [Lemma 2.1](#) we know that $|(\varphi_s, \alpha^2 \partial_s \varphi_s)| \lesssim 1$ and $\|\alpha^2 \partial_s \varphi_s\| \lesssim 1$. Thus, the first and the second bounds in [Lemma 3.1](#) imply, respectively,

$$\|D_{0151}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-1}t, \quad \|D_{0152}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2}t.$$

For D_{0153} we use the commutation relations to rewrite it as

$$D_{0153} = -i \int_0^t e^{iH_{\varphi_t} s} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (H_{\varphi_t} + M)^{-1} g_s \tilde{\psi}_s \otimes \Omega ds$$

with g_s from [\(2-6\)](#). Therefore, [Proposition 2.2](#) yields

$$\|D_{0153}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2}t.$$

To summarize, we have shown that

$$\|D_{01}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-1}(1+t). \tag{3-3}$$

Proof of Lemma 3.1. For any $\gamma \in \mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}$ and $(\Phi_k)_{k \in \mathbb{R}^3} \subset \mathcal{F}$, we use [\(3-1\)](#) to find

$$\begin{aligned} \left\langle \gamma, (-\Delta + 1)^{-\frac{1}{2}} \int_{\mathbb{R}^3} e^{ik \cdot x} u \otimes \Phi_k \frac{dk}{|k|} \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} &= \left\langle \nabla(-\Delta + 1)^{-\frac{1}{2}} \gamma, \int_{\mathbb{R}^3} \frac{ike^{ik \cdot x}}{|k|(1+|k|^2)} u \otimes \Phi_k dk \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} \\ &\quad + \left\langle (-\Delta + 1)^{-\frac{1}{2}} \gamma, \int_{\mathbb{R}^3} \frac{ike^{ik \cdot x}}{|k|(1+|k|^2)} (\nabla u) \otimes \Phi_k dk \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} \\ &\quad + \left\langle (-\Delta + 1)^{-\frac{1}{2}} \gamma, \int_{\mathbb{R}^3} \frac{e^{ik \cdot x}}{|k|(1+|k|^2)} u \otimes \Phi_k dk \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}}. \end{aligned}$$

Clearly,

$$\|\nabla(-\Delta + 1)^{-\frac{1}{2}}\gamma\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq \|\gamma\|_{\mathcal{L}^2 \otimes \mathcal{F}} \quad \text{and} \quad \|(-\Delta + 1)^{-\frac{1}{2}}\gamma\|_{\mathcal{L}^2 \otimes \mathcal{F}} \leq \|\gamma\|_{\mathcal{L}^2 \otimes \mathcal{F}},$$

so

$$\begin{aligned} & \left\| (-\Delta + 1)^{-\frac{1}{2}} \int_{\mathbb{R}^3} e^{ik \cdot x} u \otimes \Phi_k \frac{dk}{|k|} \right\|_{\mathcal{L}^2 \otimes \mathcal{F}} \\ & \lesssim \|u\|_{\mathcal{H}^1} \sup_{x \in \mathbb{R}^d} \left(\left\| \int_{\mathbb{R}^3} \frac{ike^{ik \cdot x}}{|k|(1+|k|^2)} \Phi_k dk \right\|_{\mathcal{F}} + \left\| \int_{\mathbb{R}^3} \frac{e^{ik \cdot x}}{|k|(1+|k|^2)} \Phi_k dk \right\|_{\mathcal{F}} \right). \end{aligned}$$

If $\Phi_k = b_k^* \Omega$, we use the fact that

$$\frac{1}{|k|(1+|k|^2)}, \frac{k}{|k|(1+|k|^2)} \in \mathcal{L}^2(\mathbb{R}^3)$$

to conclude that, uniformly in $x \in \mathbb{R}^3$,

$$\left\| \int_{\mathbb{R}^3} \frac{ike^{ik \cdot x}}{|k|(1+|k|^2)} b_k^* \Omega dk \right\|_{\mathcal{F}} \lesssim \alpha^{-1}, \quad \left\| \int_{\mathbb{R}^3} \frac{e^{ik \cdot x}}{|k|(1+|k|^2)} b_k^* \Omega dk \right\|_{\mathcal{F}} \lesssim \alpha^{-1}.$$

This proves the first bound in the lemma. If $\Phi_k = b^*(f)b_k^* \Omega$, one can similarly show that

$$\left\| \int_{\mathbb{R}^3} \frac{ike^{ik \cdot x}}{|k|(1+|k|^2)} b^*(f)b_k^* \Omega dk \right\|_{\mathcal{F}} \lesssim \frac{\|f\|_2}{\alpha^2}, \quad \left\| \int_{\mathbb{R}^3} \frac{e^{ik \cdot x}}{|k|(1+|k|^2)} b^*(f)b_k^* \Omega dk \right\|_{\mathcal{F}} \lesssim \frac{\|f\|_2}{\alpha^2}.$$

This proves the second bound in the lemma. \square

4. Bound on D_1

Bound on D_{111} . We recall equation (2-30) for D_{111} . In this equation, we commute $e^{ik \cdot x}$ with $e^{-iH_{\varphi_t} s}$. Thus, if we introduce the operator

$$H_{\varphi}(k) := e^{ik \cdot x} H_{\varphi} e^{-ik \cdot x} = (i\nabla_x + k)^2 + V_{\varphi} + \int_{\mathbb{R}^3} |\varphi(k)|^2 dk, \quad (4-1)$$

we obtain

$$\begin{aligned} D_{111} = & \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} e^{-iH_{\varphi_t}(k)s_1} e^{i(k+k') \cdot x} \\ & \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|} ds_1 ds. \end{aligned}$$

Controlling D_{111} is harder than controlling D_{01} because there are two slowly decaying terms $|k|^{-1}$ and $|k'|^{-1}$. The beginning of the proof, however, is similar; namely, for a large constant $M > 0$ to be specified, independent of t and α , we integrate by parts in s using

$$e^{i\tilde{H}_{\varphi_t} s} = -i(\tilde{H}_{\varphi_t} + M)^{-1} e^{-iMs} [\partial_s e^{i(\tilde{H}_{\varphi_t} + M)s}].$$

In this way we obtain

$$\begin{aligned}
D_{111} = & -i \int_0^t \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t} t} (\tilde{H}_{\varphi_t} + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} \\
& \quad \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_{t-s_1}) b_k^* b_{k'}^* \tilde{\psi}_{t-s_1} \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|} ds_1 \\
& + i \int_0^t \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t} s_1} (\tilde{H}_{\varphi_t} + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} \\
& \quad \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_0) b_k^* b_{k'}^* \tilde{\psi}_0 \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|} ds_1 \\
& + M \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (\tilde{H}_{\varphi_t} + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} \\
& \quad \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|} ds_1 ds \\
& + i \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (\tilde{H}_{\varphi_t} + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} \\
& \quad \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* b_{k'}^* [\partial_s \tilde{\psi}_s] \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|} ds_1 ds \\
& + i \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (\tilde{H}_{\varphi_t} + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} \\
& \quad \times W^*(\alpha^2 \varphi_t) [\partial_s W(\alpha^2 \varphi_s)] b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|} ds_1 ds.
\end{aligned}$$

We now use (2-13), which implies

$$\begin{aligned}
(\tilde{H}_{\varphi_t} + M)^{-1} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) &= W^*(\alpha^2 \varphi_t) (\tilde{H}_{\alpha^F} + M)^{-1} W(\alpha^2 \varphi_s) \\
&= W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) (\tilde{H}_{\varphi_s} + M)^{-1},
\end{aligned}$$

in order to commute $(\tilde{H}_{\varphi_t} + M)^{-1}$ to the right through $W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s)$. Moreover, we use Lemma A.3 to compute $\partial_s W(\alpha^2 \varphi_s)$. In this way we obtain

$$\begin{aligned}
D_{111} = & -i \int_0^t e^{i\tilde{H}_{\varphi_t} t} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) Q_1 ds \\
& + i \int_0^t e^{i\tilde{H}_{\varphi_t} s_1} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_0) Q_2 ds_1 \\
& + M \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) Q_3 ds_1 ds \\
& + i \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) Q_4 ds_1 ds \\
& + i \int_0^t \int_0^{t-s} e^{i\tilde{H}_{\varphi_t}(s+s_1)} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) Q_5 ds_1 ds
\end{aligned}$$

with

$$\begin{aligned}
Q_1 &:= (\tilde{H}_{\varphi_s} + M)^{-1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)(t-s)} e^{i(k'+k)\cdot x} b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|}, \\
Q_2 &:= (\tilde{H}_{\varphi_0} + M)^{-1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} b_k^* b_{k'}^* \tilde{\psi}_0 \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|},
\end{aligned}$$

$$\begin{aligned}
 Q_3 &:= (\tilde{H}_{\varphi_s} + M)^{-1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|}, \\
 Q_4 &:= (\tilde{H}_{\varphi_s} + M)^{-1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} b_k^* b_{k'}^* [\partial_s \tilde{\psi}_s] \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|}, \\
 Q_5 &:= (\tilde{H}_{\varphi_s} + M)^{-1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)s_1} e^{i(k'+k)\cdot x} (b^*(\alpha^2 \partial_s \varphi_s) - b(\alpha^2 \partial_s \varphi_s) \\
 &\quad + i \operatorname{Im}(\varphi_s, \alpha^2 \partial_s \varphi_s)) b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk'}{|k'|} \frac{dk}{|k|}.
 \end{aligned}$$

(Here, we suppress the dependence on t, s and s_1 in the notation of the Q_j .)

In the remainder of this section we shall show that, uniformly for $0 \leq s, s_1 \leq t \leq \alpha^2$,

$$\|Q_j\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} \quad \text{if } j = 1, 2, 3, 4, 5. \tag{4-2}$$

This will imply that

$$\|D_{111}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t(1+t). \tag{4-3}$$

Since the operator $(\tilde{H}_{\varphi_s} + M)^{-1}(-\Delta + \mathcal{N} + M)$ is *not* bounded, bounding the Q_j is rather involved. (Here \mathcal{N} was introduced in (2-41).) With the notation

$$Z_\varphi := V_\varphi + \int_{\mathbb{R}^3} |\varphi(x)|^2 dk + \int_{\mathbb{R}^3} (e^{-ik\cdot x} b_k + e^{ik\cdot x} b_k^*) \frac{dk}{|k|} + b(\varphi) + b^*(\varphi),$$

we abbreviate (2-14) as

$$\tilde{H}_\varphi = -\Delta + \mathcal{N} + Z_\varphi.$$

Defining

$$\tilde{Z}_\varphi := (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} Z_\varphi (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}},$$

we have

$$\begin{aligned}
 (\tilde{H}_\varphi + M)^{-1} &= (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} (1 + \tilde{Z}_\varphi)^{-1} (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} \\
 &= (-\Delta + \mathcal{N} + M)^{-1} - (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} (1 + \tilde{Z}_\varphi)^{-1} (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} Z_\varphi (-\Delta + \mathcal{N} + M)^{-1}.
 \end{aligned}$$

It is not difficult to see that for every $\varepsilon > 0$ and $A > 0$ there is an M such that

$$\|\tilde{Z}_\varphi\|_{\mathcal{L}^2 \otimes \mathcal{F} \rightarrow \mathcal{L}^2 \otimes F} \leq \varepsilon \tag{4-4}$$

for all φ with $\|\varphi\|_{\mathcal{L}^2} \leq A$; for details of this argument we refer to [Frank and Schlein 2014]. Thus, using the bound on $\|\varphi_s\|_{\mathcal{L}^2}$ from Lemma 2.1, we can choose M in such a way that

$$\|\tilde{Z}_{\varphi_s}\|_{\mathcal{L}^2 \otimes \mathcal{F} \rightarrow \mathcal{L}^2 \otimes F} \leq \frac{1}{2} \quad \text{for all } s > 0.$$

Therefore, the operator $1 + \tilde{Z}_{\varphi_s}$ in the above formula for $(H_{\varphi_s} + M)^{-1}$ is invertible. We use this formula to decompose

$$\begin{aligned}
 Q_1 &= \left(1 - (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} (1 + \tilde{Z}_{\varphi_s})^{-1} (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} \left(V_{\varphi_s} + \int_{\mathbb{R}^3} |\varphi_s(x)|^2 dk + b(\varphi_s) + b^*(\varphi_s) \right) \right) Q_{10} \\
 &\quad - (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} (1 + \tilde{Z}_{\varphi_s})^{-1} (Q_{11} + Q_{12})
 \end{aligned} \tag{4-5}$$

with

$$Q_{10} := (-\Delta + \mathcal{N} + M)^{-1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)(t-s)} e^{i(k'+k)\cdot x} b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk' dk}{|k'| |k|},$$

$$Q_{11} := (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} \left(\int_{\mathbb{R}^3} e^{-ik''\cdot x} b_{k''} \frac{dk''}{|k''|} \right) (-\Delta + \mathcal{N} + M)^{-1} \\ \times \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)(t-s)} e^{i(k'+k)\cdot x} b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk' dk}{|k'| |k|},$$

$$Q_{12} := (-\Delta + \mathcal{N} + M)^{-\frac{1}{2}} \left(\int_{\mathbb{R}^3} e^{ik''\cdot x} b_{k''}^* \frac{dk''}{|k''|} \right) (-\Delta + \mathcal{N} + M)^{-1} \\ \times \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)(t-s)} e^{i(k'+k)\cdot x} b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk' dk}{|k'| |k|}.$$

Using (4-4), the fact that $(-\Delta + \mathcal{N} + M)^{-\frac{1}{2}}(b(\varphi_s) + b^*(\varphi_s))$ is bounded uniformly in s , as well as the estimates $\|V_{\varphi_s}\|_{\infty} \lesssim 1$ (from (C-1) and Proposition 2.2) and $\|\varphi_s\|_2 \lesssim 1$ (from Lemma 2.1), we conclude from (4-5) that

$$\|Q_1\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \|Q_{10}\|_{\mathcal{L}^2 \otimes \mathcal{F}} + \|Q_{11}\|_{\mathcal{L}^2 \otimes \mathcal{F}} + \|Q_{12}\|_{\mathcal{L}^2 \otimes \mathcal{F}}.$$

We now bound the three terms on the right side separately.

Bound on Q_{10} . To control Q_{10} we prove an analogue of Lemma 3.1 for the case of two singularities.

Lemma 4.1. For $u \in \mathcal{H}^2(\mathbb{R}^3)$, $f \in \mathcal{L}^2(\mathbb{R}^3)$ and $s \in \mathbb{R}$,

$$\left\| (-\Delta + 1)^{-1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)s} e^{i(k+k')\cdot x} b_k^* b_{k'}^* u \otimes \Omega \frac{dk' dk}{|k'| |k|} \right\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} \|u\|_{\mathcal{H}^2}.$$

Before proving this lemma we show how to use it to bound Q_{10} . Note that, since Q_{10} involves only $b_k^* b_{k'}^* \Omega$, the operator $(-\Delta + \mathcal{N} + M)^{-1}$ in its definition can be replaced by $(-\Delta + 2\alpha^{-2} + M)^{-1}$. This observation, together with Lemma 4.1 and the uniform boundedness of $\tilde{\psi}_s$ in \mathcal{H}^2 for $s \in [0, \alpha^2]$ (see Proposition 2.2), proves that

$$\|Q_{10}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2}. \tag{4-6}$$

Proof of Lemma 4.1. We shall show that for any $\gamma \in \mathcal{L}^2(\mathbb{R}^3) \otimes \mathcal{F}$,

$$\left| \left\langle \gamma, (-\Delta + 1)^{-1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)s} e^{i(k+k')\cdot x} b_k^* b_{k'}^* u \otimes \Omega \frac{dk' dk}{|k'| |k|} \right\rangle \right| \lesssim \alpha^{-2} \|\gamma\|_{\mathcal{L}^2 \otimes \mathcal{F}} \|u\|_{\mathcal{H}^2}.$$

We integrate by parts twice in x and use (3-1) with k replaced by $k + k'$. A typical term that is obtained in this way in the inner product on the left side is

$$\left\langle e^{iH_{\varphi_t}(k)s} \partial_{x_i} \partial_{x_j} (-\Delta + 1)^{-1} \gamma, \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i(k+k')\cdot x} b_k^* b_{k'}^* u \otimes \Omega \frac{(k_i + k'_i)(k_j + k'_j) dk' dk}{|k| |k'| (1 + |k + k'|^2)^2} \right\rangle.$$

Since $\partial_{x_i} \partial_{x_j} (-\Delta + 1)^{-1}$ is bounded and $e^{iH_{\varphi_t}(k)s}$ is unitary, the vector on the left side of the inner product is bounded in norm by $\|\gamma\|_{\mathcal{L}^2 \otimes \mathcal{F}}$. We now show that the vector on the right side of the inner

product is bounded as well. We compute

$$\begin{aligned} & \left\| \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{i(k+k') \cdot x} b_k^* b_{k'}^* u \otimes \Omega \frac{(k_i + k'_i)(k_j + k'_j)}{|k||k'| (1 + |k + k'|^2)^2} dk' dk \right\|_{\mathcal{L}^2 \otimes \mathcal{F}}^2 \\ & = 2\alpha^{-4} \|u\|_2^2 \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \frac{(k_i + k'_i)^2 (k_j + k'_j)^2}{|k|^2 |k'|^2 (1 + |k + k'|^2)^4} dk' dk. \end{aligned}$$

The desired bound now follows from the fact that the double integral on the right side is finite. Other terms that arise in the integration by parts are controlled similarly and we omit the details. \square

Bound on Q_{11} . By considering the number of involved field particles, we can replace \mathcal{N} in the definition of Q_{11} by numbers and obtain

$$\begin{aligned} Q_{11} & = (-\Delta + \alpha^{-2} + M)^{-\frac{1}{2}} \left(\int_{\mathbb{R}^3} e^{-ik'' \cdot x} b_{k''} \frac{dk''}{|k''|} \right) (-\Delta + 2\alpha^{-2} + M)^{-1} \\ & \quad \times \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)(t-s)} e^{i(k'+k) \cdot x} b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk' dk}{|k'| |k|}. \end{aligned}$$

Next, by commuting $b_{k''}$ to the right,

$$\begin{aligned} Q_{11} & = \alpha^{-2} (-\Delta + \alpha^{-2} + M)^{-\frac{1}{2}} \int_{\mathbb{R}^3} ((i\nabla - k')^2 + 2\alpha^{-2} + M)^{-1} \\ & \quad \times \int_{\mathbb{R}^3} e^{-ik' \cdot x} e^{-iH_{\varphi_t}(k)(t-s)} e^{i(k'+k) \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega \frac{dk' dk}{|k'|^2 |k|} \\ & \quad + \alpha^{-2} (-\Delta + \alpha^{-2} + M)^{-\frac{1}{2}} \int_{\mathbb{R}^3} ((i\nabla - k)^2 + 2\alpha^{-2} + M)^{-1} \\ & \quad \times \int_{\mathbb{R}^3} e^{-ik \cdot x} e^{-iH_{\varphi_t}(k)(t-s)} e^{i(k'+k) \cdot x} b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk' dk}{|k'| |k|^2}. \end{aligned}$$

It remains to compute the norm of this expression. Since this is considerably easier than for Q_{12} , we omit the details and only state the final result,

$$\|Q_{11}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3}. \quad (4-7)$$

Bound on Q_{12} . In the same way as for Q_{11} , we can replace \mathcal{N} by a number, so that

$$\begin{aligned} Q_{12} & = (-\Delta + 3\alpha^{-2} + M)^{-\frac{1}{2}} \int_{\mathbb{R}^3} e^{ik'' \cdot x} b_{k''}^* (-\Delta + 2\alpha^{-2} + M)^{-1} \\ & \quad \times \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-iH_{\varphi_t}(k)(t-s)} e^{i(k'+k) \cdot x} b_k^* b_{k'}^* \tilde{\psi}_s \otimes \Omega \frac{dk' dk}{|k'| |k|}. \end{aligned}$$

Next, we commute $e^{ik'' \cdot x}$ and $e^{i(k'+k) \cdot x}$ to the right and obtain

$$\begin{aligned} Q_{12} & = \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} b_k^* b_{k'}^* b_{k''}^* e^{i(k+k'+k'') \cdot x} ((i\nabla - k - k' - k'')^2 + 3\alpha^{-2} + M)^{-\frac{1}{2}} \\ & \quad \times ((i\nabla - k - k')^2 + 2\alpha^{-2} + M)^{-1} e^{-iH_{\varphi_t}(-k')(t-s)} \tilde{\psi}_s \otimes \Omega \frac{dk'' dk' dk}{|k''| |k'| |k|}. \end{aligned}$$

We now compute the norm of this expression. For the part of the norm over \mathcal{F} , we use the fact that

$$\begin{aligned} \alpha^6 \langle \Omega, b_{k_1} b_{k_2} b_{k_3} b_{k_4}^* b_{k_5}^* b_{k_6}^* \Omega \rangle &= \delta(k_1 - k_4) \delta(k_2 - k_5) \delta(k_3 - k_6) + \delta(k_1 - k_4) \delta(k_2 - k_6) \delta(k_3 - k_5) \\ &\quad + \delta(k_1 - k_5) \delta(k_2 - k_4) \delta(k_3 - k_6) + \delta(k_1 - k_5) \delta(k_2 - k_6) \delta(k_3 - k_4) \\ &\quad + \delta(k_1 - k_6) \delta(k_2 - k_4) \delta(k_3 - k_5) + \delta(k_1 - k_6) \delta(k_2 - k_4) \delta(k_3 - k_6) \end{aligned}$$

to write

$$\|Q_{12}\|_{\mathcal{L}^2 \otimes \mathcal{F}}^2 = \alpha^{-6} (X_1 + \dots + X_6), \tag{4-8}$$

where, for instance,

$$\begin{aligned} X_1 := \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} &\left\langle e^{-iH_{\varphi_t}(-k')(t-s)} \tilde{\psi}_s, ((i\nabla - k - k' - k'')^2 + 3\alpha^{-2} + M)^{-1} \right. \\ &\left. \times ((i\nabla - k - k')^2 + 2\alpha^{-2} + M)^{-2} e^{-iH_{\varphi_t}(-k')(t-s)} \tilde{\psi}_s \right\rangle \frac{dk''}{|k''|^2} \frac{dk'}{|k'|^2} \frac{dk}{|k|^2} \end{aligned}$$

and

$$\begin{aligned} X_2 := \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} &\left\langle e^{-iH_{\varphi_t}(-k'')(t-s)} \tilde{\psi}_s, ((i\nabla - k - k' - k'')^2 + 3\alpha^{-2} + M)^{-1} \right. \\ &\left. \times ((i\nabla - k - k'')^2 + 2\alpha^{-2} + M)^{-1} ((i\nabla - k - k')^2 + 2\alpha^{-2} + M)^{-1} e^{-iH_{\varphi_t}(-k')(t-s)} \tilde{\psi}_s \right\rangle \frac{dk''}{|k''|^2} \frac{dk'}{|k'|^2} \frac{dk}{|k|^2}. \end{aligned}$$

By the Schwarz inequality we have $|X_2| \leq X_1$ and, similarly,

$$|X_j| \leq X_1 \quad \text{for all } j = 1, \dots, 6. \tag{4-9}$$

Thus it suffices to control X_1 .

We first perform the k'' integral and then the k integral. We make use of the following bounds.

Lemma 4.2. *One has the operator inequalities*

$$\int_{\mathbb{R}^3} ((i\nabla - k'')^2 + 1)^{-1} \frac{dk''}{|k''|^2} \lesssim 1, \tag{4-10}$$

$$\int_{\mathbb{R}^3} ((i\nabla_x - k)^2 + 1)^{-2} \frac{dk}{|k|^2} \lesssim (-\Delta + 1)^{-1}. \tag{4-11}$$

Before proving the lemma, let us see that they provide the desired bounds on X_1 . First, conjugating (4-10) with $e^{i(k+k') \cdot x}$ and assuming that $M + 3\alpha^2 \geq 1$, we obtain, uniformly in $k, k' \in \mathbb{R}^3$,

$$\int_{\mathbb{R}^3} ((i\nabla - k - k' - k'')^2 + 3\alpha^{-2} + M)^{-1} \frac{dk''}{|k''|^2} \lesssim 1. \tag{4-12}$$

Similarly, conjugating (4-11) with $e^{ik' \cdot x}$, we obtain, uniformly in $k' \in \mathbb{R}^3$,

$$\int_{\mathbb{R}^3} ((i\nabla_x - k - k')^2 + 2\alpha^{-2} + M)^{-2} \frac{dk}{|k|^2} \lesssim ((i\nabla - k')^2 + 1)^{-1}. \tag{4-13}$$

Inserting (4-12) and (4-13) into the definition of X_1 , we obtain

$$X_1 \lesssim \int_{\mathbb{R}^3} \left\langle e^{-iH_{\varphi_t}(-k')(t-s)} \tilde{\psi}_s, ((i\nabla - k')^2 + 1)^{-1} e^{-iH_{\varphi_t}(-k')(t-s)} \tilde{\psi}_s \right\rangle \frac{dk'}{|k'|^2}.$$

Since $(-\Delta + 1)^{-\frac{1}{2}}(H_{\varphi_t} + M)^{\frac{1}{2}}$ is bounded, uniformly in t (by [Corollary B.2](#) and [Lemma 2.1](#)), we also know that $((i\nabla - k')^2 + 1)^{-\frac{1}{2}}(H_{\varphi_t}(-k') + M)^{\frac{1}{2}}$ is bounded, uniformly in t . Thus,

$$\begin{aligned} X_1 &\lesssim \int_{\mathbb{R}^3} \left\langle e^{-iH_{\varphi_t}(-k')(t-s)} \tilde{\psi}_s, (H_{\varphi_t}(-k') + M)^{-1} e^{-iH_{\varphi_t}(-k')(t-s)} \tilde{\psi}_s \right\rangle \frac{dk'}{|k'|^2} \\ &= \int_{\mathbb{R}^3} \left\langle \tilde{\psi}_s, (H_{\varphi_t}(-k') + M)^{-1} \tilde{\psi}_s \right\rangle \frac{dk'}{|k'|^2} \\ &\lesssim \int_{\mathbb{R}^3} \left\langle \tilde{\psi}_s, ((i\nabla - k')^2 + M)^{-1} \tilde{\psi}_s \right\rangle \frac{dk'}{|k'|^2}. \end{aligned}$$

Applying [\(4-10\)](#) again, we see that the latter expression is bounded by a constant times $\|\tilde{\psi}_s\|_{\mathcal{L}^2}^2 = 1$ by [Lemma 2.1](#). This, together with [\(4-8\)](#) and [\(4-9\)](#), implies that

$$\|Q_{12}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3}. \quad (4-14)$$

Proof of [Lemma 4.2](#). We only prove [\(4-11\)](#), since the proof of [\(4-10\)](#) is similar and simpler. By applying a Fourier transform, we see that we need to prove

$$\int_{\mathbb{R}^3} ((p+k)^2 + 1)^{-2} \frac{dk}{|k|^2} \lesssim (p^2 + 1)^{-1} \quad \text{for } p \in \mathbb{R}^3.$$

We split the integral into the regions $4|k| > |p| + 1$ and $4|k| \leq |p| + 1$. In the first region we bound $|k|^{-2} \leq 16/(|p| + 1)^2$ and note that

$$\int_{\{4|k| > |p| + 1\}} ((p+k)^2 + 1)^{-2} dk \leq \int_{\mathbb{R}^3} ((p+k)^2 + 1)^{-2} dk = \int_{\mathbb{R}^3} (k^2 + 1)^{-2} dk < \infty.$$

In the second region we distinguish the cases $|p| < 1$ and $|p| \geq 1$. In the first case we bound

$$\int_{\{4|k| \leq |p| + 1\}} ((p+k)^2 + 1)^{-2} \frac{dk}{|k|^2} \leq \int_{\{4|k| \leq |p| + 1\}} \frac{dk}{|k|^2} \leq \int_{\{|k| \leq \frac{1}{2}\}} \frac{dk}{|k|^2} < \infty.$$

For $|p| \geq 1$ we note that in the second region we have $2|k| \leq |p|$ and therefore $(p+k)^2 \geq \frac{1}{4}p^2 \geq k^2$. Thus,

$$((p+k)^2 + 1)^{-2} \leq \left(\frac{1}{4}p^2 + 1\right)^{-1} (k^2 + 1)^{-1}.$$

Since $(k^2 + 1)^{-1}|k|^{-2}$ is integrable, we obtain again a bound of the required form. \square

Bounds on Q_2, \dots, Q_5 . The terms Q_2, \dots, Q_4 are controlled in exactly the same way as Q_1 . (For Q_4 we use the fact that $\|\partial_s \tilde{\psi}_s\|_{\mathcal{H}^2} \lesssim 1$ for $t \leq \alpha^2$ by [Proposition 2.2](#).) The argument for Q_5 is also similar. In fact, the term involving $\text{Im}(\varphi_s, \alpha^2 \partial_s \varphi_s)$ is controlled as before. For the term involving $b^*(\alpha^2 \partial_s \varphi_s)$ we have to prove a simple extension of [Lemma 4.1](#) where we have operators $b^*(f)b_k^*b_{k'}^*$, with $f \in \mathcal{L}^2$ (similarly as the second part in [Lemma 3.1](#)). Finally, the term involving $b(\alpha^2 \partial_s \varphi_s)$ can be commuted to the right and therefore becomes a less singular term which can be controlled already with [Lemma 3.1](#). These arguments prove [\(4-2\)](#) and complete the proof of [\(4-3\)](#).

Bound on D_{112} . The term D_{112} in (2-34) contains only one factor $|k'|^{-1}$ and can therefore be controlled essentially by the same method as D_{01} , based on Lemma 3.1. In order to create a factor of $(H_{\varphi_t} + M)^{-1}$, we integrate by parts in s_1 . This, however, will create a factor of \tilde{H}_{φ_t} in one of the terms. When dealing with D_{211} we will explain how to remove this term by integrating by parts in s . Since $\|g_{s,t}\|_\infty \lesssim \alpha^{-2}|t-s|$ and $\|\partial_s g_{s,t}\|_\infty = \|g_s\|_\infty \lesssim \alpha^{-2}$ by Proposition 2.2, this factor behaves well in the bounds. When applying Lemma 3.1 we also use $\|\partial_s \tilde{\psi}_s\|_{\mathcal{H}^1} \lesssim 1$ from Proposition 2.2; see also the remark at the beginning of Section 2B concerning the bounds on $\partial_t \tilde{\psi}_t$. Without going into details we state the final result,

$$\|D_{112}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3} t^2 (1+t). \tag{4-15}$$

Bound on D_{121} . Also the term D_{121} in (2-31) contains only one factor of $|k|^{-1}$ and can be controlled as just sketched for D_{112} and as explained in detail for D_{211} . In order to control the terms that appear when integrating by parts in s we make use of $\|\partial_s \sigma_{\tilde{\psi}_s}\|_{\mathcal{L}^2} \lesssim 1$ and $\|\partial_s \tilde{\psi}_s\|_{\mathcal{H}^1} \lesssim 1$ from Proposition 2.2 in addition to the bounds from Lemma 2.1. Moreover, we need an obvious extension of Lemma 3.1 to the case with $b^*(f_1)b^*(f_2)b_k^*$, which is proved in the same way. Combining all this, we end up with

$$\|D_{121}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t (1+t). \tag{4-16}$$

Bound on D_{122} . The term D_{122} contains no $|k|^{-1}$ term. Using $\|g_{s,t}\|_\infty \lesssim \alpha^{-2}|t-s|$ for $0 \leq s \leq t \leq \alpha^2$ by Proposition 2.2 and $\|b(\sigma_{\tilde{\psi}_s})\Omega\|_{\mathcal{F}} = \alpha^{-1}\|\sigma_{\tilde{\psi}_s}\|_2 \lesssim \alpha^{-1}$ by Lemma 2.1, we obtain immediately

$$\|D_{122}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3} t^3. \tag{4-17}$$

5. Estimation on D_2

Bound on D_{211} . We recall equation (2-32) for D_{211} . In this equation we commute $e^{-ik \cdot x}$ through $e^{-iH_{\varphi_t} s_1}$, which introduces again the operator $H_{\varphi_t}(k)$ from (4-1), and we commute b_k with b_k^* . In this way, we obtain

$$D_{211} = \alpha^{-2} \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} e^{-iH_{\varphi_t}(k)s_1} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|^2} ds_1 ds.$$

The difficulty in controlling D_{211} comes again from the k -integral. It is not enough to bound the norm of the integrand as it stands, since $|k|^{-2}$ is not integrable. Thus, we need to gain some extra decay from $e^{-iH_{\varphi_t}(k)s_1}$. To get this decay, we integrate by parts in s_1 using

$$e^{-iH_{\varphi_t}(k)s_1} = i e^{iMs_1} (H_{\varphi_t}(k) + M)^{-1} \partial_{s_1} e^{-i[H_{\varphi_t}(k)+M]s_1} \tag{5-1}$$

with a large constant $M > 0$ independent of α and t . We obtain

$$\begin{aligned} D_{211} &= i\alpha^{-2} \int_0^t \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t} t} (H_{\varphi_t}(k) + M)^{-1} e^{-iH_{\varphi_t}(k)(t-s)} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|^2} ds \\ &\quad - i\alpha^{-2} \int_0^t \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t} s} (H_{\varphi_t}(k) + M)^{-1} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|^2} ds \end{aligned}$$

$$\begin{aligned}
 & + \alpha^{-2} M \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (H_{\varphi_t}(k) + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} \\
 & \qquad \qquad \qquad \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|^2} ds_1 ds \\
 & + \alpha^{-2} \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} \tilde{H}_{\varphi_t} (H_{\varphi_t}(k) + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} \\
 & \qquad \qquad \qquad \times W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|^2} ds_1 ds \\
 & = D_{2111} + D_{2112} + D_{2113} + D_{2114},
 \end{aligned}$$

where $D_{211k}, k = 1, \dots, 4$, are naturally defined.

We first show how to deal with the terms D_{2111}, D_{2112} and D_{2113} . The term D_{2114} is harder because of the additional factor of \tilde{H}_{φ_t} .

The following lemma quantifies in which sense the operator $(H_{\varphi_t} + M)^{-1}$ leads to additional decay in k .

Lemma 5.1. For $u \in \mathcal{H}^2(\mathbb{R}^3)$,

$$\int_{\mathbb{R}^3} \|(i\nabla + k|^2 + 1)^{-1} u\|_2 \frac{dk}{|k|^2} \lesssim \|u\|_{\mathcal{H}^2}. \tag{5-2}$$

Proof. By Fourier transform, we have

$$\|(i\nabla + k|^2 + 1)^{-1} u\|_2^2 = \int_{\mathbb{R}^3} \frac{1}{(1 + |p + k|^2)^2 (1 + |p|^2)^2} (1 + |p|^2)^2 |\hat{u}(p)|^2 dp.$$

We now observe that

$$\frac{1}{(1 + |p + k|^2)^2 (1 + |p|^2)^2} \lesssim \frac{1}{(1 + |k|^2)^2}.$$

This can be proved by considering separately the regions where $|p| \leq \frac{1}{2}|k|$ and $|p| \geq \frac{1}{2}|k|$. Thus,

$$\|(i\nabla + k|^2 + 1)^{-1} u\|_2^2 \lesssim \frac{1}{(1 + |k|^2)^2} \|u\|_{\mathcal{H}^2}^2,$$

and the claimed bound follows by integration over k . □

Let us return to the terms D_{2111}, D_{2112} and D_{2113} . It follows from [Corollary B.2](#) by conjugating with the unitary $e^{ik \cdot x}$ that there is an $M > 0$ such that the operator $(H_{\varphi_t}(k) + M)^{-1} (|i\nabla + k|^2 + 1)$ is uniformly bounded in α and t . This, together with the boundedness of ψ_s in \mathcal{H}^2 for $s \in [0, \alpha^2]$ from [Proposition 2.2](#), yields

$$\int_{\mathbb{R}^3} \|(H_{\varphi_t}(k) + M)^{-1} \tilde{\psi}_s\|_2 \frac{dk}{|k|^2} \lesssim 1,$$

and therefore

$$\|D_{2111}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t, \quad \|D_{2112}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t, \quad \|D_{2113}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t^2. \tag{5-3}$$

We now turn to the term D_{2114} , which contains the operator \tilde{H}_{φ_t} . The idea is to remove this operator by integrating by parts in s using

$$\tilde{H}_{\varphi_t} e^{i\tilde{H}_{\varphi_t}s} = -i\partial_s e^{i\tilde{H}_{\varphi_t}s}. \tag{5-4}$$

This leads to

$$\begin{aligned} D_{2114} = & -i\alpha^{-2} \int_0^t \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}t} (H_{\varphi_t}(k) + M)^{-1} e^{-iH_{\varphi_t}(k)(t-s_1)} W^*(\alpha^2\varphi_t) W(\alpha^2\varphi_{s_1}) \tilde{\psi}_{s_1} \otimes \Omega \frac{dk}{|k|^2} ds_1 \\ & + i\alpha^{-2} \int_0^t \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}s_1} (H_{\varphi_t}(k) + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} W^*(\alpha^2\varphi_t) W(\alpha^2\varphi_0) \tilde{\psi}_0 \otimes \Omega \frac{dk}{|k|^2} ds_1 \\ & + i\alpha^{-2} \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (H_{\varphi_t}(k) + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} \\ & \qquad \qquad \qquad \times W^*(\alpha^2\varphi_t) W(\alpha^2\varphi_s) \partial_s \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|^2} ds_1 ds \\ & + i\alpha^{-2} \int_0^t \int_0^{t-s} \int_{\mathbb{R}^3} e^{i\tilde{H}_{\varphi_t}(s+s_1)} (H_{\varphi_t}(k) + M)^{-1} e^{-iH_{\varphi_t}(k)s_1} \\ & \qquad \qquad \qquad \times W^*(\alpha^2\varphi_t) (\partial_s W(\alpha^2\varphi_s)) \tilde{\psi}_s \otimes \Omega \frac{dk}{|k|^2} ds_1 ds. \end{aligned}$$

The first three terms on the right side can be bounded by [Lemma 5.1](#) together with the uniform boundedness in \mathcal{H}^2 of $\tilde{\psi}_s$ and $\partial_s \tilde{\psi}_s$ in $[0, \alpha^2]$ from [Proposition 2.2](#); see also the remark at the beginning of [Section 2B](#) concerning the bounds on $\partial_t \tilde{\psi}_t$. For the fourth term on the right side we use the formula [\(A-4\)](#) for $\partial_s W(\alpha^2\varphi_s)$. Then the term can be bounded by proceeding in the same way as for D_{015} and using [Lemma 5.1](#) together with the fact that $\alpha^2 \partial_s \varphi_s$ is uniformly bounded in \mathcal{L}^2 for all times by [Lemma 2.1](#). To summarize, we obtain

$$\|D_{2114}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t(1+t), \tag{5-5}$$

and, because of [\(5-3\)](#),

$$\|D_{211}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t(1+t). \tag{5-6}$$

Bound on D_{212} . The term D_{212} involves a single difficult operator $\int b_k^* e^{ik' \cdot x} |k'|^{-1} dk'$ and can be controlled using the technique from bounding D_{01} . We first integrate by parts with respect to s_1 using [\(5-1\)](#) (with $k = 0$) to create a factor of $(H_{\varphi_t} + M)^{-1}$. Using this factor we can apply [Lemma 3.1](#) as in the bound of D_{01} . In one of the terms, however, the integration by parts creates a factor \tilde{H}_{φ_t} . We remove this operator via [\(5-4\)](#) by integrating by parts in s . The factor $g_{s,t}$ and its derivative $\partial_s g_{s,t} = -g_s$ are bounded by [Proposition 2.2](#) and do not create any problems. Eventually, this shows that

$$\|D_{212}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3} t^2(1+t). \tag{5-7}$$

Bound on D_{221} . The term D_{221} appears in [\(2-33\)](#). We use $b_k b^*(\sigma_{\tilde{\psi}_s}) \Omega = \alpha^{-2} \sigma_{\tilde{\psi}_s}(k) \Omega$. By the Schwarz inequality, [\(C-2\)](#) and [Lemma 2.1](#) we have $\| |k|^{-1} \sigma_{\tilde{\psi}_s}(k) \Omega \|_1 \lesssim \| \sigma_{\tilde{\psi}_s} \|_{\mathcal{L}^2_{(1)}} \lesssim \| \psi_s \|_{H^1}^2 \lesssim 1$. From this one easily concludes that

$$\|D_{221}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t^2.$$

Bound on D_{222} . The term D_{222} appears in (2-37). Using the bound on $g_{s,t}$ from Proposition 2.2 and the fact that $b(\sigma_{\tilde{\psi}_s})\Omega$ has norm of order α^{-1} by Lemma 2.1, one obtains

$$\|D_{222}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-3} t^3.$$

6. Bounds on D_3 , D_4 and D_5

We recall that we have already controlled D_{32} , D_{42} and D_{52} in (2-39), (2-40) and (2-42). The remaining terms D_{31} , D_{41} and D_{51} have at most a single term $|k|^{-1}$ and can be bounded using the methods we have already developed. Therefore we will be rather brief.

For each of the terms D_{311} , D_{312} , D_{412} , D_{511} and D_{512} we first integrate by parts in s_1 to generate a factor of $(H_{\varphi_t} + M)^{-1}$, which allows us to apply Lemma 3.1. One of the terms, however, will involve \tilde{H}_{φ_t} , which we have to remove by integrating by parts in s . Using the bounds from Lemma 2.1 and Proposition 2.2 we obtain

$$\begin{aligned} \|D_{311}\|_{\mathcal{L}^2 \otimes \mathcal{F}} &\lesssim \alpha^{-2} t(1+t), & \|D_{312}\|_{\mathcal{L}^2 \otimes \mathcal{F}} &\lesssim \alpha^{-3} t^2(1+t), & \|D_{412}\|_{\mathcal{L}^2 \otimes \mathcal{F}} &\lesssim \alpha^{-3} t^2(1+t), \\ \|D_{511}\|_{\mathcal{L}^2 \otimes \mathcal{F}} &\lesssim \alpha^{-3} t(1+t), & \|D_{512}\|_{\mathcal{L}^2 \otimes \mathcal{F}} &\lesssim \alpha^{-4} t^2(1+t + \alpha^{-1} t^2). \end{aligned}$$

The remaining term D_{411} can be immediately bounded by

$$\|D_{411}\|_{\mathcal{L}^2 \otimes \mathcal{F}} \lesssim \alpha^{-2} t^2.$$

7. Proof of the almost orthogonality relations

7A. Proof of (2-28). We recall that

$$\langle \Omega, e^{-iH_{\varphi_t} t} D_0 \rangle_{\mathcal{F}} = \left\langle \Omega, \int_0^t e^{-iH_{\varphi_t}(t-s)} P_{\tilde{\psi}_s}^{\perp} \int_{\mathbb{R}^3} (e^{ik \cdot x} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} ds \right\rangle_{\mathcal{F}}.$$

We commute the operator b_k^* to the left and use $b_k \Omega = 0$. For the commutator we obtain from Corollary A.2 (with the definition (2-5) of $g_{s,t}$)

$$\begin{aligned} \langle \Omega, e^{-iH_{\varphi_t} t} D_0 \rangle_{\mathcal{F}} &= \left\langle \Omega, \int_0^t e^{-iH_{\varphi_t}(t-s)} P_{\tilde{\psi}_s}^{\perp} g_{s,t} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \tilde{\psi}_s \otimes \Omega ds \right\rangle_{\mathcal{F}} \\ &= \int_0^t e^{-iH_{\varphi_t}(t-s)} P_{\tilde{\psi}_s}^{\perp} g_{s,t} \tilde{\psi}_s \langle \Omega, W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) \Omega \rangle_{\mathcal{F}} ds. \end{aligned}$$

Thus,

$$\|\langle \Omega, e^{-iH_{\varphi_t} t} D_0 \rangle_{\mathcal{F}}\|_{\mathcal{L}^2} \leq t \sup_{0 \leq s \leq t} \|g_{s,t}\|_{\infty} \|\tilde{\psi}_s\|_2.$$

Thus, by the bound on $g_{s,t}$ from Proposition 2.2 and the conservation of the \mathcal{L}^2 norm of $\tilde{\psi}_s$, we obtain the claimed bound (2-28).

7B. Proof of (2-29). For $\Phi \in \mathcal{F}$, let

$$\begin{aligned} \Theta_\Phi(t) &:= \langle \tilde{\psi}_t \otimes \Phi, e^{-iH_{\varphi_t} t} D_0 \rangle_{\mathcal{L}^2 \otimes \mathcal{F}} \\ &= \left\langle \tilde{\psi}_t \otimes \Phi, \int_0^t e^{-iH_{\varphi_t}(t-s)} P_{\tilde{\psi}_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} ds \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}}. \end{aligned}$$

We shall show that

$$|\Theta_\Phi(t)| \lesssim \alpha^{-2} t^2 (1 + \alpha^{-2} t^2) \|\Phi\|_{\mathcal{F}}, \tag{7-1}$$

which by duality implies (2-29).

Our goal will be to derive an ordinary differential equation for Θ_Φ . We use the presence of the operator $P_{\tilde{\psi}_s}^\perp$ to obtain (with inner products in $\mathcal{L}^2 \otimes \mathcal{F}$)

$$\begin{aligned} \partial_t \Theta_\Phi &= \left\langle \partial_t \tilde{\psi}_t \otimes \Phi, \int_0^t e^{-iH_{\varphi_t}(t-s)} P_{\tilde{\psi}_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} ds \right\rangle \\ &\quad + \left\langle \tilde{\psi}_t \otimes \Phi, \int_0^t (\partial_t e^{-iH_{\varphi_t}(t-s)}) P_{\tilde{\psi}_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} ds \right\rangle \\ &\quad + \left\langle \tilde{\psi}_t \otimes \Phi, \int_0^t e^{-iH_{\varphi_t}(t-s)} P_{\tilde{\psi}_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} (\partial_t W^*(\alpha^2 \varphi_t)) W(\alpha^2 \varphi_s) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} ds \right\rangle. \end{aligned}$$

For the first term we use equation (2-10) for $\partial_t \tilde{\psi}_t$. In the second term, we compute, using Duhamel's formula,

$$\begin{aligned} \partial_t e^{-iH_{\varphi_t}(t-s)} &= -iH_{\varphi_t} e^{-iH_{\varphi_t}(t-s)} - i \int_0^{t-s} e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t H_{\varphi_t}) e^{-iH_{\varphi_t} s_1} ds_1 \\ &= -i(H_{\varphi_t} + (t-s)\partial_t \|\varphi_t\|_2^2) e^{-iH_{\varphi_t}(t-s)} - i \int_0^{t-s} e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t V_{\varphi_t}) e^{-iH_{\varphi_t} s_1} ds_1. \end{aligned}$$

Note that the part involving H_{φ_t} will cancel the contribution from the first term, except for part of the constant $\omega(t)$. Finally, for the third term we use Lemma A.3 and Lemma A.1 to obtain

$$\begin{aligned} \partial_t W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) &= \alpha^2 W^*(\alpha^2 \varphi_t) [b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t) + i \operatorname{Im}(\varphi_t, \partial_t \varphi_t)] W(\alpha^2 \varphi_s) \\ &= \alpha^2 W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) [b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t) + 2i \operatorname{Im}(\partial_t \varphi_t, \varphi_s) + i \operatorname{Im}(\varphi_t, \partial_t \varphi_t)] \\ &= \alpha^2 W^*(\alpha^2 \varphi_t) W(\alpha^2 \varphi_s) [b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t) + 2i \operatorname{Im}(\partial_t \varphi_t, \varphi_s - \varphi_t) + i \operatorname{Im}(\partial_t \varphi_t, \varphi_t)]. \end{aligned}$$

Putting all this into the above formula, we obtain

$$\partial_t \Theta_\Phi = M_1 + M_2 + M_3,$$

where the terms M_1 , M_2 and M_3 are defined, using the notation

$$\Phi_{s,t} := W^*(\alpha^2 \varphi_s) W(\alpha^2 \varphi_t) \Phi,$$

by

$$\begin{aligned}
 M_1(t) &:= -i \int_0^t \int_0^{t-s} \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t V_{\varphi_t}) e^{-iH_{\varphi_t} s_1} P_{\tilde{\psi}_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle ds_1 ds, \\
 M_2(t) &:= \alpha^2 \int_0^t \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s)} P_{\tilde{\psi}_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} (b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t)) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle ds, \\
 M_3(t) &:= \int_0^t m(s,t) \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s)} P_{\tilde{\psi}_s}^\perp \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle ds
 \end{aligned}$$

with

$$m(s,t) := -i(t-s) \partial_t \|\varphi_t\|_2^2 + 2i\alpha^2 \operatorname{Im}(\partial_t \varphi_t, \varphi_s - \varphi_t).$$

Since $\Theta_\Phi(0) = 0$, we conclude that

$$\Theta_\Phi(t) = \int_0^t (M_1(s) + M_2(s) + M_3(s)) ds. \quad (7-2)$$

Below we shall show that

$$|M_1(t)| \lesssim \alpha^{-3} t^2 \|\Phi\|_{\mathcal{F}}, \quad |M_2(t)| \lesssim \alpha^{-2} t \|\Phi\|_{\mathcal{F}}, \quad |M_3(t)| \lesssim \alpha^{-3} t^2 \|\Phi\|_{\mathcal{F}}. \quad (7-3)$$

Together with (7-2) this will prove (7-1) and therefore (2-29).

Bound on M_1 . Using the fact that $P_{\tilde{\psi}_s}^\perp = 1 - |\tilde{\psi}_s\rangle\langle\tilde{\psi}_s|$ (see the proof of Lemma 2.4), we have the decomposition

$$M_1 = M_{11} - M_{12},$$

where

$$M_{11}(t) := -i \int_0^t \int_0^{t-s} \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t V_{\varphi_t}) e^{-iH_{\varphi_t} s_1} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} ds_1 ds$$

and, with $\sigma_{\tilde{\psi}_s}$ from (2-2),

$$M_{12}(t) := -i \int_0^t \int_0^{t-s} \left\langle \tilde{\psi}_t, e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t V_{\varphi_t}) e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s \right\rangle_{\mathcal{L}^2} \left\langle \Phi_{s,t}, b^*(\sigma_{\tilde{\psi}_s}) \Omega \right\rangle_{\mathcal{F}} ds_1 ds.$$

The second term is easy to control. In fact, the a priori bounds from Lemma 2.1 together with $\|\partial_t V_{\varphi_t}\|_\infty \lesssim \alpha^{-2}$ from (C-8) imply

$$\left| \left\langle \tilde{\psi}_t, e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t V_{\varphi_t}) e^{-iH_{\varphi_t} s_1} \tilde{\psi}_s \right\rangle_{\mathcal{L}^2} \right| \lesssim \alpha^{-2}$$

and

$$\left| \left\langle \Phi_{s,t}, b^*(\sigma_{\tilde{\psi}_s}) \Omega \right\rangle_{\mathcal{F}} \right| \lesssim \alpha^{-1} \|\Phi\|_{\mathcal{F}}.$$

This yields a bound of the form (7-3).

We now bound the integrand in M_{11} . We have

$$\begin{aligned} & \left| \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t V_{\varphi_t}) e^{-iH_{\varphi_t}s_1} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} \right| \\ & \leq \| (H_{\varphi_t} + M)^{\frac{1}{2}} \tilde{\psi}_t \otimes \Phi_{s,t} \| \| (H_{\varphi_t} + M)^{-\frac{1}{2}} (\partial_t V_{\varphi_t}) (H_{\varphi_t} + M)^{\frac{1}{2}} \| \left\| (H_{\varphi_t} + M)^{-\frac{1}{2}} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\|. \end{aligned}$$

By [Corollary B.2](#) and an easy modification of its proof, for M sufficiently large (but independent of t and α), the operators $(H_{\varphi_t} + M)^{\pm \frac{1}{2}} (-\Delta + 1)^{\mp \frac{1}{2}}$ are both bounded uniformly in t . Therefore [Lemma 3.1](#) and the a priori bounds from [Lemma 2.1](#) yield

$$\begin{aligned} & \left| \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t V_{\varphi_t}) e^{-iH_{\varphi_t}s_1} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} \right| \\ & \lesssim \alpha^{-1} \|\tilde{\psi}_t\|_{\mathcal{H}^1} \|\Phi\|_{\mathcal{F}} \|(-\Delta + 1)^{-\frac{1}{2}} (\partial_t V_{\varphi_t}) (-\Delta + 1)^{\frac{1}{2}}\| \|\psi_s\|_{\mathcal{H}^1} \\ & \lesssim \alpha^{-1} \|\Phi\|_{\mathcal{F}} \|(-\Delta + 1)^{-\frac{1}{2}} (\partial_t V_{\varphi_t}) (-\Delta + 1)^{\frac{1}{2}}\|. \end{aligned}$$

Finally, using the fact that $\|\nabla \partial_t V_{\varphi_t}\|_{\infty} \lesssim \alpha^{-2}$ (see [\(C-8\)](#)), we obtain that the operator appearing in this bound has norm $\lesssim \alpha^{-2}$. Thus, we finally obtain

$$\left| \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s-s_1)} (\partial_t V_{\varphi_t}) e^{-iH_{\varphi_t}s_1} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} \right| \lesssim \alpha^{-3},$$

which, when integrated over s_1 and s , leads to the bound in [\(7-3\)](#).

Bound on M_2 . As for M_1 , we use $P_{\tilde{\psi}_s}^{\perp} = 1 - |\tilde{\psi}_s\rangle\langle\tilde{\psi}_s|$ to get the decomposition

$$M_2 = M_{21} - M_{22}$$

with

$$M_{21}(t) := \alpha^2 \int_0^t \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s)} \int_{\mathbb{R}^3} (e^{ik \cdot x} (b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t)) b_k^* \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle ds$$

and, with $\sigma_{\tilde{\psi}_s}$ from [\(2-2\)](#),

$$M_{22}(t) := \alpha^2 \int_0^t \langle \tilde{\psi}_t, e^{-iH_{\varphi_t}(t-s)} \tilde{\psi}_s \rangle_{\mathcal{L}^2} \langle \Phi_{s,t}, (b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t)) b^*(\sigma_{\tilde{\psi}_s}) \Omega \rangle_{\mathcal{F}} ds.$$

Once again the bound on M_{22} is straightforward. Namely, we commute $b^*(\sigma_{\tilde{\psi}_s})$ to the left through $b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t)$ and obtain

$$\langle \Phi_{s,t}, (b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t)) b^*(\sigma_{\tilde{\psi}_s}) \Omega \rangle_{\mathcal{F}} = -\langle \Phi_{s,t}, b^*(\sigma_{\tilde{\psi}_s}) b^*(\partial_t \varphi_t) \Omega \rangle_{\mathcal{F}} + \alpha^{-2} \langle \partial_t \varphi_t, \sigma_{\tilde{\psi}_s} \rangle \langle \Phi_{s,t}, \Omega \rangle_{\mathcal{F}}.$$

By similar computations as, for instance, in the bound on D_{32} and by the a priori bounds from [Lemma 2.1](#), we obtain

$$\left| \langle \Phi_{s,t}, (b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t)) b^*(\sigma_{\tilde{\psi}_s}) \Omega \rangle_{\mathcal{F}} \right| \lesssim \alpha^{-2} \|\Phi\|_{\mathcal{F}} \|\sigma_{\tilde{\psi}_s}\| \|\partial_t \varphi_t\| \lesssim \alpha^{-4} \|\Phi\|_{\mathcal{F}}.$$

By the conservation of the \mathcal{L}^2 norm of $\tilde{\psi}_t$ we conclude

$$|M_{22}(t)| \lesssim \alpha^{-2} t \|\Phi\|_{\mathcal{F}},$$

which is of the form claimed in (7-3).

We now discuss M_{21} . Again we commute b_k^* to the left through $b(\partial_t \varphi_t) - b^*(\partial_t \varphi_t)$ and obtain

$$M_{21} = M_{211} + M_{212},$$

where

$$M_{211}(t) := -\alpha^2 \int_0^t \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s)} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* b^*(\partial_t \varphi_t) \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} ds$$

and, with g_s from (2-6),

$$M_{212}(t) := \int_0^t \langle \tilde{\psi}_t, e^{-iH_{\varphi_t}(t-s)} g_s \tilde{\psi}_s \rangle_{\mathcal{L}^2} \langle \Phi_{s,t}, \Omega \rangle_{\mathcal{F}} ds.$$

Since $\|g_s\|_{\infty} \lesssim \alpha^{-2}$ by Proposition 2.2, we obtain immediately

$$|M_{212}(t)| \lesssim \alpha^{-2} t \|\Phi\|_{\mathcal{F}}.$$

To control M_{211} we bound

$$\begin{aligned} & \left| \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s)} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* b^*(\partial_t \varphi_t) \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} \right| \\ & \leq \| (H_{\varphi_t} + M)^{\frac{1}{2}} \tilde{\psi}_t \otimes \Phi_{s,t} \| \left\| (H_{\varphi_t} + M)^{-\frac{1}{2}} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* b^*(\partial_t \varphi_t) \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\|. \end{aligned}$$

As for M_{11} , we use Lemma 2.1 and Corollary B.2 (and a simple extension of its proof) to choose M large enough, but independent of t and α , so that $(H_{\varphi_t} + M)^{\pm \frac{1}{2}} (-\Delta + 1)^{\mp \frac{1}{2}}$ are both bounded uniformly in t . Therefore Lemma 3.1 and the a priori bounds from Lemma 2.1 yield

$$\begin{aligned} & \left| \left\langle \tilde{\psi}_t \otimes \Phi_{s,t}, e^{-iH_{\varphi_t}(t-s)} \int_{\mathbb{R}^3} (e^{ik \cdot x} b_k^* b^*(\partial_t \varphi_t) \tilde{\psi}_s \otimes \Omega) \frac{dk}{|k|} \right\rangle_{\mathcal{L}^2 \otimes \mathcal{F}} \right| \lesssim \alpha^{-2} \|\tilde{\psi}_t\|_{\mathcal{H}^1} \|\Phi\|_{\mathcal{F}} \|\partial_t \varphi_t\|_{\mathcal{L}^2} \|\tilde{\psi}_s\|_{\mathcal{H}^1} \\ & \lesssim \alpha^{-4} \|\Phi\|_{\mathcal{F}}. \end{aligned}$$

This, when integrated over s and multiplied by α^2 , leads to the bound in (7-3).

Bound on M_3 . The a priori bounds from Lemma 2.1 yield

$$|m(s, t)| \lesssim \alpha^{-2} |t - s|.$$

Moreover, applying Lemma 3.1 as in the bound on M_{21} we find that the absolute value of the inner product in the integral defining M_3 is bounded by a constant times $\alpha^{-1} \|\Phi\|_{\mathcal{F}}$. This yields the bound in (7-3).

This concludes the proof of (2-29).

Appendix A: Some properties of the Weyl operators

In this appendix we collect some standard properties of the Weyl operators $W(f)$ defined in (1-7) in terms of $b(f)$ and $b^*(f)$. They are well known, but we provide proofs for the sake of completeness. We recall that the commutation relations for b_k and b_k^* involve a factor α^{-2} .

Lemma A.1. *The operators b_k , b_k^* and $W(f)$ satisfy the following relations,*

$$b_k W(f) = W(f)(b_k + \alpha^{-2} f(k)) \quad \text{and} \quad b_k^* W(f) = W(f)(b_k^* + \alpha^{-2} \bar{f}(k)). \quad (\text{A-1})$$

Proof. For $t > 0$ we consider the operators

$$F_t := W(tf) = e^{t(b^*(f)-b(f))}, \quad (\text{A-2})$$

which satisfy

$$\partial_t F_t = (b^*(f) - b(f))F_t, \quad F_0 = \text{Id}.$$

Multiplying by b_k and using the commutation relations, we obtain the following equation for $b_k F_t$:

$$\partial_t b_k F_t = (b^*(f) - b(f))b_k F_t + \alpha^{-2} f(k)F_t, \quad b_k F_0 = b_k.$$

Therefore, by Duhamel's principle applied to the latter equation,

$$b_k F_t = e^{t(b^*(f)-b(f))}b_k + \alpha^{-2} f(k) \int_0^t e^{(t-s)(b^*(f)-b(f))} F_s ds.$$

Recalling the definition of F_t in (A-2), we can rewrite this as

$$b_k F_t = F_t b_k + t\alpha^{-2} f(k)F_t. \quad (\text{A-3})$$

At $t = 1$ we obtain the first identity in the lemma. The second one is proved similarly. \square

By applying Lemma A.1 twice, we obtain:

Corollary A.2.
$$\begin{aligned} [b_k^*, W^*(f)W(g)] &= -\alpha^{-2}(\bar{f}(k) - \bar{g}(k))W^*(f)W(g), \\ [b_k, W^*(f)W(g)] &= -\alpha^{-2}(f(k) - g(k))W^*(f)W(g). \end{aligned}$$

Next, we'll consider the case where f depends (differentiably) on a parameter.

Lemma A.3.

$$\partial_t W(f_t) = \frac{1}{2}\alpha^{-2}((f_t, \partial_t f_t) - (\partial_t f_t, f_t))W(f_t) + W(f_t)(b^*(\partial_t f_t) - b(\partial_t f_t)), \quad (\text{A-4})$$

$$\partial_t W(f_t) = -\frac{1}{2}\alpha^{-2}((f_t, \partial_t f_t) - (\partial_t f_t, f_t))W(f_t) + (b^*(\partial_t f_t) - b(\partial_t f_t))W(f_t). \quad (\text{A-5})$$

Proof. For $s > 0$ we consider the operators

$$F(s, t) := W(sf_t), \quad (\text{A-6})$$

which satisfy

$$\partial_s F(s, t) = (b^*(f_t) - b(f_t))F(s, t), \quad F(0, t) = \text{Id}.$$

We differentiate this equation with respect to t and obtain

$$\begin{aligned}\partial_s \partial_t F(s, t) &= (b^*(f_t) - b(f_t)) \partial_t F(s, t) + (b^*(\partial_t f_t) - b(\partial_t f_t)) F(s, t), \\ \partial_t F(0, t) &= 0.\end{aligned}$$

Therefore, by Duhamel's principle,

$$\begin{aligned}\partial_t F(s, t) &= \int_0^s e^{(b^*(f_t) - b(f_t))(s-s_1)} (b^*(\partial_t f_t) - b(\partial_t f_t)) F(s_1, t) ds_1 \\ &= \int_0^s W((s-s_1)f_t) (b^*(\partial_t f_t) - b(\partial_t f_t)) W(s_1 f_t) ds_1.\end{aligned}$$

In order to simplify the integrand we now use [Lemma A.1](#) and obtain

$$(b^*(\partial_t f_t) - b(\partial_t f_t)) W(s_1 f_t) = \alpha^{-2} W(s_1 f_t) s_1 ((f_t, \partial_t f_t) - (\partial_t f_t, f_t)) + W(s_1 f_t) (b^*(\partial_t f_t) - b(\partial_t f_t)).$$

If we insert this into the above formula for $\partial_t F(s, t)$, we obtain

$$\partial_t F(s, t) = \alpha^{-2} \frac{1}{2} s^2 W(s f_t) ((f_t, \partial_t f_t) - (\partial_t f_t, f_t)) + s W(s f_t) (b^*(\partial_t f_t) - b(\partial_t f_t)).$$

At $s = 1$, we obtain the first identity in the lemma. The second one is proved similarly. \square

Lemma A.4. For any $f, g \in \mathcal{L}^2$,

$$\langle \Omega, W^*(g) W(f) \Omega \rangle = e^{i\alpha^{-2} \text{Im}(g, f) - \alpha^{-2} \|f - g\|^2 / 2}.$$

Proof. Let $f_t := tf + (1-t)g$ and $F(t) := \langle \Omega, W^*(g) W(f_t) \Omega \rangle$. By [Lemma A.3](#), using that $\text{Im}(f_t, \partial_t f_t) = \text{Im}(f_t, f - g) = \text{Im}(g, f)$,

$$\partial_t F(t) = \langle \Omega, W^*(g) W(f_t) (b^*(f - g) + i\alpha^{-2} \text{Im}(g, f)) \Omega \rangle.$$

Next, by [Corollary A.2](#), since $(g - f_t, f - g) = -t\|f - g\|^2$,

$$W^*(g) W(f_t) b^*(f - g) = b^*(f - g) W^*(g) W(f_t) + \alpha^{-2} (g - f_t, f - g) W^*(g) W(f_t),$$

so

$$\partial_t F(t) = (-\alpha^{-2} t \|f - g\|^2 + i\alpha^{-2} \text{Im}(g, f)) F(t).$$

Since $F(0) = 1$, we conclude that

$$F(t) = e^{-\alpha^{-2} t^2 \|f - g\|^2 / 2 + i\alpha^{-2} t \text{Im}(g, f)},$$

which, at $t = 1$, gives the assertion. \square

Appendix B: The effective Schrödinger operator

In this appendix we investigate the operator and form domains of the effective Schrödinger operator H_φ from (2-12) with potential V_φ from (2-1).

Lemma B.1. *For every $A > 0$ and $\varepsilon > 0$ there is an $M > 0$ such that if $\|\varphi\| \leq A$, then for all $\psi \in \mathcal{H}^1(\mathbb{R}^3)$,*

$$\| |V_\varphi|^{\frac{1}{2}} \psi \| \leq \varepsilon \| (-\Delta + M)^{\frac{1}{2}} \psi \|$$

and for all $\psi \in \mathcal{H}^2(\mathbb{R}^3)$,

$$\| V_\varphi \psi \| \leq \varepsilon \| (-\Delta + M) \psi \|.$$

Proof. As in [Frank and Schlein 2014, Section 2.1], the Hardy–Littlewood–Sobolev inequality implies

$$\| V_\varphi \|_6 \lesssim \| \varphi \|_2. \tag{B-1}$$

This implies, by the Hölder and Sobolev inequalities,

$$\int_{\mathbb{R}^3} |V_\varphi| |\psi|^2 dx \leq \| V_\varphi \|_6 \| \psi \|_{\frac{12}{5}}^2 \lesssim \| \varphi \|_2 \| \nabla \psi \|_2^{\frac{1}{2}} \| \psi \|_2^{\frac{3}{2}}$$

and

$$\int_{\mathbb{R}^3} |V_\varphi|^2 |\psi|^2 dx \leq \| V_\varphi \|_6^2 \| \psi \|_3^2 \lesssim \| \varphi \|_2^2 \| \Delta \psi \|_2^{\frac{1}{2}} \| \psi \|_2^{\frac{3}{2}}.$$

These bounds easily imply the assertions of the lemma. □

Corollary B.2. *For every $A > 0$ there are $M > 0$ and $C > 0$ such that if $\|\varphi\|_2 \leq A$ then for all $f \in \mathcal{L}^2(\mathbb{R}^3)$*

$$\| (H_\varphi + M)^{-\frac{1}{2}} f \|_2 \leq C \| (-\Delta + 1)^{-\frac{1}{2}} f \|_2$$

and

$$\| (H_\varphi + M)^{-1} f \|_2 \leq C \| (-\Delta + 1)^{-1} f \|_2.$$

Proof. To prove the first assertion, we write

$$(H_\varphi + M)^{-1} = (-\Delta + M)^{-\frac{1}{2}} (1 + (-\Delta + M)^{-\frac{1}{2}} V_\varphi (-\Delta + M)^{-\frac{1}{2}})^{-1} (-\Delta + M)^{-\frac{1}{2}}$$

and note that according to Lemma B.1 we can choose M such that $\|\varphi\| \leq A$ implies

$$\| (-\Delta + M)^{-\frac{1}{2}} V_\varphi (-\Delta + M)^{-\frac{1}{2}} \| \leq \varepsilon^2.$$

Similarly, for the second assertion we write

$$(H_\varphi + M)^{-1} = (1 + (-\Delta + M)^{-1} V_\varphi)^{-1} (-\Delta + M)^{-1}$$

and choose M such that $\|\varphi\| \leq A$ implies $\| (-\Delta + M)^{-1} V_\varphi \| \leq \varepsilon$. □

Appendix C: Well-posedness of the Landau–Pekar equations

In this appendix we prove Lemma 2.1 and Proposition 2.2. Recall that the weighted spaces $\mathcal{L}^2_{(m)} = \mathcal{L}^2(\mathbb{R}^3; (1 + k^2)^m dk)$ were introduced in (1-11). We begin with some bounds on the coupling terms V_φ and σ_ψ introduced in (2-1) and (2-2).

Lemma C.1. *We have*

$$\|\partial^\beta V_\varphi\|_\infty \lesssim \|\varphi\|_{\mathcal{L}^2_{|\beta|+1}} \quad \text{for all } \beta \in \mathbb{N}_0^3, \quad (\text{C-1})$$

$$\|\sigma_\psi\|_{\mathcal{L}^2_{(1)}} \lesssim \|\psi\|_{\mathcal{H}^1}^2, \quad \|\sigma_\psi\|_{\mathcal{L}^2_{(3)}} \lesssim \|\psi\|_{\mathcal{H}^2}^2. \quad (\text{C-2})$$

Proof. By the Schwarz inequality,

$$|\partial^\beta V_\varphi(x)| \leq 2 \int_{\mathbb{R}^3} |k|^{|\beta|-1} |\varphi(k)| dk \leq 2 \|\varphi\|_{\mathcal{L}^2_{|\beta|+1}} \left(\int_{\mathbb{R}^3} \frac{|k|^{2(|\beta|-1)} dk}{(1+k^2)^{2(|\beta|+1)}} \right)^{\frac{1}{2}}$$

and the last integral is finite.

We have

$$\|\sigma_\psi\|_2^2 = \left\| \frac{1}{|k|} \int_{\mathbb{R}^3} |\psi(x)|^2 e^{ik \cdot x} dx \right\|_2^2 = 2\pi^2 \iint_{\mathbb{R}^3 \times \mathbb{R}^3} \frac{|\psi(x)|^2 |\psi(y)|^2}{|x-y|} dx dy.$$

By the Hardy–Littlewood–Sobolev inequality, we know this is bounded by a constant times $\|\psi\|_{\frac{5}{3}}^2 = \|\psi\|_{\frac{12}{5}}^4$, which, by the Sobolev embedding theorem, is bounded by a constant times $\|\psi\|_{\mathcal{H}^1}^4$. Moreover, by Plancherel,

$$\|\sigma_\psi\|_{\mathcal{L}^2(|k|^{2m})}^2 = \int_{\mathbb{R}^3} |k|^{2(m-1)} \left| \int_{\mathbb{R}^3} |\psi|^2 e^{ik \cdot x} dx \right|^2 dk = (2\pi)^3 (|\psi|^2, (-\Delta)^{m-1} |\psi|^2).$$

In particular, for $m = 1$ we get $\|\psi\|_{\mathcal{H}^1}^4$, which by Sobolev is controlled by $\|\psi\|_{\mathcal{H}^1}^2$. For $m = 3$, the claimed bound follows easily using $\|\psi\|_\infty \lesssim \|\psi\|_{\mathcal{H}^2}$ and again Sobolev. \square

Proof of Lemma 2.1. Local well-posedness in $\mathcal{H}^1 \times \mathcal{L}^2$ follows by a standard fixed-point argument and one sees that $\|\psi_t\|_2$ and $\mathcal{E}(\psi_t, \varphi_t)$ are conserved. One can use (B-1) and the Sobolev inequality to show that [Frank and Schlein 2014, Section 2.1],

$$\mathcal{E}(\psi, \varphi) \geq \|\nabla \psi\|_2^2 + \|\varphi\|_2^2 - C \|\varphi\|_2 \|\nabla \psi\|_2^{\frac{1}{2}} \|\psi\|_2^{\frac{3}{2}} \quad (\text{C-3})$$

for some universal constant $C > 0$. This, together with conservation of $\mathcal{E}(\psi_t, \varphi_t)$, yields global well-posedness as well as the uniform bounds (2-3).

According to (C-2) and the first bound in (2-3), we have $\|\sigma_{\psi_t}\| \lesssim \|\psi_t\|_{\mathcal{H}^1}^2 \lesssim 1$, which is the third bound in (2-4).

By equation (1-9) for φ_t we have

$$\|\alpha^2 \partial_t \varphi_t\|_2 \leq \|\varphi_t\|_2 + \|\sigma_{\psi_t}\|_2$$

and therefore, by the second bound in (2-3) and the third bound in (2-4), we obtain the first bound in (2-4).

Finally, $\varphi_t - \varphi_s = \int_s^t \partial_{s_1} \varphi_{s_1} ds_1$, so for $t > s$, by the first bound in (2-4),

$$\|\varphi_t - \varphi_s\|_2 \leq \int_s^t \|\partial_{s_1} \varphi_{s_1}\|_2 ds_1 \lesssim \alpha^{-2} |t - s|.$$

This proves the second bound in (2-4) and completes the proof of the lemma. \square

Before dealing with $\mathcal{H}^4 \times \mathcal{L}^2_{(3)}$ -regularity in Proposition 2.2, we need to establish $\mathcal{H}^2 \times \mathcal{L}^2_{(1)}$ -regularity.

Lemma C.2. *If $(\psi_0, \varphi_0) \in \mathcal{H}^2(\mathbb{R}^3) \times \mathcal{L}_{(1)}^2(\mathbb{R}^3)$, then $(\psi_t, \varphi_t) \in \mathcal{H}^2(\mathbb{R}^3) \times \mathcal{L}_{(1)}^2(\mathbb{R}^3)$ for all $t \in \mathbb{R}$ and*

$$\|\psi_t\|_{\mathcal{H}^2} \lesssim 1 + \alpha^{-2}|t|, \quad \|\varphi_t\|_{\mathcal{L}_{(1)}^2(\mathbb{R}^3)} \lesssim 1 + \alpha^{-2}|t|$$

with implicit constants depending only on the initial data. Moreover,

$$\|\partial_t \psi_t\|_{\mathcal{L}^2} \lesssim 1 + \alpha^{-2}|t|, \quad \|\partial_t \sigma \psi_t\|_{\mathcal{L}^2} \lesssim 1 + \alpha^{-2}|t|. \quad (\text{C-4})$$

If, in addition, $\varphi_0 \in \mathcal{L}_{(m)}^2(\mathbb{R}^3)$, $m = 2, 3$, then $\varphi_t \in \mathcal{L}_{(m)}^2(\mathbb{R}^3)$ for all $t \in \mathbb{R}$ and

$$\|\varphi_t\|_{\mathcal{L}_{(m)}^2(\mathbb{R}^3)} \lesssim 1 + \alpha^{-6}|t|^3.$$

Proof. By a standard fixed-point argument one can show local existence of solutions in $\mathcal{H}^2 \times \mathcal{L}_{(1)}^2$. In the following we will construct a functional, which is equivalent to the \mathcal{H}^2 norm of ψ and which grows in a controlled way as time increases. This will prove, in particular, that ψ_t belongs to \mathcal{H}^2 for all times.

We claim that for every $A > 0$ there is a constant $M > 0$ such that

$$\mathcal{E}^{(2)}(\psi, \varphi) := \|(-\Delta + V_\varphi + M)\psi\|_2^2$$

satisfies

$$\frac{1}{2}\|\psi\|_{\mathcal{H}^2} \leq (\mathcal{E}^{(2)}(\psi, \varphi))^{\frac{1}{2}} \leq \frac{3}{2}\|\psi\|_{\mathcal{H}^2} \quad (\text{C-5})$$

for all $\psi \in \mathcal{H}^2$ and all φ satisfying $\|\varphi\|_2 \leq A$. In fact, much as in the proof of [Corollary B.2](#), we have

$$\|(-\Delta + V_\varphi + M)\psi\|_2 - \|(-\Delta + M)\psi\|_2 \leq \|V_\varphi(-\Delta + M)^{-1}\| \|(-\Delta + M)\psi\|_2$$

and according to [Lemma B.1](#) we can choose M such that the first factor on the right side is less than ε for $\|\varphi\|_2 \leq A$.

According to [Lemma 2.1](#) there is an $A > 0$ (depending only on $\|\psi_0\|_{\mathcal{H}^1}$ and $\|\varphi_0\|_{\mathcal{L}^2}$) such that $\|\varphi_t\|_{\mathcal{L}^2} \leq A$ for all t . We choose M corresponding to this value of A and compute, using the equation for ψ_t ,

$$\begin{aligned} \partial_t \mathcal{E}^{(2)}(\psi_t, \varphi_t) &= 2 \operatorname{Re}((-\Delta + V_{\varphi_t} + M)\psi_t, (-\Delta + V_{\varphi_t} + M)\partial_t \psi_t) + 2 \operatorname{Re}((-\Delta + V_{\varphi_t} + M)\psi_t, (\partial_t V_{\varphi_t})\psi_t) \\ &= 2 \operatorname{Re}((-\Delta + V_{\varphi_t} + M)\psi_t, (\partial_t V_{\varphi_t})\psi_t). \end{aligned}$$

By the Schwarz and the Hölder inequalities,

$$\partial_t \mathcal{E}^{(2)}(\psi_t, \varphi_t) \leq 2(\mathcal{E}^{(2)}(\psi_t, \varphi_t))^{\frac{1}{2}} \|\partial_t V_{\varphi_t}\|_6 \|\psi_t\|_3.$$

By [\(B-1\)](#) and [Lemma 2.1](#), $\|\partial_t V_{\varphi_t}\|_6 \lesssim \|\partial_t \varphi_t\|_2 \lesssim \alpha^{-2}$, and by the Sobolev inequality and [Lemma 2.1](#), $\|\psi_t\|_3 \lesssim \|\psi_t\|_{\mathcal{H}^1} \lesssim 1$. Thus,

$$\partial_t \mathcal{E}^{(2)}(\psi_t, \varphi_t) \lesssim \alpha^{-2}(\mathcal{E}^{(2)}(\psi_t, \varphi_t))^{\frac{1}{2}},$$

which implies $(\mathcal{E}^{(2)}(\psi_t, \varphi_t))^{\frac{1}{2}} \lesssim 1 + \alpha^{-2}|t|$. According to [\(C-5\)](#), this implies the claimed bound on $\|\psi_t\|_{\mathcal{H}^2}$.

The remaining bounds are proved in a straightforward way. We have

$$\|\partial_t \psi_t\|_2 \leq \|-\Delta \psi_t\|_2 + \|V_{\varphi_t} \psi_t\|_2 \leq \|\psi_t\|_{\mathcal{H}^2} + \|V_{\varphi_t}\|_6 \|\psi_t\|_3.$$

By the bound on $\|\psi_t\|_{\mathcal{H}^2}$ together with (B-1) and the bounds from Lemma 2.1, we obtain the first bound in (C-4). Moreover,

$$\partial_t \sigma_{\psi_t} = 2|k|^{-1} \int_{\mathbb{R}^3} \operatorname{Re}(\overline{\psi_t} \partial_t \psi_t) e^{ik \cdot x} dx$$

and so, by the Hardy–Littlewood–Sobolev inequality as in (B-1),

$$\|\partial_t \sigma_{\psi_t}\|_2 \lesssim \|\psi_t \partial_t \psi_t\|_{\frac{6}{5}} \leq \|\psi_t\|_3 \|\partial_t \psi_t\|_2.$$

By the first bound in (C-4) and Lemma 2.1, we obtain the second bound in (C-4).

In order to deduce the bounds on φ_t , we use Duhamel’s formula:

$$\varphi_t(k) = e^{-it/\alpha^2} \varphi_0(k) - i\alpha^{-2} \int_0^t e^{-i(t-s)/\alpha^2} \sigma_{\psi_s}(k) ds. \quad (\text{C-6})$$

If $\varphi_0 \in \mathcal{L}_{(m)}^2$, $m = 1, 2, 3$, we deduce that $\varphi_t \in \mathcal{L}_{(m)}^2$ provided we can bound $\|\sigma_{\psi_s}\|_{\mathcal{L}_{(m)}^2}$. This quantity can be controlled by Sobolev norms of ψ_s according to (C-2). \square

Proof of Proposition 2.2. The basic strategy is the same as in the proof of Lemma C.2, except that verifying the properties of the functional is more complicated in this case. Again we do not give the details of the local existence via a fixed-point argument.

We claim that for every $A > 0$ there is a constant $M > 0$ such that

$$\mathcal{E}^{(4)}(\psi, \varphi) := \|(-\Delta + V_\varphi + M)^2 \psi\|_2^2$$

satisfies

$$\frac{1}{2} \|\psi\|_{\mathcal{H}^4} \leq (\mathcal{E}^{(4)}(\psi, \varphi))^{\frac{1}{2}} \leq \frac{3}{2} \|\psi\|_{\mathcal{H}^4} \quad (\text{C-7})$$

for all $\psi \in \mathcal{H}^4$ and all φ satisfying $\|\varphi\|_{\mathcal{L}_{(3)}^2} \leq A$. To show this, we first observe that, as in the proof of Lemma C.2,

$$\begin{aligned} & \left| \|(-\Delta + V_\varphi + M)^2 \psi\|_2 - \|(-\Delta + M)(-\Delta + V_\varphi + M)\psi\|_2 \right| \\ & \leq \|V_\varphi(-\Delta + M)^{-1}\| \|(-\Delta + M)(-\Delta + V_\varphi + M)\psi\|_2 \end{aligned}$$

and that $\|V_\varphi(-\Delta + M)^{-1}\|$ can be made arbitrarily small for $\|\varphi\|_{\mathcal{L}^2}$ bounded by choosing M large. Thus, it suffices to show that $\|(-\Delta + M)(-\Delta + V_\varphi + M)\psi\|_2$ is equivalent to $\|(-\Delta + M)^2 \psi\|_2$. We compute

$$\begin{aligned} & \left| \|(-\Delta + M)(-\Delta + V_\varphi + M)\psi\|_2 - \|(-\Delta + V_\varphi + M)(-\Delta + M)\psi\|_2 \right| \\ & \leq \|(2\nabla V_\varphi \cdot \nabla + \Delta V_\varphi)(-\Delta + M)^{-1}\| \|(-\Delta + M)\psi\|_2. \end{aligned}$$

According to (C-1), the first factor on the right side can be made arbitrarily small for $\|\varphi\|_{\mathcal{L}_{(3)}^2}$ bounded by choosing M large. We conclude by applying the argument in Lemma C.2 again to compare $\|(-\Delta + V_\varphi + M)(-\Delta + M)\psi\|_2$ to $\|(-\Delta + M)^2 \psi\|_2$. This proves the claim.

According to Lemma C.2, for every $\tau > 0$ there is an $A > 0$ (depending only on $\|\psi_0\|_{\mathcal{H}^2}$, $\|\varphi_0\|_{\mathcal{L}_{(3)}^2}$ and τ) such that $\|\varphi_t\|_{\mathcal{L}_{(3)}^2} \leq A$ for all $|t| \leq \tau\alpha^2$. We choose M corresponding to this value of A and

compute, using the equation for ψ_t ,

$$\begin{aligned} \partial_t \mathcal{E}^{(4)}(\psi_t, \varphi_t) &= 2 \operatorname{Re}((-\Delta + V_{\varphi_t} + M)^2 \psi_t, (-\Delta + V_{\varphi_t} + M)^2 \partial_t \psi_t) \\ &\quad + 2 \operatorname{Re}((-\Delta + V_{\varphi_t} + M)^2 \psi_t, (\partial_t V_{\varphi_t})(-\Delta + V_{\varphi_t} + M) \psi_t) \\ &\quad + 2 \operatorname{Re}((-\Delta + V_{\varphi_t} + M)^2 \psi_t, (-\Delta + V_{\varphi_t} + M)(\partial_t V_{\varphi_t}) \psi_t) \\ &= 4 \operatorname{Re}((-\Delta + V_{\varphi_t} + M)^2 \psi_t, (\partial_t V_{\varphi_t})(-\Delta + V_{\varphi_t} + M) \psi_t) \\ &\quad - 2 \operatorname{Re}((-\Delta + V_{\varphi_t} + M)^2 \psi_t, (2\nabla \partial_t V_{\varphi_t} \cdot \nabla + \Delta \partial_t V_{\varphi_t}) \psi_t). \end{aligned}$$

Therefore, by the Schwarz inequality,

$$\begin{aligned} &\partial_t \mathcal{E}^{(4)}(\psi_t, \varphi_t) \\ &\leq 2(\mathcal{E}^{(4)}(\psi_t, \varphi_t))^{\frac{1}{2}} (2\|\partial_t V_{\varphi_t}\|_{\infty} \|(-\Delta + V_{\varphi_t} + M)\psi_t\|_2 + 2\|\nabla \partial_t V_{\varphi_t}\|_{\infty} \|\nabla \psi_t\|_2 + \|\Delta \partial_t V_{\varphi_t}\|_{\infty} \|\psi_t\|_2). \end{aligned}$$

According to [Lemma C.2](#) and [\(C-5\)](#), all terms involving ψ_t here are bounded by a constant for $|t| \leq \tau\alpha^2$. Assume that we can prove that all terms involving φ_t here are bounded by a constant times α^{-2} for $|t| \leq \tau\alpha^2$. Then we will have shown that

$$\partial_t \mathcal{E}^{(4)}(\psi_t, \varphi_t) \lesssim \alpha^{-2} (\mathcal{E}^{(4)}(\psi_t, \varphi_t))^{\frac{1}{2}}$$

for $|t| \leq \tau\alpha^2$, which implies that $(\mathcal{E}^{(4)}(\psi_t, \varphi_t))^{\frac{1}{2}} \lesssim 1 + \alpha^{-2}|t| \lesssim 1$ for $|t| \leq \tau\alpha^2$. According to [\(C-7\)](#), this proves that $\|\psi_t\|_{\mathcal{H}^4} \lesssim 1$ for $|t| \leq \tau\alpha^2$.

Thus, it remains to prove that for all multi-indices β with $|\beta| \leq 2$,

$$\|\partial_x^\beta \partial_t V_{\varphi_t}\|_{\infty} \lesssim \alpha^{-2} \quad \text{for } |t| \leq \tau\alpha^{-2}. \tag{C-8}$$

If we insert the equation of φ_t into the definition of V_{φ_t} , we find

$$\partial_t V_{\varphi_t}(x) = -i\alpha^{-2} \int_{\mathbb{R}^3} (e^{-ik \cdot x} \varphi_t(k) - e^{ik \cdot x} \overline{\varphi_t(k)}) \frac{dk}{|k|}. \tag{C-9}$$

(Note that the contribution from σ_{ψ_t} cancels.) Using this formula, we obtain

$$\|\partial_x^\beta \partial_t V_{\varphi_t}\|_{\infty} \lesssim \alpha^{-2} \|\varphi_t\|_{\mathcal{L}_{|\beta|+1}^2}$$

in the same way as we obtained [\(C-1\)](#). This implies [\(C-8\)](#) in view of the bounds on φ_t from [Lemma C.2](#).

It is straightforward to deduce the remaining bounds claimed in the proposition. The bound on $\|\varphi_t\|_{\mathcal{L}_{(3)}^2}$ follows from [Lemma C.2](#). Because of the equation for ψ_t , we have

$$\|\partial_t \psi_t\|_{\mathcal{H}^2} \leq \|\Delta \psi_t\|_{\mathcal{H}^2} + \|V_{\varphi_t} \psi_t\|_{\mathcal{H}^2} \lesssim \|\psi_t\|_{\mathcal{H}^4} + \sum_{|\beta| \leq 2} \|\partial^\beta V_{\varphi_t}\|_{\infty} \|\psi_t\|_{\mathcal{H}^2}.$$

Using the fact that $\|\psi_t\|_{\mathcal{H}^4} \lesssim 1$ and $\|\varphi_t\|_{\mathcal{L}_{(3)}^2} \lesssim 1$, which by [\(C-1\)](#) controls $\|\partial^\beta V_{\varphi_t}\|_{\infty}$ for $|\beta| \leq 2$, we conclude that $\|\partial_t \psi_t\|_{\mathcal{H}^2} \lesssim 1$. The second bound in [\(2-8\)](#) follows from [Lemma C.2](#).

Finally, we need to prove the bounds on g_s and $g_{s,t}$. By the Schwarz inequality as in the proof of [\(C-1\)](#) together with the equation for φ_s we find

$$\|g_s\|_{\infty} \lesssim \|\partial_s \varphi_s\|_{\mathcal{L}_{(1)}^2} \leq \alpha^{-2} (\|\varphi_s\|_{\mathcal{L}_{(1)}^2} + \|\sigma_{\psi_s}\|_{\mathcal{L}_{(1)}^2}).$$

According to (C-2) and Lemma 2.1 we have $\|\sigma_{\psi_s}\|_{\mathcal{L}^2_{(1)}} \lesssim \|\psi_s\|_{H^1}^2 \lesssim 1$. Moreover, if $|t|, |s| \leq \tau\alpha^2$, then Lemma C.2 implies $\|\varphi_s\|_{\mathcal{L}^2_{(1)}} \lesssim 1$. Thus,

$$\|g_s\|_{\infty} \lesssim \alpha^{-2},$$

as claimed. Moreover, $g_{s,t} = \int_s^t g_{s_1} ds_1$, so for $t > s$

$$\|g_{s,t}\|_{\infty} \leq \int_s^t \|g_{s_1}\|_{\infty} ds_1 \lesssim \alpha^{-2}(t-s).$$

This proves (2-9). \square

Appendix D: Reduced density matrices

Here we show how the approximation of $e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega$ in Theorem 1.3 yields approximations to its reduced density matrices in Theorem 1.2. The argument relies on the following abstract lemma.

Lemma D.1. *Let \mathcal{H}_1 and \mathcal{H}_2 be Hilbert spaces; let $\Psi, \Phi \in \mathcal{H}_1 \otimes \mathcal{H}_2$ and $f \in \mathcal{H}_1$ and $g \in \mathcal{H}_2$ such that*

$$\begin{aligned} \Psi &= f \otimes g + \Phi, \\ \|f\|_{\mathcal{H}_1} &\leq C, \quad \|g\|_{\mathcal{H}_2} \leq C, \quad \|\Phi\|_{\mathcal{H}_1 \otimes \mathcal{H}_2} \leq C\varepsilon, \\ \|\langle g, \Phi \rangle_{\mathcal{H}_2}\|_{\mathcal{H}_1} &\leq C\varepsilon^2, \quad \|\langle f, \Phi \rangle_{\mathcal{H}_1}\|_{\mathcal{H}_2} \leq C\varepsilon^2 \end{aligned}$$

for some $C > 0$ and $\varepsilon > 0$. Define

$$\gamma_1 = \text{Tr}_{\mathcal{H}_2} |\Psi\rangle\langle\Psi|, \quad \gamma_2 = \text{Tr}_{\mathcal{H}_1} |\Psi\rangle\langle\Psi|.$$

Then

$$\text{Tr}_{\mathcal{H}_1} |\gamma_1 - \|g\|_{\mathcal{H}_2}^2 |f\rangle\langle f| | \leq 3C^2\varepsilon^2, \quad \text{Tr}_{\mathcal{H}_2} |\gamma_2 - \|f\|_{\mathcal{H}_1}^2 |g\rangle\langle g| | \leq 3C^2\varepsilon^2.$$

Before proving this lemma, let us use it to derive Theorem 1.2 from Theorem 1.3. We apply the lemma with $\mathcal{H}_1 = \mathcal{L}^2(\mathbb{R}^3)$, $\mathcal{H}_2 = \mathcal{F}$, $f = e^{-i\int_0^t \omega(s) ds} \psi_t$, $g = \Omega$,

$$\Psi = W^*(\alpha^2 \varphi_t) e^{-i\tilde{H}_\alpha^F t} \psi_0 \otimes W(\alpha^2 \varphi_0) \Omega, \quad \Phi = W^*(\alpha^2 \varphi_t) R(t).$$

Then Theorem 1.3 implies that the assumptions of the lemma are satisfied with $\varepsilon = \alpha^{-1}(1 + |t|)$. We have $\|f\|^2 = \|\psi_t\|^2 = \|\psi_0\|^2 = 1$, $\|g\|^2 = \|\Omega\|^2 = 1$ and $|f\rangle\langle f| = |\psi_t\rangle\langle\psi_t|$. Moreover,

$$\text{Tr}_{\mathcal{H}_2} |\Psi\rangle\langle\Psi| = \gamma_t^{\text{particle}}, \quad \text{Tr}_{\mathcal{H}_1} |\Psi\rangle\langle\Psi| = W^*(\alpha^2 \varphi_t) \gamma_t^{\text{field}} W(\alpha^2 \varphi_t).$$

Thus, the conclusion of Theorem 1.2 follows from the lemma.

We now turn to the proof of the lemma. It relies on the bound

$$\text{Tr}_{\mathcal{H}_1} |\text{Tr}_{\mathcal{H}_2} |\Psi_1\rangle\langle\Psi_2| | \leq \|\Psi_1\|_{\mathcal{H}_1 \otimes \mathcal{H}_2} \|\Psi_2\|_{\mathcal{H}_1 \otimes \mathcal{H}_2} \tag{D-1}$$

valid for any vectors $\Psi_1, \Psi_2 \in \mathcal{H}_1 \otimes \mathcal{H}_2$. For the proof of (D-1) recall the variational characterization of the trace norm,

$$\text{Tr}_{\mathcal{H}_1} |K| = \sup_{(e_j), (e'_j)} \text{Re} \sum_j \langle e_j, K e'_j \rangle_{\mathcal{H}_1},$$

where the supremum is over all orthonormal systems (e_j) and (e'_j) in \mathcal{H}_1 . Thus, if (b_k) is an orthonormal basis in \mathcal{H}_2 , then

$$\begin{aligned} \operatorname{Re} \sum_j \langle e_j, (\operatorname{Tr}_{\mathcal{H}_2} |\Psi_1\rangle\langle\Psi_2|) e'_j \rangle_{\mathcal{H}_1} &= \operatorname{Re} \sum_{j,k} \langle e_j \otimes b_k, \Psi_1 \rangle_{\mathcal{H}_1 \otimes \mathcal{H}_2} \langle \Psi_2, e'_j \otimes b_k \rangle_{\mathcal{H}_1 \otimes \mathcal{H}_2} \\ &\leq \left(\sum_{j,k} |\langle e_j \otimes b_k, \Psi_1 \rangle_{\mathcal{H}_1 \otimes \mathcal{H}_2}|^2 \right)^{\frac{1}{2}} \left(\sum_{j,k} |\langle \Psi_2, e'_j \otimes b_k \rangle_{\mathcal{H}_1 \otimes \mathcal{H}_2}|^2 \right)^{\frac{1}{2}} \\ &\leq \|\Psi_1\|_{\mathcal{H}_1 \otimes \mathcal{H}_2} \|\Psi_2\|_{\mathcal{H}_1 \otimes \mathcal{H}_2}, \end{aligned}$$

where the last inequality comes from the orthonormality of $(e_j \otimes b_k)$ and $(e'_j \otimes b_k)$. Therefore the variational characterization of the trace norm yields (D-1).

Proof. Since $\operatorname{Tr}_{\mathcal{H}_2} |f \otimes g\rangle\langle\Phi| = |f\rangle\langle(g, \Phi)_{\mathcal{H}_2}|$, we have

$$\gamma_1 - \|g\|_{\mathcal{H}_2}^2 |f\rangle\langle f| = |f\rangle\langle(g, \Phi)_{\mathcal{H}_2}| + |\langle\Phi, g\rangle_{\mathcal{H}_2}\rangle\langle f| + \operatorname{Tr}_2 |\Phi\rangle\langle\Phi|.$$

By (D-1) and the assumptions the trace norm, each one of the three operators on the right side is bounded by $C^2\varepsilon^2$. This proves the first inequality in the lemma. The second one is proved similarly. \square

Finally, we show that the α^{-2} error bound in Theorem 1.2 (for times of order one) is due to the fact that φ_t is time-dependent. The proof makes use of the fact that for arbitrary normalized vectors a and b in a Hilbert space \mathcal{H} one has

$$\operatorname{Tr}_{\mathcal{H}} |a\rangle\langle a| - |b\rangle\langle b| = 2(1 - |\langle a, b \rangle|^2)^{\frac{1}{2}}, \quad (\text{D-2})$$

as is easily verified.

Proof of Lemma 1.4. Because of Theorem 1.2, it suffices to prove that there are $\varepsilon > 0$ and $c > 0$ such that for all $|t| \leq \varepsilon$ and all $\alpha \geq 1$,

$$\operatorname{Tr}_{\mathcal{F}} |W(\alpha^2\varphi_t)\Omega\rangle\langle W(\alpha^2\varphi_t)\Omega| - |W(\alpha^2\varphi_0)\Omega\rangle\langle W(\alpha^2\varphi_0)\Omega| \geq c\alpha^{-1}|t|.$$

According to Lemma A.4 and (D-2), this is equivalent to

$$1 - e^{-\alpha^2\|\varphi_t - \varphi_0\|_2^2} = 1 - \left| \langle \Omega, W^*(\alpha^2\varphi_0)W(\alpha^2\varphi_t)\Omega \rangle \right|^2 \geq \frac{1}{4}c^2\alpha^{-2}t^2.$$

Since $\|\varphi_t - \varphi_0\|_2 \lesssim \alpha^{-2}|t|$ by Lemma 2.1, it suffices to prove that there are $\varepsilon > 0$ and $c' > 0$ such that for all $|t| \leq \varepsilon$ and all $\alpha \geq 1$,

$$\|\varphi_t - \varphi_0\|_2 \geq c'\alpha^{-2}|t|.$$

Since $\varphi_0 + \sigma_{\psi_0} \neq 0$, this will clearly follow if we can prove that for all $|t| \leq \alpha^2$ and $\alpha \geq 1$,

$$\|\varphi_t - \varphi_0 + i\alpha^{-2}t(\varphi_0 + \sigma_{\psi_0})\|_2 \leq C\alpha^{-2}t^2. \quad (\text{D-3})$$

To prove this, we use equation (1-8) for φ_t to write

$$\varphi_t - \varphi_0 = \int_0^t \partial_s \varphi_s ds = -i\alpha^{-2} \int_0^t (\varphi_s + \sigma_{\psi_s}) ds = -i\alpha^{-2}t(\varphi_0 + \sigma_{\psi_0}) + r_t$$

with

$$r_t := -i\alpha^{-2} \int_0^t \int_0^s (\partial_{s_1} \varphi_{s_1} + \partial_{s_1} \sigma \psi_{s_1}) ds_1 ds.$$

By Lemma 2.1 and Proposition 2.2, the \mathcal{L}^2 -norm of the integrand of r_t is bounded by a constant uniformly in $|s_1| \leq \alpha^2$ and $\alpha \geq 1$. This yields (D-3) and completes the proof. \square

Appendix E: Improving the result of [Frank and Schlein 2014]

We now show how the techniques from [Frank and Schlein 2014] can be extended to times $|t| = o(\alpha)$. This argument is due to an anonymous referee, whom we thank for kind permission to include it in our paper.

Proposition E.1. *Let $\varphi \in L^2(\mathbb{R}^3)$ and $\alpha_0 > 0$. Assume that $\Psi \in L^2(\mathbb{R}^3) \otimes \mathcal{F}$ satisfies*

$$\|(p^2 + \mathcal{N} + 1)^{\frac{1}{2}} \Psi\| \leq M, \quad \|(p^2 + 1)^{\frac{1}{2}} \mathcal{N} \Psi\| \leq M\alpha^{-2}.$$

Then for all $\alpha \geq \alpha_0$ and all $t \in \mathbb{R}$,

$$\|e^{-i\tilde{H}_{\alpha^2}^F t} W(\alpha^2 \varphi) \Psi - e^{-iH_{\varphi} t} W(\alpha^2 \varphi) \Psi\|^2 \leq M^2(1 + 2\alpha^{-1})(e^{C|t|/(2\alpha)} - 1),$$

where C depends only on α_0 and an upper bound on $\|\varphi\|_{\mathcal{L}^2}$.

Note that this result can be applied, in particular, to $\Psi = \psi \otimes \Omega$ with $\|\psi\|_{\mathcal{H}^1} \leq M$. We also recall that the effective Schrödinger operator H_{φ} was defined in (2-12).

Proof. Let $A(t) := \|e^{-i\tilde{H}_{\alpha^2}^F t} W(\alpha^2 \varphi) \Psi - e^{-iH_{\varphi} t} W(\alpha^2 \varphi) \Psi\|^2$. It is shown in [Frank and Schlein 2014, Proposition 9] that $A'(t) = f(t) + g(t)$ with

$$f(t) \leq CM\alpha^{-1} A(t)^{\frac{1}{2}}, \quad \int_0^T g(t) dt \leq CM^2\alpha^{-2}T,$$

where C depends only on α_0 and an upper bound on $\|\varphi\|_{\mathcal{L}^2}$. We bound $f(t) \leq \frac{1}{2}C\alpha^{-1}(A(t) + M^2)$ and therefore

$$A(T) \leq \int_0^T f(t) dt + \int_0^T g(t) dt \leq \frac{1}{2}C\alpha^{-1} \int_0^T A(t) dt + \frac{1}{2}CM^2\alpha^{-1}(1 + 2\alpha^{-1})T.$$

Thus,

$$A(T) + M^2(1 + 2\alpha^{-1}) \leq M^2(1 + 2\alpha^{-1}) + \frac{1}{2}C\alpha^{-1} \int_0^T (A(t) + M^2(1 + 2\alpha^{-1})) dt$$

and, by Gronwall's inequality, for all $t \geq 0$

$$A(t) + M^2(1 + 2\alpha^{-1}) \leq M^2(1 + 2\alpha^{-1})e^{Ct/(2\alpha)}. \quad \square$$

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Note added in proof

After this work was accepted for publication, the preprint by M. Griesemer [2016] appeared on the arXiv. This preprint studies the dynamics generated by the initial conditions given by the minimizing pair (ψ_*, φ_*) of the energy functional $\mathcal{E}(\psi, \varphi)$ under the constraint $\|\psi\| = 1$ up to times of order $o(\alpha^2)$.

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