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SIMON BORTZ, JOHN HOFFMAN, STEVE HOFMANN,
JOSÉ LUIS LUNA GARCÍA AND KAJ NYSTRÖM

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AND PARABOLIC UNIFORMLY RECTIFIABLE SETS**



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Let $E \subset \mathbb{R}^{n+1}$ be a parabolic uniformly rectifiable set. We prove that every bounded solution u to

$$\partial_t u - \Delta u = 0 \quad \text{in } \mathbb{R}^{n+1} \setminus E$$

satisfies a Carleson measure estimate condition. An important technical novelty of our work is that we develop a corona domain approximation scheme for E in terms of regular $\text{Lip}(1/2, 1)$ graph domains. This scheme has an analogous elliptic version which improves on the known results in that setting.

1. Introduction

For more than forty years, there has been significant interest in quantitative estimates for solutions of (linear) elliptic and parabolic partial differential equations in the absence of smoothness. In this area of research, the lack of smoothness presents itself in the structure or regularity of the coefficients of the operator, or in the geometry of the domain. Recently, sustained efforts in this area have provided characterizations of quantitative geometric notions (e.g., uniform rectifiability) in terms of quantitative estimates for harmonic functions [Garnett et al. 2018; Hofmann et al. 2016] and a geometric characterization of the L^p -solvability of the Dirichlet problem [Azzam et al. 2020]. This paper concerns the parabolic analogue of [Hofmann et al. 2016] and overcomes the substantial difficulty introduced by the distinguished time direction and the anisotropic scaling. To deal with this difficulty, we are forced to build appropriate approximating domains with better properties than would be enjoyed by the parabolic analogues of the chord-arc domains constructed in that paper. In particular, our construction improves on that of [Hofmann et al. 2016], even in the elliptic setting. We shall discuss these issues in more detail momentarily.

We shall prove the following.

Theorem 1.1 (a Carleson measure estimate for bounded caloric functions). *Let $n \geq 2$. Let $E \subset \mathbb{R}^{n+1}$ be a set which is uniformly rectifiable in the parabolic sense. Then for any solution to $(\partial_t - \Delta_X)u = 0$ in $\mathbb{R}^{n+1} \setminus E$ with $u \in L^\infty(\mathbb{R}^{n+1} \setminus E)$,*

$$\sup_{(t,X) \in E, r > 0} r^{-n-1} \iint_{B((t,X),r)} |\nabla u|^2 \delta(s, Y) dY ds \leq C \|u\|_{L^\infty(E^c)}^2, \quad (1.2)$$

where $\delta(s, Y) := \text{dist}((s, Y), E)$ and C depends only on the dimension and the parabolic uniformly rectifiable constants for E .

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Here and below, $\text{dist}((s, Y), E)$ is the parabolic distance from (s, Y) to the given set E , and the ball $B((t, X), r)$ is defined with respect to the parabolic metric; see (2.1) and (2.2) below.

In the case that Ω is an open set, the following holds.

Theorem 1.3. *Let $n \geq 2$. Let $\Omega \subset \mathbb{R}^{n+1}$ be an open set for which $\partial\Omega$ is uniformly rectifiable in the parabolic sense. Then for any solution to $(\partial_t - \Delta_X)u = 0$ in Ω with $u \in L^\infty(\Omega)$,*

$$\sup_{(t, X) \in E, r > 0} r^{-n-1} \iint_{B((t, X), r) \cap \Omega} |\nabla u|^2 \delta(s, Y) dY ds \leq C \|u\|_{L^\infty(\Omega)}^2,$$

where C depends only the dimension and the parabolic uniformly rectifiable constants for $\partial\Omega$. Here, $\delta(s, Y) := \text{dist}((s, Y), \partial\Omega)$, the parabolic distance to $\partial\Omega$.

The notion of parabolic uniform rectifiability was introduced in [Hofmann et al. 2003; 2004] and is defined below, but we first provide some context here. Through the works of Hofmann, Lewis, Murray and Silver [Hofmann 1995; 1997; Hofmann and Lewis 1996; 2005; Lewis and Murray 1995; Lewis and Silver 1988], it was shown that the *good* parabolic graphs for parabolic singular integrals and parabolic potential theory are regular $\text{Lip}(1/2, 1)$ graphs, that is, graphs which are $\text{Lip}(1/2, 1)$ (in time-space coordinates) and which possess extra regularity in time in the sense that a (nonlocal) half-order time derivative of the defining function of the graph is in the space of functions of parabolic bounded mean oscillation. This is in contrast to the elliptic setting, where one often views Lipschitz graphs as the *good* graphs for singular integrals and potential theory (because of [Coifman et al. 1982; 1983; Coifman and Semmes 1991; Dahlberg 1977; David 1991]), and where the BMO estimate for the gradient is an automatic consequence of Rademacher’s theorem and the inclusion of L^∞ in BMO. The definition of parabolic uniform rectifiability in [Hofmann et al. 2003; 2004] is given in terms of parabolic β numbers,¹ but we do not work with this definition directly here. Instead, we work with an equivalent notion of parabolic uniform rectifiability in terms of the existence of appropriate corona decompositions recently established in [Bortz et al. 2023; 2022a]. However, it is worth remarking that the graph of a $\text{Lip}(1/2, 1)$ function is parabolic uniformly rectifiable if and only if the function has a half-order time derivative in parabolic BMO. In contrast to the case of *elliptic* uniform rectifiability, which has reached a state of maturity that includes numerous interesting characterizations, this is not the case for parabolic uniform rectifiability. In fact, beyond [Hofmann et al. 2003; 2004], the only correct and more systematic studies of parabolic uniformly rectifiable sets can be found in [Bortz et al. 2023; 2022a].² In these works, parabolic uniform rectifiability is characterized in terms of a bilateral coronization by regular $\text{Lip}(1/2, 1)$ graphs (Lemma 2.14), and this characterization is the starting point for the analysis in this paper. In general there are many interesting open problems in this and related areas, and it should be emphasized that parabolic uniform rectifiability is *significantly* different to its elliptic counterpart; see [Bortz et al. 2022a, Observation 4.19].

¹These β numbers can be traced back to the work of P. Jones [1990].

²There are works of J. Rivera-Noriega in this area, but these articles have significant gaps or no proofs. Some of these gaps are outlined in [Bortz et al. 2023; 2022a; 2022b].

To give an idea of the methods involved in the proof of [Theorem 1.1](#), the primary novelty of our work is a *corona domain approximation scheme* ([Proposition 3.25](#)) in terms of regular $\text{Lip}(1/2, 1)$ graph domains. This is in contrast to the (elliptic) NTA domain approximations produced in [[Hofmann et al. 2016](#)] for uniformly rectifiable sets. In fact, our proof here carries over without modification to the elliptic setting,³ providing an (improved) approximation by Lipschitz domains. In [[Hofmann et al. 2016](#)] the authors use Whitney cubes to construct these NTA domains using *dyadic sawtooths* and exploiting an elliptic bilateral corona decomposition. The heuristic in the elliptic setting is that these *sawtooth domains* inherit many essential properties of the original boundary. In contrast, in the parabolic setting the analogous constructions do not necessarily inherit even the most basic properties. One of the most readily apparent difficulties in the parabolic setting comes from the fact that the natural *lower dimensional parabolic measure* can easily fail to see relatively nice sets. In particular, given a cube (with respect to the standard coordinates) in \mathbb{R}^{n+1} , two of the faces (those orthogonal to the time axis) have zero natural *parabolic surface measure*, which says that, not only does the boundary of a cube fail to be uniformly rectifiable in the parabolic sense, it fails even to have the Ahlfors–David regularity property. The method outlined in this paper circumvents this difficulty by *lifting* the graphs in the parabolic bilateral corona decomposition ([Lemma 2.14](#)) in a manner that respects the stopping time regimes and thereby produces the graph domains rather directly. We also point out that, while the analogous elliptic results (in [[Hofmann et al. 2016](#)]) proceed along the lines of *extrapolation of Carleson measures* it was later seen in [[Hofmann et al. 2019](#)] that this was unnecessary and a more direct approach is available. Therefore, upon proving [Proposition 3.25](#), the proof of [Theorem 1.1](#) proceeds as in [[Hofmann et al. 2019](#)].

Let us provide some motivation for the estimate in [Theorem 1.1](#). As remarked above, (elliptic) uniform rectifiability has been characterized by various properties of harmonic functions and among these characterizations is the elliptic version of the Carleson measure estimate in [Theorem 1.1](#); see [[Garnett et al. 2018](#); [Hofmann et al. 2016](#)]. We therefore expect that the estimate in [Theorem 1.1](#) is a significant step in characterizing parabolic uniform rectifiability by properties of caloric functions. We suspect that additional considerations and conditions will need to be made, as was the case for nonsymmetric operators in the elliptic setting [[Azzam et al. 2022](#)], in the converse, *free-boundary* direction due to the lack of self-adjointness of the heat operator. In domains that are sufficiently nice topologically, the estimate in [Theorem 1.1](#) (and its elliptic analogue) is also intimately tied to the solvability of the L^p -Dirichlet boundary value problem in the parabolic setting [[Dindoš et al. 2017](#); [Genschaw and Hofmann 2020](#)] (see [[Kenig et al. 2000](#); [2016](#)] and related work in [[Dindoš et al. 2011](#); [Hofmann and Le 2018](#); [Zhao 2018](#)] for the elliptic theory). Indeed, in the case of regular $\text{Lip}(1/2, 1)$ graph domains it is known that estimate (1.2) for bounded null-solutions to general parabolic operators of the form $\mathcal{L} = \partial_t - \text{div}_X A \nabla_X$ is equivalent to the solvability of the L^p -Dirichlet boundary value problem for some $p > 1$ [[Dindoš et al. 2017](#)] (boundary value problem means the data is prescribed on the lateral boundary). In fact, merely assuming parabolic Ahlfors–David regularity⁴ and a *backwards thickness condition* (also of Ahlfors–David regular type), the solvability of the L^p -Dirichlet boundary value problem is implied by a stronger estimate where the

³Except that the technical [Lemma 3.24](#) is no longer needed.

⁴In particular, without assuming that the domain is the region above a regular $\text{Lip}(1/2, 1)$ graph.

L^∞ -norm on the right-hand side of (1.2) is replaced by the (boundary) BMO norm of the data; see [Genschaw and Hofmann 2020]. This stronger estimate is unlikely to hold⁵ in the present setting due to the lack of (nontangential) accessibility to the boundary.

The rest of this paper is organized as follows. In Section 2 we introduce the notions and notation used throughout the paper. In Section 3 we construct approximating domains, each adapted to a particular stopping time regime in the parabolic bilateral corona decomposition, Lemma 2.14. In Section 4 we prove the main theorems of the paper (Theorems 1.1 and 1.3) using the constructions produced in Section 3. In Section 5 we discuss some possible extensions of the results here.

2. Preliminaries

In this paper, we work in \mathbb{R}^{n+1} identified with $\mathbb{R} \times \mathbb{R}^n = \{(t, X) : t \in \mathbb{R}, X \in \mathbb{R}^n\}$ and $n \in \mathbb{N}, n \geq 2$.⁶ We use the notation

$$\text{dist}(A, B) := \inf_{(t, X) \in A, (s, Y) \in B} |X - Y| + |t - s|^{1/2} \tag{2.1}$$

to denote the parabolic distance between A and B , with $A, B \subseteq \mathbb{R}^{n+1}$. We also use the notation $B((t, X), r)$ for the parabolic ball centered at (t, X) with radius $r > 0$, that is,

$$B((t, X), r) := \{(s, Y) : \text{dist}((t, X), (s, Y)) < r\}. \tag{2.2}$$

Given $E \subset \mathbb{R}^{n+1}$ we let $\text{diam}(E)$ denote the diameter, or parabolic diameter, defined with respect to the parabolic metric.

Definition 2.3 (parabolic Hausdorff measure). Given $s > 0$ we let \mathcal{H}_p^s denote the s -dimensional parabolic Hausdorff measure. More specifically, for a set $E \subset \mathbb{R}^{n+1}$ and $\epsilon > 0$ we define

$$\begin{aligned} \mathcal{H}_{p,\epsilon}^s(E) &:= \inf \left\{ \sum_i \text{diam}(E_i)^s : E \subseteq \bigcup_i E_i, \text{diam}(E_i) \leq \epsilon \right\}, \\ \mathcal{H}_p^s(E) &:= \lim_{\epsilon \rightarrow 0^+} \mathcal{H}_{p,\epsilon}^s(E) = \limsup_{\epsilon \rightarrow 0^+} \mathcal{H}_{p,\epsilon}^s(E). \end{aligned}$$

The following family of planes will be important in this work.

Definition 2.4 (t -independent planes). We say that an n -dimensional plane P in \mathbb{R}^{n+1} is t -independent if it contains a line in the t -direction. Equivalently, if \vec{v} is the normal vector to P , then $\vec{v} \cdot (1, \vec{0}) = 0$.

The following local energy (Caccioppoli) inequality holds for solutions to the heat equation.

Lemma 2.5 (Caccioppoli inequality). Let $B = B((t, X), r)$, and suppose that $u \in (1 + \alpha)B$ is a solution to $(\partial_t - \Delta_X)u = 0$ for some $\alpha > 0$. Then

$$\int_B |\nabla_X u(t, X)|^2 dX dt \lesssim r^{-2} \int_{(1+\alpha)B} |u|^2 dX dt,$$

where the implicit constant depends on the dimension and α .

⁵The elliptic analogue does not hold (in general) in the complement of uniformly rectifiable set.

⁶We apologize for the departure from the usual notation (X, t) , but we will often be working with graphs and it is convenient to have the *last* variable as the graph variable.

Definition 2.6 (Ahlfors–David regular). We say $E \subset \mathbb{R}^{n+1}$ is (parabolic) Ahlfors–David regular, written E is ADR, if it is closed and there exists a constant $C > 0$ such that

$$C^{-1}r^{n+1} \leq \mathcal{H}_p^{n+1}(B((t, X), r) \cap E) \leq Cr^{n+1}, \quad \forall(t, X) \in E, \quad r \in (0, \text{diam}(E)).$$

We will call the C of [Definition 2.6](#) the Ahlfors–David regularity constant and if a particular constant depends on the Ahlfors–David regularity constant, we will say that the constant *depends on ADR*. We will sometimes write $\sigma := \mathcal{H}_p^{n+1}|_E$ to denote the *surface measure* on E . (The underlying set defining σ will always be clear from the context.)

An ADR set E can be viewed as a space of homogeneous type, (E, dist, σ) , with homogeneous dimension $n + 1$. All such sets have a nice filtration, which we will refer to as the *dyadic cubes* on E .

Lemma 2.7 [[Christ 1990](#); [David and Semmes 1991](#); [1993](#); [Hytönen and Kairema 2012](#); [Hytönen and Martikainen 2012](#)]. Assume that $E \subset \mathbb{R}^{n+1}$ is (parabolic) ADR in the sense of [Definition 2.6](#) with constant C . Then E admits a parabolic dyadic decomposition in the sense that there exist constants $a_0 > 0$, $\gamma > 0$, and $c_* < \infty$, such that for each $k \in \mathbb{Z}$ there exists a collection of Borel sets, \mathbb{D}_k , which we will call (dyadic) cubes, such that

$$\mathbb{D}_k := \{Q_j^k \subset E : j \in \mathfrak{J}_k\},$$

where \mathfrak{J}_k denotes some (countable) index set depending on k , with the decomposition satisfying

- (i) $E = \cup_j Q_j^k$, for each $k \in \mathbb{Z}$.
- (ii) If $m \geq k$, then either $Q_i^m \subset Q_j^k$ or $Q_i^m \cap Q_j^k = \emptyset$.
- (iii) For each (j, k) and each $m < k$, there is a unique i such that $Q_j^k \subset Q_i^m$.
- (iv) $\text{diam}(Q_j^k) \leq c_*2^{-k}$.
- (v) Each Q_j^k contains $E \cap B((t_j^k, Z_j^k), a_02^{-k})$ for some $(t_j^k, Z_j^k) \in E$.
- (vi) $E(\{(t, Z) \in Q_j^k : \text{dist}((t, Z), E \setminus Q_j^k) \leq \varrho2^{-k}\}) \leq c_*\varrho^\gamma E(Q_j^k)$, for all k, j and for all $\varrho \in (0, \alpha)$.

Remark 2.8. We denote by $\mathbb{D} = \mathbb{D}(E)$ the collection of all Q_j^k , i.e.,

$$\mathbb{D} := \bigcup_k \mathbb{D}_k.$$

Given a cube $Q \in \mathbb{D}$, we set

$$\mathbb{D}_Q := \{Q' \in \mathbb{D} : Q' \subseteq Q\}.$$

For a dyadic cube $Q \in \mathbb{D}_k$, we let $\ell(Q) := 2^{-k}$, and we refer to this quantity as the size or side-length of Q . Evidently, $\ell(Q) \sim \text{diam}(Q)$ with constant of comparison depending at most on n and C . Note that (iv) and (v) of [Lemma 2.7](#) imply that for each cube $Q \in \mathbb{D}_k$, there is a point $(t_Q, X_Q) \in E$ and a ball $B((t_Q, X_Q), r)$ such that $r \approx 2^{-k} \approx \text{diam}(Q)$ and

$$E \cap B((t_Q, X_Q), r) \subset Q \subset E \cap B((t_Q, X_Q), Cr), \tag{2.9}$$

for some uniform constant C . We shall refer to the point (t_Q, X_Q) as the center of Q . Given a dyadic cube $Q \subset E$ and $K > 1$, we define the K dilate of Q as

$$KQ := \{(t, X) \in E : \text{dist}((t, X), E) < (K - 1) \text{diam}(Q)\}. \tag{2.10}$$

Throughout the paper we assume that E is uniformly rectifiable in the parabolic sense. We nominally define this notion in language that will be meaningful to those intimately familiar with the work of David and Semmes, but we will not discuss and introduce all the relevant terminology (the interested reader may consult [Bortz et al. 2022a, Definition 4]), as it will not be used in the present work. In fact, the reader can safely ignore the following definition, as parabolic uniform rectifiability is equivalent to the existence of a bilateral corona decomposition [Bortz et al. 2023, Theorem 3.3] (see Lemma 2.14 below) and the latter is the formulation of parabolic uniform rectifiability that we will actually use throughout the paper.

Definition 2.11 (uniformly rectifiable in the parabolic sense (P-UR)). We say a set $E \subset \mathbb{R}^{n+1}$ is uniformly rectifiable in the parabolic sense (P-UR) if E is ADR and satisfies the $(2, 2)$ geometric lemma with respect to t -independent planes and the measure \mathcal{H}_p^{n+1} ; see [Bortz et al. 2022a, Definition 4.1].⁷ We say that a constant depends on P-UR if it depends on the ADR and Carleson measure constant in the definition of the $(2, 2)$ geometric lemma (with respect to t -independent planes and the measure \mathcal{H}_p^{n+1}).

In order to state the bilateral corona decomposition, we need to define regular $\text{Lip}(1/2, 1)$ graphs and coherent subsets of dyadic cubes.

Definition 2.12 (regular $\text{Lip}(1/2, 1)$ graphs). We say that Γ is a regular $\text{Lip}(1/2, 1)$ graph if there exists a t -independent plane P and a function $\psi : P \rightarrow P^\perp$ such that

$$\Gamma = \{(p, \psi(p)) : p \in P\},$$

where, upon identifying P with $\mathbb{R}^n = \mathbb{R} \times \mathbb{R}^{n-1} = \{(t, x') : t \in \mathbb{R}, x' \in \mathbb{R}^{n-1}\}$, there exist constants b_1 and b_2 such that ψ has the following two properties:

- ψ is a $\text{Lip}(1/2, 1)$ function with constant bounded by b_1 , that is

$$|\psi(t, x') - \psi(s, y')| \leq b_1(|x' - y'| + |t - s|^{1/2}), \quad \forall (t, x'), (s, y') \in \mathbb{R}^n.$$

- ψ has a half-order time derivative in parabolic-BMO with parabolic-BMO norm bounded by b_2 , that is,

$$\|D_t^{1/2} \psi\|_{\text{P-BMO}(\mathbb{R}^n)} \leq b_2,$$

where P-BMO is the space of bounded mean oscillation with respect to parabolic balls (or cubes) and $D_t^{1/2} \psi(t, x')$ denotes the half-order time derivative. The half-order time derivative of ψ can be defined by the Fourier transform or by

$$D_t^{1/2} \psi(t, x') := \hat{c} \text{ p.v. } \int_{\mathbb{R}} \frac{\psi(s, x') - \psi(t, x')}{|s - t|^{3/2}} dt, \quad \forall t \in \mathbb{R}, \quad \forall x' \in \mathbb{R}^{n-1},$$

where \hat{c} is an appropriate constant.

⁷In [Bortz et al. 2022a], a different measure was used in place of \mathcal{H}_p^{n+1} , but these measures are equivalent when the set E is P-UR (with respect to either measure). See [Bortz et al. 2023, Corollary B.2].

Definition 2.13 (coherency [David and Semmes 1993]). Suppose E is a d -dimensional ADR set with dyadic cubes $\mathbb{D}(E)$. Let $\mathcal{S} \subset \mathbb{D}(E)$. We say that \mathcal{S} is *coherent* if the following conditions hold:

- (a) \mathcal{S} contains a unique maximal element $Q(\mathcal{S})$ which contains all other elements of \mathcal{S} as subsets.
- (b) If Q belongs to \mathcal{S} and if $Q \subset \tilde{Q} \subset Q(\mathcal{S})$, then $\tilde{Q} \in \mathcal{S}$.
- (c) Given a cube $Q \in \mathcal{S}$, either all of its children belong to \mathcal{S} , or none of them do.

We say that \mathcal{S} is *semicoherent* if only conditions (a) and (b) hold.

The following is the bilateral corona decomposition.

Lemma 2.14 [Bortz et al. 2023, Theorem 3.3]. Suppose that $E \subset \mathbb{R}^{n+1}$ is P-UR. Given any positive constant $\eta \ll 1$ and $K := \eta^{-1}$, there are constants $C_\eta = C_\eta(\eta, n, \text{ADR}, \text{P-UR})$ and $b_2 = b_2(n, \text{ADR}, \text{P-UR})$ and a disjoint decomposition $\mathbb{D}(E) = \mathcal{G} \cup \mathcal{B}$ satisfying the following properties:

- (1) The **good** collection \mathcal{G} is further subdivided into disjoint **stopping time regimes**, $\mathcal{G} = \cup_{\mathcal{S}^* \in \mathcal{S}} \mathcal{S}^*$ such that each such regime \mathcal{S}^* is coherent.
- (2) The **bad** cubes, as well as the maximal cubes $Q(\mathcal{S}^*)$ satisfy a Carleson packing condition:

$$\sum_{Q' \subset Q, Q' \in \mathcal{B}} \sigma(Q') + \sum_{\mathcal{S}^*: Q(\mathcal{S}^*) \subset Q} \sigma(Q(\mathcal{S}^*)) \leq C_\eta \sigma(Q), \quad \forall Q \in \mathbb{D}(E).$$

- (3) For each \mathcal{S}^* , there is a regular $\text{Lip}(1/2, 1)$ graph $\Gamma_{\mathcal{S}^*}$, where the function defining the graph has $\text{Lip}(1/2, 1)$ constant at most η (that is, $b_1 \leq \eta$) and whose half-order time derivative has P-BMO norm bounded by b_2 , such that, for every $Q \in \mathcal{S}^*$,

$$\sup_{(t, X) \in KQ} \text{dist}((t, X), \Gamma_{\mathcal{S}^*}) + \sup_{(s, Y) \in B_Q^* \cap \Gamma_{\mathcal{S}^*}} \text{dist}((s, Y), E) < \eta \text{diam}(Q), \tag{2.15}$$

where $B_Q^* := B(x_Q, K \text{diam}(Q))$.

Remark 2.16. Notice that if \mathcal{S} is any coherent subregime of \mathcal{S}^* ,⁸ then item (3) holds for every $Q \in \mathcal{S}$. Also, note that below we may insist that K is large, but this should be interpreted as taking η small.

Definition 2.17 (Whitney cubes and Whitney regions). Given an ADR set $E \subset \mathbb{R}^{n+1}$ we let $\mathcal{W}(E^c)$ be the standard (parabolic) Whitney decomposition of E^c , that is, $\mathcal{W}(E^c) = \{I_i\}$ is a collection of closed parabolic dyadic cubes⁹ with disjoint interiors, $\cup_{\mathcal{W}(E^c)} I_i = E^c$, and for each $I \in \mathcal{W}(E^c)$,

$$4 \text{diam}(I) \leq \text{dist}(4I, E) \leq \text{dist}(I, E) \leq 100 \text{diam}(I).$$

(A similar construction can be found in Lemma 3.24 below). For $\eta \ll 1 \ll K$ and $Q \in \mathbb{D}(E)$, we define

$$\mathcal{W}_Q(\eta, K) = \{I \in \mathcal{W}(E^c) : \eta^{1/4} \text{diam}(Q) \leq \text{dist}(I, E) \leq \text{dist}(I, Q) \leq K^{1/4} \text{diam}(Q)\}$$

and

$$\mathcal{W}_Q^*(\eta, K) = \{I \in \mathcal{W}(E^c) : \eta^4 \text{diam}(Q) \leq \text{dist}(I, E) \leq \text{dist}(I, Q) \leq K \text{diam}(Q)\}.$$

⁸This means $\mathcal{S} \subset \mathcal{S}^*$ and \mathcal{S} satisfies the coherency conditions in Definition 2.13.

⁹This means cubes from the collection of parabolic cubes in $\mathbb{R}^{n+1} = \mathbb{R} \times \mathbb{R}^n$ with vertices at the lattice points $2^{2k}\mathbb{Z} \times 2^k\mathbb{Z}^n$, for each $k \in \mathbb{Z}$.

Comparing volumes, we see that $\#\mathcal{W}_Q \leq C(n, \eta, K)$ (we use the notation $\#A$ to denote the cardinality of a finite set A). For $\eta \ll 1 \ll K$ and $Q \in \mathbb{D}(E)$, we set

$$U_Q(\eta, K) = \bigcup_{I \in \mathcal{W}_Q(\eta, K)} I \quad \text{and} \quad U_Q^*(\eta, K) = \bigcup_{I \in \mathcal{W}_Q^*(\eta, K)} I.$$

Remark 2.18. The reader may readily verify that the Whitney regions U_Q and U_Q^* have bounded overlaps, that is,

$$\sum_{Q \in \mathbb{D}(E)} 1_{U_Q}(t, X) + \sum_{Q \in \mathbb{D}(E)} 1_{U_Q^*}(t, X) \lesssim 1, \quad \forall (t, X) \in \mathbb{R}^{n+1},$$

where the implicit constant depends on the dimension, ADR, η , and K .

3. Domain approximation in stopping time regimes

In this section we assume that E has a bilateral corona decomposition and we fix S , a coherent subregime of a stopping time regime S^* in the bilateral corona decomposition (by Remark 2.16 the same estimates hold for S). Our goal is to construct a family of graphs that approximate the set E well in the sense of Lemma 2.14 (3) but have the additional property that they lie *above* (or on) the set E at the scale and location of the maximal cube Q_S . Other important properties of the construction will also be established including containment properties with respect to the Whitney regions defined above (see Definition 2.17). In the sequel we will often insist on further smallness of η depending on dimension and the ADR constant for E . Compared to [Hofmann et al. 2016], the constructions outlined in this section are the main novelties of this paper.

Let $Q_S := Q(S)$ be the maximal cube in the coherent subregime under consideration. Recall that $S \subseteq S^*$ and that there exists a regular $\text{Lip}(1/2, 1)$ graph, Γ_{S^*} , such that Lemma 2.14 (3) holds for S^* and hence also for S . Without loss of generality we may assume that the t -independent plane over which $\Gamma := \Gamma_{S^*}$ is defined is the plane $\mathbb{R}^n \times \{0\}$. Let $f : \mathbb{R}^n \rightarrow \mathbb{R}$ be the regular $\text{Lip}(1/2, 1)$ function that defines Γ_{S^*} , that is,

$$\Gamma := \Gamma_{S^*} = \{(t, x', f(t, x')) : (t, x') \in \mathbb{R}^n\}.$$

We define the \mathbb{R}^{n+1} -valued function

$$F(t, x') = (t, x', f(t, x')).$$

Inspired by [David and Semmes 1991], we define the *stopping time distance* $d : \mathbb{R}^{n+1} \rightarrow \mathbb{R}$ by¹⁰

$$d_S[(t, X)] = \inf_{Q \in \mathcal{S}} [\text{dist}((t, X), Q) + \text{diam}(Q)].$$

Given $\alpha \in [\frac{7}{8}, \frac{31}{32}]$ we introduce

$$g_\alpha(t, x') := f(t, x') + \eta^\alpha d[F(t, x')] \quad \text{and} \quad G_\alpha(t, x') := (t, x', g_\alpha(t, x')).$$

As $\alpha \in [\frac{7}{8}, \frac{31}{32}]$, below we will drop the subscript α and all constants will be independent of α . As we have fixed S , we will also drop the subscript S from d_S .

¹⁰Note that we take the stopping time distance in the *subregime*.

We first prove that g is $\text{Lip}(1/2, 1)$.

Lemma 3.1. *If $\eta^{7/8} \leq \frac{1}{2}$, then g is a $\text{Lip}(1/2, 1)$ function with constant less than $3\eta^\alpha$, and the function*

$$G(t, x') := (t, x', g(t, x'))$$

satisfies

$$\frac{1}{2}d[F(t, x')] \leq d[G(t, x')] \leq 2d[F(t, x')].$$

Proof. Note first that d is $\text{Lip}(1/2, 1)$ (on \mathbb{R}^{n+1}) with constant no more than 1, that is,

$$|d[(t, X)] - d[(s, Y)]| \leq \text{dist}((t, X), (s, Y)).$$

This follows from the fact that d is the infimum of nonnegative $\text{Lip}(1/2, 1)$ functions with constant 1. Using this we see that

$$\begin{aligned} |g(t, x') - g(s, y')| &\leq |f(t, x') - f(s, y')| + \eta^\alpha |d[(t, x', f(t, x'))] - d[(s, y', f(s, y'))]| \\ &\leq \eta[|t - s|^{1/2} + |x' - y'|] + \eta^\alpha[|t - s|^{1/2} + |x' - y'| + |f(t, x') - f(s, y')|] \\ &\leq 3\eta^\alpha[|t - s|^{1/2} + |x' - y'|]. \end{aligned}$$

To deduce the inequalities involving $d[G(t, x')]$ and $d[F(t, x')]$ we consider two cases. If $d[F(t, x')] = 0$, then $G(t, x') = F(t, x')$ so that $d[G(t, x')] = 0$. Otherwise, $d[F(t, x')] > 0$, and using that d is $\text{Lip}(1/2, 1)$ with constant 1, we have

$$|d[F(t, x')] - d[G(t, x')]| \leq \text{dist}(F(t, x'), G(t, x')) = |f(t, x') - g(t, x')| \leq \eta^\alpha d[F(t, x')] \leq \frac{1}{2}d[F(t, x')].$$

From this we easily obtain

$$\frac{1}{2}d[F(t, x')] \leq d[G(t, x')] \leq 2d[F(t, x')]. \quad \square$$

We will use the following elementary lemma several times.

Lemma 3.2. *If Γ' is the graph of a $\text{Lip}(1/2, 1)$ function φ with $\text{Lip}(1/2, 1)$ -norm less than $\frac{1}{2}$, then*

$$\frac{1}{2}|x_n - \varphi(t, x')| \leq \text{dist}((t, X), \Gamma') \leq |x_n - \varphi(t, x')|,$$

for all $(t, X) = (t, x', x_n)$.

Proof. The inequality on the right-hand side is trivial. To prove the inequality on the left-hand side, we can, after a translation, assume that $(t, x', \varphi(t, x')) = (0, 0, 0)$. Furthermore, we can without loss of generality assume that $x_n \geq 0$ (the case $x_n < 0$ is treated in the same way). Then $|x_n - \varphi(0, 0)| = x_n$. If $(s, y') \in \mathbb{R}^n$ satisfies $|y'| + |s|^{1/2} > x_n$, then

$$\text{dist}((t, X), (s, y', \varphi(s, y'))) \geq |y'| + |s|^{1/2} \geq x_n.$$

If $(s, y') \in \mathbb{R}^n$ satisfies $|y'| + |s|^{1/2} \leq x_n$, then $|\varphi(s, y')| \leq \frac{1}{2}x_n$ and hence

$$\text{dist}((t, X), (s, y', \varphi(s, y'))) \geq |x_n - \varphi(s, y')| \geq (1 - \frac{1}{2})x_n = \frac{1}{2}x_n.$$

These estimates prove the lemma. □

We will need the following properties of the stopping time distance.

Lemma 3.3. *Let $A > 1$. If $(t, X) \in \mathbb{R}^{n+1}$ satisfies $0 < 2d[(t, X)] \leq A \operatorname{diam}(Q_S)$, then there exists $Q^* \in \mathcal{S}$ such that*

$$\operatorname{dist}((t, X), Q^*) \leq 2d[(t, X)] \leq A \operatorname{diam}(Q^*) \leq C_{n, \operatorname{ADR}} d[(t, X)]. \tag{3.4}$$

If $d[(t, X)] = 0$, then there exists, for every $\epsilon \in (0, A \operatorname{diam}(Q_S))$, $Q_\epsilon \in \mathcal{S}$ such that

$$\operatorname{dist}((t, X), Q_\epsilon) \leq \epsilon < A \operatorname{diam}(Q_\epsilon) \leq C_{n, \operatorname{ADR}} \epsilon. \tag{3.5}$$

Proof. We start with proving (3.4). By definition there exists $Q \in \mathcal{S}$ such that

$$\operatorname{dist}((t, X), Q) + \operatorname{diam}(Q) \leq 2d[(t, X)].$$

Let $Q^* \in \mathcal{S}$ be the smallest cube satisfying $Q \subseteq Q^* \subseteq Q_S$ such that

$$A \operatorname{diam}(Q^*) \geq 2d[(t, X)]. \tag{3.6}$$

Such a cube exists because Q_S is a *candidate*. Notice that since Q^* contains Q , $\operatorname{dist}((t, X), Q^*) \leq \operatorname{dist}((t, X), Q) \leq 2d[(t, X)]$ which proves the first inequality in (3.4). The second inequality in (3.4) holds by the choice of Q^* . To see that the last inequality holds, we first note that if $Q^* = Q$, then $\operatorname{diam}(Q^*) = \operatorname{diam}(Q) \leq 2d[(t, X)]$ and we are done. Otherwise, the child of Q^* containing Q , namely Q' , fails to satisfy (3.6) and hence

$$A \operatorname{diam}(Q^*) \lesssim_{n, \operatorname{ADR}} A \operatorname{diam}(Q') \leq 2d[(t, X)].$$

Since $A > 1$, we have that $\operatorname{diam}(Q^*) \lesssim_{n, \operatorname{ADR}} d[(t, X)]$ (with the implicit constant independent of A). This proves (3.4).

To verify (3.5), note that by definition there exists $Q \in \mathcal{S}$ such that

$$\operatorname{dist}((t, X), Q) + \operatorname{diam}(Q) \leq \epsilon \leq A \operatorname{diam}(Q_S).$$

This allows us to repeat the argument above to produce Q_ϵ . □

Lemma 3.7. *If $(t, X) \in B((t_{Q_S}, X_{Q_S}), \frac{1}{4}K \operatorname{diam}(Q_S)) \cap E$ with $(t, X) = (t, x', x_n)$, then*

$$\operatorname{dist}((t, X), \Gamma) \lesssim \eta d[(t, X)] \quad \text{and} \quad |x_n - f(t, x')| \lesssim \eta d[(t, X)].$$

Here the implicit constants depend only on the dimension and ADR .

Proof. The second inequality follows from the first and Lemma 3.2. If $d[(t, X)] = 0$, then Lemma 3.3 gives that for $n \in \mathbb{N}$, we have $(t, X) \in K Q_{1/n}$ for some $Q_{1/n} \in \mathcal{S}$ with $\operatorname{diam}(Q_{1/n}) \approx 1/n$. Then using Lemma 2.14(3) we have $\operatorname{dist}((t, X), \Gamma) \lesssim 1/n$ for all n and hence $(t, X) \in \Gamma$. This proves the lemma in the case $d[(t, X)] = 0$.

Now assume $d[(t, X)] > 0$ and note that $2d[(t, X)] < (K - 1) \operatorname{diam}(Q_S)$ if $K > 6$. Applying Lemma 3.3, there exists Q^* such that

$$\operatorname{dist}((t, X), Q^*) \leq (K - 1) \operatorname{diam}(Q^*) \lesssim d[(t, X)].$$

Then [Lemma 2.14](#) (3) gives the desired inequality

$$\text{dist}((t, X), \Gamma) \leq \eta \text{diam}(Q^*) \lesssim \eta d[(t, X)]. \quad \square$$

Let

$$\Gamma^+ := \{(t, x', g(t, x')) : (t, x') \in \mathbb{R}^n\}$$

denote the graph of g . We first prove that we did not lose too much by modifying f and that, in fact, E lies below Γ^+ (near Q_S).

Lemma 3.8. *If $(t, X) \in B((t_{Q_S}, X_{Q_S}), \frac{1}{4}K \text{diam}(Q_S)) \cap E$ with $(t, X) = (t, x', x_n)$, then*

(a) $\frac{1}{8}\eta^\alpha d[(t, X)] \leq \text{dist}((t, X), \Gamma^+) \leq 3\eta^\alpha d[(t, X)]$, and

(b) $x_n \leq g(t, x') - \frac{1}{4}\eta^\alpha d[(t, X)]$.

Proof. If $d[(t, X)] = 0$, then $d[F(t, x')] = 0$ by [Lemma 3.7](#), and

$$(t, X) = (t, x', f(t, x')) = (t, x', g(t, x')).$$

This implies (a) and (b).

Assume $d[(t, X)] > 0$. [Lemma 3.7](#) yields the estimate

$$|x_n - f(t, x')| \leq C\eta d[(t, X)]. \tag{3.9}$$

If $C\eta < \frac{1}{2}$, then following the lines of the proof of [Lemma 3.1](#), we have that

$$\frac{1}{2}d[(t, X)] \leq d[F(t, x')] \leq 2d[(t, X)]. \tag{3.10}$$

Thus, by definition of g ,

$$g(t, x') - x_n = \eta^\alpha d[F(t, x')] + (f(t, x') - x_n) \geq \frac{1}{2}\eta^\alpha d[(t, X)] - C\eta d[(t, X)] \geq \frac{1}{4}\eta^\alpha d[(t, X)],$$

provided $C\eta \leq \frac{1}{4}\eta^\alpha$. This proves (b), and when combined with [Lemma 3.2](#), it gives the lower bound in (a).

To verify the upper bound in (a), we use (3.9) and (3.10) to write

$$|g(t, x') - x_n| \leq \eta^\alpha d[F(t, x')] + |f(t, x') - x_n| \leq 2\eta^\alpha d[(t, X)] + C\eta d[(t, X)] \leq 3\eta^\alpha d[(t, X)]. \quad \square$$

We remind the reader that we have previously defined certain Whitney regions (see [Definition 2.17](#)). We now investigate how these Whitney regions interact with the graphs we are constructing. First we need to see how they interact with the original graph Γ . As in the elliptic setting [[Hofmann et al. 2016](#)], we have the following.

Lemma 3.11. *If $Q \in \mathcal{S}$ and $I \in \mathcal{W}_Q$, then I is either above or below Γ (it does not meet Γ). Moreover, we have the estimate*

$$\text{dist}(I, \Gamma) \geq \eta^{1/2} \text{diam}(Q).$$

Proof. The first statement, about the cubes being above or below the graph, follows from the estimate. Suppose for the sake of contradiction that there exists $I \in \mathcal{W}_Q$, $Q \in \mathcal{S}$ such that $\text{dist}(I, \Gamma) < \eta^{1/2} \text{diam}(Q)$,

and let $(s, Y) \in \Gamma$ be such that $\text{dist}((s, Y), I) \leq \eta^{1/2} \text{diam}(Q)$. By construction, $\text{dist}((t, Z), (t_Q, X_Q)) \lesssim K^{1/4} \text{diam}(Q)$ for all $(t, Z) \in I$ and hence

$$\text{dist}((s, Y), (t_Q, X_Q)) \leq \eta^{1/2} \text{diam}(Q) + CK^{1/4} \text{diam}(Q) \lesssim K^{1/4} \text{diam}(Q).$$

By Lemma 2.14 (3), $\text{dist}((s, Y), E) \leq \eta \text{diam}(Q)$. Choosing $(t_0, Z_0) \in I$ such that $\text{dist}((t_0, Z_0), (s, Y)) = \text{dist}((s, Y), I) \leq \eta^{1/2} \text{diam}(Q)$, we have that

$$\begin{aligned} \text{dist}(I, E) &\leq \text{dist}((t_0, Z_0), (s, Y)) + \text{dist}((s, Y), E) \\ &\leq \eta^{1/2} \text{diam}(Q) + \eta \text{diam}(Q) \leq 2\eta^{1/2} \text{diam}(Q) < \eta^{1/4} \text{diam}(Q), \end{aligned}$$

provided $\eta^{1/4} < \frac{1}{2}$. This violates the assumption that $I \in \mathcal{W}_Q$. □

In light of Lemma 3.11, for $Q \in \mathcal{S}$ we have that $\mathcal{W}_Q = \mathcal{W}_Q^+ \cup \mathcal{W}_Q^-$, where \mathcal{W}_Q^+ is the collection of Whitney cubes above Γ and \mathcal{W}_Q^- is the collection of Whitney cubes below Γ . We then define

$$U_Q^\pm := \bigcup_{I \in \mathcal{W}_Q^\pm} I.$$

The following lemma says that U_Q^+ still lies above Γ^+ and, when $(t, X) \in U_Q^+$, the distance from (t, X) to Γ^+ is roughly the distance to E .

Lemma 3.12. *Let $Q \in \mathcal{S}$. If η is sufficiently small, then U_Q^+ lies above Γ^+ and*

$$x_n - g(t, x') \geq \frac{1}{2} \text{dist}((t, X), \Gamma), \quad \forall (t, X) = (t, x', x_n) \in U_Q^+. \tag{3.13}$$

Moreover,

$$\text{dist}((t, X), \Gamma^+) \approx \text{dist}((t, X), E), \quad (t, X) \in U_Q^+, \tag{3.14}$$

where the implicit constants depend on dimension, ADR , η , and K .

Proof. Recall that $\eta = K^{-1}$. Let $(t, X) = (t, x', x_n) \in I$ for some $I \in \mathcal{W}_Q^+$. As $\text{dist}((t, X), (t_Q, X_Q)) \lesssim K^{1/4} \text{diam}(Q)$ and $Q \in \mathcal{S}$, we have

$$\text{dist}((t, X), \Gamma) \lesssim (K^{1/4} + \eta) \text{diam}(Q) \lesssim K^{1/4} \text{diam}(Q). \tag{3.15}$$

Using Lemma 3.2,

$$|(t, X) - F(t, x')| \lesssim K^{1/4} \text{diam}(Q),$$

and therefore $\text{dist}(F(t, x'), Q) \lesssim K^{1/4} \text{diam}(Q)$. It follows that $d[F(t, x')] \lesssim K^{1/4} \text{diam}(Q)$, and using Lemma 3.11,

$$x_n - f(t, x') \geq \text{dist}((t, X), \Gamma) \geq \eta^{1/2} \text{diam}(Q) \gtrsim \eta^{1/2} K^{-1/4} d[F(t, x')] \approx \eta^{3/4} d[F(t, x')].$$

By the definition of g and the fact that $\alpha \geq \frac{7}{8}$, Lemma 3.2 implies that

$$\begin{aligned} x_n - g(t, x') &= x_n - f(t, x') - \eta^\alpha d[F(t, x')] \geq \frac{1}{2}(x_n - f(t, x')) \\ &\geq \frac{1}{2} \text{dist}((t, X), \Gamma) \geq \frac{1}{2} \eta^{1/2} \text{diam}(Q), \end{aligned} \tag{3.16}$$

where the next-to-last inequality yields (3.13), and where we have used Lemma 3.11 in the last step. In particular, U_Q^+ lies above Γ^+ . Using (3.15) and the last inequality in (3.16), and then the properties of the Whitney cubes in \mathcal{W}_Q , we have

$$\text{dist}((t, X), \Gamma) \approx_{\eta, K} \text{diam}(Q) \approx_{\eta, K} \text{dist}((t, X), E).$$

Combining (3.16) and the last displayed estimate and then using Lemma 3.2, we obtain

$$\text{dist}((t, X), \Gamma^+) \geq \frac{1}{2}(x_n - g(t, x')) \geq \frac{1}{4} \text{dist}((t, X), \Gamma) \approx_{\eta, K} \text{dist}((t, X), E)$$

and

$$\text{dist}((t, X), \Gamma^+) \leq x_n - g(t, x') \leq x_n - f(t, x') \leq 2 \text{dist}((t, X), \Gamma) \approx \text{dist}((t, X), E). \quad \square$$

We also require that close to Q_S , the region above Γ^+ shall be contained in a collection of Whitney regions associated to $Q \in \mathcal{S}$. This can be done using the Whitney regions U_Q^* .

Lemma 3.17. *Suppose $(t, X) = (t, x', x_n)$ satisfies $x_n > g(t, x')$ and*

$$(t, X) \in B((t_{Q_S}, X_{Q_S}), \frac{1}{32}K \text{diam}(Q_S)).$$

Then

$$\text{dist}((t, X), E) \geq \text{dist}((t, X), \Gamma^+), \tag{3.18}$$

and there exists $Q^* \in \mathcal{S}$ such that $(t, X) \in U_{Q^*}^*$.

Proof. Let (t, X) be as above. By Lemma 3.8, we see that $d[(t, X)] \neq 0$. To prove (3.18), we note that if (s, Y) is the closest point to (t, X) in E , then

$$\text{dist}((t, X), (s, Y)) \leq \text{dist}((t, X), (t_{Q_S}, X_{Q_S})) < \frac{1}{32}K \text{diam}(Q_S).$$

Thus $\text{dist}((s, Y), (t_{Q_S}, X_{Q_S})) < \frac{1}{16}K \text{diam}(Q_S)$, so in particular,

$$(s, Y) \in B((t_{Q_S}, X_{Q_S}), \frac{1}{4}K \text{diam}(Q_S)) \cap E.$$

By Lemma 3.8 (b), (s, Y) lies below Γ^+ and hence the line segment between (s, Y) and (t, X) meets Γ^+ . This proves (3.18).

To prove the existence of $Q^* \in \mathcal{S}$ such that $(t, X) \in U_{Q^*}^*$, we break the proof into cases.

Case 1: $x_n - g(t, x') \geq \eta^3 d[(t, X)]$. In this case, by (3.18) and Lemma 3.2,

$$\text{dist}((t, X), E) \geq \text{dist}((t, X), \Gamma^+) \geq \frac{1}{2}\eta^3 d[(t, X)].$$

Since $d[(t, X)] \leq (\frac{1}{32}K + 1) \text{diam}(Q_S) < (\frac{1}{2}(K - 1)) \text{diam}(Q_S)$, we use Lemma 3.3 to produce Q^* with

$$\text{dist}((t, X), Q^*) \leq (K - 1) \text{diam}(Q^*) \lesssim_{n, \text{ADR}} d[(t, X)].$$

Thus,

$$(t, X) \in B((t_{Q^*}, X_{Q^*}), K \text{diam}(Q^*))$$

and

$$\text{dist}((t, X), E) \gtrsim \eta^3 d[(t, X)] \gtrsim \eta^3 (K - 1) \text{diam}(Q^*) \approx \eta^2 \text{diam}(Q),$$

where the implicit constants depend on the dimension and ADR. Letting $I \in \mathcal{W}$ be such that $(t, X) \in I$, it follows that $I \in \mathcal{W}_Q^*$, provided η is sufficiently small.

Case 2: $x_n - g(t, x') \leq \eta^3 d[(t, X)]$. In this case, note that

$$\text{dist}((t, X), G(t, x')) \leq x_n - g(t, x') \leq \eta^3 d[(t, X)]. \tag{3.19}$$

Thus, since d is Lipschitz with norm 1 with respect to $\text{dist}(\cdot)$, we have, for $\eta^3 < \frac{1}{2}$,

$$\frac{1}{2}d[(t, X)] \leq d[G(t, x')] \leq 2d[(t, X)]. \tag{3.20}$$

In particular, $d[G(t, x')] > 0$. Notice then that

$$d[G(t, x')] \leq 2d[(t, X)] \leq \left(\frac{1}{16}K + 2\right) \text{diam}(Q_S) \leq \frac{1}{8}(K - 1) \text{diam}(Q_S),$$

provided that K is large enough, and [Lemma 3.3](#) then yields $Q^* \in \mathcal{S}$ such that

$$\text{dist}(G(t, x'), Q^*) \leq \frac{1}{4}(K - 1) \text{diam}(Q^*) \approx d[G(t, x')]. \tag{3.21}$$

Combining the latter estimate with [\(3.19\)](#) and [\(3.20\)](#), we see that

$$(t, X) \in B((t_{Q^*}, X_{Q^*}), \frac{1}{2}(K - 1) \text{diam}(Q^*)).$$

Claim 3.22. *For η chosen small enough ($\eta^2 < \frac{1}{2}$ will suffice at this stage),*

$$\text{dist}(G(t, x'), E) \geq \eta^2 d[G(t, x')].$$

Taking the claim for granted momentarily, by [\(3.19\)](#), [\(3.20\)](#), and [\(3.21\)](#), we have

$$\begin{aligned} \text{dist}((t, X), E) &\geq \text{dist}(G(t, x'), E) - \text{dist}((t, X), G(t, x')) \\ &\geq \text{dist}(G(t, x'), E) - \eta^3 d[(t, X)] \gtrsim \eta^2 d[(t, X)] \\ &\approx \eta^2 [G(t, x')] \gtrsim \eta^2 (K - 1) \text{diam}(Q^*) \approx \eta \text{diam}(Q^*), \end{aligned}$$

and the lemma is proved. It remains to prove [Claim 3.22](#).

Proof of Claim 3.22. Let $(s, Y) = (s, y', y_n) \in E$ be such that

$$\text{dist}(G(t, x'), (s, Y)) = \text{dist}(G(t, x'), E).$$

Assume, for the sake of obtaining a contradiction, that

$$\text{dist}(G(t, x'), (s, Y)) < \eta^2 d[G(t, x')].$$

Then for $\eta^2 < \frac{1}{2}$, since d is Lipschitz with norm 1 with respect to $\text{dist}(\cdot)$, we have

$$d[G(t, y')] \leq 2d[(s, Y)].$$

Hence, under the current assumption that $\text{dist}(G(t, x'), (s, Y)) < \eta^2 d[G(t, x')]$,

$$|y' - x'| + |t - s|^{1/2} < 2\eta^2 d[(s, Y)].$$

Since g is $\text{Lip}(1/2, 1)$ with constant $3\eta^\alpha$ (in particular less than 1),

$$\text{dist}(G(t, x'), G(s, y')) \lesssim \eta^2 d[(s, Y)].$$

Thus,

$$|y_n - g(s, y')| = \text{dist}((s, Y), G(s, y')) \lesssim \eta^2 d[(s, Y)],$$

which contradicts the conclusion of [Lemma 3.8](#), provided that

$$(s, Y) \in B((t_{Q_S}, X_{Q_S}), \frac{1}{4}K \text{diam}(Q_S)). \tag{3.23}$$

Indeed, the latter is true, as we now show. Recall that by hypothesis

$$(t, X) \in B((t_{Q_S}, X_{Q_S}), \frac{1}{32}K \text{diam}(Q_S)).$$

Moreover, in the scenario of Case 2,

$$\text{dist}((t, X), G(t, x')) = x_n - g(t, x') < \eta^3 d[(t, X)] \leq \text{diam}(Q_S),$$

and therefore

$$\text{dist}(G(t, x'), (t_{Q_S}, X_{Q_S})) < \frac{1}{16}K \text{diam}(Q_S).$$

Since (s, Y) is the closest point on E to $G(t, x')$, it must be that

$$(s, Y) \in B((t_{Q_S}, X_{Q_S}), \frac{1}{8}K \text{diam}(Q_S)).$$

In particular, [\(3.23\)](#) holds, and this proves the claim. □

Our next goal is to produce a regular version of the graphs we have constructed above. The vehicle for this regularization is the following lemma.

Lemma 3.24. *Let $h : \mathbb{R}^n \rightarrow \mathbb{R}$, $h(t, x') \geq 0$, be a $\text{Lip}(1/2, 1)$ function with $\text{Lip}(1/2, 1)$ constant (at most) 1. There exists a function $H : \mathbb{R}^n \rightarrow \mathbb{R}$ such that*

(1) $c_1 h(t, x') \leq H(t, x') \leq c_2 h(t, x')$ for all $(t, x') \in \mathbb{R}^n$.

(2) If $Z = \{(t, x') : h(t, x') = 0\}$, then

$$h(t, x')^{2m-1} |\partial_t^m H(t, x')| + h(t, x')^{m-1} |\nabla_{x'}^m H(t, x')| \leq c_{n,m}, \quad \forall (t, x') \in Z^c, \quad m \in \mathbb{N}.$$

(3) $H \in \text{Lip}(1/2, 1)$ with constant less than c_3 .

Here c_1 , c_2 , and c_3 depend on the dimension alone and $c_{n,m}$ depends on the dimension and m . Moreover, H enjoys the estimate

$$\|D_t^{1/2} H\|_{\text{P-BMO}} \leq c_4,$$

where c_4 depends only on the dimension.

The proof has many *standard* elements (if one knows where to look), but is a little lengthy. The proof can be found in the [Appendix](#).

Now we are ready to create our regularized graph. Let $h(t, x') := \frac{1}{2}d[F(t, x')]$, and let $H(t, x')$ be the function provided by Lemma 3.24.¹¹ We define two functions

$$\psi_{\eta, S}^{\pm}(t, x) := f(t, x) \pm \eta^{15/16}H(t, x').$$

We hope that it is clear to the reader that the function

$$g_{\alpha}^{-}(t, x') := f(t, x') - \eta^{\alpha}d[F(t, x')]$$

has properties analogous to those of $g_{\alpha}(t, x')$ except that g_{α}^{-} is below f and E , that the cubes in \mathcal{W}_Q^{-} are below the graph of g_{α}^{-} , etc. We next deduce that $\psi_{\eta, S}^{\pm}$ has the same properties as the functions g_{α} (and g_{α}^{-}), enumerated below.

Proposition 3.25. *Let E be uniformly rectifiable in the parabolic sense. Let $\mathbb{D}(E) = \mathcal{G} \cup \mathcal{B}$, $\mathcal{G} = \cup_{S^* \in \mathcal{S}} S^*$, and $\{\Gamma_{S^*}\}_{S^* \in \mathcal{S}}$ be the bilateral corona decomposition of E given by Lemma 2.14, with constants $\eta \ll 1$, $K = \eta^{-1}$, and b_2 . Let M_0 be the constant from Lemma 4.1 below, with $\tilde{b}_1 = 2$ and $\tilde{b}_2 = 1 + b_2$. If η is sufficiently small, depending only on the dimension and ADR, then the following holds.*

Let $S^* \in \mathcal{S}$. Then for every coherent subregime S of S^* , there is a t -independent plane P_S and two regular parabolic graphs Γ_S^{\pm} over P_S given by the functions $\psi_{\eta, S}^{\pm}$, with

$$\|\psi_{\eta, S}^{\pm}\|_{\text{Lip}(1/2, 1)} \leq C_n \eta^{15/16} \quad \text{and} \quad \|D_t^{1/2} \psi_{\eta, S}^{\pm}\|_{\text{P-BMO}} \leq (1 + b_2),$$

with the following properties (in the coordinates given by $P_S \oplus P_S^{\perp}$):

(1) If $Q \in S$, then \mathcal{W}_Q has a disjoint decomposition $\mathcal{W}_Q = \mathcal{W}_Q^+ \cup \mathcal{W}_Q^-$, and if we let $U_Q^{\pm} := \cup_{I \in \mathcal{W}_Q^{\pm}} I$, then

$$U_Q^{\pm} \subseteq B((t_{Q_S}, X_{Q_S}), K^{3/4} \text{diam}(Q_S)) \cap \{\pm x_n > \pm \psi_{\eta, S}^{\pm}(t, x')\}.$$

Here the notation $\{\pm x_n > \pm \psi_{\eta, S}^{\pm}(t, x')\}$ means

$$\{(t, X) = (t, x', x_n) : \pm x_n > \pm \psi_{\eta, S}^{\pm}(t, x')\}.$$

In particular,

$$\bigcup_{Q \in S} U_Q^{\pm} \subseteq B((t_{Q_S}, X_{Q_S}), K^{3/4} \text{diam}(Q_S)) \cap \{\pm x_n > \pm \psi_{\eta, S}^{\pm}(t, x')\}.$$

(2) If $(t, X) \in \cup_{Q \in S} U_Q^{\pm}$, then

$$\text{dist}((t, X), E) \approx_{\eta} \text{dist}((t, X), \Gamma_S^{\pm}).$$

(3) If $(t, X) \in B((t_{Q_S}, X_{Q_S}), \frac{1}{32}K \text{diam}(Q_S)) \cap \{\pm x_n > \pm \psi_{\eta, S}^{\pm}(t, x')\}$, then

$$\text{dist}((t, X), E) \geq \text{dist}((t, X), \Gamma_S^{\pm}).$$

(4) $B((t_{Q_S}, X_{Q_S}), \frac{1}{32}K \text{diam}(Q_S)) \cap \{\pm x_n > \pm \psi_{\eta, S}^{\pm}(t, x')\} \subseteq \cup_{Q \in S} U_Q^*$.

¹¹ See the proof of Lemma 3.1, from which one can easily deduce that $d[F(t, x')]$ has $\text{Lip}(1/2, 1)$ norm less than $1 + \eta$ and hence $h(t, x')$ has $\text{Lip}(1/2, 1)$ -norm less than 1. This allows one to apply Lemma 3.24.

(5) *There exist $(t_S^\pm, X_S^\pm) \in \Gamma_S^\pm$ such that*

$$B((t_{Q_S}, X_{Q_S}), M_0 K^{3/4} \text{diam}(Q_S)) \subset B((t_S^\pm, X_S^\pm), K^{7/8} \text{diam}(Q_S)) \tag{3.26}$$

and

$$B((t_S^\pm, X_S^\pm), K^{7/8} \text{diam}(Q_S)) \subset B((t_{Q_S}, X_{Q_S}), \frac{1}{32} K \text{diam}(Q_S)). \tag{3.27}$$

Proof. Let $\psi_{\eta,S}^\pm$ be as constructed before the statement of the proposition. Both $\psi_{\eta,S}^\pm$ are $\text{Lip}(1/2, 1)$ with constant less than $C\eta^{15/16}$ because f is $\text{Lip}(1/2, 1)$ with constant less than η and H has $\text{Lip}(1/2, 1)$ -norm less than $c_3 = c_3(n)$; see [Lemma 3.24](#). Similarly,

$$\|D_t^{1/2} \psi_{\eta,S}^\pm\|_{\text{P-BMO}} \leq \|D_t^{1/2} f\|_{\text{P-BMO}} + \eta^{15/16} \|D_t^{1/2} H\|_{\text{P-BMO}} \leq \|D_t^{1/2} f\|_{\text{P-BMO}} + c_4 \eta^{15/16} < b_2 + 1, \tag{3.28}$$

provided η is sufficiently small. We define Γ_S^\pm to be the graphs of $\psi_{\eta,S}^\pm$, respectively.

We claim that

$$g_{7/8}(t, x') \geq \psi_{\eta,S}^+(t, x') \geq g_{31/32}(t, x') \tag{3.29}$$

and

$$g_{7/8}^-(t, x') \leq \psi_{\eta,S}^-(t, x') \leq g_{31/32}^-(t, x'). \tag{3.30}$$

Indeed, these inequalities are a result of the fact that $\eta^{7/8} \gg \eta^{15/16} \gg \eta^{31/32}$ when η is very small along with the properties of H in relation to d . For example,

$$\psi_{\eta,S}^+(t, x') - g_{31/32}(t, x') = \eta^{15/16} H(t, x') - \eta^{31/32} d[F(t, x')].$$

Using [Lemma 3.24](#) we have

$$d[F(t, x')] = 2h(t, x') \approx_n H(t, x').$$

Since the constants are independent of η the second inequality in [\(3.29\)](#) follows. The other inequalities are treated similarly.

With [\(3.29\)–\(3.30\)](#) at hand, properties (1) and (2) can be deduced directly from [Lemmas 3.11](#) and [3.12](#). Note that to prove property (2), we observe that $(t, X) \in U_Q$ implies that (t, X) is above the graphs of both $g_{7/8}$ and $g_{31/32}$. Similarly, properties (3) and (4) can be deduced from [\(3.29\)](#) (or [\(3.30\)](#)) and [Lemma 3.17](#): to prove (3) and (4) in, e.g., the context of $\psi_{\eta,S}^+(t, x')$, we simply observe that if (t, X) is above $\psi_{\eta,S}^+(t, x')$, then it is also above $g_{31/32}$.

To prove (5), let (s, Y) be the closest point on Γ_S to (t_{Q_S}, X_{Q_S}) and observe from [Lemma 2.14](#) (3) that

$$\text{dist}((t_{Q_S}, X_{Q_S}), (s, Y)) \leq \eta \text{diam}(Q_S).$$

As $(s, Y) = (s, y', y_n) = F(s, y')$, we have that

$$H(s, y') \approx_n d[F(s, y')] \leq \text{diam}(Q_S) + \eta \text{diam}(Q_S) < 2 \text{diam}(Q_S),$$

where we have used the properties of H given by [Lemma 3.24](#). Then by definition,

$$\text{dist}((t_{Q_S}, X_{Q_S}), (s, y', \psi_{\eta,S}^\pm(s, y'))) \leq \text{dist}((t_{Q_S}, X_{Q_S}), (s, Y)) + \eta^{15/16} H(s, y') \lesssim_n \eta^{15/16} \text{diam}(Q_S).$$

Setting $(t_S^\pm, X_S^\pm) = (s, y', \psi_{\eta,S}^\pm(s, y'))$ and taking η sufficiently small (and hence K sufficiently large), we arrive at (5). □

4. Carleson measure estimates: Proof of Theorems 1.1 and 1.3

Before we get into the details of proving Theorems 1.1 and 1.3, we point out that the domains we produced in Proposition 3.25 support (a local version of) the Carleson measure estimate.

Lemma 4.1 [Hofmann and Lewis 2005, Lemma A.2]. *Let \tilde{b}_1, \tilde{b}_2 be fixed nonnegative constants. Let $\varphi(t, x')$ be a regular $\text{Lip}(1/2, 1)$ function, with $\text{Lip}(1/2, 1)$ constant \tilde{b}_1 , such that $\|D_t^{1/2}\varphi\|_{\text{P-BMO}} \leq \tilde{b}_2$. Let*

$$\Omega^+ = \{(t, X) = (t, x', x_n) : x_n > \varphi(t, x')\}.$$

Then there exist $M_0 = M_0(n, \tilde{b}_1, \tilde{b}_2) > 1$ and $c_5 = c_5(n, \tilde{b}_1, \tilde{b}_2)$, such that if u is a bounded solution to $(\partial_t - \Delta_x)u = 0$ in

$$\Omega^+((t_0, X_0), M_0r) := B((t_0, X_0), M_0r) \cap \Omega^+,$$

for some $(t_0, X_0) \in \partial\Omega$, then

$$\iint_{B((t_0, X_0), r) \cap \Omega^+} |\nabla u(s, Y)|^2 \tilde{\delta}(s, Y) dY ds \leq c_5 r^{n+1} \|u\|_{L^\infty(\Omega^+((t_0, X_0), M_0r))}^2. \tag{4.2}$$

Here $\tilde{\delta}(s, Y) = \text{dist}((s, Y), \partial\Omega^+)$. An analogous statement holds for bounded solutions to $(\partial_t - \Delta_x)u = 0$ in

$$\Omega^-((t_0, X_0), M_0r) := B((t_0, X_0), M_0r) \cap \Omega^-,$$

where

$$\Omega^- = \{(t, X) = (t, x', x_n) : x_n < \varphi(t, x')\}.$$

Proof. The lemma is a consequence of [Hofmann and Lewis 2005, Lemma A.2], henceforth abbreviated [HL, A.2]. However, to reduce the proof of the lemma to [HL, A.2] one has to note two things. First, by using the parabolic version of the Dahlberg–Kenig–Stein pullback, the operator $(\partial_t - \Delta)$ is transformed to an operator of the form treated in [HL, A.2] in the upper half-space. Furthermore, $\Omega^+((t_0, X_0), M_0r)$ is transformed into a region containing a Carleson region of size roughly M_0r . Second, while stated for solutions in the upper half-space, [HL, A.2] uses only that u is a solution in a Carleson region. \square

Now we are ready to prove Theorem 1.1.

Proof of Theorem 1.1. Let E be uniformly rectifiable in the parabolic sense and let u be a bounded solution to $(\partial_t - \Delta_x)u = 0$ in E^c . We may assume that $\|u\|_{L^\infty(E^c)} \neq 0$ since the conclusion of Theorem 1.1 holds trivially if $\|u\|_{L^\infty(E^c)} = 0$. Let

$$v := \frac{u}{\|u\|_{L^\infty(E^c)}}.$$

Then $\|v\|_{L^\infty(E^c)} = 1$ and it clearly suffices to prove the theorem with v in place of u .

For each $Q \in \mathbb{D}(E)$ we set

$$\beta_Q = \iint_{U_Q} |\nabla_X v|^2 \delta(s, Y) dY ds.$$

We first reduce the proof of the theorem to a statement concerning the β_Q .

Claim 4.3. *If there exists C (independent of v) such that*

$$\sum_{Q \subseteq Q_0} \beta_Q \leq C\sigma(Q_0), \quad \forall Q_0 \in \mathbb{D}(E), \tag{4.4}$$

then there exists C' such that

$$\sup_{(t,X) \in E, r > 0} r^{-n-1} \iint_{B((t,X),r)} |\nabla_X v|^2 \delta(s, Y) dY ds \leq C'. \tag{4.5}$$

In particular, to prove the theorem it is enough to verify (4.4)

Sketch of Proof of Claim 4.3. To prove that (4.4) implies (4.5), we select a collection $\{Q_0^i\}_i \subset \mathbb{D}(E)$ such that, for each i , $\text{diam}(Q_0^i) \approx \kappa r$, and such that the collection has uniformly bounded cardinality depending only on n and ADR. Furthermore, $B((t, X), r) \cap E \subset \cup_i Q_0^i$. Choosing κ large enough depending only on allowable parameters, we have that $B((t, X), r) \setminus E \subset \cup_i \cup_{Q \subset Q_0^i} U_Q$. We can then apply (4.4) to each $Q_0 = Q_0^i$. We omit the routine details. \square

We have now reduced everything to the setting of our dyadic machinery and we are almost ready to begin employing the constructions in Proposition 3.25. Notice that these constructions are only likely to be helpful when bounding a β_Q when Q is a *good* cube. That is why the following claim is important when handling the *bad* cubes.

Claim 4.6. *There exists a constant A depending only the dimension, K, η , and ADR such that*

$$\beta_Q \leq A\sigma(Q).$$

Sketch of Proof of Claim 4.6. The claim follows readily from Lemma 2.5 and ADR as in [Hofmann et al. 2016]. We omit the details. \square

We now prove (4.4). Fix $Q_0 \in \mathbb{D}(E)$. If $Q_0 \in S^*$ for some $S^* \in \mathcal{S}$ we let $S = S^* \cap \mathbb{D}_{Q_0}$ and note that S is a coherent subregime of S^* with maximal cube Q_0 . \mathbb{D}_{Q_0} has the disjoint decomposition

$$\mathbb{D}_{Q_0} = \{Q \in \mathcal{B} : Q \subseteq Q_0\} \cup \left(\bigcup_{S^* : Q(S^*) \subset Q_0} S^* \right) \cup S, \tag{4.7}$$

where $S = \emptyset$ if Q_0 is not in a stopping time regime (i.e., if Q_0 is a bad cube). By Lemma 2.14 (2) and Claim 4.6,

$$\sum_{\substack{Q \subseteq Q_0 \\ Q \in \mathcal{B}}} \beta_Q \leq C \sum_{\substack{Q \subseteq Q_0 \\ Q \in \mathcal{B}}} \sigma(Q) \leq CC_{\eta,K} \sigma(Q_0). \tag{4.8}$$

Let us suppose, for the moment, that we can show that

$$\sum_{Q \in S^*} \beta_Q \leq C\sigma(Q(S^*)), \tag{4.9}$$

for all S^* such that $Q(S^*) \subset Q_0$, and that

$$\sum_{Q \in S} \beta_Q \leq C\sigma(Q(S)) = C\sigma(Q_0), \tag{4.10}$$

if Q_0 is in some stopping time regime. Then, by Lemma 2.14 (2),

$$\sum_{Q \in \mathcal{S}} \beta_Q + \sum_{\mathcal{S}^*: Q(\mathcal{S}^*) \subset Q_0} \sum_{Q \in \mathcal{S}^*} \beta_Q \lesssim \sigma(Q_0) + \sum_{\mathcal{S}^*: Q(\mathcal{S}^*) \subset Q_0} \sigma(Q(\mathcal{S}^*)) \lesssim (C_{\eta, K} + 1)\sigma(Q_0).$$

Combining this estimate with (4.8) and using the decomposition of \mathbb{D}_{Q_0} in (4.7) proves (4.4) and hence the theorem. Thus it suffices to verify (4.9) and (4.10). In the following we only prove (4.9) as the only change needed when proving (4.10) is to change \mathcal{S}^* to \mathcal{S} .

To prove (4.9) we use Proposition 3.25. Fix \mathcal{S}^* such that $Q(\mathcal{S}^*) \subset Q$, and let $\psi_{\eta, \mathcal{S}^*}^\pm$ be the functions from Proposition 3.25. Then by Proposition 3.25 (1), if $Q \in \mathcal{S}^*$, then $U_Q = U_Q^+ \cup U_Q^-$ and so $\beta_Q = \beta_Q^+ + \beta_Q^-$, where

$$\beta_Q^\pm := \iint_{U_Q^\pm} |\nabla_X v|^2 \delta(s, Y) dY ds.$$

Clearly it is enough to show that

$$\sum_{Q \in \mathcal{S}^*} \beta_Q^\pm \leq C\sigma(Q(\mathcal{S}^*)).$$

We prove the estimate for the sum of the β_Q^+ leaving the straightforward modification needed to handle the sum of the β_Q^- to the interested reader. Moreover, since the $\{U_Q\}$, and hence the $\{U_Q^\pm\}$, have bounded overlap it is enough to prove that

$$\iint_{\cup_{Q \in \mathcal{S}^*} U_Q^+} |\nabla_X v|^2 \delta(s, Y) dY ds \leq C\sigma(Q(\mathcal{S}^*)). \tag{4.11}$$

Let

$$\begin{aligned} \Omega_{\mathcal{S}^*}^+ &= B((t_{\mathcal{S}^*}^+, X_{\mathcal{S}^*}^+), K^{7/8} \text{diam}(Q_{\mathcal{S}^*})) \cap \{x_n > \psi_{\eta, \mathcal{S}^*}^+(t, x')\}, \\ \tilde{\Omega}_{\mathcal{S}^*}^+ &= B((t_{Q_{\mathcal{S}^*}}, X_{Q_{\mathcal{S}^*}}), K^{3/4} \text{diam}(Q_{\mathcal{S}^*})) \cap \{x_n > \psi_{\eta, \mathcal{S}^*}^+(t, x')\}, \\ \widehat{\Omega}_{\mathcal{S}^*}^+ &= B((t_{Q_{\mathcal{S}^*}}, X_{Q_{\mathcal{S}^*}}), M_0 K^{3/4} \text{diam}(Q_{\mathcal{S}^*})) \cap \{x_n > \psi_{\eta, \mathcal{S}^*}^+(t, x')\}, \end{aligned}$$

where we recall that we use the coordinates $P_{\mathcal{S}^*} \oplus P_{\mathcal{S}^*}^\perp$ and that the notation $\{x_n > \psi_{\eta, \mathcal{S}^*}^+(t, x')\}$ means

$$\{(t, x', x_n) : x_n > \psi_{\eta, \mathcal{S}^*}^+(t, x')\}.$$

We note that $\widehat{\Omega}_{\mathcal{S}^*}^+ \subset \Omega_{\mathcal{S}^*}^+$ by (3.26). Proposition 3.25 (4) and (3.27) ensure that $\Omega_{\mathcal{S}^*}^+$ is an open subset of E^c and hence v is a solution in $\Omega_{\mathcal{S}^*}^+$. Applying Lemma 4.1 we have

$$\iint_{\tilde{\Omega}_{\mathcal{S}^*}^+} |\nabla_X v|^2 \tilde{\delta}(s, Y) dY ds \lesssim \text{diam}(Q(\mathcal{S}^*))^{n+1} \approx \sigma(Q(\mathcal{S}^*)), \tag{4.12}$$

where $\tilde{\delta}(s, Y) = \text{dist}((s, Y), \Gamma_{\mathcal{S}^*}^+)$ and $\Gamma_{\mathcal{S}^*}^+$ is the graph of $\psi_{\mathcal{S}^*}^+$. Note that if $Q \in \mathcal{S}^*$, we have $U_Q^+ \in \tilde{\Omega}_{\mathcal{S}^*}^+$ by Proposition 3.25 (1). Moreover, by Proposition 3.25 (2), we have $\tilde{\delta}(s, Y) \approx \delta(s, Y)$ in $\cup_{Q \in \mathcal{S}^*} U_Q^+$. Thus,

$$\begin{aligned} \iint_{\cup_{Q \in \mathcal{S}^*} U_Q^+} |\nabla_X v|^2 \delta(s, Y) dY ds &\approx \iint_{\cup_{Q \in \mathcal{S}^*} U_Q^+} |\nabla_X v|^2 \tilde{\delta}(s, Y) dY ds \\ &\leq \iint_{\tilde{\Omega}_{\mathcal{S}^*}^+} |\nabla_X v|^2 \tilde{\delta}(s, Y) dY ds \lesssim \sigma(Q(\mathcal{S}^*)), \end{aligned} \tag{4.13}$$

where we used (4.12) in the last inequality. This proves (4.11), and the proof of the theorem is complete. \square

The proof of [Theorem 1.3](#) is nearly identical, the only difference being that in this case one needs to take $\mathcal{W}(\Omega)$, a Whitney decomposition of Ω , instead of $\mathcal{W}(E^c)$. Modulo the following remark we leave the details to the interested reader.

Remark 4.14. As in the elliptic setting, [Theorem 1.3](#) does *not* require the corkscrew condition. On the other hand, the converse of the *elliptic* version of [Theorem 1.3](#) [[Garnett et al. 2018](#); [Azzam et al. 2022](#)] requires the additional assumption of interior corkscrews. Note that when carrying out the proof of [Theorem 1.3](#), without the corkscrew assumption it may be the case that the Whitney regions U_Q are empty for some cubes $Q \in \mathbb{D}(\partial\Omega)$, but this does not affect the analysis above.

5. Further remarks

In this section we make some remarks concerning possible extensions and consequences of [Theorem 1.1](#) and the constructions in [Proposition 3.25](#). These extensions and consequences can be proved, or, we expect that they can be proved, largely using the tools already developed in the elliptic setting. Again, we believe that the main novelty of this paper is the approximation scheme, that is, [Proposition 3.25](#).

The first observation is that solutions to the heat equation are (locally) smooth and that t -derivatives of solutions are, in fact, solutions. This allows one to produce a Caccioppoli-type inequality for the t -derivative which, in turn, allows one to improve the Carleson measure estimate in [Theorems 1.1](#) and [1.3](#) to one that includes the t -derivative. In particular, under the same hypotheses as [Theorem 1.1](#), the estimate

$$\sup_{(t,X) \in E, r > 0} r^{-n-1} \iint_{B((t,X),r)} (|\nabla u|^2 + \delta(s, Y)^2 |\partial_s u|^2) \delta(s, Y) dY ds \leq C \|u\|_{L^\infty(E^c)}^2 \tag{5.1}$$

holds with a constant C depending only the dimension and the parabolic uniformly rectifiable constants for E .

The second observation is that the proof of [Theorem 1.1](#) uses essentially three properties of u :

- (i) $u \in L^\infty(E^c)$,
- (ii) the Caccioppoli’s inequality of [Lemma 2.5](#), and
- (iii) the local square function estimate stated in [Lemma 4.1](#).

If one wants to extend the validity of [Theorem 1.1](#) to more general parabolic equations in divergence form,

$$\mathcal{L} = \partial_t - \operatorname{div}_X A(t, X) \nabla_X,$$

where A is an $n \times n$ uniformly elliptic matrix, then some regularity conditions on the coefficients need to be imposed in order to guarantee property (iii). A natural sufficient condition is the parabolic analogue of the *Kenig–Pipher condition*.¹² More specifically, this means that A satisfies

$$|\nabla_X A(s, Y)| \delta(s, Y), |\partial_t A(s, Y)| \delta^2(s, Y) \in L^\infty(\mathbb{R}^{n+1} \setminus E),$$

¹²In fact, in [[Hofmann and Lewis 2005](#), Lemma A.2], a slightly more general class of coefficients is permitted.

where $\delta(s, Y) = \text{dist}((s, Y), E)$, and that there exists a constant M such that

$$\begin{aligned} \sup_{(t, X) \in E, r > 0} r^{-n-1} \iint_{B((t, X), r)} |\nabla_X A(s, Y)|^2 \delta(s, Y) dY ds &\leq M, \\ \sup_{(t, X) \in E, r > 0} r^{-n-1} \iint_{B((t, X), r)} |\partial_t A(s, Y)|^2 \delta^3(s, Y) dY ds &\leq M. \end{aligned} \tag{5.2}$$

In particular, our results apply to this class of coefficients.

A final observation is that it seems likely that some form of ϵ -approximability [Hofmann et al. 2016; 2019] should hold in this parabolic setting along with the corresponding quantitative Fatou theorem [Bortz and Hofmann 2020]. In fact, it may be more reasonable to use the dyadic constructions from [Hofmann et al. 2016] in proving these results. Indeed, our construction here would provide some of the necessary initial estimates (Theorem 1.1), but it seems easier to deduce (parabolic) BV estimates using dyadic cubes. We also mention that it is natural to use the estimate (5.1) and to prove ϵ -approximability via a Carleson measure estimate which includes the t -derivative of the approximator. Note that this estimate was not included in [Rivera-Noriega 2003] and therefore the proof used in [Rivera-Noriega 2003, Proposition 4.3] is valid only if one works with a vertical version of the nontangential counting function \mathcal{N} (or by a spatially nontangential version based on time-slice cones with t fixed), rather than a fully nontangential version.

Appendix: Proof of Lemma 3.24

The idea is to follow Stein’s construction of the regularized distance [Stein 1970, Chapter VI, §1 & §2] and to combine this with ideas from some of the estimates produced in [David and Semmes 1991].

Proof of Lemma 3.24. Let $d := n + 1$ and $Z = \{(t, x') : h(t, x') = 0\}$. We note that Z is closed since h is continuous. We set $H(t, x') = 0$ for all $(t, x') \in Z$.

We need to define H off of Z , and following [David and Semmes 1991], we begin by producing a Whitney-type decomposition of Z^c with respect to h .¹³ For each $(t, x') \in Z^c$, we let $I_{(t, x')}$ be the largest closed (parabolic) dyadic cube containing (t, x') satisfying

$$\text{diam}(I_{(t, x')}) \leq \frac{1}{20} \inf_{(\tau, z') \in I_{(t, x')}} h(\tau, z').$$

Recall that the diameter is defined with respect to the parabolic metric. To see that such a cube exists, set $r = \frac{1}{2}h(t, x')$ and note, as h is $\text{Lip}(1/2, 1)$ with constant 1, that $h(\tau, z') \geq r$ in $B((t, x'), r)$. Therefore, any cube which contains (t, x') and which has diameter less than $\frac{1}{20}r$ is a *candidate* for $I_{(t, x')}$. We conduct this construction for each $(t, x') \in Z^c$, and we enumerate the resulting maximal cubes (without repetition) as $\{I_i\}_{i \in \Lambda}$. We note that

$$10 \text{diam}(I_i) \leq h(t, x') \leq 60 \text{diam}(I_i), \quad \forall (t, x') \in 10I_i, \tag{A.1}$$

where κI is the parabolic dilation of I by a factor of κ . Indeed, if $(t, x') \in 10I_i$, then (t, x') is at most a (parabolic) distance of $10 \text{diam}(I_i)$ from a point in I_i , and hence, using the selection criterion for I_i and

¹³In contrast to the usual Whitney decomposition, in which h is the distance to a closed set, the present version remains valid even in the case that Z is empty.

the $\text{Lip}(1/2, 1)$ condition for h ,

$$h(t, x') \geq \min_{(\tau, z') \in I_i} h(\tau, z') - 10 \text{diam}(I_i) \geq 10 \text{diam}(I_i).$$

To prove the upper bound in (A.1), we note that if I is the parent of I_i , then I fails the selection criteria. Hence there exists $(\tau, z') \in I$ such that $h(\tau, z') < 20 \text{diam}(I) = 40 \text{diam}(I_i)$, and as h is $\text{Lip}(1/2, 1)$ with constant 1 and $\text{dist}((\tau, z'), (t, x')) \leq 20 \text{diam}(I_i)$, it follows that

$$h(t, x') \leq h(\tau, z') + 20 \text{diam}(I_i) \leq 60 \text{diam}(I_i).$$

Using (A.1), we have that

$$\frac{1}{6} \text{diam}(I_j) \leq \text{diam}(I_i) \leq 6 \text{diam}(I_j) \tag{A.2}$$

whenever $10I_i \cap 10I_j \neq \emptyset$. By comparing volumes, it follows that the $\{10I_j\}$ have bounded overlap, with a constant depending on the dimension alone, that is,

$$\sum_{i \in \Lambda} 1_{I_i}(t, x') \leq N, \quad \forall (t, x') \in \mathbb{R}^n, \tag{A.3}$$

where $N = N(n)$.

Let $Q_0 = \{(t, x') \in \mathbb{R}^n : |x'|_\infty \leq \frac{1}{2}, |t| < \frac{1}{4}\}$ be the unit parabolic cube in \mathbb{R}^n . Let $\varphi \in C_0^\infty(3Q_0)$, with $0 \leq \varphi \leq 1$ and $\varphi \equiv 1$ on $2Q_0$. Clearly the $\text{Lip}(1/2, 1)$ constant of φ is bounded. For each $i \in \Lambda$, let (t_i, x'_i) be the center of I_i and $\ell(I_i)$ be the parabolic side length of I_i , that is, $\ell(I_i) = \frac{1}{2}r_i$ and

$$I_i = \{(t, x') : |x - x_i|_\infty < r_i, |t - t_i| < r_i^2\}.$$

For $i \in \Lambda_i$ we set

$$\varphi_i(t, x') = \varphi\left(\frac{t - t_i}{\ell(I_i)^2}, \frac{x - x_i}{\ell(I_i)}\right).$$

Then $0 \leq \varphi_i \leq 1$, φ_i is supported in $3I_i$, $\varphi_i \equiv 1$ on $2I_i$, φ_i is $\text{Lip}(1/2, 1)$ with constant less than $\ell(I_i)^{-1} \approx_n \text{diam}(I_i)^{-1}$, and (on all of \mathbb{R}^n),

$$\ell(I_i)^{2m} |\partial_t^m \varphi_i| + \ell(I_i) |\nabla_{x'} \varphi_i| \approx_{n,m} \text{diam}(I_i)^{2m} |\partial_t^m \varphi_i| + \text{diam}(I_i) |\nabla_{x'} \varphi_i| \lesssim \tilde{c}_{n,m}.$$

We now define

$$H(t, x') := \sum_{i \in \Lambda} \text{diam}(I_i) \varphi_i(t, x').$$

Using (A.1) we see that $3I_i$ does not meet Z for any $i \in \Lambda$, and hence $H(t, x') = 0$ for all $(t, x') \in Z$. By construction, if $(t, x') \in Z^c$ then $(t, x') \in I_j$ for some $j \in \Lambda$, and as $\varphi_j(t, x') = 1$, using also (A.1), we have that

$$H(t, x') \geq \text{diam}(I_j) \geq \frac{1}{60} h(t, x').$$

This proves the lower bound in (1). To prove the upper bound in (1) we have, by (A.2) and (A.3),

$$H(t, x') \leq \sum_{i: 3I_i \cap 3I_j \neq \emptyset} \text{diam}(Q_j) \leq 6N \text{diam}(Q_i) \leq \frac{3}{5} N h(t, x'),$$

where we used (A.1) in the last inequality. Having proved (1), we see that the proof of (2) is very similar. For instance, using the bounds for the t -derivatives of φ_i , if $(t, x') \in Z^c$, then $(t, x') \in I_j$ for some j , and hence

$$|\partial_t^m H(t, x')| \leq \tilde{c}_{n,m} \sum_{i:3I_i \cap 3I_j \neq \emptyset} \text{diam}(I_i) \text{diam}(I_i)^{-2m} \lesssim CN \text{diam}(I_j)^{-2m+1} \approx h(t, x')^{-2m+1}.$$

The bound for $|\nabla_{x'}^m H|$ has a similar proof. Finally, to see that H is $\text{Lip}(1/2, 1)$, we have

$$\begin{aligned} |H(t, x') - H(s, y')| &\leq \sum_{i:(t,x') \in 3I_i} \text{diam}(I_i) |\varphi_i(t, x') - \varphi(s, y')| + \sum_{i:(s,y') \in 3I_i} \text{diam}(I_i) |\varphi_i(t, x') - \varphi(s, y')| \\ &\leq 2c'N[|t - s|^{1/2} + |x' - y'|], \end{aligned}$$

where we used that φ_i is a $\text{Lip}(1/2, 1)$ function with constant $c' \text{diam}(I_i)^{-1}$.

Now we get to the heart of the matter, that is, proving the half-order in time regularity of H (this part is not in Stein’s book, but rather draws a great deal of inspiration from [David and Semmes 1991]). By the results in [Hofmann et al. 2003, pp. 370–373], it suffices to verify the Carleson measure estimate

$$\tilde{v}(s, y', \rho) \leq c'_4 \rho^{n+1}, \quad \forall (s, y') \in \mathbb{R}^n, \quad \rho > 0, \tag{A.4}$$

where

$$\tilde{v}(s, y', \rho) := \int_0^\rho \iint_{B((s,y'),\rho)} \hat{\gamma}(\tau, z', r)^2 d\sigma(\tau, z') \frac{dr}{r},$$

where $d\sigma(\tau, z') = \sqrt{1 + |\nabla_{z'} H(\tau, z')|^2} dz' dt$ and

$$\hat{\gamma}(\tau, z', r) := \inf_L \left[r^{-d} \iint_{B((\tau,z'),r)} \left(\frac{H(t, x') - L(x')}{r} \right)^2 d\sigma(t, x') \right]^{1/2},$$

where the infimum is taken over all affine functions L of x' only, and we recall that $d = n + 1$.

The idea behind the proof of the estimate (A.4) is as follows. If the scale r is large with respect to $h(\tau, z')$, then H is well approximated by just the linear function 0. If the scale is small with respect to $h(\tau, z')$, then, necessarily, $h(\tau, z') > 0$ and H is flat (below the scale r) by the derivative estimates (2) and therefore we can approximate H by its x' -tangent plane.

Now let us begin the proof of (A.4). Fix $(s, y') \in \mathbb{R}^n$ and $\rho > 0$. Set $h_\rho(t, x') := \min\{\frac{1}{60}h(t, x'), \rho\}$, and write

$$\begin{aligned} \tilde{v}(s, y', \rho) &= \int_0^\rho \iint_{B((s,y'),\rho)} \hat{\gamma}(\tau, z', r)^2 d\sigma(\tau, z') \frac{dr}{r} \\ &= \iint_{B((s,y'),\rho)} \int_{h_\rho(\tau,z')}^\rho \hat{\gamma}(\tau, z', r)^2 \frac{dr}{r} d\sigma(\tau, z') + \iint_{B((s,y'),\rho)} \int_0^{h_\rho(\tau,z')} \hat{\gamma}(\tau, z', r)^2 \frac{dr}{r} d\sigma(\tau, z') \\ &= T_1 + T_2. \end{aligned}$$

Let us handle term T_2 first. For (τ, z') and r in the domain of integration, $r > 0$ and $h(\tau, z') \geq 60r$. In particular, $(\tau, z') \in I_j$ for some $j \in \Lambda$, and for all such j it holds that $I_j \cap B((s, y'), \rho) \neq \emptyset$ and $r \leq \text{diam}(I_j)$

(by the right-hand estimate in (A.1)). Thus we have

$$T_2 \leq \sum_{j \in \tilde{\Lambda}} \iint_{I_j \cap B((s, y'), \rho)} \int_0^{\min\{\text{diam}(I_j), \rho\}} \hat{\gamma}(\tau, z', r)^2 \frac{dr}{r} d\sigma(\tau, z'), \tag{A.5}$$

where $\tilde{\Lambda} = \{j \in \Lambda : B((s, y'), \rho) \cap I_j \neq \emptyset\}$. Fix $j \in \tilde{\Lambda}$, $(\tau, z') \in I_j$, and $r \leq \text{diam}(I_j)$. Using the affine function

$$L_{(\tau, z')} (x') = H(\tau, z') + \nabla_{z'} H(\tau, z') \cdot (x' - z'),$$

we find by Taylor's theorem, Lemma 3.24 (2) (already proved above), and (A.1) that

$$\begin{aligned} \hat{\gamma}(\tau, z', r)^2 &\leq r^{-d} \iint_{B((\tau, z'), r)} \left(\frac{|H(t, x') - L_{(\tau, z')} (x')|}{r} \right)^2 d\sigma(t, x') \\ &\lesssim r^2 \sup_{B(x_{I_j}, 2 \text{diam}(I_j))} [|\partial_t H|^2 + |\nabla^2 H|^2] \lesssim r^2 \sup_{B(x_{I_j}, 2 \text{diam}(I_j))} h^{-2}(t, x') \\ &\lesssim r^2 \text{diam}(I_j)^{-2} \lesssim r^2 (\min\{\text{diam}(I_j), \rho\})^{-2}, \end{aligned} \tag{A.6}$$

where $x_{I_j} = (t_{I_j}, x'_{I_j})$ is the center of I_j . Using (A.6) and (A.5) we obtain

$$\begin{aligned} T_2 &\lesssim \sum_{j \in \tilde{\Lambda}} \iint_{I_j \cap B((s, y'), \rho)} \int_0^{\min\{\text{diam}(I_j), \rho\}} r^2 (\min\{\text{diam}(I_j), \rho\})^{-2} \frac{dr}{r} d\sigma(\tau, z') \\ &\lesssim \sum_{j \in \tilde{\Lambda}} \iint_{I_j \cap B((s, y'), \rho)} 1 d\sigma(\tau, z') \lesssim \sum_{j \in \tilde{\Lambda}} \sigma(I_j \cap B((s, y'), \rho)) \lesssim \sigma(B((s, y'), \rho)) \lesssim \rho^d, \end{aligned}$$

as desired.

Having obtained the desired bound for T_2 , we turn our attention to T_1 . For $(\tau, z') \in \mathbb{R}^n$ and $r > 0$, set

$$\Lambda'(\tau, z', r) := \{i \in I_i \cap B((\tau, z'), r) \neq \emptyset\}.$$

Note that in term T_1 , we have $r \in (\frac{1}{60}h(\tau, z'), \rho)$, so that $h(\hat{\tau}, \hat{z}') < 61r$ for all $(\hat{\tau}, \hat{z}') \in B((\tau, z'), r)$ because h has $\text{Lip}(1/2, 1)$ constant 1. Hence, by (A.1), we have $\text{diam}(I_i) \leq 7r \leq 7\rho$ for all $i \in \Lambda'(\tau, z', r)$. In particular, since $(\tau, z') \in B((s, y'), \rho)$,

$$\bigcup_{i \in \Lambda'(\tau, z', r)} I_i \subset B((s, y'), 10\rho). \tag{A.7}$$

For $(\tau, z') \in B((s, y'), \rho)$, with $r \in (\frac{1}{60}h(\tau, z'), \rho)$, we plug $L = 0$ into the definition of $\hat{\gamma}$ and use property (1) (which we have already established) along with (A.1) to see that

$$\begin{aligned} \hat{\gamma}(\tau, z', r)^2 &\leq r^{-d} \iint_{B((\tau, z'), r)} \left(\frac{H(t, x')}{r} \right)^2 d\sigma(t, x') \lesssim \sum_{i \in \Lambda'(\tau, z', r)} r^{-d} \iint_{I_i} \text{diam}(I_i)^2 r^{-2} d\sigma(t, x') \\ &\lesssim \sum_{i \in \Lambda'(\tau, z', r)} \left(\frac{\text{diam}(I_i)}{r} \right)^{d+2} \lesssim \sum_{i \in \Lambda'(\tau, z', r)} \left(\frac{\text{diam}(I_i)}{r} \right)^{d+1}, \end{aligned} \tag{A.8}$$

where we used the fact that $\text{diam}(I_i) \lesssim r$ in the estimate on the last line. Thus, if we let

$$\Lambda_0 := \{i \in \Lambda : I_i \subset B((s, y'), 10\rho), \text{diam}(I_i) \leq 7\rho\},$$

then using (A.7), the definition of $\Lambda'(\tau, z', r)$, and again using the fact that $\text{diam}(I_i) \leq 7r$ for $i \in \Lambda'(\tau, z', r)$, we obtain

$$\begin{aligned} T_1 &\leq \iint_{B((s, y'), \rho)} \int_{h(\tau, z')/60}^\rho \sum_{i \in \Lambda'(\tau, z', r)} \text{diam}(I_i)^{d+1} \frac{dr}{r^{d+2}} d\sigma(\tau, z') \\ &\lesssim \sum_{i \in \Lambda_0} \text{diam}(I_i)^{d+1} \int_{\text{diam}(I_i)/7}^\rho \int_{\text{dist}((\tau, z'), I_i) < r} 1 d\sigma(\tau, z') \frac{dr}{r^{d+2}} \\ &\lesssim \sum_{i \in \Lambda_0} \text{diam}(I_i)^{d+1} \int_{\text{diam}(I_i)/7}^\infty \frac{dr}{r^2} \lesssim \sum_{i \in \Lambda_0} \text{diam}(I_i)^d \lesssim \rho^d. \end{aligned}$$

This yields the desired Carleson measure estimate and concludes the proof of the lemma. \square

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References

- [Azzam et al. 2020] J. Azzam, S. Hofmann, J. M. Martell, M. Mourougolou, and X. Tolsa, “Harmonic measure and quantitative connectivity: geometric characterization of the L^p -solvability of the Dirichlet problem”, *Invent. Math.* **222**:3 (2020), 881–993. [MR](#) [Zbl](#)
- [Azzam et al. 2022] J. Azzam, J. Garnett, M. Mourougolou, and X. Tolsa, “Uniform rectifiability, elliptic measure, square functions, and ε -approximability via an ACF monotonicity formula”, *Int. Math. Res. Not.* **2022** (online publication June 2022).
- [Bortz and Hofmann 2020] S. Bortz and S. Hofmann, “Quantitative Fatou theorems and uniform rectifiability”, *Potential Anal.* **53**:1 (2020), 329–355. [MR](#) [Zbl](#)
- [Bortz et al. 2022a] S. Bortz, J. Hoffman, S. Hofmann, J. L. Luna-Garcia, and K. Nyström, “Coronizations and big pieces in metric spaces”, *Ann. Inst. Fourier (Grenoble)* **72**:5 (2022), 2037–2078. [MR](#) [Zbl](#)
- [Bortz et al. 2022b] S. Bortz, J. Hoffman, S. Hofmann, J. L. Luna-Garcia, and K. Nyström, “On big pieces approximations of parabolic hypersurfaces”, *Ann. Fenn. Math.* **47**:1 (2022), 533–571. [MR](#) [Zbl](#)
- [Bortz et al. 2023] S. Bortz, J. Hoffman, S. Hofmann, J. L. Luna-Garcia, and K. Nyström, “Corona decompositions for parabolic uniformly rectifiable sets”, *J. Geom. Anal.* **33**:3 (2023), art. id. 96. [MR](#) [Zbl](#)
- [Christ 1990] M. Christ, “A $T(b)$ theorem with remarks on analytic capacity and the Cauchy integral”, *Colloq. Math.* **60**:2 (1990), 601–628. [MR](#) [Zbl](#)
- [Coifman and Semmes 1991] R. Coifman and S. Semmes, “ L^2 estimates in nonlinear Fourier analysis”, pp. 79–95 in *Harmonic analysis* (Sendai, Japan, 1990), edited by S. Igari, Springer, 1991. [MR](#) [Zbl](#)
- [Coifman et al. 1982] R. R. Coifman, A. McIntosh, and Y. Meyer, “L’intégrale de Cauchy définit un opérateur borné sur L^2 pour les courbes lipschitziennes”, *Ann. of Math. (2)* **116**:2 (1982), 361–387. [MR](#) [Zbl](#)
- [Coifman et al. 1983] R. R. Coifman, G. David, and Y. Meyer, “La solution des conjectures de Calderón”, *Adv. Math.* **48**:2 (1983), 144–148. [MR](#) [Zbl](#)

- [Dahlberg 1977] B. E. J. Dahlberg, “Estimates of harmonic measure”, *Arch. Ration. Mech. Anal.* **65**:3 (1977), 275–288. [MR](#) [Zbl](#)
- [David 1991] G. David, *Wavelets and singular integrals on curves and surfaces*, Lecture Notes in Math. **1465**, Springer, 1991. [MR](#) [Zbl](#)
- [David and Semmes 1991] G. David and S. Semmes, *Singular integrals and rectifiable sets in \mathbb{R}^n : beyond Lipschitz graphs*, Astérisque **193**, Soc. Math. France, Paris, 1991. [MR](#) [Zbl](#)
- [David and Semmes 1993] G. David and S. Semmes, *Analysis of and on uniformly rectifiable sets*, Math. Surv. Monogr. **38**, Amer. Math. Soc., Providence, RI, 1993. [MR](#) [Zbl](#)
- [Dindoš et al. 2011] M. Dindoš, C. Kenig, and J. Pipher, “BMO solvability and the A_∞ condition for elliptic operators”, *J. Geom. Anal.* **21**:1 (2011), 78–95. [MR](#) [Zbl](#)
- [Dindoš et al. 2017] M. Dindoš, S. Petermichl, and J. Pipher, “BMO solvability and the A_∞ condition for second order parabolic operators”, *Ann. Inst. H. Poincaré C Anal. Non Linéaire* **34**:5 (2017), 1155–1180. [MR](#) [Zbl](#)
- [Garnett et al. 2018] J. Garnett, M. Mourougolou, and X. Tolsa, “Uniform rectifiability from Carleson measure estimates and ε -approximability of bounded harmonic functions”, *Duke Math. J.* **167**:8 (2018), 1473–1524. [MR](#) [Zbl](#)
- [Genschaw and Hofmann 2020] A. Genschaw and S. Hofmann, “A weak reverse Hölder inequality for caloric measure”, *J. Geom. Anal.* **30**:2 (2020), 1530–1564. [MR](#) [Zbl](#)
- [Hofmann 1995] S. Hofmann, “A characterization of commutators of parabolic singular integrals”, pp. 195–210 in *Fourier analysis and partial differential equations* (Miraflores de la Sierra, Spain, 1992), edited by J. García-Cuerva et al., CRC, Boca Raton, FL, 1995. [MR](#) [Zbl](#)
- [Hofmann 1997] S. Hofmann, “Parabolic singular integrals of Calderón-type, rough operators, and caloric layer potentials”, *Duke Math. J.* **90**:2 (1997), 209–259. [MR](#) [Zbl](#)
- [Hofmann and Le 2018] S. Hofmann and P. Le, “BMO solvability and absolute continuity of harmonic measure”, *J. Geom. Anal.* **28**:4 (2018), 3278–3299. [MR](#) [Zbl](#)
- [Hofmann and Lewis 1996] S. Hofmann and J. L. Lewis, “ L^2 solvability and representation by caloric layer potentials in time-varying domains”, *Ann. of Math. (2)* **144**:2 (1996), 349–420. [MR](#) [Zbl](#)
- [Hofmann and Lewis 2005] S. Hofmann and J. L. Lewis, “The L^p Neumann problem for the heat equation in non-cylindrical domains”, *J. Funct. Anal.* **220**:1 (2005), 1–54. [MR](#) [Zbl](#)
- [Hofmann et al. 2003] S. Hofmann, J. L. Lewis, and K. Nyström, “Existence of big pieces of graphs for parabolic problems”, *Ann. Fenn. Math.* **28**:2 (2003), 355–384. [MR](#) [Zbl](#)
- [Hofmann et al. 2004] S. Hofmann, J. L. Lewis, and K. Nyström, “Caloric measure in parabolic flat domains”, *Duke Math. J.* **122**:2 (2004), 281–346. [MR](#) [Zbl](#)
- [Hofmann et al. 2016] S. Hofmann, J. M. Martell, and S. Mayboroda, “Uniform rectifiability, Carleson measure estimates, and approximation of harmonic functions”, *Duke Math. J.* **165**:12 (2016), 2331–2389. [MR](#) [Zbl](#)
- [Hofmann et al. 2019] S. Hofmann, J. M. Martell, and S. Mayboroda, “Transference of scale-invariant estimates from Lipschitz to non-tangentially accessible to uniformly rectifiable domains”, preprint, 2019. [arXiv 1904.13116](#)
- [Hytönen and Kairema 2012] T. Hytönen and A. Kairema, “Systems of dyadic cubes in a doubling metric space”, *Colloq. Math.* **126**:1 (2012), 1–33. [MR](#) [Zbl](#)
- [Hytönen and Martikainen 2012] T. Hytönen and H. Martikainen, “Non-homogeneous Tb theorem and random dyadic cubes on metric measure spaces”, *J. Geom. Anal.* **22**:4 (2012), 1071–1107. [MR](#) [Zbl](#)
- [Jones 1990] P. W. Jones, “Rectifiable sets and the traveling salesman problem”, *Invent. Math.* **102**:1 (1990), 1–15. [MR](#) [Zbl](#)
- [Kenig et al. 2000] C. Kenig, H. Koch, J. Pipher, and T. Toro, “A new approach to absolute continuity of elliptic measure, with applications to non-symmetric equations”, *Adv. Math.* **153**:2 (2000), 231–298. [MR](#) [Zbl](#)
- [Kenig et al. 2016] C. Kenig, B. Kirchheim, J. Pipher, and T. Toro, “Square functions and the A_∞ property of elliptic measures”, *J. Geom. Anal.* **26**:3 (2016), 2383–2410. [MR](#) [Zbl](#)
- [Lewis and Murray 1995] J. L. Lewis and M. A. M. Murray, *The method of layer potentials for the heat equation in time-varying domains*, Mem. Amer. Math. Soc. **545**, Amer. Math. Soc., Providence, RI, 1995. [MR](#) [Zbl](#)

- [Lewis and Silver 1988] J. L. Lewis and J. Silver, “Parabolic measure and the Dirichlet problem for the heat equation in two dimensions”, *Indiana Univ. Math. J.* **37**:4 (1988), 801–839. [MR](#) [Zbl](#)
- [Rivera-Noriega 2003] J. Rivera-Noriega, “Absolute continuity of parabolic measure and area integral estimates in non-cylindrical domains”, *Indiana Univ. Math. J.* **52**:2 (2003), 477–525. [MR](#) [Zbl](#)
- [Stein 1970] E. M. Stein, *Singular integrals and differentiability properties of functions*, Princeton Math. Ser. **30**, Princeton Univ. Press, 1970. [MR](#) [Zbl](#)
- [Zhao 2018] Z. Zhao, “BMO solvability and A_∞ condition of the elliptic measures in uniform domains”, *J. Geom. Anal.* **28**:2 (2018), 866–908. [MR](#) [Zbl](#)

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SIMON BORTZ: sbortz@ua.edu

Department of Mathematics, University of Alabama, Tuscaloosa, AL, United States

JOHN HOFFMAN: jlh82b@mail.missouri.edu

Department of Mathematics, University of Missouri at Columbia, Columbia, MO, United States

STEVE HOFMANN: hofmanns@missouri.edu

Department of Mathematics, University of Missouri at Columbia, Columbia, MO, United States

JOSÉ LUIS LUNA GARCÍA: lunagaj@mcmaster.ca

Department of Mathematics & Statistics, McMaster University, Hamilton, Canada

KAJ NYSTRÖM: kaj.nystrom@math.uu.se

Department of Mathematics, Uppsala University, Uppsala, Sweden

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Strong semiclassical limits from Hartree and Hartree–Fock to Vlasov–Poisson equations	891
LAURENT LAFLECHE and CHIARA SAFFIRIO	
Marstrand–Mattila rectifiability criterion for 1-codimensional measures in Carnot groups	927
ANDREA MERLO	
Finite-time blowup for a Navier–Stokes model equation for the self-amplification of strain	997
EVAN MILLER	
Eigenvalue bounds for Schrödinger operators with random complex potentials	1033
OLEG SAFRONOV	
Carleson measure estimates for caloric functions and parabolic uniformly rectifiable sets	1061
SIMON BORTZ, JOHN HOFFMAN, STEVE HOFMANN, JOSÉ LUIS LUNA GARCÍA and KAJ NYSTRÖM	