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We obtain an improvement of the Beckner inequality $||f||_2^2 - ||f||_p^2 \le (2-p) ||\nabla f||_2^2$ valid for $p \in [1, 2]$ and the Gaussian measure. Our improvement is essential for the intermediate case $p \in (1, 2)$, and moreover, we find the natural extension of the inequality for any real p.

1. Introduction

1.1. The history of the problem. The Poincaré inequality [Nash 1958] for the standard Gaussian measure

$$d\gamma_n = \frac{e^{-|x|^2/2}}{\sqrt{(2\pi)^n}} dx$$

states that

$$\int_{\mathbb{R}^n} f^2 \, d\gamma_n - \left(\int_{\mathbb{R}^n} f \, d\gamma_n \right)^2 \le \int_{\mathbb{R}^n} |\nabla f|^2 \, d\gamma_n \tag{1}$$

for any smooth bounded function $f : \mathbb{R}^n \to \mathbb{R}$. Later William Beckner [1989] generalized (1) for any real power $p, 1 \le p \le 2$, as follows:

$$\int_{\mathbb{D}^n} f^p \, d\gamma_n - \left(\int_{\mathbb{D}^n} f \, d\gamma_n \right)^p \le \frac{p(p-1)}{2} \int_{\mathbb{D}^n} f^{p-2} |\nabla f|^2 \, d\gamma_n \tag{2}$$

for any smooth bounded $f: \mathbb{R}^n \to (0, \infty)$. We caution the reader that in [Beckner 1989], inequality (2) was formulated in a slightly different but equivalent form (see Theorem 1, inequality (3) in that paper). It should be also mentioned that in the case p=2, inequality (2) does coincide with (1) for all $f \geq 0$ but it does not imply the Poincaré inequality for the functions taking the negative values, especially when $\int_{\mathbb{R}^n} f \, d\gamma_n = 0$. If $p \to 1+$ then (2) provides us with log-Sobolev inequality (see [Beckner 1989]). In general, the constant p(p-1)/2 is sharp in the right-hand side of (2), as can be seen for n=1 on the test functions $f(x) = e^{\varepsilon x}$ by sending $\varepsilon \to 0$.

Later Beckner's inequality (2) was studied by many mathematicians for different measures, in different settings and for different spaces as well. We refer the reader to [Arnold et al. 2007; Da Pelo et al. 2016;

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Barthe et al. 2006; 2007; Barthe and Roberto 2003; Bobkov and Tetali 2003; Bobkov and Götze 1999; Boucheron et al. 2005; Chafaï 2004; Wang 2005; Latała and Oleszkiewicz 2000; Kolesnikov 2007].

An analysis done in [Ivanisvili and Volberg 2015c] indicates that the right-hand side of (2) can be improved. In the present paper we address this issue: what is the precise estimate of the difference given in the left-hand side of (2), and can the requirement $p \in [1, 2]$ be avoided by slightly changing the right-hand side of (2)?

We give complete answers to these questions. For example, if $p = \frac{3}{2}$ we will obtain an improvement in Beckner's inequality (2):

$$\int_{\mathbb{R}^n} f^{3/2} \, d\gamma_n - \left(\int_{\mathbb{R}^n} f \, d\gamma_n \right)^{3/2} \le \int_{\mathbb{R}^n} \left(f^{3/2} - \frac{1}{\sqrt{2}} (2f - \sqrt{f^2 + |\nabla f|^2}) \sqrt{f + \sqrt{f^2 + |\nabla f|^2}} \right) d\gamma_n. \tag{3}$$

The left-hand side of (3) coincides with the left-hand side of (2) for $p = \frac{3}{2}$, but the right-hand side of (3) is strictly smaller than the right-hand side in (2). Indeed, notice that we have the *pointwise* inequality

$$x^{3/2} - \frac{1}{\sqrt{2}}(2x - \sqrt{x^2 + y^2})\sqrt{x + \sqrt{x^2 + y^2}} \le \frac{3}{8}x^{-1/2}y^2 \quad \text{for all } x, y \ge 0,$$
 (4)

which follows from the homogeneity, i.e., take x = 1, and the rest is a direct computation which follows by introducing a new variable

$$u := \sqrt{1 + \sqrt{1 + y^2}}.$$

As one can see, the improvement of Beckner's inequality (2) is essential. Indeed, if $y \to \infty$ then the right-hand side of (4) increases as y^2 whereas the left-hand side of (4) increases as $y^{3/2}$. Also notice that if $x \to 0$ then the difference of both sides of (4) tends to infinity. The only place where the quantities in (4) are comparable is when $y/x \to 0$.

1.2. *Main results.* Let k be a real parameter. Let $H_k(x)$ be the Hermite function which satisfies the Hermite differential equation

$$H_k'' - xH_k' + kH_k = 0, \quad x \in \mathbb{R}, \tag{5}$$

and which grows relatively slowly, that is, $H_k(x) = x^k + o(x^k)$ as $x \to +\infty$. If k is a nonnegative integer then H_k is the *probabilists*' Hermite polynomial of degree k with the leading coefficient 1; for example, $H_0(x) = 1$, $H_1(x) = x$, $H_2(x) = x^2 - 1$, etc. In general, for arbitrary $k \in \mathbb{R}$ one should think that H_k is the analytic extension of the Hermite polynomials in k (existence and many other properties will be mentioned in Section 2).

For $k \in \mathbb{R}$, let R_k be the rightmost zero of $H_k(x)$ (see Lemma 7). If $k \le 0$ then we set $R_k = -\infty$. Define $F_k(x)$ as

$$F_k\left(\left|\frac{H_k'(q)}{H_k(q)}\right|\right) = \frac{H_{k+1}(q)}{H_k^{1+1/k}(q)} \quad \text{for } q \in (R_k, \infty).$$
 (6)

We will see in the next section that $F_k \in C^2([0, \infty))$ is well defined and $F_k(0) = 1$. Moreover, if k > -1 then F_k will be a decreasing concave function, and if k < -1 then F_k will be an increasing convex function.

One may observe that

$$F_1(y) = 1 - y^2$$
, $F_2(y) = \frac{1}{\sqrt{2}}(2 - \sqrt{1 + y^2})\sqrt{1 + \sqrt{1 + y^2}}$.

If k = 0 then definition (6) should be understood in the limiting sense as

$$F_{\exp}(H_{-1}(q)) = q \exp\left(\alpha - \int_{1}^{q} H_{-1}(s) \, ds\right) \quad \text{for all } q \in \mathbb{R},\tag{7}$$

where

$$\alpha = \int_{1}^{\infty} \left(H_{-1}(s) - \frac{1}{s} \right) ds \approx -0.266 \dots$$
 (8)

Theorem 1. For any $p \in \mathbb{R} \setminus [0, 1]$ and any smooth bounded $f \geq 0$ with $\int_{\mathbb{R}^n} f^p d\gamma_n < \infty$ we have

$$\int_{\mathbb{R}^n} f^p F_{1/(p-1)} \left(\frac{|\nabla f|}{f} \right) d\gamma_n \le \left(\int_{\mathbb{R}^n} f \, d\gamma_n \right)^p. \tag{9}$$

The inequality is reversed if $p \in (0, 1)$.

Proposition 2. We have

$$1 - \frac{p(p-1)}{2}t^2 \le F_{1/(p-1)}(t) \quad \text{for all } t \ge 0, \ p \in (1,2].$$
 (10)

It remains to notice that estimate (10) applied to (9) immediately gives (2).

The improvement will be essential when $t \to \infty$. For example, it will become clear in the next section that as $t \to \infty$ we have

$$F_{1/(p-1)}(t) \sim -t^p \left(H'_{1/(p-1)}(R_{1/(p-1)})\right)^{1-p} \quad \text{for } p > 1.$$
 (11)

Another immediate application of Theorem 1 is the following corollary.

Corollary 3. For any $p \in (1, 2]$ and any smooth bounded $f \ge 0$ we have

$$\int_{\mathbb{R}^n} f^p \, d\gamma_n - \left(\int_{\mathbb{R}^n} f \, d\gamma_n \right)^p \le \left(H'_{1/(p-1)}(R_{1/(p-1)}) \right)^{1-p} \int_{\mathbb{R}^n} |\nabla f|^p \, d\gamma_n. \tag{12}$$

Estimate (12) will follow by showing that, for any $y \ge 0$, the map

$$x \to x^p - x^p F_{1/(1-p)} \left(\frac{y}{x}\right) \tag{13}$$

is decreasing for x > 0, and the limit $x \to 0$ gives (12) by (11).

Appearance of the roots of Hermite functions in (12) seems quite unexpected, especially when these estimates are obtained on the Hamming cube. For example, in [Ivanisvili and Volberg 2016] we were able to extend (12) to the Hamming cube but for a particular power $p = \frac{3}{2}$:

$$\mathbb{E}f^{3/2} - (\mathbb{E}f)^{3/2} \le \frac{1}{\sqrt{2}} \mathbb{E}|\nabla f|^{3/2}, \quad f: \{-1, 1\}^n \to \mathbb{R}_+. \tag{14}$$

We refer the reader to that paper for the notations, and we notice that the result announced there is a counterpart of (9) for $p = \frac{3}{2}$ on the Hamming cube, where the identity $x^{3/2}F_2(y/x) = \Re(x+iy)^{3/2}$

was used. Next, let $A \subset \{-1, 1\}^n$, and let $w_A(x)$ denote the number of edges containing x between the set A and its complement. Clearly $w_A(x)$ lives on the *boundary* of the set A: $w_A(x) = 4|\nabla \mathbb{1}_A|^2$. If A has cardinality 2^{n-1} then the classical edge-isoperimetric inequality [Harper 1966] states that $\sum_{x \in \{-1,1\}^n} w_A(x) \ge 2^n$. On the other hand, taking $f = \mathbb{1}_A$ in (14) gives

$$\sum_{x \in \{-1,1\}^n} w_A(x)^{3/4} \ge (2 - \sqrt{2})2^n,$$

which is a new edge-isoperimetric inequality and does not follow from the classical one.

Theorem 1 generates several inequalities. If $p \to 1+$ then (9) gives the log-Sobolev inequality. If p=2 then (9) provides us with the Poincaré inequality. If $p \to \pm \infty$ then we obtain a new Sobolev inequality:

Corollary 4. For any smooth bounded f we have

$$\int_{\mathbb{R}^n} \exp(f) F_{\exp}(|\nabla f|) d\gamma_n \le \exp\left(\int_{\mathbb{R}^n} f d\gamma_n\right),$$

where F_{exp} is defined in (7).

Finally if $p \to 0$ we obtain a new "negative log-Sobolev" inequality:

Corollary 5. For any smooth bounded $f \ge 0$ with $\int_{\mathbb{R}^n} \ln f \, d\gamma_n > -\infty$ we have

$$\int_{\mathbb{R}^n} -\ln f \, d\gamma_n + \ln \left(\int_{\mathbb{R}^n} f \, d\gamma_n \right) \leq \int_{\mathbb{R}^n} -F_{-\ln} \left(\frac{|\nabla f|}{f} \right) d\gamma_n,$$

where $F_{-\ln}(t)$ is defined as

$$F_{-\ln}\left(\frac{H_{-2}(x)}{H_{-1}(x)}\right) = \int_{1}^{x} H_{-1}(s) \, ds - c + \ln H_{-1}(x), \quad x \in \mathbb{R}.$$

All these estimates extend to uniformly log-concave probability measures in the following sense (for the proof see Section 3).

Corollary 6. Let $d\mu = e^{-U} dx$ be a probability measure, where $\text{Hess } U \ge R \cdot \text{Id for some } R > 0$. For any $p \in \mathbb{R} \setminus [0, 1]$ and any smooth bounded $f \ge 0$ with $\int_{\mathbb{R}^n} f^p d\mu < \infty$ we have

$$\int_{\mathbb{R}^n} f^p F_{1/(p-1)} \left(\frac{|\nabla f|}{f\sqrt{R}} \right) d\mu \le \left(\int_{\mathbb{R}^n} f \, d\mu \right)^p. \tag{15}$$

The inequality is reversed if $p \in (0, 1)$.

The limiting cases of (15) when $p \to \pm \infty$ and $p \to 0$ should be understood in the sense of functions $F_{\rm exp}$ and $F_{-\ln}$ as in Corollary 4 and Corollary 5.

To summarize, the current paper provides us with estimates of Φ -entropy (see [Chafaï 2004])

$$\mathbf{Ent}_{\gamma_n}^{\Phi}(f) := \int_{\mathbb{R}^n} \Phi(f) \, d\gamma_n - \Phi\left(\int_{\mathbb{R}^n} f \, d\gamma_n\right)$$

for the following examples:

- $\Phi(x) = x^p$ for $p \in \mathbb{R} \setminus [0, 1]$ using Theorem 1.
- $\Phi(x) = -x^p$ for $p \in (0, 1)$ using Theorem 1.
- $\Phi(x) = e^x$ using Corollary 4, or by taking $p \to \pm \infty$ in Theorem 1.
- $\Phi(x) = -\ln x$ using Corollary 5, or by taking $p \to 0$ in Theorem 1.
- $\Phi(x) = x \ln x$ by taking $p \to 1$ in Theorem 1.

2. The proof of the theorem

The proof of the theorem amounts to checking that the real-valued function

$$M(x, y) = x^p F_k\left(\frac{y}{x}\right), \quad k = \frac{1}{1-p},\tag{16}$$

defined on $[\varepsilon, \infty) \times [0, \infty)$ for any $\varepsilon > 0$, obeys a necessary smoothness condition, has a boundary condition $M(x, 0) = x^p$ and satisfies the partial differential inequality

$$\begin{pmatrix} M_{xx} + M_y/y & M_{xy} \\ M_{xy} & M_{yy} \end{pmatrix} \le 0, \tag{17}$$

with reversed inequality in (17) if $p \in (0, 1)$. Then by Theorem 1 in [Ivanisvili and Volberg 2015c] we obtain that

$$\int_{\mathbb{R}^n} f^p F_k\left(\frac{|\nabla f|}{f}\right) d\gamma_n = \int_{\mathbb{R}^n} M(f, |\nabla f|) d\gamma_n \le M\left(\int_{\mathbb{R}^n} f d\gamma_n, 0\right) = \left(\int_{\mathbb{R}^n} f d\gamma_n\right)^p$$

for any smooth bounded $f \ge \varepsilon$, which is the statement of the theorem we want to prove (except we need to justify the passage to the limit $\varepsilon \to 0$ and this will be done later). Notice that the inequality is reversed if $p \in (0, 1)$; indeed, in this case we should work with -M(x, y) instead of M(x, y).

Next we will need some tools regarding the Hermite functions H_k .

2.1. Properties of Hermite functions. H_k can be defined (see [Hayman and Ortiz 1975]) by

$$H_k(x) = -\frac{2^{-k/2}\sin(\pi k) \Gamma(k+1)}{2\pi} \sum_{n=0}^{\infty} \frac{\Gamma((n-k)/2)}{n!} (-x\sqrt{2})^n,$$
 (18)

or in terms of the confluent hypergeometric functions (see [Durand 1975]) by

$$H_k(x) = \sqrt{\frac{2^k}{\pi}} \left[\cos\left(\frac{\pi k}{2}\right) \Gamma\left(\frac{k+1}{2}\right) {}_{1}F_{1}\left(-\frac{k}{2}, \frac{1}{2}; \frac{x^2}{2}\right) + t\sqrt{2} \sin\left(\frac{\pi k}{2}\right) \Gamma\left(\frac{k}{2} + 1\right) {}_{1}F_{1}\left(\frac{1-k}{2}, \frac{3}{2}; \frac{x^2}{2}\right) \right]. \tag{19}$$

If k is a nonnegative integer then one should understand (18) and (19) in the limiting sense. Notice the recurrence properties

$$H'_k(x) = kH_{k-1}(x),$$
 (20)

$$H_{k+1}(x) = xH_k(x) - H'_k(x).$$
 (21)

These properties follow from (18) and the fact that $\Gamma(z+1) = z\Gamma(z)$.

We also notice that

$$H_k(x) := e^{x^2/4} D_k(x),$$
 (22)

where $D_k(x)$ is the parabolic cylinder function; i.e., it is the solution of the equation

$$D_k'' + \left(k + \frac{1}{2} - \frac{x^2}{4}\right)D_k = 0.$$

Since $H_k(x)$ is an entire function in x and k (see [Temme 2015] for the parabolic cylinder function), sometimes it will be convenient to write H(x, k) instead of $H_k(x)$. The precise asymptotic for $x \to +\infty$, x > 0 and any $k \in \mathbb{R}$ is given by

$$H_k(x) \sim x^k \cdot \sum_{n=0}^{\infty} (-1)^n \frac{(-k)_{2n}}{n! (2x^2)^n}.$$
 (23)

Here $(a)_n = 1$ if n = 0, and $(a)_n = a(a+1)\cdots(a+n-1)$ if n > 0. When $x \to -\infty$ we have

$$H_k(x) \sim |x|^k \cos(k\pi) \sum_{n=0}^{\infty} (-1)^n \frac{(-k)_{2n}}{n!(2x^2)^n} + \frac{\sqrt{2\pi}}{\Gamma(-k)} |x|^{-k-1} e^{x^2/2} \sum_{n=0}^{\infty} \frac{(1+k)_{2n}}{n!(2x^2)^n}.$$
 (24)

We refer the reader to [Temme 2015; Olver et al. 2010]. For instance, for (23) we can use the asymptotic formula (12.9.1) in [Olver et al. 2010] for the parabolic cylinder function. To verify (24) we can express $H_k(-x)$ as a linear combination of two parabolic cylinder functions but with argument x instead of -x, see (12.2.15) in [Olver et al. 2010], and then we can use (12.9.1) and (12.9.2) in the same paper.

Next we will need the result of Elbert and Muldoon [1999] which describes the behavior of the real zeros of $H_k(x)$ for any real k.

Lemma 7. For $k \le 0$, the function $H_k(x)$ has no real zeros, and it is positive on the real axis. For $n < k \le n+1$, $n=0,1,\ldots$, the function $H_k(x)$ has n+1 real zeros. Each zero is increasing function of k on its interval of definition.

The proof of the lemma is Theorem 3.1 in [Elbert and Muldoon 1999]. It is explained in the paper that as k passes through each nonnegative integer n a new leftmost zero appears at $-\infty$, while the rightmost zero passes through the largest zero of $H_k(x)$. They also include more precise information about the asymptotic behavior of the zeros as $k \to \infty$.

Further we will need Turán's inequality for H_k for any real k.

Lemma 8. We have Turán's inequality:

$$H_k^2(x) - H_{k-1}(x)H_{k+1}(x) > 0 \quad \text{for all } k \in \mathbb{R}, \ x \ge L_k,$$
 (25)

where L_k denotes the leftmost zero of H_k . If $k \le 0$ then $L_k = -\infty$.

The lemma is known as Turán's inequality when k is a nonnegative integer. Unfortunately we could not find the reference in the case when k is different from a positive integer; therefore we decided to include the proof of the lemma.

The following is borrowed from [Madhava Rao and Thiruvenkatachar 1949].

Proof. Take $f(x) = e^{-x^2/2} (H_k^2(x) - H_{k-1}(x)H_{k+1}(x))$. Asymptotic formulas (23) and (24) imply that

$$\lim_{x \to +\infty} f(x) = 0 \quad \text{for all } k \in \mathbb{R},$$

$$f(x) \sim \sqrt{2\pi} |x| > 0 \quad \text{for } x \to -\infty, \ k = 0,$$

$$2e^{-x^2/2}$$
(26)

$$f(x) \sim \frac{2\pi e^{x^2/2}}{\Gamma(-k)\Gamma(-k+1)} |x|^{-2k-2} \quad \text{for } x \to -\infty, \ k \notin \{0\} \cup \mathbb{N}.$$

On the other hand, notice that by (20) and (21) we have

$$f'(x) = -e^{-x^2/2}H_kH_{k-1}. (27)$$

If $k \le 0$ then by Lemma 7 we have f' < 0, and because of the conditions $f(-\infty) = +\infty$ and $f(\infty) = 0$ we obtain that f > 0 on \mathbb{R} . To verify the statement for k > 0 we notice that

$$f''(x) = e^{-x^2/2} (H_k^2 - kH_{k-1}^2).$$
(28)

Now we notice that if $H_k(c) = 0$ then $H_{k-1}(c) \neq 0$. Indeed, assume to the contrary that $H_{k-1}(c) = 0$. Then by (20) we have $H'_k(c) = 0$ and by (5) we obtain $H''_k(c) = 0$, and again taking derivatives in (20) we obtain that $H_{k-2}(c) = 0$. Repeating this process we obtain that $H_{k-N}(c) = 0$ for any large integer N > 0. But this contradicts Lemma 7.

Thus by (27) and (28) we obtain that c is a local minimum of f if and only if $H_{k-1}(c) = 0$. Then $f(c) = e^{-x^2/2}H_k^2(c) > 0$. Finally we obtain that $f: [L_k, \infty) \to \mathbb{R}$ is positive on its local minimum points, $f(\infty) = 0$ and $f(L_k) > 0$ (because H_{k-1} and H_{k+1} have opposite signs at zeros of H_k by (21)). Therefore f > 0 on $[L_k, \infty) \to \mathbb{R}$ and the lemma is proved.

Remark 9. If $k \in \mathbb{N}$ then H_k is the probabilists' Hermite polynomial of degree k, so f(x) will be even and inequality (25) will hold for all $x \in \mathbb{R}$, which confirms the classical Turán's inequality. However, if k > 0 but $k \notin \mathbb{N}$ then (25) fails when $x \to -\infty$; see (26).

Finally the next corollary together with Lemma 7 implies that

$$\left| \frac{H_k'}{H_k} \right| = \operatorname{sign}(k) \frac{H_k'(q)}{H_k(q)}$$

is positive and decreasing for $q \in (R_k, \infty)$ and $k \in \mathbb{R} \setminus \{0\}$.

Corollary 10. For any $x \ge L_k$ and any $k \in \mathbb{R} \setminus \{0\}$ we have

$$\operatorname{sign}[(H_k')^2 - H_k H_k''] = \operatorname{sign}(k).$$

Proof. The proof follows from Lemma 8 and the identity

$$k(H_k^2 - H_{k-1}H_{k+1}) = (H_k')^2 - H_k H_k''$$
(29)

by
$$(5)$$
, (20) and (21) .

2.2. Checking the partial differential inequality. Let p = 1 + 1/k. Further we assume $k \neq 0, -1$. Define $F = F_k$ as in the Introduction:

$$F(t) = \frac{H_{k+1}(q)}{H_k^{1+1/k}(q)}, \quad \text{where } \left| \frac{H_k'(q)}{H_k(q)} \right| = t, \ q \in (R_k, \infty), \ t \in (0, \infty).$$
 (30)

Notice that by Corollary 10, the function

$$\left| \frac{H_k'(q)}{H_k(q)} \right| = \operatorname{sign}(k) \frac{H_k'(q)}{H_k(q)}$$

is positive decreasing in q for $q \in (R_k, \infty)$; moreover, by (23) and (20) we have $H'_k(q)/H_k(q) \sim k/q$ when $q \to +\infty$. From the same asymptotic formulas it follows that when $t \to 0+$ we have

$$F(t) = 1 - \frac{p(p-1)}{2}t^2 + O(t^4).$$

Therefore F is a well-defined function and $F \in C^2([0, \infty))$.

Take a positive $\varepsilon > 0$ and define M(x, y) as in (16):

$$M(x, y) := x^p F\left(\frac{y}{x}\right) \quad \text{for } y \ge 0, \ x > \varepsilon > 0.$$
 (31)

Clearly $M(x, \sqrt{y}) \in C^2([\varepsilon, \infty) \times [0, \infty))$. By Theorem 1 in [Ivanisvili and Volberg 2015c] we have the inequality

$$\int_{\mathbb{R}^n} M(f, |\nabla f|) \, d\gamma_n \le M\left(\int_{\mathbb{R}^n} f \, d\gamma_n, 0\right) \tag{32}$$

for all smooth bounded $f \ge \varepsilon$ if (17) holds. In terms of F (see (31)) condition (17) takes the form

$$tFF''p(p-1) + F'F'' - t(p-1)^2(F')^2 \ge 0$$
, i.e., the determinant of (17) is nonnegative, (33)

$$F''(t+t^3) + F'(2t^2 + 1 - 2pt^2) + Fp(p-1)t \le 0$$
, i.e., the trace of (17) is nonpositive, (34)

where t = y/x is the argument of F. In fact we will show that we have equality in (33) instead of inequality; therefore the sign of (17) will depend on the sign of the trace in (34). We will see that inequality (34) will be reversed for $p \in (0, 1)$.

From (30), (29), (20) and (21) we obtain

$$F'(t) = -\frac{k+1}{|k|} \frac{1}{H_{\nu}^{1/k}},\tag{35}$$

$$F''(t) = \frac{F'}{|k|} \frac{H_k H_{k-1}}{H_k^2 - H_{k+1} H_{k-1}},$$
(36)

$$F(t) = -\frac{|k|}{k+1} \frac{H_{k+1}}{H_k} F'. \tag{37}$$

If we plug (36) and (37) into (33) we obtain that the left-hand side of (33) is zero. If we plug (36) and (37) into (34) we obtain

left-hand side of (34) =
$$\left[\frac{(kH_{k-1}^2 - H_k^2 + H_{k-1}H_{k+1})^2 + H_{k-1}^2H_k^2}{H_k^2(H_k^2 - H_{k+1}H_{k-1})} \right] F'.$$

Thus the sign of left-hand side of (34) coincides with the sign of F', which coincides with sign(-(k+1)). The condition $p \in \mathbb{R} \setminus [0, 1]$ implies that k > -1 and therefore (17) holds. The condition $p \in (0, 1)$ implies that k < -1 and therefore the inequality in (17) is reversed.

Thus we have obtained (32) for smooth bounded functions $f \ge \varepsilon$. Next we claim that for an arbitrary smooth bounded $f \ge 0$ with $\int_{\mathbb{R}^n} f^p d\gamma_n < \infty$, we can apply the inequality to $f_\varepsilon := f + \varepsilon$ and send ε to 0 in (9). Indeed, it follows from (6) and (23) that as $t \to \infty$ we have

$$F(t) \sim \begin{cases} -t^{1+1/k} (H'_k(R_k))^{-1/k} & \text{for } k > 0, \\ \operatorname{sign}(-1-k) \left(\frac{e^{t^2/2} \sqrt{2\pi}}{t|\Gamma(-1-k)|} \right)^{-1/k} & \text{for } k < 0, \ k \neq -1. \end{cases}$$
(38)

Thus for p > 1, that is, k > 0, the claim about the limit follows from the estimate $|F(t)| \le C_1 + C_2 t^p$ together with the Lebesgue dominated convergence theorem.

If p < 0, that is, $k \in (-1, 0)$, we rewrite (9) in a standard way as

$$\int_{\mathbb{R}^n} f_{\varepsilon}^p \, d\gamma_n - \left(\int_{\mathbb{R}^n} f_{\varepsilon} \, d\gamma_n \right)^p \le \int_{\mathbb{R}^n} f_{\varepsilon}^p \left(1 - F\left(\frac{|\nabla f|}{f_{\varepsilon}} \right) \right) d\gamma_n. \tag{39}$$

Since f is bounded, $f \ge 0$ and $\int_{\mathbb{R}^n} f^p d\gamma_n < \infty$, there is no issue with the left-hand side of (39) when $\varepsilon \to 0$. For the right-hand side of (39) we notice that the function $x^p(1 - F(y/x))$ is nonnegative and decreasing in x. Then the claim follows from the monotone convergence theorem. The nonnegativity follows from the observation that F(0) = 1 and F' < 0 (see (35) where we have k > -1). The monotonicity follows from (6), (35), (20) and the straightforward computations

$$\frac{\partial}{\partial x}(x^p(1-F(y/x))) = x^{p-1}(p-pF(t)+tF'(t)) = x^{p-1}p\left[1-\frac{q}{H_t^{1/k}(q)}\right],\tag{40}$$

where

$$|k|\frac{H_{k-1}(q)}{H_k(q)} = t = \frac{y}{x}$$

and $q \in (R_k, \infty)$. The last expression in (40) is negative because

$$1 \ge F(t) = \frac{H_{k+1}}{H_k^{1+1/k}} = \frac{qH_k - kH_{k-1}}{H_k^{1+1/k}} > \frac{q}{H_k^{1/k}}.$$

Finally if $p \in (0, 1)$, that is, k < -1, we have the opposite inequality in (39). In this case the situation is absolutely the same as for $k \in (-1, 0)$ except now we should consider the function $x^p(F(y/x) - 1)$, which is nonnegative and decreasing in x; see (40). This finishes the proof of the theorem.

Now let us show Proposition 2. Since F(0) = 1, it is enough to show a stronger inequality, namely $F' + p(p-1)t \ge 0$. From (35) and the fact that k > 1 since $p \in (1, 2)$, we obtain that it is enough to

show the inequality

$$-\frac{p}{H_k^{1/k}} + p(p-1)\frac{H_k'}{H_k} \ge 0 \quad \text{for all } k \ge 1, \ q \in (R_k, \infty).$$

Using (20) and p = 1 + 1/k we notice that the inequality can be rewritten as $1 \ge H_k(q)/H_{k-1}^{k/(k-1)}(q)$ for all $q \in (R_k, \infty)$. To verify the last inequality recall that F(0) = 1 and F'(t) < 0. Therefore $F(t) \le 1$. We also recall the definition of F(t); see (30). It follows that $1 \ge F = H_{k+1}/H_k^{1+1/k}$ for all k > 0. The last inequality is the same as

$$1 \ge \frac{H_k(q)}{H_{k-1}^{k/(k-1)}(q)} \quad \text{for all } q \in (R_k, \infty), \ k > 1.$$
(41)

This finishes the proof of Proposition 2.

To verify Corollary 3 we only need to prove the monotonicity of the map (13) for $p \in (1, 2]$, that is, $k \ge 1$, and the rest will follow from (38). If k = 1 there is nothing to prove; therefore we assume k > 1. By (40) it is enough to show that $L(q) := H_k^{1/k}(q) - q \le 0$ for $q \in (R_k, \infty)$. The growth condition (24) on H_k implies that $\lim_{q \to \infty} L(q) = 0$. If $L'(q) \ge 0$ then we are done. Using (20) we notice that $L'(q) \ge 0$ is equivalent to (41), which was already proved.

2.3. *Proof of Corollaries 4 and 5.* Notice that as $y \to 0$ we have

$$F_{\text{exp}}(y) = 1 - \frac{y^2}{2} + O(y^4)$$
 and $F_{-\ln}(y) = -\frac{y^2}{2} + O(y^4)$.

One can check that

$$M_{\exp}(x, y) := e^x F_{\exp}(y), \quad M_{\exp}(x, 0) = e^x, \quad M_{\exp}(x, \sqrt{y}) \in C^2(\mathbb{R} \times \mathbb{R}_+),$$

 $M_{-\ln}(x, y) := -\ln(x) + F_{-\ln}\left(\frac{y}{x}\right), \quad M_{-\ln}(x, 0) = -\ln x, \quad x > 0,$

and $M_{-\ln}(x, \sqrt{y}) \in C^2([\varepsilon, \infty) \times \mathbb{R}^+)$ for any $\varepsilon > 0$. By straightforward computations we notice that if we set $\psi(q) = \alpha - \int_1^q H_{-1}(s) ds$ then using the identity $1 = qH_{-1}(q) + H_{-2}(q)$ we obtain

$$F_{\exp}(H_{-1}) = qe^{\psi}, \quad F'_{\exp}(H_{-1}) = -e^{\psi} \quad \text{and} \quad F''_{\exp}(H_{-1}) = -\frac{H_{-1}}{H_{-2}}.$$

Similarly we compute that

$$F'_{-\ln}\left(\frac{H_{-2}}{H_{-1}}\right) = -H_{-1}$$
 and $F''_{-\ln}\left(\frac{H_{-2}}{H_{-1}}\right) = -\frac{H_{-2}H_{-1}^2}{H_{-1}^2 - H_{-2}}.$

Next one notices that $M_{\rm exp}$ and $M_{-\ln}$ satisfy (17) (in fact the determinant of (17) is zero). Then by Theorem 1 in [Ivanisvili and Volberg 2015c] we obtain the corollaries. The passage to the limit for $M_{-\ln}(x,y)$ when $\varepsilon \to 0$ follows from the monotone convergence theorem. Indeed, we notice that $-F_{-\ln}(y/x) \ge 0$ is decreasing in x. We apply Corollary 5 to $f_{\varepsilon} = f + \varepsilon$ and send $\varepsilon \to 0$.

2.3.1. How we guessed the functions M_{exp} and $M_{-\ln}$. One may ask how to find the functions M_{exp} and $M_{-\ln}$. To find M_{exp} we should apply (9) to functions $f = e^{g/p}$, where g is some fixed function. Then (9) takes the form

$$\int_{\mathbb{R}^n} e^g F_{1/(p-1)} \left(\frac{|\nabla g|}{p} \right) d\gamma_n \le \left(\int_{\mathbb{R}^n} e^{g/p} \, d\gamma_n \right)^p. \tag{42}$$

Now we take $p \to \infty$. The right-hand side of (42) tends to $\exp(\int_{\mathbb{R}^n} g \, d\gamma_n)$. For the left-hand side of (42) we should compute the limit

$$F_{\exp}(t) := \lim_{p \to \infty} F_{1/(p-1)} \left(\frac{t}{p} \right) = \lim_{p \to \infty} F_{1/(p-1)} \left(\frac{t}{p-1} \right) = \lim_{k \to 0+} F_k(tk).$$

In fact all equalities can be justified by direct calculations using the fact that $H_k(x) = H(x, k)$ is the entire function of x and k; see [Temme 2015] for the parabolic cylinder function and formula (22).

It is clear that $F_{\text{exp}}(0) = 1$. Next if we take $k \to 0+$ in (6) we obtain

$$\lim_{k\to 0+} F_k\left(\left|\frac{H_k'}{H_k}\right|\right) = \lim_{k\to 0+} F_k\left(k\frac{H_{k-1}}{H_k}\right) = \lim_{k\to 0+} F_k\left(k\frac{H_{k-1}}{H_0}\right) = F_{\mathrm{exp}}(H_{k-1}).$$

On the other hand, for the right-hand side of (6) we have

$$\lim_{k \to 0+} \frac{H_{k+1}(q)}{H_k^{1+1/k}} = q \lim_{k \to 0+} H_k^{-1/k}.$$

Here we have used $H_0(q) = 1$ and $H_1(q) = q$. Thus it remains to find $\lim_{k \to 0+} H_k^{-1/k}$. If we take the derivative in k of (20) we obtain $H_{xk}(x,k) = H(x,k-1) + kH_k(x,k)$ (here subindices denote partial derivatives). Now taking k = 0 we obtain $H_{xk}(x,0) = H(x,-1)$. Thus $H_k(x,0)$ is an antiderivative of $H(x,-1) = H_{-1}$. So

$$\lim_{k \to 0+} H_k^{-1/k} = \lim_{k \to 0+} \exp\left(-\frac{1}{k} \ln(1 + kH_k(x, 0) + o(k))\right) = \exp\left(-\int H_{-1}(s) \, ds\right).$$

Finally we obtain

$$F_{\exp}(H_{-1}(q)) = q \exp\left(C - \int_{1}^{q} H_{-1}\right). \tag{43}$$

In order to satisfy the condition $F_{\text{exp}}(0) = 1$, the constant c must be chosen as $C = \int_1^\infty (H_{-1} - 1/s) \, ds$; indeed send $q \to \infty$ in (43). This gives Corollary 4. It is worth mentioning that we have also obtained

$$H_k(x,0) = \int_1^x H_{-1}(s) ds - \alpha;$$

see (8).

To find $M_{-\ln}$, let $F(x, k) := F_k(x)$. Let $F_k(x, k)$ denote the partial derivative in k of F(x, k). If we send $p \to 0$, p < 0 in (9) and compare the terms of order p we obtain

$$\int_{\mathbb{R}^n} \left(\ln f - F_k \left(\frac{|\nabla f|}{f}, -1 \right) \right) d\gamma_n \ge \ln \left(\int_{\mathbb{R}^n} f \, d\gamma_n \right).$$

where

It remains to find the function $F_k(x, -1)$. Let us equate terms of order (k + 1) as $k \to -1$, k < -1 in

$$F\left(\frac{|H_x(x,k)|}{H(x,k)},k\right) = \frac{H(x,k+1)}{H(x,k)^{1+1/k}}.$$

Straightforward computation shows that

$$F_k\left(\frac{H_{-2}(x)}{H_{-1}(x)}, -1\right) = H_k(x, 0) + \ln H_{-1}(x) = \int_1^x H_{-1}(s) \, ds - \alpha + \ln H_{-1}(x),$$

$$\alpha = \int_1^\infty \left(H_{-1}(s) - \frac{1}{s}\right) ds.$$

3. Concluding remarks

The reader may wonder how we guessed the choice (16). Of course it was not a random guess. Function (16) is the best possible in the sense that the determinant of (17) is identically zero:

$$M_{yy}\left(M_{xx} + \frac{M_y}{y}\right) - M_{xy}^2 = 0,$$

$$M(x, 0) = x^p \quad \text{for } x \ge 0.$$

$$(44)$$

Initially this was the way we started looking for M(x, y) as the solution of the Monge–Ampère equation with a drift (44). By a proper change of variables, the equation reduces to the backwards heat equation (see [Ivanisvili and Volberg 2015c] for more details where the connection with the theory of exterior differential systems of R. Bryant et al. [1991] was exploited)

$$u_{xx} + u_t = 0, \tag{45}$$

$$u(x, 0) = Cx^{p/(p-1)}$$
 for $x > 0$. (46)

One can notice that the Hermite polynomials do satisfy (45) and (46) when p/(p-1) is a positive integer. In general, one should invoke Hermite functions and this is the reason for the appearance of these functions in our theorem.

Another possibility is to assume that M(x, y) should be homogeneous of degree p, which forces M to have the form (31) for some F. Next setting h = F/F' and further by a subtle change of variables, one obtains Hermite differential equation (5).

Nevertheless, for the formal proof of Theorem 1 we do not need to go through the details. We have M(x, y) defined by (16) and we just need to check that it satisfies the desired properties.

That M(x, y) satisfies (17) makes it possible to extend Theorem 1 in a semigroup setting for uniformly log-concave probability measures. Indeed, let $d\mu = e^{-U} dx$, where Hess $U \ge R \cdot \text{Id}$, R > 0. Let $L = \Delta - \nabla U \cdot \nabla$ and $P_t = e^{tL}$ be the semigroup with generator L; see [Ivanisvili and Volberg 2015c; Bakry et al. 2014].

Corollary 11. For any $p \in \mathbb{R} \setminus [0, 1]$ and any smooth bounded $f \geq 0$ with $\int_{\mathbb{R}^n} f^p d\mu < \infty$ we have

$$P_t \left[f^p F_{1/(p-1)} \left(\frac{|\nabla f|}{f \sqrt{R}} \right) \right] \le (P_t f)^p F_{1/(p-1)} \left(\frac{|\nabla P_t f|}{P_t f \sqrt{R}} \right).$$

The inequality is reversed if $p \in (0, 1)$.

Proof. Notice that $\widetilde{M}(x, y) = M(x, y/\sqrt{R})$ satisfies (17). Now it remains to use inequality (2.3) from [Ivanisvili and Volberg 2015c].

By taking $t \to \infty$ and using the fact that $|\nabla P_t f| \le e^{-tR} P_t |\nabla f|$, we obtain Corollary 6.

Finally we would like to mention that having characterization (17) of functional inequalities (32) makes our approach to problem (9) systematic. Very similar *local estimates* happen to rule some *global inequalities*. We refer the reader to our recent papers on this subject [Ivanisvili and Volberg 2015a;2015b; Ivanisvili 2016].

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