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**REGULARITY FOR NONLOCAL EQUATIONS
WITH LOCAL NEUMANN BOUNDARY CONDITIONS**

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We establish fine results on the boundary behavior of solutions to nonlocal equations in $C^{k,\gamma}$ domains which satisfy local Neumann conditions on the boundary. Such solutions typically blow up at the boundary like $v \asymp \text{dist}^{s-1}$ and are sometimes called large solutions. In this setup we prove optimal regularity results for the quotients v/dist^{s-1} , depending on the regularity of the domain and on the data of the problem. The results of this article will be important in a forthcoming work on nonlocal free boundary problems.

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1. Introduction

The study of nonlocal operators of the form

$$Lv(x) = \text{p.v.} \int_{\mathbb{R}^n} (v(x) - v(x+h))K(h) dh, \tag{1-1}$$

where $K : \mathbb{R}^n \rightarrow [0, \infty]$ is a kernel satisfying for some $s \in (0, 1)$

$$K(h) = \frac{K(h/|h|)}{|h|^{n+2s}}, \quad 0 < \lambda \leq K(\theta) \leq \Lambda \quad \text{for all } \theta \in \mathbb{S}^{n-1}, \quad K(h) = K(-h) \tag{1-2}$$

has been an important area of research in analysis and probability for the past 30 years. Operators L of the type (1-1)–(1-2) arise naturally as generators of $2s$ -stable Lévy processes, and are used to model different kinds of real-world phenomena involving long range interactions, e.g., in mathematical finance and in physics. From a PDE perspective, it is of particular interest to study the effect of the nonlocality of L on the regularity of solutions to nonlocal equations. By now, the question of *interior* regularity of solutions

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is fairly well-understood, and there are several important works in this context, such as [Caffarelli and Silvestre 2009; 2011a; 2011b; Silvestre 2006; Bass and Levin 2002; Kassmann 2009; Di Castro et al. 2014; 2016; Barrios et al. 2014; Ros-Oton and Serra 2016b].

A much more delicate question is the one of *boundary* regularity of solutions to nonlocal problems. Previous works have mostly focused on nonlocal Poisson problems, given as

$$\begin{cases} Lv = f & \text{in } \Omega, \\ v = 0 & \text{in } \mathbb{R}^n \setminus \Omega. \end{cases} \quad (1-3)$$

The nonlocal Poisson problem (1-3) arises naturally as the Euler–Lagrange equation of a nonlocal energy minimization problem and can therefore be studied via variational methods, but also via nonvariational methods. For (1-3) it was proved (see [Ros-Oton and Serra 2014; Grubb 2015]) that weak solutions satisfy $v \in C^s(\bar{\Omega})$, once $\partial\Omega \in C^{1,\gamma}$ and $f \in L^\infty(\Omega)$. The C^s regularity of solutions is optimal, as one can see from the explicit example (see [Gettoor 1961; Landkof 1972; Dyda 2012])

$$(-\Delta)^s(1 - |x|^2)_+^s = c_{n,s} > 0 \quad \text{in } B_1, \quad (1-4)$$

which also remains valid for L satisfying (1-1)–(1-2) (see [Ros-Oton 2016]). However, it turns out that once the domain, the kernel, and the data are regular enough, also the quotient v/d^s will be regular, yielding a fine description of the behavior of the solution v at the boundary. The best known result in the literature, establishing optimal boundary regularity of weak solutions of (1-3) in terms of the regularity of the domain and the data was shown in [Ros-Oton and Serra 2017; Abatangelo and Ros-Oton 2020; Grubb 2015] (see also [Ros-Oton and Serra 2016a; 2016b; Abels and Grubb 2023]) and reads as

$$\partial\Omega \in C^{k+1,\gamma}, \quad f \in C^{k+\gamma-s}(\bar{\Omega}) \quad \implies \quad \frac{v}{d^s} \in C^{k,\gamma}(\bar{\Omega}) \quad \text{for all } k \in \mathbb{N} \cup \{0\}, \quad \gamma \in (0, 1). \quad (1-5)$$

All the previously mentioned results on the nonlocal Poisson problem (1-3) address weak solutions for which one can prove that they must remain bounded in $\bar{\Omega}$ (see [Servadei and Valdinoci 2014; Korvenpää et al. 2016]). However, explicit computations reveal that there also exist pointwise solutions of (1-3), which explode at the boundary of the domain behaving asymptotically like d^{s-1} . The following most prominent example goes back to a work by Hmissi [1994] (see also [Bogdan 1999, Example 1, p. 239; Bogdan et al. 2009, Example 3.3; Dyda 2012]):

$$(-\Delta)^s(1 - |x|^2)_+^{s-1} = 0 \quad \text{in } B_1. \quad (1-6)$$

The example (1-6) has initiated the conceptual study of boundary blow-up for solutions to nonlocal equations (see [Grubb 2014; 2015; 2018; 2023; Abatangelo 2015; 2017; Abatangelo et al. 2023; Chan et al. 2021]). In this theory, solutions such as (1-6) are sometimes called “large solutions”. Due to the explosion at the boundary, the above function cannot be a weak solution, and clearly violates (1-5).

In order to have a unified framework which also allows for singular behavior at the boundary, it is necessary to keep track of the boundary behavior of the solution, or more precisely to prescribe somehow the behavior of the quotient v/d^{s-1} . In this spirit, the following Neumann problem, which was introduced

in [Grubb 2014] (see also [Grubb 2018; 2023]), can be seen as a generalization of (1-3)

$$\begin{cases} Lv = f & \text{in } \Omega, \\ v = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \partial_\nu \left(\frac{v}{d^{s-1}} \right) = g & \text{on } \partial\Omega, \end{cases} \tag{1-7}$$

where $\nu(x_0) \in \mathbb{S}^{n-1}$ denotes the inner unit normal at $x_0 \in \partial\Omega$. The problem (1-7) is a natural *nonlocal Neumann problem* with inhomogeneous Neumann data g , and one can show that the problem is well-posed in suitable function spaces, at least if the domain is C^∞ (see [Grubb 2014]). Moreover, the solutions blow up at every boundary point where v/d^{s-1} does not vanish.

Remark 1.1. The functions in (1-4) and (1-6), are both solutions to (1-7), with $g \equiv 1$ and $f = c_{n,s}$ and with $g \equiv (s - 1)2^{s-2}$ and $f = 0$, respectively, in case $\Omega = B_1$.

The Neumann condition in (1-7) is purely local in nature in the sense that it is imposed only on the topological boundary $\partial\Omega$. Therefore, (1-7) is conceptually completely different from the nonlocal Neumann problem introduced in [Du et al. 2012; Dipierro et al. 2017] (see also [Alves and Torres Ledesma 2020; Vondraček 2021; Audrito et al. 2023; Foghem and Kassmann 2024; Grube and Hensiek 2024]). It is also of entirely different nature than [Barles et al. 2014a; 2014b; Bogdan et al. 2003; Chen and Kim 2002], where local boundary conditions are imposed, but instead the operator is changed, depending on the domain.

Main result. The aforementioned regularity results (1-5) from [Ros-Oton and Serra 2017; Abatangelo and Ros-Oton 2020] do not apply to (1-7) since solutions are in general not continuous and might even explode at the boundary. However, it is natural to expect fine regularity results for the quotients v/d^{s-1} depending on the regularity of the domain and the data.

When Ω is C^∞ and $K|_{\mathbb{S}^{n-1}}$ is C^∞ , the regularity theory for (1-7) was developed by Grubb [2014] using an approach via pseudodifferential operators.

Our goal in this work is twofold: to establish sharp boundary regularity estimates for (1-7) in $C^{k,\gamma}$ domains, and at the same time to prove them for the first time as localized estimates in $\Omega \cap B_2$. This is new even for C^∞ domains, and it is crucial for our application to free boundary problems.

Our main result is the following:

Theorem 1.2. *Let $L, K, s, \lambda, \Lambda$ be as in (1-1)–(1-2). Let $k \in \mathbb{N}$, $\gamma \in (0, 1)$ with $\gamma \neq s$, and $\Omega \subset \mathbb{R}^n$ be a $C^{k+1,\gamma}$ domain, and $K \in C^{2k+2\gamma+3}(\mathbb{S}^{n-1})$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to*

$$\begin{cases} Lv = f & \text{in } \Omega \cap B_2, \\ v = 0 & \text{in } B_2 \setminus \Omega, \\ \partial_\nu \left(\frac{v}{d^{s-1}} \right) = g & \text{on } \partial\Omega \cap B_2, \end{cases}$$

where $\nu : \partial\Omega \rightarrow \mathbb{S}^{n-1}$ is the normal vector of Ω , and $f \in C(\Omega) \cap \mathcal{X}(\Omega \cap B_2)$, $g \in C^{k-1+\gamma}(\partial\Omega \cap B_2)$,

$$\mathcal{X}(\Omega \cap B_2) = \begin{cases} d^{s-\gamma} L^\infty(\Omega \cap B_2) & \text{if } k + \gamma \leq 2s, \\ C^{k-2s+\gamma}(\Omega \cap B_2) & \text{if } k + \gamma > 2s. \end{cases} \tag{1-8}$$

Then, it holds that $v/d^{s-1} \in C_{\text{loc}}^{k+\gamma}(\bar{\Omega} \cap B_2)$ and

$$\left\| \frac{v}{d^{s-1}} \right\|_{C^{k,\gamma}(\bar{\Omega} \cap B_1)} \leq c \left(\left\| \frac{v}{d^{s-1}} \right\|_{L^\infty(\Omega \cap B_2)} + \|v\|_{L_{2s}^1(\mathbb{R}^n \setminus B_2)} + \|f\|_{\mathcal{X}(\Omega \cap B_2)} + \|g\|_{C^{k-1+\gamma}(\partial\Omega \cap B_2)} \right),$$

for some $c > 0$, which only depends on $n, s, \lambda, \Lambda, k, \gamma, \Omega$, and $\|K\|_{C^{2k+2\gamma+3}(\mathbb{S}^{n-1})}$.

For the definition of $L_{2s}^1(\mathbb{R}^n)$ and the notion of viscosity solutions, we refer to Section 2.

The regularity we obtain for v/d^{s-1} depending on the regularity of the domain Ω and the data f, g is expected to be optimal. For f and g , this is an immediate consequence of interior Schauder theory (see [Ros-Oton and Serra 2016b]), and the order of the equation. For the regularity of the domain, we observe that our results align with the ones in [Abatangelo and Ros-Oton 2020] once $v \in C(\bar{\Omega} \cap B_2)$. We obtain results with regularity assumptions on K that are expected to be optimal in case Ω is a half-space (see Theorem 1.7). As in [Grubb 2014], we rule out the case $\gamma = s$. The result is expected to be false in this case. It corresponds to proving Schauder-type regularity estimates of integer order.

Another key advantage of our approach is that it allows for localized results in $\Omega \cap B_2$. Nonetheless, if $\Omega \subset B_2$, and v is a solution to (1-7), by application of the maximum principle (see Lemma 3.4) to the estimate in Theorem 1.2 we can obtain the following bound which is purely in terms of f and g :

$$\left\| \frac{v}{d^{s-1}} \right\|_{C^{k,\gamma}(\bar{\Omega})} \leq c(\|f\|_{\mathcal{X}(\Omega)} + \|g\|_{C^{k-1+\gamma}(\partial\Omega)}).$$

Thus, we have the following generalization of (1-5) to solutions of (1-7):

$$\begin{aligned} \partial\Omega \in C^{k+1,\gamma}, \quad f \in C^{k-2s+\gamma}(\bar{\Omega}), \quad g \in C^{k-1+\gamma}(\partial\Omega) \\ \implies \frac{v}{d^{s-1}} \in C^{k,\gamma}(\bar{\Omega}) \quad \text{for all } k \in \mathbb{N}, \gamma \in (0, 1). \end{aligned} \quad (1-9)$$

A weak maximum principle and nonlocal problems with local Dirichlet conditions. The example (1-6) of a nontrivial s -harmonic function that vanishes outside B_1 implies that the Poisson problem (1-3) for the fractional Laplacian is ill-posed even in the homogeneous case. Therefore, maximum principles are usually established under an additional assumption on the boundary behavior of the solution, ruling out “large” solutions such as (1-6) (see [Silvestre 2007; Servadei and Valdinoci 2014; Felsinger et al. 2015; Jarohs and Weth 2019; Feulefack and Jarohs 2023; Fernández-Real and Ros-Oton 2024a]). Note that a similar phenomenon occurs for local equations, where any constant function is a pointwise solution inside the solution domain.

In this paper, we prove the following nonlocal weak maximum principle, which allows for solutions that blow up at the boundary.

Proposition 1.3. *Let $L, K, s, \lambda, \Lambda$ be as in (1-1)–(1-2). Let $\gamma > 0$ and $\Omega \subset \mathbb{R}^n$ be a $C^{1,\gamma}$ domain. Let $v \in L_{2s}^1(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to*

$$\begin{cases} Lv \geq 0 & \text{in } \Omega, \\ v \geq 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \frac{v}{d^{s-1}} \geq 0 & \text{on } \partial\Omega. \end{cases}$$

Then, $v \geq 0$.

The condition $v/d^{s-1} \geq 0$ in Proposition 1.3 includes solutions that blow up at the boundary, such as (1-6). Previously, maximum principles including large solutions have been established in [Abatangelo 2015; Grube and Hensiek 2023; Liu and Zhuo 2025; Li and Liu 2023]. Proposition 1.3 extends these results to general $2s$ -stable integrodifferential operators, and to $C^{1,\gamma}$ domains, respectively.

Recall that a natural way to make the nonlocal Poisson problem (1-3) well-posed is to impose Neumann boundary conditions as in (1-7). Another way would be to prescribe the limit of the quotient v/d^{s-1} directly, which leads to the following nonlocal problem with local Dirichlet data, which was introduced independently in [Grubb 2014; Abatangelo 2015]:

$$\begin{cases} Lv = f & \text{in } \Omega, \\ v = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \frac{v}{d^{s-1}} = h & \text{on } \partial\Omega. \end{cases} \quad (1-10)$$

The weak maximum principle in Proposition 1.3 implies that the problems (1-10) and (1-3) are equivalent, when $h \equiv 0$. Thus, (1-10) can be seen as an inhomogeneous nonlocal Dirichlet problem.

Another contribution of this article is the following Schauder-type boundary regularity estimate for solutions to nonlocal equations with local Dirichlet data:

Theorem 1.4. *Let $L, K, s, \lambda, \Lambda$ be as in (1-1)–(1-2). Let $k \in \mathbb{N}, \gamma \in (0, 1)$ with $\gamma \neq s$, and $\Omega \subset \mathbb{R}^n$ be a $C^{k+1,\gamma}$ domain, and $K \in C^{2k+2\gamma+3}(\mathbb{S}^{n-1})$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to (1-10) with $f \in C(\Omega) \cap \mathcal{X}(\Omega)$ and $h \in C^{k+\gamma}(\partial\Omega)$, where \mathcal{X} is as in (1-8).*

Then, it holds that $v/d^{s-1} \in C^{1+\gamma}_{\text{loc}}(\bar{\Omega})$, and

$$\left\| \frac{v}{d^{s-1}} \right\|_{C^{k,\gamma}(\bar{\Omega})} \leq c(\|f\|_{\mathcal{X}(\Omega)} + \|h\|_{C^{k+\gamma}(\partial\Omega)})$$

for some $c > 0$, which only depends on $n, s, \lambda, \Lambda, k, \gamma, \Omega$, and $\|K\|_{C^{2k+2\gamma+3}(\mathbb{S}^{n-1})}$.

We refer to [Grubb 2015; 2023] for similar results in the framework of pseudodifferential operators.

Note that (1-10) can always be reduced to the homogeneous problem (1-3). In fact, if Ω and h are regular enough, one can extend h to a smooth function in $\bar{\Omega}$ and consider $w := v - d^{s-1}h$. Then, w solves the homogeneous problem (1-3) with a new right-hand side $\tilde{f} = f - L(d^{s-1}h)$. Since $L(d^{s-1}h)$ has good regularity properties (see Corollary 2.5), we can prove Theorem 1.4, by application of the results in [Ros-Oton and Serra 2017; Abatangelo and Ros-Oton 2020].

Strategy of the proof: regularity for nonlocal problems with local Neumann data. Since the nonlocal problem with inhomogeneous local Dirichlet data (1-10) can always be reduced to the homogeneous problem (1-3) for which the boundary regularity theory was already established (see [Ros-Oton and Serra 2017; Abatangelo and Ros-Oton 2020]), the proof of Theorem 1.4 is rather simple.

In sharp contrast to that, for the Neumann problem (1-7) there is no cheap way to obtain the boundary regularity results in Theorem 1.2 from the existing theory. In fact, it is already highly nontrivial to establish Hölder continuity of the quotient v/d^{s-1} up to the boundary (see Theorem 1.6 below).

Our proof of Theorem 1.6 goes in *three main steps*.

First, we establish a weak maximum principle for solutions to the Neumann problem (1-7).

Proposition 1.5. *Let $L, K, s, \lambda, \Lambda$ be as in (1-1)–(1-2). Let $\gamma > 0$, $\Omega \subset \mathbb{R}^n$ be a $C^{2,\gamma}$ domain, and $K \in C^{5+2\gamma}(\mathbb{S}^{n-1})$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to*

$$\begin{cases} Lv \geq 0 & \text{in } \Omega, \\ v \geq 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \partial_\nu \left(\frac{v}{b_\Omega} \right) \leq 0 & \text{on } \partial\Omega, \end{cases}$$

where b_Ω is defined in (3-2). Then $v \geq 0$.

This result seems to be the first maximum principle for nonlocal problems with local Neumann boundary conditions in the literature. We believe it to be of independent interest and refer to Lemma 3.4 for a corresponding L^∞ bound in the case of inhomogeneous data. The function b can be thought of as a special regularized distance function taken to the power $s - 1$. We stress that the result is no longer true if the function b is replaced by \tilde{d}^{s-1} , where \tilde{d} is another regularized distance function. In fact, Proposition 1.5 holds true for the function in (1-6) if $b = (1 - |\cdot|)_+^{s-1}$, but fails if we replace b by the regularized distance $\tilde{d} = (1 - |\cdot|^4)$.

The proof of Proposition 1.5 follows from a nonlocal Hopf-type lemma for solutions to the inhomogeneous Dirichlet problem (1-10) (see Lemma 3.3), which in turn follows from the weak maximum principle in Proposition 1.3. All of these results rely heavily on explicit barriers for (1-10) in $C^{1,\gamma}$ domains that are adapted to the geometry of the domain and blow up at the boundary like d^{s-1} . These barriers can be seen as perturbations of (1-6), or rather of 1D solutions such as

$$L(x_n)_+^{s-1} = 0 \quad \text{in } \{x_n > 0\}. \tag{1-11}$$

Note that (1-11) follows simply by differentiating the equation

$$L(x_n)_+^s = 0 \quad \text{in } \{x_n > 0\}.$$

The previous identity is a classical fact for nonlocal operators (1-1)–(1-2) (see [Fernández-Real and Ros-Oton 2024a, Lemma 2.6.2]).

The *second* main step in the proof of Theorem 1.6 is to establish Hölder continuity of order α , for $\alpha \in (0, 1)$ small enough, up to the boundary of v/d^{s-1} for solutions to (1-7) in $C^{1,\gamma}$ domains.

Theorem 1.6. *Let $L, K, s, \lambda, \Lambda$ be as in (1-1)–(1-2). Let $\gamma \in (0, 1)$, $\Omega \subset \mathbb{R}^n$ be a $C^{2,\gamma}$ domain, and $K \in C^{5+2\gamma}(\mathbb{S}^{n-1})$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to*

$$\begin{cases} Lv = f & \text{in } \Omega \cap B_2, \\ v = 0 & \text{in } B_2 \setminus \Omega, \\ \partial_\nu \left(\frac{v}{d^{s-1}} \right) = g & \text{on } \partial\Omega \cap B_2, \end{cases}$$

with $f \in C(\Omega \cap B_2)$ and $g \in C(\partial\Omega \cap B_2)$. Then, there exists $\alpha_0 > 0$, such that when $d^{s-\alpha} f \in L^\infty(\Omega \cap B_2)$ for some $\alpha \in (0, \alpha_0]$, and it holds that $v/d^{s-1} \in C^\alpha_{\text{loc}}(\bar{\Omega} \cap B_2)$, and

$$\left\| \frac{v}{d^{s-1}} \right\|_{C^\alpha(\bar{\Omega} \cap B_1)} \leq c \left(\left\| \frac{v}{d^{s-1}} \right\|_{L^\infty(\Omega \cap B_2)} + \|v\|_{L^1_{2s}(\mathbb{R}^n \setminus B_2)} + \|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_2)} + \|g\|_{L^\infty(\partial\Omega \cap B_2)} \right),$$

where $c > 0$ and α_0 depend only on $n, s, \lambda, \Lambda, \gamma$, and the $C^{2,\gamma}$ radius of Ω .

The proof of Theorem 1.6 uses the weak maximum principle in Proposition 1.5 and the interior weak Harnack inequality, to establish a weak Harnack inequality for v/d^{s-1} at the boundary (see Lemma 4.1). This allows us to deduce a so called “growth lemma” for v/d^{s-1} , stating that v/d^{s-1} must be large pointwise in a ball centered at the boundary, if v/d^{s-1} was large in a measure-theoretic sense in a ball away from the boundary. Such growth lemma allows to establish oscillation decay for v/d^{s-1} at the boundary, and to deduce the Hölder estimate in Theorem 1.6. A similar proof for the classical Laplacian can be found in [Lian and Zhang 2023].

Once the boundary Hölder estimate is shown, we can establish the higher order boundary regularity in Theorem 1.2 via a blow-up argument. This is the *third*, and last step of the proof. Theorem 1.6 is crucial in order to deduce uniform convergence of the blow-up sequence.

The blow-up argument follows the scheme in [Abatangelo and Ros-Oton 2020] and relies on a Liouville theorem in the half-space with local Neumann data (see Theorem 5.1). However, major modifications have to be made in most of the steps due to the boundary blow-up of solutions. For instance, we need to show the following new result (see Corollary 2.5):

$$\partial\Omega \in C^{k+1,\gamma} \implies L(d^{s-1}) \in C^{k-1+\gamma-s}(\bar{\Omega}) \quad \text{if } k + \gamma > 1 + s.$$

Moreover, the presence of a Neumann boundary condition complicates some of the arguments, such as the proof of a stability result for viscosity solutions (see Lemma 2.13). Finally, as in [Abatangelo and Ros-Oton 2020] we need to make use of a suitable notion of nonlocal equations up to a polynomial (see [Dipierro et al. 2019; 2022]) in order to account for solutions that grow too fast at infinity (see Definition 2.8).

Applications to free-boundary problems. We end the discussion of the main results of this article by shedding some light on a, perhaps unexpected, connection between nonlocal problems with local Neumann boundary data and free boundary problems. This connection is a main motivation for us to study (1-7). Let us explain this phenomenon in the particular case of the fractional Laplacian.

The nonlocal one-phase free boundary problem, which was introduced in [Caffarelli et al. 2010] (see also [Ros-Oton and Weidner 2024]), deals with the minimization of the functional

$$\mathcal{I}(w) := \iint_{(B_1^c \times B_1^c)^c} (w(x) - w(y))^2 \frac{dy dx}{|x - y|^{n+2s}} + M|\{w > 0\} \cap B_1| \tag{1-12}$$

for some $M > 0$ and with prescribed values of w in $\mathbb{R}^n \setminus B_1$. One can show (see [Caffarelli et al. 2010; Fernández-Real and Ros-Oton 2024b]) that local minimizers of (1-12) are $C^s(B_1)$ and that they are viscosity solutions to

$$\begin{cases} (-\Delta)^s w = 0 & \text{in } \Omega \cap B_1, \\ w = 0 & \text{in } B_1 \setminus \Omega, \\ \frac{w}{d^s} = c_{n,s}M & \text{on } \partial\Omega \cap B_1, \end{cases} \tag{1-13}$$

where $c_{n,s} > 0$ is a constant and $\Omega := \{w > 0\}$. An important question in the theory is to determine the regularity of the free boundary $\partial\Omega$ near so called “regular points”. These are the points $x_0 \in \partial\Omega \cap B_1$ for

which blow-ups of w are half-space solutions, i.e., (up to rotations and multiplicative constants)

$$\frac{w(x_0 + rx)}{r^s} \rightarrow w_0(x) := (x_n)_+^s \quad \text{locally uniformly.}$$

One can show using the extension for $(-\Delta)^s$ (see [De Silva and Roquejoffre 2012; De Silva and Savin 2012; De Silva et al. 2014]) that once a sequence (w_ε) of viscosity solutions (1-13) is “ ε -close” to the half-space solution w_0 in the sense that

$$(x_n - \varepsilon)_+^s \leq w_\varepsilon(x) \leq (x_n + \varepsilon)_+^s,$$

then it holds, as $\varepsilon \searrow 0$, that

$$\frac{w_\varepsilon(x) - (x_n)_+^s}{\varepsilon} \rightarrow (x_n)_+^{s-1} u(x),$$

where u solves the so called “linearized problem”

$$\begin{cases} (-\Delta)^s((x_n)_+^{s-1} u) = 0 & \text{in } \{x_n > 0\} \cap B_1, \\ \partial_n u = 0 & \text{on } \{x_n = 0\} \cap B_1. \end{cases} \tag{1-14}$$

Hence, $(x_n)_+^{s-1} u$ is a solution to a nonlocal problem with local Neumann data (1-7) in the half-space, and it explodes at the boundary $\{x_n = 0\} \cap B_1$.

In order to establish regularity results for the free boundary $\Omega = \{w > 0\}$ near regular points, it is an important step to establish boundary regularity results for the solution to the linearized problem. For (1-14) this was done in [De Silva and Roquejoffre 2012; De Silva and Savin 2012; De Silva et al. 2014], using the Caffarelli–Silvestre extension.

In the light of this connection, our main result Theorem 1.2 also makes a contribution to the theory of the nonlocal one-phase problem (1-12), and provides a completely new proof of the regularity for (1-14), even in the case of the fractional Laplacian.

We end this discussion by stating a variant of Theorem 1.2 in the special case $\Omega = \{x_n > 0\}$. This result holds true under assumptions on the regularity of K which are expected to be optimal, and it will be helpful in the study of the nonlocal one-phase free boundary problem (1-13) with respect to general nonlocal operators (1-1)–(1-2), which we plan to investigate in a future work (see [Ros-Oton and Weidner 2025]).

Theorem 1.7. *Let $L, K, s, \lambda, \Lambda$ be as in (1-1)–(1-2). Let $k \in \mathbb{N}, \gamma \in (0, 1)$ with $\gamma \neq s$. Let*

$$u \in C(\{x_n \geq 0\} \cap B_2) \quad \text{with } (x_n)_+^{s-1} u \in L_{2s}^1(\mathbb{R}^n)$$

be a viscosity solution to

$$\begin{cases} L((x_n)_+^{s-1} u) = f & \text{in } \{x_n > 0\} \cap B_2, \\ \partial_n u = g & \text{on } \partial\{x_n = 0\} \cap B_2. \end{cases}$$

with $f \in C(\{x_n > 0\} \cap B_2) \cap \mathcal{X}(\{x_n > 0\} \cap B_2)$, $g \in C^{k-1+\gamma}(\{x_n = 0\} \cap B_2)$, and $K \in C^{k-2s+\gamma}(\mathbb{S}^{n-1})$ if $k + \gamma > 2s$, where \mathcal{X} is as in (1-8). Then, it holds that

$$\|u\|_{C^{k,\gamma}(\{x_n \geq 0\} \cap B_1)} \leq c \left(\|u\|_{L^\infty(\{x_n > 0\} \cap B_2)} + \|(x_n)_+^{s-1} u\|_{L_{2s}^1(\mathbb{R}^n \setminus B_2)} + \|f\|_{\mathcal{X}(\{x_n > 0\} \cap B_2)} + \|g\|_{C^{k-1+\gamma}(\{x_n = 0\} \cap B_2)} \right)$$

for some $c > 0$, which only depends on $n, s, \lambda, \Lambda, k, \gamma$, and (if $k + \gamma > 2s$) also on $\|K\|_{C^{k-2s+\gamma}(\mathbb{S}^{n-1})}$.

Finally, we make the following remark.

Remark 1.8. The following two problems are equivalent if $v \in C(\bar{\Omega} \cap B_2)$, i.e., if solutions do not blow up on $\partial\Omega \cap B_2$:

$$\left\{ \begin{array}{l} Lv = f \quad \text{in } \Omega \cap B_2, \\ v = 0 \quad \text{in } B_2 \setminus \Omega, \\ \partial_\nu \left(\frac{v}{d^{s-1}} \right) = g \quad \text{on } \partial\Omega \cap B_2, \end{array} \right. \Leftrightarrow \left\{ \begin{array}{l} Lv = f \quad \text{in } \Omega \cap B_2, \\ v = 0 \quad \text{in } B_2 \setminus \Omega, \\ \frac{v}{d^s} = g \quad \text{on } \partial\Omega \cap B_2. \end{array} \right.$$

Indeed, since $v \equiv 0$ in $B_2 \setminus \Omega$, it holds for any $x_0 \in \partial\Omega \cap B_2$ that

$$\partial_\nu \left(\frac{v}{d^{s-1}} \right) = \lim_{x \rightarrow x_0} \frac{\frac{v}{d^{s-1}}(x) - \lim_{z \rightarrow x_0} \frac{v}{d^{s-1}}(z)}{d(x)} = \lim_{x \rightarrow x_0} \frac{v}{d^s}(x).$$

Recall that the second problem is satisfied by minimizers to the nonlocal one-phase problem (1-13). Moreover, the above problem is the nonlocal counterpart of the over-determined Serrin’s problem whenever $\Omega \subset B_2$ (see for instance [Fall and Jarohs 2015; Soave and Valdinoci 2019; Biswas and Jarohs 2020; Dipierro et al. 2023]).

Organization of the paper. In Section 2 we introduce the notion of viscosity solutions to (1-7) and give some preliminary lemmas. Among them are already several new results of independent interest, such as the construction of explicit barriers exploding at the boundary (see Section 2.3), an analysis of the regularity of $L(d^{s-1})$ in terms of the regularity of the domain (see Corollary 2.5), and a stability result for viscosity solutions (see Lemma 2.13). In Section 3 we prove maximum principles for solutions to nonlocal problems with local Dirichlet and Neumann data (see Propositions 1.3 and 1.5). Section 4 is devoted to the proof of the Hölder estimate up to the boundary (see Theorem 1.6). In Section 5 we prove a Liouville theorem in the half-space (see Theorem 5.1), and in Section 6 we carry out a blow-up argument to prove our main result, Theorem 1.2. Finally, Section 7 contains the proof of the regularity for the inhomogeneous Dirichlet problem (see Theorem 1.4).

2. Preliminaries

In this section, we give several important definitions, such as the definitions of viscosity solutions to (1-7). In Section 2.2 we establish the regularity of $L(d^{s-1})$ depending on the regularity of the domain and in Section 2.3 we use these results to construct barrier functions. In Section 2.4, we introduce the notion of nonlocal equations satisfied up to a polynomial, and in Section 2.5 we establish stability of viscosity solutions and prove that the sum of two viscosity solutions is again a viscosity solution.

From now on, we denote by $\mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$ the class of operators (1-1) with kernels satisfying (1-2). Moreover, whenever we say $K \in C^\alpha(\mathbb{S}^{n-1})$ for some $\alpha > 0$, we mean that $\|K\|_{C^\alpha(\mathbb{S}^{n-1})} \leq \Lambda$. Sometimes, we denote the class of operators (1-1) satisfying (1-2) and $K \in C^\alpha(\mathbb{S}^{n-1})$ by $\mathcal{L}_s^{\text{hom}}(\lambda, \Lambda, \alpha)$.

Moreover, given an open, bounded domain $\Omega \subset \mathbb{R}^n$ with $\partial\Omega \in C^\beta$ for some $\beta > 1$, $d := d_\Omega : \mathbb{R}^n \rightarrow [0, \infty)$ will denote the regularized distance which satisfies $d \in C^\infty(\Omega) \cap C^\beta(\bar{\Omega})$ and $d \equiv 0$ in $\mathbb{R}^n \setminus \Omega$. Crucially, we have $\text{dist}(\cdot, \Omega) \leq d \leq C \text{dist}(\cdot, \Omega)$ in \mathbb{R}^n , i.e., the topological distance and the regularized distance are

pointwise comparable. We will often use the fact that $|D^k d| \leq cd^{\beta-k}$ (see [Fernández-Real and Ros-Oton 2024a, Definition 2.7.5]). Throughout this article, we will define $d^{s-1} \equiv 0$ in $\mathbb{R}^n \setminus \Omega$.

In the following, whenever $x_0 \in \partial\Omega$, we write $v/d^{s-1}(x_0) := \lim_{\Omega \ni x \rightarrow x_0} v/d^{s-1}(x)$.

2.1. Function spaces and solution concepts. Let us introduce the following function space:

$$L^1_\alpha(\mathbb{R}^n) := \left\{ u : \|u\|_{L^1_\alpha(\mathbb{R}^n)} := \int_{\mathbb{R}^n} |u(y)|(1+|y|)^{-n-\alpha} dy < \infty \right\}, \quad \alpha > 0.$$

Typically, we will use the previous definition with $\alpha = 2s$. We are now in a position to give the notion of viscosity solution to (1-7).

Definition 2.1 (viscosity solution). Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{1,\gamma}$. By $\nu \in \mathbb{S}^{n-1}$, we denote the inner normal vector to $\partial\Omega$.

(i) We say that $v \in C(\Omega) \cap L^1_{2s}(\mathbb{R}^n)$ is a viscosity subsolution to

$$Lv = f \quad \text{in } \Omega \cap B_1, \quad (2-1)$$

where $f \in C(\Omega \cap B_1)$, if for any $x \in \Omega \cap B_1$ and any neighborhood $N_x \subset \Omega$ of x it holds that

$$L\phi(x) \leq f(x) \quad \text{for all } \phi \in C^2(N_x) \cap L^1_{2s}(\mathbb{R}^n) \quad \text{such that } v(x) = \phi(x), \quad \phi \geq v. \quad (2-2)$$

We say that v is a viscosity supersolution to (2-2) if (2-2) holds true for $-v$ and $-f$ instead of v and f . Moreover, v is a viscosity solution to (2-2), if it is a viscosity subsolution and a viscosity supersolution.

(ii) For any function $b \in L^1_{2s}(\mathbb{R}^n)$ with $b/d^{s-1} \in C^1(\bar{\Omega})$ we say that $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ is a viscosity subsolution to

$$\partial_\nu(v/b) = g \quad \text{on } \partial\Omega \cap B_1,$$

where $g \in C(\partial\Omega \cap B_1)$, if for any $x \in \partial\Omega \cap B_1$ and any neighborhood $N_x \subset \bar{\Omega} \cap B_1$ of x it holds that

$$\partial_\nu\phi(x) \leq g(x) \quad \text{for all } \phi \in C^2(N_x) \cap L^\infty(\bar{\Omega}) \quad \text{such that } v/b(x) = \phi(x), \quad \phi \leq v/b. \quad (2-3)$$

We say that v is a viscosity supersolution to (2-3) if (2-3) holds true for $-v$ and $-g$ instead of v and g . Moreover, v is a viscosity solution to (2-3), if it is a viscosity subsolution and a viscosity supersolution.

Clearly, if in (i) $Lv(x)$, or if in (ii) $\partial_\nu(v/d^{s-1})(x) = \lim_{\Omega \ni y \rightarrow x} (v/d^{s-1})(y)$ exists in the strong sense, then the notions of viscosity solutions coincide with the ones for strong solutions (see [Fernández-Real and Ros-Oton 2024a, Lemma 3.4.13]).

2.2. Nonlocal operators and the distance function. The goal of this subsection is to establish several lemmas on the regularity of $L(d^{s-1})$ depending on the regularity of Ω . Lemma 2.3 will help us to establish barriers in $C^{1,\gamma}$ domains and Corollary 2.5 is crucial for domains that are more regular.

The following lemma is a slight modification of [Fernández-Real and Ros-Oton 2024a, Lemma B.2.4].

Lemma 2.2. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be a bounded Lipschitz domain with Lipschitz constant L and $C^{0,1}$ radius $\rho_0 > 0$. Let $x_0 \in \Omega$ with $\rho := d_\Omega(x_0)$, $\gamma > -1$ and $\gamma < \beta$. Then,*

$$\int_{\Omega \setminus B_{\rho/2}} d_\Omega^\gamma(x_0 + y) |y|^{-n-\beta} dy \leq C(1 + \rho^{\gamma-\beta})$$

for some constant $C > 0$, depending only on $n, \gamma, \beta, \rho_0, L$, and, if $\gamma > 0$ or $\beta \leq 0$ also on $\text{diam}(\Omega)$.

Proof. We assume that $x_0 = 0$. By [Fernández-Real and Ros-Oton 2024a, Lemma B.2.4], there exists $\kappa > 0$ such that for any $t \in (0, \kappa)$,

$$\mathcal{H}^{n-1}(\{d = t\} \cap (B_{2^{j+1}\rho} \setminus B_{2^j\rho})) \leq C(2^j\rho)^{n-1}. \tag{2-4}$$

Note that

$$\int_{(\Omega \setminus B_{\rho/2}) \cap \{d \geq \kappa\}} d^\gamma(y) |y|^{-n-\beta} dy \leq (\text{diam}(\Omega)^\gamma \mathbb{1}_{\{\gamma > 0\}} + \kappa^\gamma \mathbb{1}_{\{\gamma \leq 0\}}) \int_{(\Omega \setminus B_{\rho/2}) \cap \{d \geq \kappa\}} |y|^{-n-\beta} dy \leq c$$

for some constant $c > 0$ depending on κ and, if $\gamma > 0$ or $\beta \leq 0$ also on $\text{diam}(\Omega)$, independent of ρ . The independence of ρ is trivial if $\kappa \leq 2\rho$ since then $\Omega \setminus B_{\rho/2} \subset \Omega \setminus B_{\kappa/4}$, and otherwise, it follows from the fact that $B_r \cap \{d \geq \kappa\} = \emptyset$ once $r \leq \kappa/2 \leq \kappa - \rho$ (recall that $d(0) = \rho$), so also in this case, we can replace the domain of integration by $\Omega \setminus B_{\kappa/2}$. Moreover, using (2-4) and the coarea formula,

$$\begin{aligned} \int_{(\Omega \setminus B_{\rho/2}) \cap \{d \leq \kappa\}} d^\gamma(y) |y|^{-n-\beta} dy &\leq c \sum_{j \geq 1} \left((2^j\rho)^{-n-\beta} \int_{(B_{2^{j+1}\rho} \setminus B_{2^j\rho}) \cap \{d \leq \kappa\}} d^\gamma(y) |\nabla d(y)| dy \right) \\ &\leq c \sum_{j \geq 1} \left((2^j\rho)^{-n-\beta} \int_0^{\min\{2^j\rho, \kappa\}} t^\gamma \left[\int_{(B_{2^{j+1}\rho} \setminus B_{2^j\rho}) \cap \{d=t\}} d\mathcal{H}^{n-1}(y) \right] dt \right) \\ &\leq c \sum_{j \geq 1} ((2^j\rho)^{-\beta+\gamma}) \leq c\rho^{\gamma-\beta} \end{aligned}$$

for some $c > 0$, where we used that $\gamma - \beta < 0$. □

The following lemma will be of central importance for the proof of Lemmas 2.6 and 2.7.

Lemma 2.3. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{1,\gamma}$ for some $\gamma > 0$. Then, for any $\delta \in (0, s)$, there exists $c_1 > 0$, depending only on $n, s, \lambda, \Lambda, \Omega, \gamma, \delta$, and the $C^{1,\gamma}$ radius of Ω , such that*

$$|L(d^{s-1})| \leq c_1 d^{\delta\gamma-s-1} \quad \text{in } \Omega.$$

Moreover, for any $\varepsilon \in (0, s)$, there exist $c_2, c_3 > 0$ depending only on $n, s, \lambda, \Lambda, \gamma, \varepsilon$, and the $C^{1,\gamma}$ radius of Ω , such that

$$-L(d^{s-1+\varepsilon}) \leq -c_2 d^{\varepsilon-s-1} + c_3 \quad \text{in } \Omega.$$

The first claim follows in a similar way as [Fernández-Real and Ros-Oton 2024a, Proposition B.2.1].

Proof. We let $x_0 \in \Omega$ and write $\rho = d(x_0)$. Then, we let

$$l(x) = (d(x_0) + \nabla d(x_0) \cdot (x - x_0))_+$$

and observe that

$$L(l^{s-1}) = 0 \quad \text{in } \{l > 0\},$$

as a consequence of $L(l^s) = 0$ and $\nabla l^s = s l^{s-1} \nabla l = s \nabla d(x_0) l^{s-1}$. Next, we claim that

$$|d^{s-1} - l^{s-1}|(x_0 + y) \leq \begin{cases} C \rho^{s+\gamma-3} |y|^2 & \text{in } B_{\rho/2}, \\ C |y|^{(1+\gamma)\delta} |d^{s-1-\delta}(x_0 + y) + l^{s-1-\delta}(x_0 + y)| & \text{in } \mathbb{R}^n \setminus B_{\rho/2}. \end{cases} \quad (2-5)$$

From here, we can compute

$$\begin{aligned} |L(d^{s-1})(x_0)| &= |L(d^{s-1} - l^{s-1})(x_0)| \\ &\leq C \rho^{s+\gamma-3} \int_{B_{\rho/2}} |y|^{2-n-2s} \, dy \\ &\quad + C \int_{(x_0+\Omega) \setminus B_{\rho/2}} |y|^{-n-2s+(1+\gamma)\delta} |d^{s-1-\delta}(x_0 + y) + l^{s-1-\delta}(x_0 + y)| \, dy \\ &\leq C(1 + \rho^{\gamma-s-1} + \rho^{\gamma\delta-s-1}), \end{aligned}$$

where we applied Lemma 2.2 to d and to l with $s-1-\delta =: \gamma < \beta := 2s - (1+\gamma)\delta$ (choosing $\gamma \in (0, s)$ so small that $\beta > 0$), in order to estimate the third integral. Since this estimate implies the first result, it remains to verify the claim (2-5). In case $x \in B_{\rho/2}(x_0)$, we estimate

$$|d^{s-1} - l^{s-1}|(x) \leq |d - l|(x) \|d^{s-2} + l^{s-2}\|_{L^\infty(B_{\rho/2}(x_0))} \leq c \|D^2 d\|_{L^\infty(B_{\rho/2}(x_0))} |x_0 - x|^2 \rho^{s-2} \leq C \rho^{s+\gamma-3} |y|^2.$$

Here, we used that $|D^2 d| \leq C d^{-1+\gamma}$ by [Fernández-Real and Ros-Oton 2024a, Lemma B.0.1] and that $l \geq c\rho$ in $B_{\rho/2}(x_0)$. The latter statement follows since by the $C^{1,\gamma}$ regularity of d , it must be

$$|d(x) - d(x_0) - \nabla d(x_0) \cdot (x - x_0)| \leq C \rho^{1+\gamma} \quad \text{for all } x \in B_{\rho/2}(x_0),$$

due to Taylor's formula, and therefore $d(x)$ and ρ are comparable in $B_{\rho/2}(x_0)$, which yields for small enough ρ for some $c > 0$,

$$l(x) \geq d(x_0) + \nabla d(x_0) \cdot (x - x_0) \geq d(x) - C \rho^{1+\gamma} \geq c \rho > 0 \quad \text{for all } x \in B_{\rho/2}(x_0).$$

We can always assume that $\rho > 0$ is small, since otherwise, the result follows by the regularity of d^{s-1} away from the boundary of Ω .

Next, for $x \in \mathbb{R}^n \setminus B_{\rho/2}(x_0)$, we make use of the following algebraic inequality, which follows from the C^δ regularity of the function $t \mapsto t^{s-1-\delta}$ in $[\min\{a, b\}, \max\{a, b\}]$,

$$|a^{s-1} - b^{s-1}| \leq c |a - b|^\delta |a^{s-1-\delta} + b^{s-1-\delta}| \quad \text{for all } a, b > 0,$$

for any $\delta \in (0, s)$ and some $c > 0$, depending only on s, δ , which allows us to estimate

$$\begin{aligned} |d^{s-1}(x) - l^{s-1}(x)| &\leq c |d(x) - l(x)|^\delta |d^{s-1-\delta}(x) + l^{s-1-\delta}(x)| \\ &\leq c |x_0 - x|^{(1+\gamma)\delta} |d^{s-1-\delta}(x) + l^{s-1-\delta}(x)|, \end{aligned}$$

where we used that by [Fernández-Real and Ros-Oton 2024a, Lemma B.2.2] it holds

$$|d(x) - l(x)| \leq C |x_0 - x|^{1+\gamma}.$$

This proves the first claim.

Now, we turn to the proof of the second result. First, we observe that by similar arguments as in the first part of the proof, we obtain

$$|d^{\varepsilon+s-1} - l^{\varepsilon+s-1}|(x_0 + y) \leq \begin{cases} C\rho^{\varepsilon+s+\gamma-3}|y|^2 & \text{in } B_{\rho/2}, \\ C|y|^{(1+\gamma)\delta}||d^{\varepsilon+s-1-\delta}(x_0 + y) + l^{\varepsilon+s-1-\delta}(x_0 + y)| & \text{in } \mathbb{R}^n \setminus B_{\rho/2}, \end{cases}$$

and therefore

$$|L(d^{\varepsilon+s-1} - l^{\varepsilon+s-1})(x_0)| \leq C(1 + \rho^{\varepsilon+\gamma-s-1} + \rho^{\varepsilon+\gamma\delta-s-1}).$$

We claim that for any $e \in \mathbb{S}^{n-1}$ it holds that

$$\begin{cases} L((x \cdot e)_+^{\varepsilon+s-1}) = c_e(x \cdot e)_+^{\varepsilon-s-1} & \text{in } \{x \cdot e > 0\}, \\ (x \cdot e)_+^{s-1+\varepsilon} = 0 & \text{in } \{x \cdot e \leq 0\} \end{cases} \quad (2-6)$$

for some $c_e \in [c_-, c_+]$, where $c_+ > c_- > 0$ depend only on n, s, λ, Λ . Once we have shown the claim (2-6), we can conclude the proof, since it implies

$$\begin{aligned} -L(d^{\varepsilon+s-1})(x_0) &\leq -L(l^{\varepsilon+s-1})(x_0) + |L(d^{\varepsilon+s-1} - l^{\varepsilon+s-1})(x_0)| \\ &\leq -c\rho^{\varepsilon-s-1} + C(1 + \rho^{\varepsilon+\gamma-s-1} + \rho^{\varepsilon+\gamma\delta-s-1}) \leq -c\rho^{\varepsilon-s-1} + C. \end{aligned}$$

Hence, it remains to prove (2-6). By the $2s$ -homogeneity of L we can apply [Fernández-Real and Ros-Oton 2024a, Lemmas B.1.5 and 1.10.3(iii)] and deduce

$$L((x \cdot e)_+^{\varepsilon+s-1}) = c_1(-\Delta)_{\mathbb{R}}^s(x \cdot e)_+^{\varepsilon+s-1} = c_1c_2(x \cdot e)_+^{\varepsilon-s-1}$$

for some constant $c_1 > 0$ and where c_2 is given by, see [Fall and Ros-Oton 2022, Lemma 2.4],

$$c_2 = (-\Delta)_{\mathbb{R}}^s(l_+^{\varepsilon+s-1})(1) = \frac{\Gamma(s + \varepsilon)}{\Gamma(-s + \varepsilon)} \frac{\sin(\pi(-1 + \varepsilon))}{\sin(\pi(-1 - s + \varepsilon))} > 0.$$

This concludes the proof. □

The following lemma is crucial in the proofs of Lemma 2.13, and in Section 6. It follows by differentiating the corresponding results in [Abatangelo and Ros-Oton 2020].

Lemma 2.4. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $k \in \mathbb{N}$, and $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{k+1,\gamma}$ for some $\gamma \in (0, 1)$ with $\gamma \neq s$, and $0 \in \partial\Omega$. Assume that $K \in C^{2k+2\gamma+3}(\mathbb{S}^{n-1})$. Let*

$$\eta \in C^{k,\gamma}(\overline{\Omega \cap B_1}) \cap C^\infty(\Omega \cap B_1).$$

Then, there exists $c > 0$, depending only on $n, s, \lambda, \Lambda, \Omega, \gamma, k$, such that the following holds true:

(i) *If $k = 1$ and $\gamma < s$, then*

$$|L(d^{s-1}(\nabla d)\eta)| \leq c(|\cdot| + |\eta(0)| + |\nabla\eta(0)|)d^{\gamma-s} \quad \text{in } \overline{\Omega} \cap B_{1/2}.$$

(ii) *If $k \geq 2$ or $\gamma > s$, then*

$$[L(d^{s-1}(\nabla d)\eta)]_{C^{k-1-s+\gamma}(\overline{\Omega \cap B_{1/2}})} \leq c(|\cdot| + |\eta(0)| + |\nabla\eta(0)|).$$

(iii) *If $k + \gamma > 2s$, then we have for any $x_0 \in \Omega \cap B_{1/2}$*

$$[L(d^{s-1}(\nabla d)\eta)]_{C^{k+\gamma-2s}(B_{d(x_0)/2}(x_0))} \leq c(|\cdot| + |\eta(0)| + |\nabla\eta(0)|)d^{s-1}(x_0).$$

Proof. By [Abatangelo and Ros-Oton 2020, Corollary 2.3] (see also [Kukuljan 2021, Corollary 3.9] for $i > k + 1$), we deduce that

$$|D^i L(d^s \eta)| \leq c(|\cdot| + |\eta(0)|)d^{k+\gamma-s-i} \quad \text{in } \bar{\Omega} \cap B_{1/2} \quad \text{for all } i \in \mathbb{N}. \quad (2-7)$$

By [Abatangelo and Ros-Oton 2020, Theorem 2.2] and the choice of ψ in the proof of [Abatangelo and Ros-Oton 2020, Corollary 2.3], it follows that the assumption $\eta \in C^{k,\gamma}(\bar{\Omega} \cap \bar{B}_1) \cap C^\infty(\Omega \cap B_1)$ is sufficient for (2-7) to hold true. Let us now prove (i) and assume that $k = 1$ and $\gamma < s$. Then, since $D^i \eta \in C^\infty(\mathbb{R}^n)$, another application of [Abatangelo and Ros-Oton 2020, Corollary 2.3] yields

$$\|L(d^s D^i \eta)\|_{C^{1+\gamma-s}(\bar{\Omega} \cap B_{1/2})} \leq C \quad \text{for } i \in \{1, 2\}.$$

Since $\nabla(d^s \eta) = s d^{s-1}(\nabla d)\eta + d^s \nabla \eta$, a combination of the previous two estimates with $i = 1$ implies

$$|L(d^{s-1}(\nabla d)\eta)| \leq s^{-1}|\nabla L(d^s \eta)| + s^{-1}|L(d^s \nabla \eta)| \leq c(|\cdot| + |\eta(0)| + |\nabla \eta(0)|)d^{\gamma-s} \quad \text{in } \bar{\Omega} \cap B_{1/2},$$

which yields the result in (i).

To see (ii) and (iii), we observe first that by application of (2-7), we have for any $i \in \mathbb{N}$

$$|D^i L(d^s(\nabla \eta))| \leq c(|\cdot| + |\nabla \eta(0)|)d^{k+\gamma-s-i} \quad \text{in } \bar{\Omega} \cap B_{1/2}.$$

Next, by differentiation, we obtain

$$D^{i+1}(d^s \eta) = s D^i(d^{s-1}(\nabla d)\eta) + D^i(d^s \nabla \eta).$$

Thus, altogether for every $i \in \mathbb{N}$

$$\begin{aligned} |D^i(L(d^{s-1}(\nabla d)\eta))| &\leq s^{-1}|D^{i+1}L(d^s \eta)| + s^{-1}|D^i L(d^s(\nabla \eta))| \\ &\leq c(|\cdot| + |\eta(0)|)d^{k+\gamma-s-(i+1)} + c(|\cdot| + |\nabla \eta(0)|)d^{k+\gamma-s-i} \\ &\leq c(|\cdot| + |\eta(0)| + |\nabla \eta(0)|)d^{k+\gamma-s-i-1} \quad \text{in } \bar{\Omega} \cap B_{1/2}. \end{aligned} \quad (2-8)$$

To conclude the proof of (ii), let $x_0 \in \Omega \cap B_{1/2}$, and note that if $k \geq 2$ or $\gamma > s$, then (2-8) applied with $i = k - 1 + \lceil \gamma - s \rceil$ implies

$$\begin{aligned} [L(d^{s-1}(\nabla d)\eta)]_{C^{k-1-s+\gamma}(B_{d(x_0)/2}(x_0))} &\leq \sup_{x,y \in B_{d(x_0)/2}(x_0)} \frac{\|D^{k-1+\lceil \gamma-s \rceil}(L(d^{s-1}(\nabla d)\eta))\|_{L^\infty(B_{d(x_0)/2}(x_0))}}{|x-y|^{\gamma-s-\lceil \gamma-s \rceil}} \\ &\leq c(|\cdot| + |\eta(0)| + |\nabla \eta(0)|)d^{\gamma-s-\lceil \gamma-s \rceil}(x_0)d^{s-\gamma+\lceil \gamma-s \rceil}(x_0) \\ &\leq c(|\cdot| + |\eta(0)| + |\nabla \eta(0)|), \end{aligned}$$

where we used that $\gamma < 1$. From here, a covering argument (see [Fernández-Real and Ros-Oton 2024a, Lemma A.1.4]) yields the desired regularity estimate in $\bar{\Omega} \cap B_{1/2}$.

To prove (iii), note that if $k + \gamma > 2s$, then (2-8) applied with $i = k + \lceil \gamma - s \rceil$ implies

$$\begin{aligned} [L(d^{s-1}(\nabla d)\eta)]_{C^{k+\gamma-2s}(B_{d(x_0)/2}(x_0))} &\leq \sup_{x,y \in B_{d(x_0)/2}(x_0)} \frac{\|D^{k+\lceil \gamma-s \rceil}(L(d^{s-1}(\nabla d)\eta))\|_{L^\infty(B_{d(x_0)/2}(x_0))}}{|x-y|^{\gamma-2s-\lceil \gamma-s \rceil}} \\ &\leq c(|\cdot| + |\eta(0)| + |\nabla\eta(0)|)d^{\gamma-s-1-\lceil \gamma-s \rceil}(x_0)d^{2s-\gamma+\lceil \gamma-s \rceil}(x_0) \\ &\leq c(|\cdot| + |\eta(0)| + |\nabla\eta(0)|)d^{s-1}(x_0), \end{aligned}$$

where we used that $\gamma < 1$. This implies (iii), and we conclude the proof. \square

As a corollary, we obtain the following result:

Corollary 2.5. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $k \in \mathbb{N}$, and $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{k+1,\gamma}$ for some $\gamma \in (0, 1)$ with $\gamma \neq s$, and $0 \in \partial\Omega$. Assume that $K \in C^{2k+2\gamma+3}(\mathbb{S}^{n-1})$. Let*

$$\eta \in C^{k,\gamma}(\overline{\Omega \cap B_1}) \cap C^\infty(\Omega \cap B_1).$$

Then, there exists $c > 0$, depending only on $n, s, \lambda, \Lambda, \Omega, \gamma, k$, such that the following holds true:

(i) *If $k = 1$ and $\gamma < s$, then*

$$|L(d^{s-1}\eta)| \leq c\|\eta\|_{C^1(\overline{\Omega \cap B_{1/2}})}d^{\gamma-s} \quad \text{in } \overline{\Omega \cap B_{1/2}}.$$

(ii) *If $k \geq 2$ or $\gamma > s$, then*

$$[L(d^{s-1}\eta)]_{C^{k-1-s+\gamma}(\overline{\Omega \cap B_{1/2}})} \leq c(\|\eta\|_{C^1(\overline{\Omega \cap B_{1/2}})} + \|\eta\|_{C^{k-1+s+\gamma}(\Omega \cap B_1)}).$$

(iii) *If $k + \gamma > 2s$, then we have for any $x_0 \in \Omega \cap B_{1/2}$*

$$[L(d^{s-1}\eta)]_{C^{k+\gamma-2s}(B_{d(x_0)/2}(x_0))} \leq c(\|\eta\|_{C^1(\overline{\Omega \cap B_{1/2}})} + \|\eta\|_{C^{k+\gamma}(\Omega \cap B_1)})d^{s-1}(x_0).$$

Proof. There exist $N \in \mathbb{N}$ and $\delta > 0$, $v_i \in \mathbb{S}^{n-1}$, $x_i \in \partial\Omega \cap B_1$, depending only on Ω , such that $\partial_{v_i}d \geq \frac{1}{2}$ in $\overline{\Omega \cap B_\delta(x_i)}$, for $i \in \{1, \dots, N\}$, and such that

$$\{x \in \overline{\Omega \cap B_{1/2}} : d(x) \leq \delta/2\} \subset \bigcup_{i=1}^N B_\delta(x_i).$$

Then, by application of Lemma 2.4(i) to $\eta := (\partial_{v_i}d)^{-1}\eta \in C^{k,\gamma}(\overline{\Omega \cap B_\delta(x_i)}) \cap C^\infty(\Omega \cap B_\delta(x_i))$, we deduce that for any $i \in \{1, \dots, N\}$

$$|L(d^{s-1}\eta)| \leq c(|\cdot| + |\eta(x_i)| + |\nabla\eta(x_i)|)d^{\gamma-s} \quad \text{in } \overline{\Omega \cap B_{\delta/2}(x_i)}.$$

Thus, we have proved (i) in $\overline{\Omega \cap B_{1/2} \cap \{d(x) \leq \delta/2\}}$. The result in $\overline{\Omega \cap B_{1/2} \cap \{d(x) > \delta/2\}}$ is immediate from the regularity of K (see [Fernández-Real and Ros-Oton 2024a, Lemma 2.2.6]).

The proofs of (ii) and (iii) follow from Lemma 2.4 in an analogous way. \square

2.3. Barriers with boundary blow-up. Let us construct barrier functions that are suitable for establishing maximum principles for solutions that blow up at the boundary. We establish a subsolution and a supersolution in the following two lemmas.

Lemma 2.6. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{1,\gamma}$ for some $\gamma > 0$. Then, for any $l \in \mathbb{R}$, $\varepsilon \in (0, \min\{s, 1-s\})$, and $M > 0$ there exists $\phi_l \in C^\infty(\Omega)$ such that*

$$\begin{cases} L\phi_l \leq -d^{\varepsilon-s-1} - M & \text{in } \Omega, \\ \phi_l = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \phi_l/d^{s-1} = l & \text{on } \partial\Omega. \end{cases}$$

Moreover, if $l \geq 0$, then there exists $\delta \in (0, 1)$, depending only on $n, s, \lambda, \Lambda, \text{diam}(\Omega), \varepsilon, M$, such that $\phi_l \geq 0$ in $\Omega \cap \{d \leq \delta\}$. And if $l < 0$, then for M large enough, depending on $n, s, \lambda, \Lambda, \text{diam}(\Omega), \varepsilon$, it holds that $\phi_l \leq 0$ in Ω .

Proof. Let $\varepsilon \in (0, s)$ and $N > 1$ to be chosen small and large, respectively, later. We set

$$\phi_l(x) := ld^{s-1}(x) - d^{s-1+\varepsilon}(x) - N\mathbb{1}_\Omega(x).$$

Then, by Lemma 2.3,

$$L\phi_l \leq c_1ld^{\delta\gamma-s-1} - c_2d^{\varepsilon-s-1} + c_3 - NL\mathbb{1}_\Omega.$$

Since $L\mathbb{1}_\Omega \geq 0$, by taking any $\delta \in (0, s)$ and then $\varepsilon < \delta\gamma$, we see that there exists $\eta > 0$, depending on $s, l, \varepsilon, M, \delta, \gamma$, such that

$$L\phi_l \leq -d^{\varepsilon-s-1} - M \quad \text{in } \Omega \cap \{d < \eta\}.$$

Moreover, there exists $c_4 > 0$, depending on $\text{diam}(\Omega)$, such that $L\mathbb{1}_\Omega \geq c_4$ in $\Omega \cap \{d \geq \eta\}$. Thus, choosing $N = Mc_4^{-1}$, we deduce that

$$L\phi_l \leq -d^{\varepsilon-s-1} - M \quad \text{in } \Omega,$$

as desired. The remaining properties of ϕ_l follow immediately from its construction. \square

Lemma 2.7. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{1,\gamma}$ for some $\gamma > 0$. Then, there is $c_1 > 0$, depending only on n, s, λ, Λ , such that for any $l \in \mathbb{R}$, $\varepsilon \in (0, \min\{s\gamma, 1-s\})$, and $M > 0$ there exists $\psi_l \in C^\infty(\Omega)$ such that*

$$\begin{cases} L\psi_l \geq c_1d^{\varepsilon-s-1} + M & \text{in } \Omega, \\ \psi_l = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \psi_l/d^{s-1} = l & \text{on } \partial\Omega. \end{cases}$$

Moreover, for any $M > 0$, if $l > 0$ is large enough, depending only on $n, s, \lambda, \Lambda, \text{diam}(\Omega), \varepsilon$, it holds that $\psi_l \geq 0$ in Ω .

Moreover, for any $\varepsilon \in (0, s)$, there is $\tilde{\psi} \in C^s(\bar{\Omega})$ such that for some $c_2 > 0$

$$\begin{cases} L\tilde{\psi} \geq d^{\varepsilon-s} & \text{in } \Omega, \\ \tilde{\psi} = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \tilde{\psi}/d^{s-1} = 0 & \text{on } \partial\Omega, \\ \partial_\nu(\tilde{\psi}/d^{s-1}) \leq c_2 & \text{on } \partial\Omega. \end{cases}$$

Proof. Let $l \in \mathbb{R}$ and $\varepsilon \in (0, \min\{s\gamma, 1 - s\})$. The proof is similar to the one of Lemma 2.6. We set

$$\psi_l(x) = ld^{s-1}(x) + d^{s-1+\varepsilon}(x) + C_2 \mathbb{1}_{\overline{\Omega}}(x)$$

and since $\varepsilon < s\gamma$, we can choose $\delta \in (0, s)$ such that $\varepsilon < \delta\gamma$ and take $C_2 > 0$ to be chosen later. By Lemma 2.3,

$$L\psi_l \geq -c_1 ld^{\delta\gamma-s-1} + c_2 d^{\varepsilon-s-1} - c_3 + c_4 C_2 \quad \text{in } \Omega,$$

for some constants $c_1, c_2, c_3, c_4 > 0$, depending only on $n, s, \lambda, \Lambda, \delta$, and the $C^{1,\gamma}$ radius of Ω . Thus, if we choose $C_2 > 0$ large enough, depending on $M, l, c_1, c_2, c_3, c_4, \varepsilon, \text{diam}(\Omega)$, then we deduce

$$L\psi_l \geq cd^{\varepsilon-s-1} + M \quad \text{in } \Omega.$$

Finally, we observe that upon choosing $l > 0$ large enough, depending only on $\varepsilon, \text{diam}(\Omega)$, we have

$$\psi_l \geq ld^{s-1} + d^{s-1+\varepsilon} \geq 0 \quad \text{in } \Omega.$$

For the second claim, we recall from [Fernández-Real and Ros-Oton 2024a, Lemma B.2.6] that for any $\varepsilon \in (0, s)$, there exist $N > 0$ and $c_1 > 0$ such that

$$L(-Nd^{s+\varepsilon}) \geq d^{\varepsilon-s} - c_1 \quad \text{in } \Omega.$$

Let $\tilde{\psi}_2 \in L^\infty(\Omega)$ be the solution to the Dirichlet problem

$$\begin{cases} L\tilde{\psi}_2 = c_1 & \text{in } \Omega, \\ \tilde{\psi}_2 = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \end{cases}$$

and observe that by the boundary regularity theory from [Fernández-Real and Ros-Oton 2024a], it holds that $\tilde{\psi}_2 \in C^s(\overline{\Omega})$, and hence for $\tilde{\psi} := -Nd^{s+\varepsilon} - \tilde{\psi}_2$, we obtain

$$\frac{\tilde{\psi}}{d^{s-1}} = 0 \quad \text{on } \partial\Omega.$$

Therefore, for some $c_2 > 0$,

$$|\partial_\nu(\tilde{\psi}/d^{s-1})| = (1-s)|\tilde{\psi}/d^s| \leq c_2 \quad \text{on } \partial\Omega,$$

as desired. □

2.4. Nonlocal equations up to a polynomial. We will need the following definition of nonlocal equations that hold true up to a polynomial. It was introduced in [Dipierro et al. 2019] for the fractional Laplacian and the theory was extended in [Dipierro et al. 2022] to general nonlocal operators (see also [Abatangelo and Ros-Oton 2020]).

Definition 2.8. For $k \in \mathbb{N}$, a bounded domain $\Omega \subset \mathbb{R}^n$, $f \in C(\Omega)$, and $K \in C^{k-1+\delta}(\mathbb{S}^{n-1})$ for some $\delta > 0$, we say that a function $u \in C(\Omega) \cap L^1_{2s+k}(\mathbb{R}^n)$ solves in the viscosity sense

$$Lu \stackrel{k}{=} f \quad \text{in } \Omega,$$

if there exist polynomials $(p_R)_{R>1} \in \mathcal{P}_{k-1}$ of degree $k-1$, and functions $(f_R)_{R>1}$ such that

$$\begin{aligned} L(u\mathbb{1}_{B_R}) &= f_R + p_R \quad \text{in } \Omega \quad \text{for all } R > \text{diam}(\Omega), \\ \|f_R - f\|_{L^\infty(\Omega)} &\rightarrow 0 \quad \text{as } R \rightarrow \infty. \end{aligned}$$

Remark 2.9. • In case $k=0$, we set $\mathcal{P}_{0-1} = \mathcal{P}_{-1} = \{0\}$. Then, $Lu \stackrel{0}{=} f$ is equivalent to $Lu = f$ (see [Dipierro et al. 2022, Corollary 2.13]).

- Instead of $K \in C^k(\mathbb{S}^{n-1})$, here we only assume $K \in C^{k-1+\delta}(\mathbb{S}^{n-1})$ for some $\delta > 0$. It is easy to see that all the arguments in [Dipierro et al. 2022] remain valid under this weaker assumption. We decided to make this change in order to have optimal assumptions on K in Theorem 1.7.
- As in [Abatangelo and Ros-Oton 2020], we assume uniform convergence $f_R \rightarrow f$. This is slightly different from [Dipierro et al. 2019], where pointwise convergence was assumed.

The following lemma is a slight improvement of [Abatangelo and Ros-Oton 2020, Lemma 3.6] (see also [Ros-Oton et al. 2025, Lemma 8.1]) in the sense that the estimate involves a weighted L^1 norm instead of a weighted L^∞ norm.

Lemma 2.10. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $u \in C(B_1)$ be a viscosity solution to*

$$Lu = f \quad \text{in } B_1.$$

Then, the following hold true:

- (i) *Let $\beta \in (0, 2s]$ if $s \neq \frac{1}{2}$ and $\beta \in (0, 1)$ if $s = \frac{1}{2}$. If $f \in C(B_1)$ and $u \in L^1_{2s}(\mathbb{R}^n)$, then it holds that $u \in C^\beta_{\text{loc}}(B_1)$ and*

$$\|u\|_{C^\beta(B_{1/2})} \leq c \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s}} dy + \|f\|_{L^\infty(B_1)} \right)$$

for some $c > 0$, depending only on $n, s, \lambda, \Lambda, \beta$.

- (ii) *If $f \in C^\alpha(B_1)$ for some $\alpha > 0$ such that $2s + \alpha \notin \mathbb{N}$, $K \in C^\alpha(\mathbb{S}^{n-1})$, and $u \in L^1_{2s+\alpha}(\mathbb{R}^n)$, then $u \in C^{2s+\alpha}_{\text{loc}}(B_1)$ and*

$$\|u\|_{C^{2s+\alpha}(B_{1/2})} \leq c \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + [f]_{C^\alpha(B_1)} \right)$$

for some $c > 0$, depending only on $n, s, \lambda, \Lambda, \alpha$.

Remark 2.11. From the proof it is apparent, that Lemma 2.10(ii) remains true if $Lu \stackrel{k}{=} f$ for $k < \alpha$.

Proof. Let us first show (ii) in case $\alpha < 1$. Let us define $v = u\mathbb{1}_{B_1}$. We claim that v solves $Lv = \tilde{f}$ in $B_{3/4}$ for some $\tilde{f} \in C^{2s+\alpha}(B_{3/4})$ with

$$\|\tilde{f}\|_{C^\alpha(B_{3/4})} \leq C \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + [f]_{C^\alpha(B_{3/4})} \right). \quad (2-9)$$

To prove it, first, we observe that for any $h \in B_{3/4}$ and $x \in B_{3/4-|h|}$, using that $K \in C^\alpha(\mathbb{S}^{n-1})$,

$$\begin{aligned} |L(u\mathbb{1}_{\mathbb{R}^n \setminus B_1})(x) - L(u\mathbb{1}_{\mathbb{R}^n \setminus B_1})(x+h)| &\leq |h|^\alpha \int_{\mathbb{R}^n \setminus B_1} |u(y)| \frac{|K(x-y) - K(x+h-y)|}{|h|^\alpha} dy \\ &\leq c|h|^\alpha \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy. \end{aligned}$$

Therefore, since we can add and subtract constants to \tilde{f} without affecting the left-hand side of the next estimate, for any $h \in B_{3/4}$ it holds that

$$\|\tilde{f} - \tilde{f}(\cdot + h)\|_{L^\infty(B_{3/4-|h|})} \leq C \left(\operatorname{osc}_{B_{3/4}} f + |h|^\alpha \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy \right). \quad (2-10)$$

From here, we deduce that there exist $g \in L^\infty(B_{3/4})$ and a constant p such that $\tilde{f} = g + p$. By construction we have

$$\|\tilde{g}\|_{L^\infty(B_{3/4})} \leq C \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + \operatorname{osc}_{B_{3/4}} f \right).$$

We split $v = v_1 + v_2$, where v_1 and v_2 are solutions to

$$\begin{cases} Lv_1 = \tilde{g} & \text{in } B_{3/4}, \\ v_1 = v & \text{in } \mathbb{R}^n \setminus B_{3/4}, \end{cases} \quad \begin{cases} Lv_2 = p & \text{in } B_{3/4}, \\ v_2 = 0 & \text{in } \mathbb{R}^n \setminus B_{3/4}, \end{cases}$$

and note that the existence of v_1, v_2 follows from [Fernández-Real and Ros-Oton 2024a, Theorem 3.2.27]). Then, by the maximum principle (see [Fernández-Real and Ros-Oton 2024a, Corollary 3.2.22]) we deduce that

$$\|v_1\|_{L^\infty(B_{3/4})} \leq C \left(\|v\|_{L^\infty(\mathbb{R}^n \setminus B_{3/4})} + \|\tilde{g}\|_{L^\infty(B_{3/4})} \right) \leq C \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + \operatorname{osc}_{B_{3/4}} f \right).$$

Hence,

$$\|v_2\|_{L^\infty(B_{3/4})} \leq \|u\|_{L^\infty(B_{3/4})} + \|v_1\|_{L^\infty(B_{3/4})} \leq C \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + \operatorname{osc}_{B_{3/4}} f \right).$$

Then, by [Abatangelo and Ros-Oton 2020, Lemma 3.7], we deduce

$$\|p\|_{L^\infty(B_{3/4})} \leq C \|v_2\|_{L^\infty(B_{3/4})} \leq C \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + \operatorname{osc}_{B_{3/4}} f \right).$$

Altogether, we have shown

$$\|\tilde{f}\|_{L^\infty(B_{3/4})} \leq \|\tilde{g}\|_{L^\infty(B_{3/4})} + \|p\|_{L^\infty(B_{3/4})} \leq C \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + \operatorname{osc}_{B_{3/4}} f \right).$$

Finally, as a direct consequence of (2-10), we deduce

$$[\tilde{f}]_{C^\alpha(B_{3/4})} \leq C \left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + [f]_{C^\alpha(B_{3/4})} \right),$$

which yields the claim (2-9).

Thus, by application of the interior regularity estimate [Fernández-Real and Ros-Oton 2024a, Theorem 2.4.1] to v , we obtain

$$\begin{aligned} \|u\|_{C^{2s+\alpha}(B_{1/2})} &= \|v\|_{C^{2s+\alpha}(B_{1/2})} \leq c(\|v\|_{L^\infty(\mathbb{R}^n)} + \|\tilde{f}\|_{C^\alpha(B_{3/4})}) \\ &\leq c\left(\|u\|_{L^\infty(B_1)} + \int_{\mathbb{R}^n \setminus B_1} \frac{|u(y)|}{|y|^{n+2s+\alpha}} dy + [f]_{C^\alpha(B_{3/4})}\right), \end{aligned}$$

as desired. This proves (ii) in case $\alpha < 1$. The case $\alpha \geq 1$ goes in the same way by considering higher order incremental quotients in the arguments above. Statement (i) was proved in [Fernández-Real and Ros-Oton 2024a, Theorem 2.4.3]. The L^∞ norm can be replaced by the $L^1_{2s}(\mathbb{R}^n)$ norm by the same truncation argument we employed above. \square

Next, we provide a lemma stating that equations up to a polynomial can be differentiated in the same way as classical nonlocal equations. This lemma will be used in the proof of Lemma 5.3.

Lemma 2.12. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $k \in \mathbb{N}$, $f \in C^1(B_1)$, and $K \in C^{k+\delta}(\mathbb{S}^{n-1})$ for some $\delta > 0$. Let $u \in C(B_1) \cap L^1_{2s+k+\delta}(\mathbb{R}^n)$ with $\partial_i u \in L^1_{2s+k-1+\delta}(\mathbb{R}^n)$. Then, it holds that*

$$Lu \stackrel{k+1}{=} f \quad \text{in } B_1,$$

and $\partial_i f_R \rightarrow \partial_i f$, if and only if

$$L(\partial_i u) \stackrel{k}{=} \partial_i f \quad \text{in } B_1.$$

Proof. Let us assume first that $\partial_i u \in L^1_{2s+k-1+\delta}(\mathbb{R}^n)$, and assume that $Lu \stackrel{k+1}{=} f$ in B_1 . Then, there exist polynomials $p_R \in \mathcal{P}_k$ and functions $f_R \in L^\infty(B_1)$ with $f_R \rightarrow f$ such that

$$L(u\mathbb{1}_{B_R}) = f_R + p_R \quad \text{in } B_1.$$

Let us now consider difference quotients $D_i^h u(x) = \frac{u(x+e_i h) - u(x)}{|h|}$ and compute

$$L(D_i^h u\mathbb{1}_{B_R}) = L(D_i^h(u\mathbb{1}_{B_R})) - L(uD_i^h \mathbb{1}_{B_R}) = D_i^h f_R + D_i^h p_R - L(uD_i^h \mathbb{1}_{B_R}),$$

where, by following the proof of [Dipierro et al. 2022, Theorem 2.1], we can decompose

$$-L(uD_i^h \mathbb{1}_{B_R})(x) = \int_{\mathbb{R}^n \setminus B_3} u D_i^h \mathbb{1}_{B_R}(y) K(x-y) dy = d_{R,h}(x) + g_{R,h}(x)$$

for functions $g_{R,h}$ such that

$$g_{R,h}(x) = \int_{\mathbb{R}^n} D_i^h(\mathbb{1}_{B_R})(y) u(y) \psi(x, y) dy = - \int_{B_R} (D_{-h}^i u(y) \psi(x, y) + u(y) D_{-h}^i \psi(x, y)) dy$$

for some function $\psi : B_1 \times (\mathbb{R}^n \setminus B_3) \rightarrow \mathbb{R}$ such that

$$\sup_{x \in B_1} \psi(x, y) \leq C \sup_{x \in B_1} (1 + |x - y|)^{-(n+2s+k-1+\delta)}, \quad \sup_{x \in B_1} |\nabla_y \psi(x, y)| \leq C \sup_{x \in B_1} (1 + |x - y|)^{-(n+2s+k+\delta)},$$

and polynomials $d_{R,h} \in \mathcal{P}_{k-1}$ with

$$d_{R,h}(x) = \sum_{|\alpha| \leq k-1} \kappa_{\alpha,h} x^\alpha, \quad \kappa_{\alpha,h} = c_\alpha \int_{B_R} D_{-h}^i [u(y) \partial_x^\alpha K(x-y)] dy, \quad c_\alpha \in \mathbb{R}.$$

Clearly, it holds that $g_{R,h} \rightarrow g_R$, and $d_{R,h} \rightarrow d_R$, as $R \rightarrow \infty$, where

$$g_R(x) = \int_{B_R} \partial_i [u(y)\psi(x, y)] dy, \quad d_R(x) = \sum_{|\alpha| \leq k-1} \kappa_\alpha x^\alpha, \quad \kappa_\alpha = c_\alpha \int_{B_R} \partial_i [u(y)\partial_x^\alpha K(x-y)] dy.$$

For the convergence $g_{R,h} \rightarrow g_R$ we are using that for any $x \in B_1$,

$$\begin{aligned} & |\mathbb{1}_{B_R}(y)(D_{-h}^i u(y)\psi(x, y) + \mathbb{1}_{B_R}(y)u(y)D_{-h}^i \psi(x, y))| \\ & \leq C|\partial_i u(y)| \sup_{x \in B_1} (1 + |x-y|)^{-(n+2s+k-1+\delta)} + C|u(y)| \sup_{x \in B_1} (1 + |x-y|)^{-(n+2s+k+\delta)} \in L^1(\mathbb{R}^n) \end{aligned}$$

and dominated convergence. Moreover, from integrating by parts, we see that it holds for any $x \in B_1$

$$\begin{aligned} \int_2^\infty R^{-1} |g_R^{(2)}(x)| dR & \leq \int_2^\infty R^{-1} \int_{\partial B_R} |u(y)||x-y|^{-(n+2s+k-1+\delta)} dy dR \\ & \leq c \int_{B_2^c} |u(y)||y|^{-(n+2s+k+\delta)} dy < \infty, \end{aligned}$$

which implies that $g_R(x) \rightarrow 0$, as $R \rightarrow \infty$, uniformly in x .

Altogether, we have shown

$$L(\partial_i u \mathbb{1}_{B_R}) = \lim_{h \rightarrow 0} D_i^h f_R + \lim_{h \rightarrow 0} D_i^h p_R + \lim_{h \rightarrow 0} d_{R,h} + \lim_{h \rightarrow 0} g_{R,h} = \partial_i f + \partial_i p_R + d_R + g_R,$$

which implies that

$$L(\partial_i u \mathbb{1}_{B_R}) \stackrel{k}{=} f \quad \text{in } B_1,$$

as desired.

Let us now show the other implication, i.e., assume that $L(\partial_i u) \stackrel{k}{=} \partial_i f$ in B_1 . Then, by [Dipierro et al. 2022] we observe that there are $F_R : \mathbb{R}^n \rightarrow \mathbb{R}$, and $P_R \in \mathcal{P}_k$ such that

$$L(u \mathbb{1}_{B_R}) = F_R + P_R \quad \text{in } B_1.$$

Clearly, by the same arguments as above, we have

$$L((D_i^h u) \mathbb{1}_{B_R}) = D_i^h L(u \mathbb{1}_{B_R}) - L(u D_i^h \mathbb{1}_{B_R}) = D_i^h F_R + D_i^h P_R + d_{R,h} + g_{R,h}$$

with $D_i^h P_R + d_{R,h} \in \mathcal{P}_{k-1}$ and $g_{R,h} \rightarrow g_R$, as $h \rightarrow 0$ with $g_R \rightarrow 0$, as $R \rightarrow \infty$. Thus, by the stability for viscosity solutions up to a polynomial (see [Abatangelo and Ros-Oton 2020, Lemma 3.5]), we have that

$$f_R + p_R = L(\partial_i u \mathbb{1}_{B_R}) = \partial_i F_R + \partial_i P_R + d_R + g_R,$$

where $d_R, p_R, \partial_i P_R \in \mathcal{P}_{k-1}$. Hence, after integrating the previous identity in x_i and denoting $\tilde{F}_R(x) = \int_{-\infty}^{x_i} (f_R - g_R)(x', y_i) dy$, we can deduce

$$F_R = \tilde{F}_R + \tilde{P}_R \quad \text{in } B_1,$$

where $\tilde{P}_R \in \mathcal{P}_k$ is such that $\partial_i \tilde{P}_R = p_R - d_R - \partial_i P_R$. Then, since $f_R \rightarrow f$ and $g_R \rightarrow 0$, as $R \rightarrow \infty$, we deduce that $\tilde{F}_R \rightarrow F$, where $\partial_i F = f$, and the proof is complete. \square

2.5. Two lemmas on viscosity solutions. In this section, we prove two auxiliary lemmas for viscosity solutions to nonlocal equations with local Neumann boundary data, namely a stability result, and that sums of viscosity subsolutions are again viscosity subsolutions. Both results are standard for nonlocal equations in the interior of the solution domain (see [Fernández-Real and Ros-Oton 2024a]). However, since we consider equations at the boundary, where solutions satisfy a Neumann condition in the viscosity sense, both results require a proof. Both proofs heavily rely on the interaction of nonlocal operators with the distance function and the results in Section 2.2.

First, we prove a stability result, which will be crucial in the blow-up argument of our proof of the higher boundary regularity.

Lemma 2.13. *Let $k \in \mathbb{N} \cup \{0\}$, $\gamma \in (0, 1)$ with $\gamma \neq s$, and $\Omega_j \subset \mathbb{R}^n$ be open, bounded domains with $\partial\Omega_j \in C^{2,\gamma}$ such that $0 \in \partial\Omega_j$, $v_0 = e_n$ for any $j \in \mathbb{N}$, and such that the $C^{2,\gamma}$ radii of Ω_j and $\text{diam}(\Omega_j)$ are uniformly bounded. Given a sequence $r_j \searrow 0$, we set $\tilde{\Omega}_j = r_j^{-1}\Omega_j$ and $\tilde{d}_j := d_{\tilde{\Omega}_j}$. Let $v_j \in L^1_{2s+k}(\mathbb{R}^n)$ with $v_j/\tilde{d}_j^{s-1} \in C(\tilde{\Omega}_j)$ be viscosity solutions to*

$$\begin{cases} L_j v_j \stackrel{k}{=} f_j & \text{in } \tilde{\Omega}_j \cap B_1, \\ v_j = 0 & \text{in } \mathbb{R}^n \setminus \tilde{\Omega}_j, \\ \partial_\nu(v_j/\tilde{d}_j^{s-1}) = g_j & \text{on } \partial\tilde{\Omega}_j \cap B_1, \end{cases}$$

where $f_j \in C(\tilde{\Omega}_j \cap B_1)$, $g_j \in C(\partial\tilde{\Omega}_j \cap B_1)$, and $(L_j)_j \subset \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda, k-1+\alpha)$ for some $\alpha > 0$. Moreover, assume that there are $v \in L^1_{2s+k}(\mathbb{R}^n)$ with $v/(x_n)_+^{s-1} \in C(\{x_n \geq 0\})$, $f \in C(\{x_n > 0\} \cap B_1)$, $g \in C(\{x_n = 0\} \cap B_1)$, $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda, k-1+\alpha)$, and $\varepsilon_j \searrow 0$, $q_j \in \mathcal{P}_k$ such that

$$\begin{aligned} v_j/\tilde{d}_j^{s-1} &\rightarrow v/(x_n)_+^{s-1} && \text{in } L^\infty_{\text{loc}}(B_1), \\ v_j &\rightarrow v && \text{in } L^1_{2s+k}(\mathbb{R}^n), \\ |f_j - p_j - f| &\rightarrow 0 && \text{in } L^\infty_{\text{loc}}(B_1 \cap \{x_n > 0\}), \\ |g_j - q_j - g|(x) &\leq c\varepsilon_j \rightarrow 0 && \text{for all } x \in \partial\tilde{\Omega}_j \cap B_1, \\ K_j &\rightarrow K && \text{in } C^{k-1+\alpha}(\mathbb{S}^{n-1}). \end{aligned}$$

Then, there exists $q \in \mathcal{P}_k$ such that v is a viscosity solution to

$$\begin{cases} Lv \stackrel{k}{=} f & \text{in } B_1 \cap \{x_n > 0\}, \\ v = 0 & \text{in } \mathbb{R}^n \setminus \{x_n > 0\}, \\ \partial_n(v/(x_n)_+^{s-1}) = g + q & \text{on } B_1 \cap \{x_n = 0\}. \end{cases}$$

If $k = 0$, the same result holds with $\tilde{d}_j^{s-\gamma} f_j \in L^\infty(\tilde{\Omega}_j \cap B_1)$ and $(x_n)_+^{s-\gamma} f \in L^\infty(\{x_n > 0\} \cap B_1)$.

Proof. Let us define $u_j := v_j/\tilde{d}_j^{s-1}$ and $u := v/(x_n)_+^{s-1}$. Since $u_j \rightarrow u$ in $L^\infty_{\text{loc}}(B_1)$ it follows that $v_j = \tilde{d}_j^{s-1}u_j \rightarrow (x_n)_+^{s-1}u = v$ in $L^\infty_{\text{loc}}(B_1 \cap \{x_n > 0\})$. This property is enough to use the stability of viscosity solutions from [Fernández-Real and Ros-Oton 2024a, Proposition 3.2.12] to v_j and v . The higher order version which we require here follows from [Abatangelo and Ros-Oton 2020, Lemma 3.5]. Since $v_j \rightarrow v$ in $L^1_{2s}(\mathbb{R}^n)$, we also have that $v = 0$ in $B_1 \setminus \{x_n > 0\}$. Consequently, it only remains to prove the convergence of the Neumann boundary condition.

To do so, let $x_0 \in B_1 \cap \{x_n = 0\}$. In case $k \geq 1$, we first truncate v and v_j in $B_2(x_0)$ and apply [Abatangelo and Ros-Oton 2020, Lemma 3.6] to obtain the equations satisfied by $v \mathbb{1}_{B_2}(x_0)$ and $v_j \mathbb{1}_{B_2}(x_0)$. In order not to over-complicate the notation, let us denote the truncations still by v and v_j and the corresponding source terms by f and f_j . Then, let $\phi \in C^2(B_r(x_0))$ for some $r \in (0, 1)$ with $\phi \leq u$ in $B_r(x_0)$, $\phi(x_0) = u(x_0)$, and $\phi \equiv u$ in $\mathbb{R}^n \setminus \overline{B_r(x_0)}$ be a test function. Given $\delta \in (0, 1)$, $\eta \in (0, \gamma)$, we define now

$$\psi^{(\delta)}(x) = -\delta \mathbb{1}_{B_r(x_0)}(x)[(x_n)_+ - (x_n)_+^{1+\eta}], \quad \psi_j^{(\delta)}(x) = -\delta \mathbb{1}_{B_r(x_0)}(x)[\tilde{d}_j(x) - \tilde{d}_j^{1+\eta}(x)].$$

There exist $C > 0$ and $\varepsilon \in (0, r/2)$, independent of δ, j , such that

$$L_j(\tilde{d}_j^{s-1} \psi_j^{(\delta)}) \leq -C \delta \tilde{d}_j^{\eta-s} \quad \text{in } \tilde{\Omega}_j \cap B_\varepsilon(x_0). \tag{2-11}$$

This is due to [Fernández-Real and Ros-Oton 2024a, Proposition B.2.1, Lemma B.2.6, Corollary B.2.8], and since $\tilde{\Omega}_j \cap B_R(x_0)$ and the respective $C^{2,\gamma}$ -radii of $\tilde{\Omega}_j$ are uniformly bounded. Indeed, the aforementioned results yield the existence of $\varepsilon_0 > 0$ such that

$$-\delta L_j(\tilde{d}_j^{s-1}[\tilde{d}_j - \tilde{d}_j^{1+\eta}]) \leq -c_1 \delta \tilde{d}_j^{\eta-s} \quad \text{in } \tilde{\Omega}_j \cap B_{\varepsilon_0}(x_0).$$

Moreover, one computes by scaling from $\tilde{\Omega}_j$ to Ω_j , denoting $d_j = d_{\Omega_j}$, and applying Lemma 2.2,

$$\begin{aligned} -\delta L_j(\tilde{d}_j^{s-1}[\tilde{d}_j - \tilde{d}_j^{1+\eta}]) \mathbb{1}_{\mathbb{R}^n \setminus B_r(x_0)} &\leq c \delta \int_{\tilde{\Omega}_j \setminus B_r(x_0)} (\tilde{d}_j^s(y) + \tilde{d}_j^{s+\eta}(y)) |y|^{-n-2s} \, dy \\ &\leq c \delta \int_{\Omega_j \setminus B_{r_j}(x_0)} (r_j^s d_j^s(y) + r_j^{s-\eta} d_j^{s+\eta}(y)) |y|^{-n-2s} \, dy \\ &\leq c_2 \delta (1 + r^{-s} + r^{\eta-s}) \quad \text{in } \tilde{\Omega}_j \cap B_{r/2}(x_0), \end{aligned} \tag{2-12}$$

where $c_2 > 0$ might depend on $\text{diam}(\Omega_j)$, which we assumed to be bounded, but not on j . Thus, by combination of the previous two computations, we deduce (2-11) upon choosing $\varepsilon < \varepsilon_0$ if necessary.

Moreover, it is immediate by construction that

$$\psi^{(\delta)} \leq 0 \quad \text{in } \mathbb{R}^n. \tag{2-13}$$

Next, we set $\phi^{(\delta)} := \phi + \psi^{(\delta)}$. For any $\delta > 0$, it still holds $\phi^{(\delta)} \leq u$ by (2-13), and $\phi^{(\delta)}(x_0) = u(x_0)$, however $u - \phi^{(\delta)}$ has a strict minimum at x_0 in $\overline{B_r(x_0)}$.

It suffices to prove for any $\delta > 0$ small enough,

$$\partial_n \phi^{(\delta)}(x_0) \leq g(x_0) + q(x_0), \tag{2-14}$$

since then it follows that $\partial_n \phi(x_0) = \partial_n \phi^{(\delta)}(x_0) + \delta \leq g(x_0) + q(x_0) + \delta$, and we obtain the desired result upon taking the limit $\delta \searrow 0$.

Let us now construct test functions $\phi_j^{(\delta)}$ for any $j \in \mathbb{N}$ as

$$\phi_j^{(\delta)} = \begin{cases} u_j + \psi_j^{(\delta)} & \text{in } \mathbb{R}^n \setminus \overline{B_r(x_0)}, \\ \phi + c_j + \psi_j^{(\delta)} & \text{in } \overline{B_r(x_0)}, \end{cases}$$

where

$$c_j = \min\{c \in \mathbb{R} : \phi + c + \psi_j^{(\delta)} \leq u_j \quad \text{in } \overline{B_r(x_0)}\}.$$

Since $\psi_j^{(\delta)} \rightarrow \psi^{(\delta)}$ (lower half-relaxed limits) in $\overline{B_r(x_0)}$, we obtain that $c_j \rightarrow 0$ and there exist $x_j \in \overline{B_r(x_0)}$ with $x_j \rightarrow x_0$ such that $\phi_j^{(\delta)}(x_j) = u_j(x_j)$ and $\phi_j^{(\delta)} \leq u_j$ by [Fernández-Real and Ros-Oton 2024a, Lemma 3.2.10 and proof of Proposition 3.2.12].

Next, we argue that $x_j \in \partial\tilde{\Omega}_j \cap B_1$. Without loss of generality, we can assume that $x_j \in B_\varepsilon(x_0)$ upon taking $j \in \mathbb{N}$ large enough. In fact, if $x_j \in \tilde{\Omega}_j \cap B_\varepsilon(x_0)$, then we can compute using Corollary 2.5(i), and (2-11),

$$\begin{aligned} L_j(\tilde{d}_j^{s-1}\phi_j^{(\delta)})(x_j) &= L_j(\tilde{d}_j^{s-1}\phi\mathbb{1}_{B_r(x_0)})(x_j) + c_j L_j(\tilde{d}_j^{s-1}\mathbb{1}_{B_r(x_0)})(x_j) \\ &\quad + L_j(v_j\mathbb{1}_{\mathbb{R}^n \setminus B_r(x_0)})(x_j) + L_j(\tilde{d}_j^{s-1}\psi_j^{(\delta)})(x_j) \\ &\leq L_j(\tilde{d}_j^{s-1}\phi)(x_j) + c_j L_j(\tilde{d}_j^{s-1})(x_j) + L_j(\tilde{d}_j^{s-1}u_j\mathbb{1}_{\mathbb{R}^n \setminus B_r(x_0)})(x_j) \\ &\quad - c_j L_j(\tilde{d}_j^{s-1}\mathbb{1}_{\mathbb{R}^n \setminus B_r(x_0)})(x_j) + c\|v_j\|_{L_{2s}^1(\mathbb{R}^n)} - C\delta\tilde{d}_j^{\eta-s}(x_j) \\ &\leq C_r\tilde{d}_j^{\gamma-s}(x_j) - C\delta\tilde{d}_j^{\eta-s}(x_j) \end{aligned} \tag{2-15}$$

for some constant $C_r > 0$, depending also on $\|v_j\|_{L_{2s}^1(\mathbb{R}^n)}$ and $\|v\|_{L_{2s}^1(\mathbb{R}^n)}$. To estimate the fourth term in the last estimate, we used an argument similar to (2-12), namely

$$\begin{aligned} -L_j(\tilde{d}_j^{s-1}\mathbb{1}_{\mathbb{R}^n \setminus B_r(x_0)})(x_j) &\leq c \int_{\tilde{\Omega}_j \setminus B_r(x_0)} \tilde{d}_j^{s-1}(y)|y|^{-n-2s} \, dy \\ &\leq c \int_{\Omega_j \setminus B_{rr_j}(x_0)} r_j^{s+1} d_j^{s-1}(y)|y|^{-n-2s} \, dy \leq cr^{-s-1} =: c_r \end{aligned}$$

for some $c_r > 0$, where we applied Lemma 2.2. Let us now recall that $\eta < \gamma$.

Hence, upon making $\varepsilon > 0$ even smaller, we can have in $\tilde{\Omega}_j \cap B_\varepsilon(x_0)$,

$$C\delta\tilde{d}_j^{\eta-\gamma} > (C_r + \mathbb{1}_{\{k=0\}}\|\tilde{d}_j^{s-\gamma} f_j\|_{L^\infty(\tilde{\Omega}_j \cap B_1)} + \mathbb{1}_{\{k \geq 1\}}\|f_j\|_{L^\infty(\tilde{\Omega}_j \cap B_1)}).$$

Then it holds that

$$L_j(\tilde{d}_j^{s-1}\phi_j^{(\delta)})(x_j) < -\mathbb{1}_{\{k=0\}}\tilde{d}_j^{\gamma-s}(x_j)\|\tilde{d}_j^{s-\gamma} f_j\|_{L^\infty(\tilde{\Omega}_j \cap B_1)} - \mathbb{1}_{\{k \geq 1\}}\|f_j\|_{L^\infty(\tilde{\Omega}_j \cap B_1)} < f_j(x_j). \tag{2-16}$$

However, by construction, $\tilde{d}_j^{s-1}\phi_j^{(\delta)}$ is a valid test function for the equation that is satisfied for $\tilde{d}_j^{s-1}u_j = v_j$ at x_j . Since we assumed that $x_j \in \tilde{\Omega}_j \cap B_1$, it must hold that $L_j(\tilde{d}_j^{s-1}\phi_j^{(\delta)})(x_j) \geq f_j(x_j)$, which contradicts (2-16).

Therefore, it must be $x_j \in \partial\tilde{\Omega}_j \cap B_1$, as we claimed before. Thus, by the boundary condition

$$\partial_{\nu_{x_j}} \phi_j^{(\delta)}(x_j) \leq g_j(x_j).$$

Passing this inequality to the limit, and using the uniform convergence $|g_j - q_j - g| \rightarrow 0$, $\nu_{x_j} \rightarrow \nu_0 = e_n$, and $\tilde{\Omega}_j \rightarrow \{x_n > 0\}$, we obtain

$$\partial_n(\phi^{(\delta)}(x_0) - q(x_0)) \leq g(x_0),$$

where $q \in \mathcal{P}_k$ is the limit of the sequence of polynomials $(q_j)_j$. Thus, we have $\partial_n \phi^{(\delta)}(x_0) \leq g(x_0) + q(x_0)$, i.e., (2-14), as desired. This concludes the proof. \square

Second, we prove that the difference of two viscosity solutions is again a viscosity subsolution.

Lemma 2.14. *Let $k \in \mathbb{N} \cup \{0\}$, $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{2,\gamma}$ for some $\gamma > 0$. Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda, k - 1 + \alpha)$ for some $\alpha > 0$. Let $v, w \in L_{2s}^1(\mathbb{R}^n)$ with $v/d^{s-1}, w/d^{s-1} \in C(\bar{\Omega})$ be viscosity solutions to*

$$\begin{cases} Lv \stackrel{k}{=} f_1 & \text{in } \Omega, \\ v = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ v/d^{s-1} = g_1 & \text{on } \partial\Omega, \end{cases} \quad \begin{cases} Lw \stackrel{k}{=} f_2 & \text{in } \Omega, \\ w = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ w/d^{s-1} = g_2 & \text{on } \partial\Omega, \end{cases}$$

for some $f_1, f_2 \in C(\Omega)$ and $g_1, g_2 \in C(\partial\Omega)$. Then, $v - w$ is a viscosity solution to

$$\begin{cases} L(v - w) \stackrel{k}{=} f_1 - f_2 & \text{in } \Omega, \\ v - w = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ (v - w)/d^{s-1} = g_1 - g_2 & \text{on } \partial\Omega. \end{cases}$$

Proof. We will only demonstrate the proof in case $k = 0$. The general case follows immediately by combining the arguments with Definition 2.8. For the nonlocal equation, the result follows for instance from [Fernández-Real and Ros-Oton 2024a, Lemma 3.4.14]. For the boundary condition, one can proceed as follows. First, we define the sup- and inf-convolutions (see [Fernández-Real and Ros-Oton 2024a, Lemma 3.2.16]),

$$\begin{aligned} (v/d^{s-1})_\varepsilon(x) &:= \inf_{\bar{D}} \left(\frac{v}{d^{s-1}}(z) + \frac{|x-z|^2}{\varepsilon} \right) \quad \text{for all } x \in \bar{D}, & (v/d^{s-1})_\varepsilon(x) &= \frac{v}{d^{s-1}}(x) \quad \text{for all } x \in \mathbb{R}^n \setminus D, \\ (w/d^{s-1})^\varepsilon &= \sup_{\bar{D}} \left(\frac{w}{d^{s-1}}(z) - \frac{|x-z|^2}{\varepsilon} \right) \quad \text{for all } x \in \bar{D}, & (w/d^{s-1})^\varepsilon(x) &= \frac{w}{d^{s-1}}(x) \quad \text{for all } x \in \mathbb{R}^n \setminus D, \end{aligned}$$

with $D \subset \Omega$ open, bounded such that $\bar{D} \cap \partial\Omega \neq \emptyset$. In analogy to [Fernández-Real and Ros-Oton 2024a, Proposition 3.2.17], we claim that for any $x \in \partial\Omega \cap \bar{D}$ it holds in the viscosity sense that

$$\partial_v(v/d^{s-1})_\varepsilon(x) \leq g_1(x) + \delta_\varepsilon, \quad \partial_v(w/d^{s-1})^\varepsilon \geq g_2(x) + \delta^\varepsilon, \quad (2-17)$$

where $\delta_\varepsilon, \delta^\varepsilon \rightarrow 0$, as $\varepsilon \rightarrow 0$. Once (2-17) is proven, since $(v/d^{s-1})_\varepsilon$ and $-(w/d^{s-1})^\varepsilon$ are both semiconcave, we have that at any point $x \in \partial\Omega \cap \bar{D}$, where $(v/d^{s-1})_\varepsilon - (w/d^{s-1})^\varepsilon$ can be touched by a paraboloid from below, the functions $(v/d^{s-1})_\varepsilon$ and $-(w/d^{s-1})^\varepsilon$ must be in $C^{1,1}$. Hence, by the linearity of ∂_v , and due to (2-17) it must hold that

$$\partial_v((v/d^{s-1})_\varepsilon - (w/d^{s-1})^\varepsilon)(x) \leq g_1(x) - g_2(x) + \delta_\varepsilon - \delta^\varepsilon \rightarrow g_1(x) - g_2(x) \quad \text{as } \varepsilon \rightarrow 0.$$

Thus, by the stability for viscosity solutions (which was provided in a significantly more general framework in Lemma 2.13), we deduce that $\partial_v((v - w)/d^{s-1}) \leq (g_1 - g_2)$ in the viscosity sense. In a similar way, one can prove $\partial_v((v - w)/d^{s-1}) \geq (g_1 - g_2)$, and thus, we obtain the desired result.

Thus, it remains to give a proof of (2-17). To see it, for any test function $\phi \in C^2(B_r(x_0))$ touching $(v/d^{s-1})_\varepsilon$ from below at $x_0 \in \partial\Omega \cap \bar{D}$, we define

$$\psi^{(\delta)}(x) = -\delta \mathbb{1}_{B_1(x_0)}(x)[d(x) - d^{1+\eta}(x)]$$

for some $\eta \in (0, \gamma)$ and observe that $\phi^{(\delta)} = \phi + \psi^{(\delta)}$ is still a valid test function, touching $(v/d^{s-1})_\varepsilon$ (strictly) from below, at x_0 . Then, there exists $x_\varepsilon \in \bar{D}$ with $x_\varepsilon \in B_{c\varepsilon}(x_0)$ for some $c > 0$, depending only on the oscillation of v/d^{s-1} , such that $\phi^{(\delta)}(\cdot + x_0 - x_\varepsilon) - \varepsilon^{-1}|x_0 - x_\varepsilon|^2$ touches v/d^{s-1} from below at x_ε . Indeed, from the definition of $(v/d^{s-1})_\varepsilon$ we deduce that there exist $x_\varepsilon \in \bar{D}$ with $x_\varepsilon \rightarrow x_0$ such that

$$\frac{v}{d^{s-1}}(x_0) \geq (v/d^{s-1})_\varepsilon(x_0) = \frac{v}{d^{s-1}}(x_\varepsilon) + \frac{|x_0 - x_\varepsilon|^2}{\varepsilon}.$$

Hence, the rate of convergence $x_\varepsilon \rightarrow x_0$ only depends on the oscillation of v/d^{s-1} . Then, since $\phi^{(\delta)}$ is a valid test function, we deduce that for any $x \in D$,

$$\phi^{(\delta)}(x + x_0 - x_\varepsilon) \leq (v/d^{s-1})_\varepsilon(x + x_0 - x_\varepsilon) \leq \frac{v}{d^{s-1}}(x) + \frac{|x_0 - x_\varepsilon|^2}{\varepsilon}$$

if $\varepsilon > 0$ is so small that $x + x_0 - x_\varepsilon \in D$. Since the aforementioned inequality becomes an equality in case $x = x_\varepsilon$, we deduce that indeed, $\phi^{(\delta)}(\cdot + x_0 - x_\varepsilon) - \varepsilon^{-1}|x_0 - x_\varepsilon|^2$ touches v/d^{s-1} from below at x_ε , as claimed.

We observe that $x_\varepsilon \notin \Omega$ since otherwise one would get a contradiction with the nonlocal equation satisfied by v , in the exact same way as in the proof of (2-16), if $\varepsilon > 0$ is small enough. Thus, $x_\varepsilon \in \partial\Omega \cap \bar{D}$, and from the boundary condition satisfied by v , it follows $\partial_\nu \phi^{(\delta)}(x_0) \leq g_1(x_\varepsilon)$. Thus, by the definition of $\phi^{(\delta)}$, we have $\partial_\nu \phi(x_0) = \partial_\nu \phi^{(\delta)}(x_0) + \delta \leq g_1(x_\varepsilon) + \delta$ for any $\delta > 0$. Thus, sending $\delta \rightarrow 0$ and recalling that $x_\varepsilon \rightarrow x_0$, as $\varepsilon \rightarrow 0$, this proves the first statement in (2-17) with $\delta_\varepsilon = g_1(x_\varepsilon) - g_1(x_0)$. Analogously, one proves the second claim in (2-17). \square

3. Nonlocal maximum principles with local Dirichlet and Neumann conditions

In this section, we establish weak maximum principles for nonlocal equations with local Dirichlet and Neumann data (see Propositions 1.3 and 1.5).

First, we establish a weak maximum principle for solutions to the inhomogeneous Dirichlet problem in (1-10) (see Proposition 1.3). Its proof goes by sliding the barrier subsolution ϕ from Lemma 2.6 underneath v from below.

Proof of Proposition 1.3. By assumption on v , we have that $v/d^{s-1} \in C(\bar{\Omega})$ with $v/d^{s-1} \geq 0$ on $\partial\Omega$. Let $z \in \partial\Omega$ be such that $\min_{\partial\Omega} v/d^{s-1} = v/d^{s-1}(z) =: l \geq 0$. Let $\varepsilon \in (0, s)$ and $M > 1$ to be chosen later, and recall the subsolution $\phi_l \in C(\Omega)$ from Lemma 2.6. We define

$$c_0 := \inf\{c \in \mathbb{R} : \phi_l/d^{s-1} - c \leq v/d^{s-1} \text{ in } \bar{\Omega}\}.$$

Since also $\phi_l/d^{s-1} \in C(\bar{\Omega})$, the above set is nonempty and $c_0 < \infty$. In fact, recalling the definition of ϕ_l , it must be

$$c_0 \leq \|v/d^{s-1}\|_{L^\infty(\bar{\Omega})} + \|(\phi_l)_+/d^{s-1}\|_{L^\infty(\bar{\Omega})} \leq \|v/d^{s-1}\|_{L^\infty(\bar{\Omega})} + l + c|\text{diam}(\Omega)|^\varepsilon, \tag{3-1}$$

which is independent of M . Moreover, since $\phi_l/d^{s-1}(z) = l = v/d^{s-1}(z)$, we have that $c_0 \geq 0$. Then, in particular, we have

$$\phi_l/d^{s-1} - c_0 \leq v/d^{s-1} \quad \text{in } \mathbb{R}^n, \quad \text{and} \quad \phi_l/d^{s-1}(x_0) - c_0 = v/d^{s-1}(x_0) \quad \text{for some } x_0 \in \bar{\Omega}.$$

In case $x_0 \in \Omega$, we have

$$\phi_l - c_0d^{s-1} - v \leq 0 \quad \text{in } \mathbb{R}^n \quad \text{and} \quad (\phi_l - c_0d^{s-1} - v)(x_0) = 0,$$

so it must be

$$0 \leq L(\phi_l - c_0d^{s-1} - v)(x_0) \leq L\phi_l(x_0) - c_0L(d^{s-1})(x_0) \leq -d^{\varepsilon-s-1}(x_0) - M + (l + c_0)cd^{\delta\gamma-s-1}(x_0),$$

where we used Lemma 2.6 and that $|L(d^{s-1})| \leq cd^{\delta\gamma-s-1}$ for any $\delta \in (0, s)$ by Lemma 2.3. Next, we fix any $\delta \in (0, s)$, and take $\varepsilon < \delta\gamma$ and M so large, depending only on $c_0, l, \text{diam}(\Omega)$ (but not on x_0), such that

$$-d^{\varepsilon-s-1}(x_0) - M + (l + c_0)cd^{\delta\gamma-s-1}(x_0) < 0.$$

Since c_0 is independent of M (see (3-1)), we obtain a contradiction. Thus, it must be $x_0 \in \partial\Omega$, which by construction yields that $c_0 = 0$, and therefore $\phi_l \leq v$ in Ω . Since $l \geq 0$, by Lemma 2.6, there exists $\delta > 0$ such that $\phi_l \geq 0$ in $\Omega \cap \{d \leq \delta\}$. Therefore, v is a viscosity solution to

$$\begin{cases} Lv \geq 0 & \text{in } \Omega \cap \{d > \delta\}, \\ v \geq 0 & \text{in } \mathbb{R}^n \setminus (\Omega \cap \{d > \delta\}). \end{cases}$$

Since $v \in C(\overline{\Omega \cap \{d > \delta\}})$, we can apply the maximum principle for viscosity solutions to v (see [Fernández-Real and Ros-Oton 2024a, Lemma 3.2.19]) and deduce that $v \geq 0$ in \mathbb{R}^n , as desired. \square

In particular, we have the following comparison principle:

Lemma 3.1. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{1,\gamma}$ for some $\gamma > 0$. Let $v, b \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1}, b/d^{s-1} \in C(\bar{\Omega})$ be viscosity solutions to*

$$\begin{cases} Lv \geq f & \text{in } \Omega, \\ v/d^{s-1} \geq 0 & \text{on } \partial\Omega, \end{cases} \quad \begin{cases} Lb \leq f & \text{in } \Omega, \\ b \leq v & \text{in } \mathbb{R}^n \setminus \Omega, \\ b/d^{s-1} \leq 0 & \text{on } \partial\Omega, \end{cases}$$

for some $f \in C(\Omega)$. Then, $v \geq b$ in \mathbb{R}^n .

Proof. Since by [Fernández-Real and Ros-Oton 2024a, Lemma 3.4.13] $w = v - b$ is a viscosity solution to $Lw \geq 0$ in Ω such that $w/d^{s-1} \geq 0$ on $\partial\Omega$, and $w \geq 0$ in $\mathbb{R}^n \setminus \Omega$, it satisfies the assumptions of Proposition 1.3. An application of this result concludes the proof. \square

As an application, we have the following version of a nonlocal Hopf lemma for viscosity solutions. The proof follows in the same way as [Fernández-Real and Ros-Oton 2024a, Proposition 2.6.6], where the Hopf lemma was proved for bounded solutions.

Lemma 3.2. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{1,\gamma}$ for some $\gamma > 0$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ satisfy, in the viscosity sense,*

$$\begin{cases} Lv = f \geq 0 & \text{in } \Omega, \\ v \geq 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ v/d^{s-1} \geq 0 & \text{on } \partial\Omega, \end{cases}$$

for some $f \in C(\Omega)$. Then, either $v \equiv 0$ in Ω , or

$$v(x) \geq C \left(\inf_{\{x \in \Omega: \text{dist}(x, \partial\Omega) \geq \delta\}} v \right) d^s(x) \quad \text{in } \Omega$$

for some $C, \delta > 0$, which depend only on $n, s, \lambda, \Lambda, \gamma, \text{diam}(\Omega)$, and the $C^{1,\gamma}$ radius of Ω .

Proof. First, by the weak maximum principle for viscosity solutions with boundary blow-up (see Proposition 1.3), we have $v \geq 0$ in \mathbb{R}^n . In order to deduce $v > 0$ in case $v \not\equiv 0$, one uses the nonlocal weak Harnack inequality (see [Fernández-Real and Ros-Oton 2024a, Theorem 3.3.1]). Then, we use the subsolution ϕ from [Fernández-Real and Ros-Oton 2024a, Corollary B.2.8] which satisfies

$$\begin{cases} L\phi \leq -1 & \text{in } N_\delta, \\ \max\{d^s, \delta^{-1}\} \geq \phi \geq \delta d^s & \text{in } \mathbb{R}^n, \end{cases}$$

for some $\delta > 0$ and where $N_\delta = \{0 < d < \delta\}$. Let us define

$$c_* = \min\{v(x) : x \in \Omega \setminus N_\delta\} > 0.$$

Then, we have

$$c_*\delta L\phi \leq Lv \quad \text{in } N_\delta \quad \text{and} \quad c_*\delta\phi \leq v \quad \text{in } \mathbb{R}^n \setminus N_\delta.$$

Hence, by the comparison principle in Lemma 3.1, we deduce that $c_*\delta\phi \leq v$ in \mathbb{R}^n , which implies the desired result. \square

Given a $C^{1,\gamma}$ domain $\Omega \subset \mathbb{R}^n$, let us now consider functions $b : \mathbb{R}^n \rightarrow \mathbb{R}$, which arise as the solution to the Dirichlet problem

$$\begin{cases} Lb = f_b & \text{in } \Omega, \\ b_\Omega = e_b & \text{in } \mathbb{R}^n \setminus \Omega, \\ b_\Omega/d^{s-1} = g_b & \text{on } \partial\Omega, \end{cases} \tag{3-2}$$

for some $f_b \geq 0$ with $f_b \not\equiv 0$, $e_b \geq 0$, and $g_b \geq 0$. With the maximum principle (see Proposition 1.3) at hand, the existence of b can be established using standard techniques. For well-posedness results in case $L = (-\Delta)^s$, we refer to [Abatangelo 2015]. Moreover, by Proposition 1.3, we have $b \geq 0$ in Ω , and by the same argument as in the proof of (7-1), we have $b/d^{s-1} \in L^\infty(\Omega)$. Moreover, if $\partial\Omega \in C^{2,\gamma}$ and f_b, e_b, g_b are smooth, then by Theorem 1.4, we have $b_\Omega/d^{s-1} \in C^{1,\gamma}(\bar{\Omega})$, and $\partial_\nu(b/d^{s-1})$ exists in the classical sense.

In the following, we will denote by b_Ω the solution to (3-2) with $f_b = g_b = 1$ and $e_b = 0$.

As a corollary of the previous results, we obtain the following pointwise formulation of a nonlocal Hopf lemma for solutions with boundary blow-up.

Lemma 3.3. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{1,\gamma}$ for some $\gamma > 0$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ satisfy, in the viscosity sense,*

$$\begin{cases} Lv \geq 0 & \text{in } \Omega, \\ v \geq 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ v/d^{s-1} = g & \text{on } \partial\Omega, \end{cases}$$

for some $g \in C(\partial\Omega)$. Let $x_0 \in \partial\Omega$ be such that $\min_{\bar{\partial\Omega}} g = g(x_0) \leq 0$. Then, either $v \equiv 0$, or we have that in the viscosity sense

$$\partial_v(v/b)(x_0) > 0$$

for any b as in (3-2) with $b/d^{s-1} = 1$ on $\partial\Omega \cap (\{g < 0\} \cup \{x_0\})$.

In particular, Lemma 3.3 implies that for the regularized distance d ,

$$\partial_v(v/d^{s-1})(x_0) = \partial_v(v/b)(x_0) + g(x_0)\partial_v(b/d^{s-1})(x_0) > g(x_0)\partial_v(b/d^{s-1})(x_0).$$

We stress that the sign of the right-hand side depends on the choice of the regularized distance d .

Proof. Since $g(x_0) \leq 0$ we have by the construction of b in (3-2)

$$\begin{aligned} L(v - g(x_0)b) &\geq -g(x_0) \geq 0 && \text{in } \Omega, \\ v - g(x_0)b &\geq 0 && \text{in } \mathbb{R}^n \setminus \Omega, \\ (v - g(x_0)b)/d^{s-1} &= g - g(x_0) \geq 0 && \text{on } \partial\Omega \cap \{g < 0\}, \\ (v - g(x_0)b)/d^{s-1} &\geq g \geq 0 && \text{on } \partial\Omega \cap \{g \geq 0\}. \end{aligned}$$

Thus, an application of Lemma 3.2 to $v - g(x_0)b$ yields that either $v - g(x_0)b \equiv 0$ in Ω , or

$$v - g(x_0)b \geq cd^s \quad \text{near } x_0. \tag{3-3}$$

We cannot have $v - g(x_0)b \equiv 0$, unless $g(x_0) = 0$ (in which case $v \equiv v - g(x_0)b \equiv 0$), since then

$$Lv = g(x_0)Lb \leq g(x_0) < 0 \quad \text{in } \Omega,$$

a contradiction. Thus, unless $v \equiv 0$, we have (3-3), and we compute, using that $b \geq 0$ and $(b/d^{s-1})(x_0) = 1$,

$$\partial_v(v/b)(x_0) = \lim_{x \rightarrow x_0} \frac{\frac{v(x)}{b(x)} - g(x_0)}{d(x)} = \lim_{x \rightarrow x_0} \frac{v(x) - g(x_0)b(x)}{b(x)d(x)} \geq c \lim_{x \rightarrow x_0} \frac{d^{s-1}(x)}{b(x)} = c > 0.$$

If the limit in the previous estimate does not exist, we need to interpret the boundary condition in the viscosity sense, i.e., take any smooth ψ with $\psi(x_0) = (v/d^{s-1})(x_0) = g(x_0)$ and $\psi \geq v/d^{s-1}$. Then, the limit $\partial_v\psi(x_0)$ exists, and an analogous computation as above yields $\partial_v\psi(x_0) \geq c > 0$, i.e., $\partial_v(v/b)(x_0) > 0$ in the viscosity sense. □

Finally, we are in a position to prove the main result of this section, a maximum principle for nonlocal equations with local Neumann conditions.

Lemma 3.4. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{2,\gamma}$ for some $\gamma > 0$ and $K \in C^{5+2\gamma}(\mathbb{S}^{n-1})$. Let $\Gamma \subset \partial\Omega$, $v \in L_{2s}^1(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ satisfy, in the viscosity sense,*

$$\begin{cases} Lv \geq f & \text{in } \Omega, \\ v \geq 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \partial_\nu(v/b) \leq g & \text{on } \partial\Omega \setminus \Gamma, \\ v/b \geq 0 & \text{on } \partial\Omega \cap \Gamma \end{cases}$$

for some $f \in C(\Omega)$ with $d^{s+1-\varepsilon} f \in L^\infty(\Omega)$ for some $\varepsilon \in (0, s]$, and $g \in C(\partial\Omega)$. Here, b is as in (3-2) with $b/d^{s-1} = 1$ on $\partial\Omega \setminus \Gamma'$ for some $\Gamma' \Subset \Gamma$. Then, there exists $c > 0$, depending only on $n, s, \lambda, \Lambda, \gamma, \varepsilon$, and the $C^{2,\gamma}$ radius of Ω and $\text{diam}(\Omega)$, such that

$$v/d^{s-1} \geq -c\|d^{s-\varepsilon} f\|_{L^\infty(\Omega)} - c\|g\|_{L^\infty(\partial\Omega \setminus \Gamma)} \quad \text{in } \Omega.$$

Proof. The case $f \geq 0$ and $g \leq 0$ follows from the Hopf lemma (see Lemma 3.3). In fact, since $v/b \in C(\partial\Omega)$, there exists $x_0 \in \partial\Omega$ with $\min_{\partial\Omega}(v/b) = (v/b)(x_0)$. If $(v/b)(x_0) \geq 0$, then we have that $v/d^{s-1} \geq 0$ on $\partial\Omega$. Otherwise, $(v/b)(x_0) < 0$, and then by assumption it must be $x_0 \in \partial\Omega \setminus \Gamma$. However, in this case Lemma 3.3 implies that either $v \equiv 0$, (in which case we are done), or $\partial_\nu(v/b)(x_0) > 0$, which contradicts $g(x_0) \leq 0$. Thus, we must have $v/d^{s-1} \geq 0$ on $\partial\Omega$. However, by the weak maximum principle (see Proposition 1.3), this implies $v \geq 0$, as desired.

Now, we explain how to get the result with general f, g . To do so, let $\tilde{\psi}_1$ be the solution to

$$\begin{cases} L\tilde{\psi}_1 = 0 & \text{in } \Omega, \\ \tilde{\psi}_1 = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \tilde{\psi}_1/d^{s-1} = h & \text{on } \partial\Omega, \end{cases}$$

for some smooth function h which satisfies $0 \leq h \leq 1$, and is such that $h = 1$ on $\partial\Omega \setminus \Gamma$, and $h = 0$ in $\partial\Omega \cap \Gamma'$.

From Lemma 3.3, we deduce that $\partial_\nu(\tilde{\psi}_1/b) < 0$ on $\partial\Omega \setminus \Gamma$. Since $\partial\Omega \in C^{2,\gamma}$, by Theorem 1.4 we have that $\partial_\nu(\tilde{\psi}_1/b) \in C^\gamma(\partial\Omega)$, and therefore, there is $c_0 > 0$ such that

$$\partial_\nu(\tilde{\psi}_1/b) \leq -c_0 < 0 \quad \text{on } \partial\Omega \setminus \Gamma.$$

Moreover, let us denote by $\tilde{\psi}_2$ the function $\tilde{\psi}$ from the second claim of Lemma 2.7, which satisfies for some $c_2 > 0$,

$$\begin{cases} L\tilde{\psi}_2 \geq d^{\varepsilon-s} & \text{in } \Omega, \\ \tilde{\psi}_2 = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \tilde{\psi}_2/d^{s-1} = 0 & \text{on } \partial\Omega, \\ \partial_\nu(\tilde{\psi}_2/b) \leq c_2 & \text{on } \partial\Omega. \end{cases}$$

Hence, if we take $M = c_0^{-1}(c_2 + 1) > 0$ and define $\tilde{\psi} := M\tilde{\psi}_1 + \tilde{\psi}_2$, we obtain

$$\begin{cases} L\tilde{\psi} \geq d^{\varepsilon-s} & \text{in } \Omega, \\ \tilde{\psi} = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \tilde{\psi}/d^{s-1} = Mh & \text{on } \partial\Omega, \\ \partial_\nu(\tilde{\psi}/b) \leq -1 & \text{on } \partial\Omega \setminus \Gamma. \end{cases}$$

We apply the previous argument with v replaced by

$$w = v + (\|d^{s-\varepsilon} f\|_{L^\infty(\Omega)} + \|g\|_{L^\infty(\partial\Omega \setminus \Gamma)})\tilde{\psi}.$$

Then, we have that in the viscosity sense,

$$\begin{cases} Lw \geq f + d^{\varepsilon-s} (\|d^{s-\varepsilon} f\|_{L^\infty(\Omega)} + \|g\|_{L^\infty(\partial\Omega \setminus \Gamma)}) \geq 0 & \text{in } \Omega, \\ w \geq v \geq 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ \partial_\nu(w/b) \leq g - (\|d^{s-\varepsilon} f\|_{L^\infty(\Omega)} + \|g\|_{L^\infty(\partial\Omega \setminus \Gamma)}) \leq 0 & \text{on } \partial\Omega \setminus \Gamma, \\ w/b \geq Mh \geq 0 & \text{on } \partial\Omega \cap \Gamma. \end{cases}$$

Altogether, by the same argument as at the beginning of the proof, we have $w \geq 0$ in Ω . Let us now observe that by construction and the same argument as in the proof of (7-1) we have

$$\tilde{\psi} \leq Cd^{s-1} \quad \text{in } \Omega$$

for some $C > 0$. Therefore, we obtain

$$v \geq -\tilde{\psi} (\|d^{s-\varepsilon} f\|_{L^\infty(\Omega)} + \|g\|_{L^\infty(\partial\Omega \setminus \Gamma)}) \geq -Cd^{s-1} (\|d^{s-\varepsilon} f\|_{L^\infty(\Omega)} + \|g\|_{L^\infty(\partial\Omega \setminus \Gamma)}) \quad \text{in } \Omega,$$

as desired. □

Proof of Proposition 1.5. This is a special case of Lemma 3.4. □

4. Hölder estimates up to the boundary

The previous maximum principle for nonlocal equations with local Neumann conditions (see Lemma 3.4) puts us in a position to establish a Harnack inequality for solutions to (1-7) at the boundary, which will eventually lead to the Hölder regularity estimate in Theorem 1.6.

To prove it, we adapt some of the ideas in [Lian and Zhang 2023] to the framework of solutions to nonlocal problems which blow up at the boundary.

For $\delta > 0$, let us define $\Omega_\delta = \{x \in \Omega : \text{dist}(x, \partial\Omega) \geq \delta\}$.

Lemma 4.1. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $0 \in \partial\Omega$ and $\partial\Omega \in C^{2,\gamma}$ for some $\gamma > 0$ and $K \in C^{5+2\gamma}(\mathbb{S}^{n-1})$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to*

$$\begin{cases} Lv \geq f & \text{in } \Omega \cap B_1, \\ v \geq 0 & \text{in } \mathbb{R}^n, \\ \partial_\nu(v/b_\Omega) \leq g & \text{on } \partial\Omega \cap B_1 \end{cases}$$

for some $f \in C(\Omega \cap B_1)$ with $d^{s-\alpha} f \in L^\infty(\Omega \cap B_1)$ for some $\alpha \in (0, s]$, and $g \in C(\overline{\partial\Omega \cap B_1})$. Assume that $0 \in \partial\Omega$. Then,

$$\int_{\Omega_{1/2} \cap B_1} (v/b_\Omega) \, dx \leq c \inf_{\Omega \cap B_{\eta^{-1}}} (v/b_\Omega) + c(\|d^{s-\alpha} f_-\|_{L^\infty(\Omega \cap B_1)} + \|g_+\|_{L^\infty(\partial\Omega \cap B_1)}),$$

where $\eta > 1$ and $c > 0$ depend only on $n, s, \lambda, \Lambda, \gamma, \alpha$, and the $C^{2,\gamma}$ radius of Ω . Here, b_Ω is defined as in (3-2).

Proof. The interior weak Harnack inequality for viscosity supersolutions (see [Fernández-Real and Ros-Oton 2024a, Theorem 3.3.1]) applied with v implies

$$\int_{\Omega_{1/2} \cap B_1} v(x) \, dx \leq c \inf_{x \in \Omega_{1/2} \cap B_1} v(x) + c\|f\|_{L^\infty(\Omega_{1/2} \cap B_1)},$$

where $c > 0$ depends on $n, s, \lambda, \Lambda, \eta, \alpha$. Moreover, since $b \asymp c > 0$ in $\Omega_{1/2} \cap B_1$, it follows for $u := v/b$ by Lemma 3.2 that

$$\int_{\Omega_{1/2} \cap B_1} u(x) \, dx \leq c \inf_{x \in \Omega_{1/2} \cap B_1} u(x) + c\|d^{s-\alpha} f\|_{L^\infty(\Omega_{1/2} \cap B_1)}.$$

Thus, it remains to show

$$\inf_{x \in \Omega_{1/2} \cap B_1} u(x) \leq c \inf_{x \in \Omega \cap B_{\eta^{-1}}} u(x) + c(\|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_1)} + \|g\|_{L^\infty(\partial\Omega \cap B_1)}). \quad (4-1)$$

Since $v \geq 0$, by the weak Harnack inequality, either $v \equiv 0$ in $\Omega_{1/2} \cap B_1$, or $\inf_{\Omega_{1/2} \cap B_1} v > 0$. Therefore, without loss of generality, we can assume that $\inf_{\Omega_{1/2} \cap B_1} v = 1$.

To prove (4-1), let us take a set $D \subset \mathbb{R}^n$ with $\partial D \in C^{2,\gamma}$ such that

$$\Omega \cap B_{1/2} \subset D \subset \Omega \cap B_1,$$

Let w be a function such that

$$\left\{ \begin{array}{ll} Lw = 0 & \text{in } D, \\ w \leq 1 & \text{in } (\Omega_{1/2} \cap B_1) \setminus D, \\ w = 0 & \text{in } \mathbb{R}^n \setminus (D \cup (\Omega_{1/2} \cap B_1)), \\ \partial_\nu(w/b_\Omega) \geq 0 & \text{on } \partial D \cap B_{2\eta^{-1}}, \\ w/b_\Omega \leq 0 & \text{on } \partial D \setminus B_{2\eta^{-1}}, \\ w/b_\Omega \geq c_1 & \text{in } \Omega \cap B_{\eta^{-1}}. \end{array} \right.$$

We construct w as follows. Let $h : \partial D \rightarrow \mathbb{R}$ and $e : \mathbb{R}^n \setminus D \rightarrow \mathbb{R}$ be smooth functions such that for some $\eta < \frac{1}{8}$

$$h = \begin{cases} 0 & \text{on } \partial D \setminus B_{2\eta^{-1}}, \\ c_1 & \text{on } \partial D \cap B_{\eta^{-1}}, \end{cases} \quad e = \begin{cases} 1 & \text{in } T, \\ 0 & \text{in } \mathbb{R}^n \setminus (D \cup (\Omega_{1/2} \cap B_1)), \end{cases}$$

where $T \Subset (\Omega_{1/2} \cap B_1) \setminus D$, $0 \leq h \leq c_1$, and $0 \leq e \leq 1$. We let w be the solution to

$$\left\{ \begin{array}{ll} Lw = 0 & \text{in } D, \\ w = e & \text{in } \mathbb{R}^n \setminus D, \\ w/b_D = h & \text{on } \partial D. \end{array} \right.$$

Then, we can show that $\partial_\nu(w/b_\Omega) \geq C > 0$ in $\partial D \cap B_{2\eta^{-1}}$ (for any given $C > 0$) by making $c_1 > 0$ small enough. Indeed, if w_1 solves the Dirichlet problem with boundary data zero and exterior data e , then by the Hopf lemma (see Lemma 3.2), we have since $w_1/d_D^s \in C^{1,\gamma}(\bar{D})$ by the boundary regularity results in [Abatangelo and Ros-Oton 2020], and since $\partial D \cap B_{2\eta^{-1}} \Subset \partial\Omega$,

$$\begin{aligned} \partial_\nu(w_1/b_\Omega) &= \partial_\nu(w_1/d_D^{s-1})(d_D^{s-1}/b_\Omega) + (w_1/d_D^{s-1})\partial_\nu(d_D^{s-1}/b_\Omega) \\ &= \partial_\nu(w_1/d_D^s)d_D + (w_1/d_D^s)\partial_\nu(d_D) = w_1/d_D^s \geq c_0 > 0 \quad \text{on } \partial D \cap B_{2\eta^{-1}}. \end{aligned}$$

Moreover, if w_2 solves the Dirichlet problem with boundary data h and exterior data zero, we get from Theorem 1.4 that $|\partial_\nu(w_2/b_\Omega)| \leq c_3c_1$ in $B_{2\eta^{-1}}$ for some $c_3 > 0$. Hence, choosing $c_1 > 0$ small enough, we deduce the claim for $w = (C/c_0)w_1 + w_2$.

Thus, we have by construction, and using that $\inf_{\Omega_{1/2} \cap B_1} v = 1$, and $w \asymp d_D^s$ near $\partial D \setminus \partial\Omega$,

$$\left\{ \begin{array}{ll} L(v - w) \geq f & \text{in } D, \\ v - w \geq 0 & \text{in } \mathbb{R}^n \setminus D, \\ \partial_\nu((v - w)/b_\Omega) \leq g & \text{on } \partial D \cap B_{2\eta^{-1}}, \\ (v - w)/b_\Omega \geq 0 & \text{on } \partial D \setminus B_{2\eta^{-1}}. \end{array} \right.$$

Note that b_Ω satisfies

$$\left\{ \begin{array}{ll} Lb_\Omega \geq 0 & \text{in } D, \\ Lb_\Omega \neq 0 & \text{in } D, \\ b_\Omega \geq 0 & \text{in } \mathbb{R}^n \setminus D, \\ b_\Omega/d_D^{s-1} = 1 & \text{on } \partial D \cap B_{4\eta^{-1}}, \\ b_\Omega/d_D^{s-1} \geq 0 & \text{on } \partial D. \end{array} \right.$$

Since $(\partial D \cap B_{4\eta^{-1}}) \ni (\partial D \cap B_{2\eta^{-1}})$, we can apply the maximum principle for the Neumann problem Lemma 3.4 with $\Gamma = \partial D \setminus B_{2\eta^{-1}}$ and $b = b_\Omega$, and deduce

$$(v - w)/b_\Omega \geq -c\|d^{s-\alpha} f_-\|_{L^\infty(D \cap B_1)} - c\|g_+\|_{L^\infty(\partial D \cap B_1)} \quad \text{in } D \cap B_1.$$

Since, by construction, we also have

$$w/b_\Omega \geq c_1 = c_1 \inf_{\Omega_{1/2} \cap B_1} v \geq c_2 \inf_{\Omega_{1/2} \cap B_1} u \quad \text{in } \Omega \cap B_{\eta^{-1}},$$

for some $c_2 > 0$, since $b_\Omega \asymp c > 0$ in $\Omega_{1/2} \cap B_1$, we deduce

$$\begin{aligned} v/b_\Omega &= (w + v - w)/b_\Omega \\ &\geq c_2 \inf_{\Omega_{1/2} \cap B_1} u - c\|d^{s-\alpha} f_-\|_{L^\infty(D \cap B_1)} - c\|g_+\|_{L^\infty(\partial D \cap B_1)} \\ &\geq c_2 \inf_{\Omega_{1/2} \cap B_1} u - c\|d^{s-\alpha} f_-\|_{L^\infty(\Omega \cap B_1)} - c\|g_+\|_{L^\infty(\partial\Omega \cap B_1)} \quad \text{in } \Omega \cap B_{\eta^{-1}}, \end{aligned}$$

where we used $D \cap B_1 \subset \Omega \cap B_1$. Hence, we obtain (4-1), as desired. □

As a corollary of the previous weak Harnack inequality at the boundary, we obtain a growth lemma.

Lemma 4.2. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{2,\gamma}$ for some $\gamma > 0$ and $K \in C^{5+2\gamma}(\mathbb{S}^{n-1})$. Let $\eta > 1$ be as in Lemma 4.1. Assume that $x_0 \in \partial\Omega$ and let $0 < R \leq 1$. Let $v \in L_{2s}^1(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to*

$$\left\{ \begin{array}{ll} Lv \geq f & \text{in } \Omega \cap B_R(x_0), \\ \partial_\nu(v/b_\Omega) \leq g & \text{on } \partial\Omega \cap B_R(x_0), \\ v \geq 0 & \text{in } B_R(x_0), \\ v \geq b_\Omega(1 - \eta^{j\beta}) & \text{in } B_{\eta^j R}(x_0) \cap \Omega \text{ for all } j \geq 1, \\ v \geq (1 - \eta^{j\beta}) & \text{in } B_{\eta^j R}(x_0) \setminus \Omega \text{ for all } j \geq 1, \\ |\Omega_{R/4} \cap B_{R/2}(x_0) \cap \{v/b_\Omega \geq \frac{1}{4}\}| \geq \frac{1}{2} |\Omega_{R/4} \cap B_{R/2}(x_0)| \end{array} \right.$$

for some $f \in C(\Omega \cap B_R(x_0))$ with $d^{s-\alpha} f \in L^\infty(\Omega \cap B_R(x_0))$ for some $\alpha \in (0, s]$, and $g \in C(\overline{\partial\Omega \cap B_R(x_0)})$. Then, there exist $\delta > 0$, and $\beta \in (0, 1)$, depending only on $n, s, \lambda, \Lambda, \gamma, \alpha$, and the $C^{2,\gamma}$ radius of Ω , such that

$$\inf_{\Omega \cap B_{\eta^{-1}R}(x_0)} (v/b_\Omega) + R^{1+\alpha} \|d^{s-\alpha} f_-\|_{L^\infty(\Omega \cap B_R(x_0))} + R \|g_+\|_{L^\infty(\partial\Omega \cap B_R(x_0))} \geq \delta.$$

Proof. Let us assume without loss of generality that $x_0 = 0$. The proof follows from an application of the weak Harnack inequality (see Lemma 4.1) to v_+ . It is slightly involved due to the appearance of the tail term.

Indeed, we have

$$Lv_+(x) \geq f(x) - \int_{\mathbb{R}^n \setminus B_R} v_-(y) K(x-y) dy =: \tilde{f}(x),$$

where we used that by assumption, $v \geq 0$ in B_R . Then, we obtain from Lemma 4.1 (after scaling), using the last assumption and setting $u := v/b_\Omega$,

$$\inf_{\Omega \cap B_{\eta^{-1}R}} u + R^{1+\alpha} \|d^{s-\alpha} \tilde{f}_-\|_{L^\infty(\Omega \cap B_{R/2})} + R \|g_+\|_{L^\infty(\partial\Omega \cap B_{R/2})} \geq c_0 \int_{\Omega_{R/4} \cap B_{R/2}} u dx \geq \frac{c_0}{8}, \tag{4-2}$$

where $c_0 > 0$ is the constant from the weak Harnack inequality.

Next, we estimate $\|d^{s-\alpha} \tilde{f}\|_{L^\infty(\Omega \cap B_R)}$. To do so, we apply a similar reasoning as in the proof of Lemma 2.2. First, we recall that for any $x \in B_{R/2}$, there exists $\kappa > 0$ such that for any $t \in (0, \kappa)$,

$$\mathcal{H}^{n-1}(\{d = t\} \cap B_{\eta^j R} \setminus B_{\eta^{j-1} R}) \leq C(\eta^j R)^{n-1},$$

where $C > 0$ depends only on n and the $C^{2,\gamma}$ radius of Ω (we refer to [Fernández-Real and Ros-Oton 2024a, Lemma B.2.4] for a reference of this fact). Next, we observe that by the coarea formula, and since $0 \leq b_\Omega \leq C d^{s-1}$,

$$\begin{aligned} & R^{1+s} \int_{\Omega \setminus B_R} v_-(y) K(x-y) dy \\ & \leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) R^{1+s} \int_{\Omega \cap (B_{\eta^j R} \setminus B_{\eta^{j-1} R})} d^{s-1}(y) |y|^{-n-2s} dy \\ & \leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) R^{1+s} \left((\eta^j R)^{-n-2s} \int_{(B_{\eta^j R} \setminus B_{\eta^{j-1} R}) \cap \{d \leq \kappa\}} d^{s-1}(y) |\nabla d(y)| dy + \kappa^{s-1} (\eta^j R)^{-2s} \right) \end{aligned}$$

$$\begin{aligned} &\leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) R^{1+s} \left((\eta^j R)^{-n-2s} \int_0^{\min\{\eta^j R, \kappa\}} t^{s-1} \left(\int_{(B_{\eta^j R} \setminus B_{\eta^{j-1} R}) \cap \{d=t\}} d\mathcal{H}^{n-1}(y) \right) dt + (\eta^j R)^{-2s} \right) \\ &\leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) R^{1+s} \left((\eta^j R)^{-1-s} + (\eta^j R)^{-2s} \right) \leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) \eta^{-2sj} \end{aligned}$$

for some $c > 0$, depending only on $n, s, \lambda, \Lambda, \kappa, C, \eta$, where we also used that $R \leq 1$. Similarly,

$$\begin{aligned} R^{1+s} \int_{(\mathbb{R}^n \setminus \Omega) \setminus B_R} v_-(y) K(x-y) dy &\leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) R^{1+s} \int_{\mathbb{R}^n \cap (B_{\eta^j R} \setminus B_{\eta^{j-1} R})} |y|^{-n-2s} dy \\ &\leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) R^{1+s} (\eta^j R)^{-2s} \leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) \eta^{-2js}, \end{aligned}$$

where we used that $R^{1-s} \leq 1$, and $c > 0$ depends only on n, s, Λ . Therefore, we obtain

$$R^{1+s} \int_{\mathbb{R}^n \setminus B_R} d^{s-1}(y) v_-(y) K(x-y) dy \leq c \sum_{j \geq 1} (1 - \eta^{j\beta}) \eta^{-2js}.$$

Since this quantity vanishes as $\beta > 0$ goes to zero, we can make the whole expression smaller than $c_0/16$, which implies, by recalling the definition of \tilde{f} ,

$$\begin{aligned} R^{1+\alpha} \|d^{s-\alpha} \tilde{f}_-\|_{L^\infty(\Omega \cap B_{R/2})} &\leq R^{1+\alpha} \|d^{s-\alpha} f_-\|_{L^\infty(\Omega \cap B_{R/2})} + R^{1+\alpha} \left\| \int_{\Omega \setminus B_R} v_-(y) K(\cdot - y) dy \right\|_{L^\infty(B_{R/2})} \\ &\leq R^{1+\alpha} \|d^{s-\alpha} f_-\|_{L^\infty(\Omega \cap B_R)} + \frac{c_0}{16}, \end{aligned}$$

and therefore by the estimate (4-2)

$$\inf_{\Omega \cap B_{\eta^{-1}R}} u + R^{1+\alpha} \|d^{s-\alpha} f_-\|_{L^\infty(\Omega \cap B_R)} + R \|g_+\|_{L^\infty(\partial\Omega \cap B_R)} \geq \frac{c_0}{8} - \frac{c_0}{16} = \frac{c_0}{16},$$

as desired. □

We are now in a position to prove the boundary Hölder regularity.

Lemma 4.3. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{2,\gamma}$ for some $\gamma > 0$ and $K \in C^{5+2\gamma}(\mathbb{S}^{n-1})$. Assume that $x_0 \in \partial\Omega$ and let $0 < R \leq 1$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to*

$$\begin{cases} Lv = f & \text{in } \Omega \cap B_R(x_0), \\ v = 0 & \text{in } B_R(x_0) \setminus \Omega, \\ \partial_\nu(v/b_\Omega) = g & \text{on } \partial\Omega \cap B_R(x_0) \end{cases}$$

for some $f \in C(\Omega \cap B_R(x_0))$ and $g \in C(\overline{\partial\Omega \cap B_R(x_0)})$. Then, there exist $c > 0$, and $\alpha_0 \in (0, 1)$, depending only on $n, s, \lambda, \Lambda, \gamma$, and the $C^{2,\gamma}$ radius of Ω , such that if $d^{s-\alpha} f \in L^\infty(\Omega \cap B_R(x_0))$ for some $\alpha \in (0, \alpha_0]$, then it holds that

$$\begin{aligned} [v/d^{s-1}]_{C^\alpha(\Omega \cap B_{R/2}(x_0))} &\leq cR^{-\alpha} \left(\|v/d^{s-1}\|_{L^\infty(\Omega)} + \|v\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + R^{1+\alpha} \|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_R(x_0))} + R \|g\|_{L^\infty(\partial\Omega \cap B_R(x_0))} \right). \end{aligned}$$

Proof. Let us assume without loss of generality that $x_0 = 0$. We will prove the desired result in two steps. Let us denote by $\eta > 1$ the constant from Lemma 4.2.

Step 1. We claim that for any $k \in \mathbb{N}$,

$$\operatorname{osc}_{B_{\eta^{-k}R}}(v/b_\Omega) \leq c\eta^{-\alpha k} \left(\|v/d^{s-1}\|_{L^\infty(\Omega)} + \|v\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + R^{1+\alpha} \|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_R)} + R^\alpha + R \|g\|_{L^\infty(\partial\Omega \cap B_R)} \right),$$

for some constant $c > 0$, depending only on $n, s, \lambda, \Lambda, \gamma$, and the $C^{2,\gamma}$ radius of Ω . To prove it, we set $\alpha_0 := \min\{\beta, \gamma s, 1 - s[-\log_\eta(1 - \delta'/2)]\}$, and $\delta := 1 - \eta^{-\alpha_0}$, where δ', β, η are the constants from Lemma 4.2. This yields

$$(1 - \delta) = \eta^{-\alpha_0}, \quad \alpha_0 \leq \min\{\beta, \gamma s, 1 - s\}, \quad \delta \leq \delta'/2. \quad (4-3)$$

Let us set $u = v/b_\Omega$, take $\alpha \in (0, \alpha_0]$, and

$$M := 4\delta^{-1}c_1 \left(\|v/d^{s-1}\|_{L^\infty(\Omega)} + \|v\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + R^{1+\alpha} \|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_R(x_0))} + R^\alpha + R \|g\|_{L^\infty(\partial\Omega \cap B_R(x_0))} \right),$$

where $c_1 > 0$ denotes the constant c_1 from Lemma 2.3.

The claim of Step 1 will follow immediately, once we construct an increasing sequence $(m_k)_k$ and a decreasing sequence $(M_k)_k$ such that for any $k \in \mathbb{N}$,

$$m_k \leq u \leq M_k \quad \text{in } B_{\eta^{-k}R}, \quad (4-4)$$

$$M_k - m_k = M\eta^{-\alpha k}. \quad (4-5)$$

We prove (4-4) and (4-5) by induction. Setting $m_0 = -(\delta/2c_1)M$, $M_0 = (\delta/2c_1)M$, we obtain the desired results for $k = 0$. Let us now assume that (4-4) and (4-5) hold true for any $j \leq k - 1$.

We will now prove it for k . Clearly, one of the following two options always holds true:

$$\begin{aligned} \left| \Omega_{\eta^{-(k-1)R/4}} \cap B_{\eta^{-(k-1)R/2}} \cap \left\{ u \geq \frac{1}{2}(M_{k-1} + m_{k-1}) \right\} \right| &\geq \frac{1}{2} \left| \Omega_{\eta^{-(k-1)R/4}} \cap B_{\eta^{-(k-1)R/2}} \right|, \\ \left| \Omega_{\eta^{-(k-1)R/4}} \cap B_{\eta^{-(k-1)R/2}} \cap \left\{ u \geq \frac{1}{2}(M_{k-1} + m_{k-1}) \right\} \right| &\leq \frac{1}{2} \left| \Omega_{\eta^{-(k-1)R/4}} \cap B_{\eta^{-(k-1)R/2}} \right|. \end{aligned}$$

In the first case, and in the second case, we define

$$w = \frac{v - (b_\Omega + \mathbb{1}_{\mathbb{R}^n \setminus \Omega} \mathbb{1}_{\{m_{k-1} < 0\}})m_{k-1}}{M_{k-1} - m_{k-1}}, \quad w = \frac{(b_\Omega + \mathbb{1}_{\mathbb{R}^n \setminus \Omega} \mathbb{1}_{\{M_{k-1} > 0\}})M_{k-1} - v}{M_{k-1} - m_{k-1}}, \quad \text{respectively.}$$

Let us assume that we are in the first case. The proof of the second case goes via the same arguments, and we will skip it. Let us verify that w satisfies the assumptions of Lemma 4.2. First, if $u(x) \geq \frac{1}{2}(M_{k-1} + m_{k-1})$ for some $x \in \Omega$, it follows that

$$\frac{w}{b_\Omega}(x) = \frac{u(x) - m_{k-1}}{M_{k-1} - m_{k-1}} \geq \frac{u(x) - m_{k-1}}{M_{k-1} - m_{k-1}} \geq \frac{\frac{M_{k-1} + m_{k-1}}{2} - m_{k-1}}{M_{k-1} - m_{k-1}} = \frac{1}{2},$$

Thus, as an immediate consequence of being in the first case, we get

$$\left| \Omega_{\eta^{-(k-1)R/4}} \cap B_{\eta^{-(k-1)R/2}} \cap \left\{ \frac{w}{b_\Omega} \geq \frac{1}{2} \right\} \right| \geq \frac{1}{2} \left| \Omega_{\eta^{-(k-1)R/4}} \cap B_{\eta^{-(k-1)R/2}} \right|.$$

Moreover, by (4-4) (for $k - 1$), we have

$$w = \frac{v - b_\Omega m_{k-1}}{M_{k-1} - m_{k-1}} \geq 0 \quad \text{in } B_{\eta^{-(k-1)}R} \cap \Omega.$$

Nonnegativity of w in $B_{\eta^{-(k-1)}R} \setminus \Omega$ follows by assumption and construction. We obtain

$$|L(b_\Omega + \mathbb{1}_{\mathbb{R}^n \setminus \Omega})| \leq c_1 \quad \text{in } \Omega,$$

and therefore $d^{s-\alpha} L(b_\Omega + \mathbb{1}_{\mathbb{R}^n \setminus \Omega}) \in L^\infty(\Omega \cap B_R)$. Then, by (4-5) (for $k - 1$) we have

$$Lw = \frac{f - L(b_\Omega + \mathbb{1}_{\mathbb{R}^n \setminus \Omega} \mathbb{1}_{\{m_{k-1} < 0\}}) m_{k-1}}{M_{k-1} - m_{k-1}} \geq \frac{f - c_1 m_{k-1}}{M_{k-1} - m_{k-1}} \quad \text{in } \Omega \cap B_R. \quad (4-6)$$

Moreover, clearly

$$\partial_\nu(w/b_\Omega) = \frac{g - \partial_\nu(b_\Omega/b_\Omega) m_{k-1}}{M_{k-1} - m_{k-1}} = \frac{g}{M_{k-1} - m_{k-1}} \quad \text{on } \partial\Omega \cap B_R.$$

It remains to verify the fourth and fifth assumption of Lemma 4.2. Let us first consider $j \leq k - 1$. In that case, for any $x \in B_{\eta^{-(k-1)+j}R} \cap \Omega$ it holds by (4-4) and (4-5) that

$$\begin{aligned} \frac{w}{b_\Omega}(x) &= \frac{u(x) - m_{k-1}}{M_{k-1} - m_{k-1}} \geq \frac{m_{k-j-1} - m_{k-1}}{M_{k-1} - m_{k-1}} \\ &\geq \frac{M_{k-1} - M_{k-j-1} + m_{k-j-1} - m_{k-1}}{M_{k-1} - m_{k-1}} = 1 - \frac{M_{k-j-1} - m_{k-j-1}}{M_{k-1} - m_{k-1}} = 1 - \eta^{\alpha j}. \end{aligned}$$

Clearly, for any $x \in B_{\eta^{-(k-1)+j}R} \setminus \Omega$ and in case $m_{k-1} < 0$, by the same arguments as above, using (4-4), we have

$$w(x) = \frac{v(x) - m_{k-1}}{M_{k-1} - m_{k-1}} \geq \frac{m_{k-j-1} - m_{k-1}}{M_{k-1} - m_{k-1}} \geq 1 - \eta^{\alpha j}.$$

If however $m_{k-1} \geq 0$, then we can use that $v = 0$ in $B_R \setminus \Omega$. Moreover, if $j > k - 1$ we compute for $x \in B_{\eta^{-(k-1)+j}R} \cap \Omega$,

$$\begin{aligned} \frac{w}{b_\Omega}(x) &= \frac{u(x) - m_{k-1}}{M_{k-1} - m_{k-1}} \geq \frac{m_0 - m_{k-1}}{M_{k-1} - m_{k-1}} \\ &\geq \frac{(M_{k-1} - m_{k-1}) - (M_0 - m_0)}{M_{k-1} - m_{k-1}} = 1 - \eta^{\alpha(k-1)} \geq 1 - \eta^{\alpha j}. \end{aligned}$$

Finally, for $x \in B_{\eta^{-(k-1)+j}R} \setminus \Omega$, again by the same arguments as above, and using that $v \geq m_0$ by construction, we have

$$w(x) = \frac{v(x) - m_{k-1} \mathbb{1}_{\{m_{k-1} < 0\}}}{M_{k-1} - m_{k-1}} \geq \frac{m_0 - m_{k-1}}{M_{k-1} - m_{k-1}} \geq 1 - \eta^{\alpha j}.$$

Consequently, all assumptions of Lemma 4.2 are satisfied for w with radius $\eta^{-(k-1)}R$. Thus, we deduce from Lemma 4.2 and the choice of δ ,

$$\begin{aligned} u - m_{k-1} &= (M_{k-1} - m_{k-1}) \frac{w}{b_\Omega} \\ &\geq 2\delta(M_{k-1} - m_{k-1}) - (\eta^{-(k-1)}R)^{1+\alpha} (\|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_{\eta^{-(k-1)}R})} + c_1|m_{k-1}|) \\ &\quad - (\eta^{-(k-1)}R) \|g\|_{L^\infty(\partial\Omega \cap B_{\eta^{-(k-1)}R})} \quad \text{in } \Omega \cap B_{\eta^{-k}R}. \end{aligned}$$

Moreover, by (4-3), the choice of M , (4-5), and the estimate $|m_{k-1}| \leq M_0 = (\delta/2c_1)M$, we estimate

$$\begin{aligned} &(\eta^{-(k-1)}R)^{1+\alpha} (\|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_{\eta^{-(k-1)}R})} + c_1|m_{k-1}|) + (\eta^{-(k-1)}R) \|g\|_{L^\infty(\partial\Omega \cap B_{\eta^{-(k-1)}R})} \\ &\leq \eta^{-\alpha(k-1)} \delta M = \delta(M_{k-1} - m_{k-1}). \end{aligned}$$

Therefore, we deduce

$$m_k := \delta(M_{k-1} - m_{k-1}) + m_{k-1} \leq u \leq M_{k-1} =: M_k \quad \text{in } \Omega \cap B_{\eta^{-k}R},$$

which proves (4-4) for k . Equation (4-5) for k follows from (4-3). The proof of Step 1 is complete.

Step 2. Now that we have established the claim of Step 1, let us show how to conclude the proof. Let us take $x, y \in B_{R/2}$. We define $k \in \mathbb{N}$ as

$$\inf\{k \in \mathbb{N} : |x - y| \geq \eta^{-k}(R/2)\}.$$

Then, $|x - y| \leq \eta^{-k+1}(R/2)$ and by Step 1, it holds that

$$\begin{aligned} \frac{|u(x) - u(y)|}{|x - y|^\alpha} &\leq \eta^{k\alpha}(R/2)^{-\alpha} \operatorname{osc}_{B_{\eta^{-k+1}(R/2)}} u \\ &\leq cR^{-\alpha} (\|u\|_{L^\infty(\Omega)} + \|v\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + R^{1+\alpha} \|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_R)} + R^{1+\alpha} + R \|g\|_{L^\infty(\partial\Omega \cap B_R)}). \end{aligned}$$

We can omit the additional summand $+R^{1+\alpha}$ by an additional scaling and normalization argument, i.e., by assuming that $R = 1$ and $\|u\|_{L^\infty(\Omega)} + \|v\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + \|d^{s-\alpha} f\|_{L^\infty(\Omega \cap B_1)} + \|g\|_{L^\infty(\partial\Omega \cap B_1)} = 1$, applying the previous estimate, and rescaling to general R . This concludes the proof after using that by Theorem 1.4 it holds that $b_\Omega/d^{s-1} \in C^\alpha(\Omega \cap B_{R/2}(x_0))$. \square

We are now in a position to deduce the boundary Hölder regularity estimate in $C^{1,\gamma}$ domains.

Proof of Theorem 1.6. Note that

$$\partial_\nu(v/b_\Omega) = \partial_\nu(v/d^{s-1}) - \partial_\nu(b_\Omega/d^{s-1})(v/d^{s-1}),$$

and recall that $|\partial_\nu(b_\Omega/d^{s-1})| \leq C$. Hence, we can apply Lemma 4.3 (with $R = \frac{1}{2}$ and varying $x_0 \in \partial\Omega$). Combining it with the interior regularity results from [Fernández-Real and Ros-Oton 2024a, Theorem 2.4.3], and a covering argument, we deduce the desired result. In order to produce the tail-term in the estimate, we employ a truncation argument in the same way as in the proof of Corollary 4.4. \square

We end this section with a boundary Hölder regularity estimate for solutions that are defined up to a polynomial and might grow fast at infinity.

Corollary 4.4. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $k \in \mathbb{N} \cup \{0\}$. Let $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{2,\gamma}$ for some $\gamma > 0$ and $K \in C^{5+2\gamma}(\mathbb{S}^{n-1})$. Let $f \in C(\Omega \cap B_4)$, $g \in C(\overline{\partial\Omega \cap B_4})$, and v with $v/d^{s-1} \in C(\overline{\Omega})$ be a viscosity solution to*

$$\begin{cases} Lv \stackrel{k}{=} f & \text{in } \Omega \cap B_4, \\ v = 0 & \text{in } B_4 \setminus \Omega, \\ \partial_\nu(v/d^{s-1}) = g & \text{on } \partial\Omega \cap B_4. \end{cases}$$

Then, there exists $\alpha_0 > 0$ such that if for some $\alpha \in (0, \alpha_0]$ we have $d^{s+1-\alpha} f \in L^\infty(\Omega \cap B_4)$, then the following holds true: If $k = 0$, and $v \in L^1_{2s}(\mathbb{R}^n)$, then $v/d^{s-1} \in C^\alpha_{\text{loc}}(\overline{\Omega \cap B_4})$, and

$$\left\| \frac{v}{d^{s-1}} \right\|_{C^\alpha(\overline{\Omega \cap B_4})} \leq c \left(\left\| \frac{v}{d^{s-1}} \right\|_{L^\infty(\Omega \cap B_4)} + \|v\|_{L^1_{2s}(\mathbb{R}^n \setminus B_4)} + \|d^{s+1-\alpha} f\|_{L^\infty(\Omega \cap B_4)} + \|g\|_{L^\infty(\partial\Omega \cap B_4)} \right).$$

If $k \in \mathbb{N}$, $K \in C^{k-1+\delta}(\mathbb{S}^{n-1})$, and $v \in L^1_{2s+k-1+\delta}(\mathbb{R}^n)$ for some $\delta > 0$, then $v/d^{s-1} \in C^\alpha_{\text{loc}}(\overline{\Omega \cap B_4})$, and

$$\left\| \frac{v}{d^{s-1}} \right\|_{C^\alpha(\overline{\Omega \cap B_4})} \leq c \left(\left\| \frac{v}{d^{s-1}} \right\|_{L^\infty(\Omega \cap B_4)} + \|v\|_{L^1_{2s+k-1+\delta}(\mathbb{R}^n \setminus B_4)} + \|d^{s+1-\alpha} f\|_{L^\infty(\Omega \cap B_4)} + \|g\|_{L^\infty(\partial\Omega \cap B_4)} \right),$$

where $c > 0$ and α_0 , depend only on $n, s, \lambda, \Lambda, \gamma, k, \delta$, and the $C^{2,\gamma}$ radius of Ω .

Proof. In case $k = 0$, the proof follows by a truncation argument. Indeed, let us define $w = v\mathbb{1}_{B_4}$ and observe that

$$Lw = f - L(v\mathbb{1}_{\mathbb{R}^n \setminus B_4}) =: \tilde{f} \quad \text{in } \Omega \cap B_2.$$

Moreover, we can estimate

$$\|d^{s+1-\alpha} \tilde{f}\|_{L^\infty(\Omega \cap B_2)} \leq c \|d^{s-1+\alpha} f\|_{L^\infty(\Omega \cap B_2)} + c \|v\|_{L^1_{2s}(\mathbb{R}^n \setminus B_4)}.$$

Thus, the desired result follows immediately by application of Theorem 1.6 to w , using that $v = 0$ in $B_4 \setminus \Omega$, by assumption.

Let now $k \in \mathbb{N}$. Again, we define $w = v\mathbb{1}_{B_4}$, but this time, since the equation only holds up to a polynomial, we obtain for any $R > 4$,

$$Lw = f_R - L(v\mathbb{1}_{B_R \setminus B_4}) + p_R =: \tilde{f} \quad \text{in } \Omega \cap B_3,$$

where $f_R \rightarrow f$ in $d^{s+1-\alpha} L^\infty(\Omega \cap B_3)$, as $R \rightarrow \infty$, and $p_R \in \mathcal{P}_{k-1}$. As in the proof of Lemma 2.10 (see also [Abatangelo and Ros-Oton 2020, Lemma 3.6; Kukuljan 2021, Lemma 4.8]), taking difference quotients of order $k-1+\delta$ of the equation for w , and using crucially that $K \in C^{k-1+\delta}(\mathbb{S}^{n-1})$, we can find a polynomial $p \in \mathcal{P}_{\lfloor k-1+\delta \rfloor}$ and h with $d^{s+1-\alpha} h \in L^\infty(\Omega \cap B_3)$ such that

$$\begin{cases} Lw = h + p & \text{in } \Omega \cap B_3, \\ w = 0 & \text{in } \mathbb{R}^n \setminus (\Omega \cap B_4), \end{cases}$$

and moreover, h satisfies the estimate

$$\|d^{s+1-\alpha} h\|_{L^\infty(\Omega \cap B_3)} \leq C \left(\|d^{s+1-\alpha} f\|_{L^\infty(\Omega \cap B_3)} + \|v\|_{L^1(\mathbb{R}^n \setminus B_4)}^{-n-2s-(k-1+\delta)} \right). \quad (4-7)$$

Next, let us take a bounded domain $D \subset \mathbb{R}^n$ with $\partial D \in C^{1,\gamma}$ such that $\Omega \cap B_2 \subset D \subset \Omega \cap B_3$. Moreover, we find w_1, w_2 such that $w_1/d_D^{s-1}, w_1/d_D^{s-1} \in C(\bar{D})$ and $w = w_1 + w_2$ satisfying

$$\begin{cases} Lw_1 = h & \text{in } D, \\ w_1 = w & \text{in } \mathbb{R}^n \setminus D, \\ w_1/d_D^{s-1} = v/d_D^{s-1} & \text{on } \partial D, \end{cases} \quad \text{and} \quad \begin{cases} Lw_2 = p & \text{in } D, \\ w_2 = 0 & \text{in } \mathbb{R}^n \setminus D, \\ w_2/d_D^{s-1} = 0 & \text{on } \partial D. \end{cases}$$

The existence of $w_2 \in L^\infty(\mathbb{R}^n)$ follows from [Fernández-Real and Ros-Oton 2024a, Theorem 3.2.27], and we obtain $w_2/d_D^s \in C^\gamma(\bar{D})$ from [Fernández-Real and Ros-Oton 2024a, Theorem 2.7.1], which yields $w_2/d_D^{s-1} = d_D(w_2/d_D^s) \in C^\gamma(\bar{D})$ since $\partial D \in C^{1,\gamma}$. Then, we can define $w_1 := w - w_2$. We claim that

$$\|w_1/d_D^{s-1}\|_{L^\infty(D)} \leq c(\|v/d_\Omega^{s-1}\|_{L^\infty(\Omega \cap B_4)} + \|d_\Omega^{s+1-\alpha}h\|_{L^\infty(\Omega \cap B_3)}). \quad (4-8)$$

To see this, let us recall the function ψ_1 (with respect to D) from Lemma 2.7, and observe that by Lemma 2.3, we can take it in such a way that

$$L(\psi_1 + d_\Omega^{s-1}) \geq c_0 d_D^{\alpha-s-1} \quad \text{in } D \quad (4-9)$$

for some $c_0 > 0$. Moreover, recall $\psi_1/d_D^{s-1} = 1$ on ∂D . Then, let us define

$$\begin{aligned} \Psi(x) = c_1 \psi_1(x) (\|v/d_D^{s-1}\|_{L^\infty(\partial D)} + \|d_D^{s+1-\alpha}h\|_{L^\infty(D)} + \|w/d_\Omega^{s-1}\|_{L^\infty(\Omega \cap B_4)}) \\ + c_1 d_\Omega^{s-1}(x) \|w/d_\Omega^{s-1}\|_{L^\infty(\Omega \cap B_4)}, \end{aligned}$$

where $c_1 := \max\{c_0^{-1}, 1\}$, and observe that by (4-9) we have

$$\begin{cases} Lw_1 \leq L\Psi & \text{in } D, \\ w_1 \leq \Psi & \text{in } \mathbb{R}^n \setminus D, \\ w_1/d_D^{s-1} \leq \Psi/d_D^{s-1} & \text{on } \partial D, \end{cases}$$

which, recalling that $\psi_1 \leq c_1 d_D^{s-1}$ in D , and $d_D \leq d_\Omega$, as well as the definition of D , imply that

$$\begin{aligned} \frac{w_1}{d_D^{s-1}} &\leq c_2 c_1 (\|v/d_D^{s-1}\|_{L^\infty(\partial D)} + \|d_D^{s+1-\alpha}h\|_{L^\infty(D)} + \|w/d_\Omega^{s-1}\|_{L^\infty(\Omega \cap B_4)}) + c_2 c_1 \frac{d_\Omega^{s-1}}{d_D^{s-1}} \|w/d_\Omega^{s-1}\|_{L^\infty(\Omega \cap B_4)} \\ &\leq c(\|v/d_\Omega^{s-1}\|_{L^\infty(\partial \Omega \cap B_3)} + \|d_\Omega^{s+1-\alpha}h\|_{L^\infty(\Omega \cap B_3)} + \|v/d_\Omega^{s-1}\|_{L^\infty(\Omega \cap B_4)}), \end{aligned}$$

which yields our claim in (4-8). As a direct consequence of (4-8), we deduce

$$\begin{aligned} \|w_2/d_D^{s-1}\|_{L^\infty(D)} &\leq \|w/d_D^{s-1}\|_{L^\infty(D)} + \|w_1/d_D^{s-1}\|_{L^\infty(D)} \\ &\leq c(\|v/d_\Omega^{s-1}\|_{L^\infty(\Omega \cap B_4)} + \|d_\Omega^{s+1-\alpha}h\|_{L^\infty(\Omega \cap B_3)}). \end{aligned} \quad (4-10)$$

Finally, we claim that

$$\|p\|_{L^\infty(D)} \leq c \|w_2/d_D^{s-1}\|_{L^\infty(D)}. \quad (4-11)$$

Once we show (4-11), then the proof is complete after combination of (4-11), (4-10), (4-7), and application of the boundary Hölder regularity estimate (see Theorem 1.6) to w in Ω , as in the case $k = 0$. We prove

(4-11) by contradiction. Suppose there are sequences $(L_j)_j, (w_j)_j, (p_j)_j$ with

$$\|p_j\|_{L^\infty(D)} = 1, \quad \text{and} \quad \begin{cases} L_j w_j = p_j & \text{in } D, \\ w_j = 0 & \text{in } \mathbb{R}^n \setminus D, \\ w_j/d_\Omega^{s-1} = 0 & \text{on } \partial D, \\ \lim_{j \rightarrow \infty} \|w_j/d_D^{s-1}\|_{L^\infty(D)} = 0. \end{cases}$$

Then, up to subsequences, it holds that $L_{j_m} \rightharpoonup L, w_{j_m}/d_D^{s-1} \rightarrow u_0$ in $L^\infty(D)$ for some $u_0 \in L^\infty(D), p_{j_m} \rightarrow p_0$ in $L^\infty(D)$. While the first convergence statement follows from [Abatangelo and Ros-Oton 2020, Lemma 3.7], the second convergence statement follows from Theorem 1.6 and the Arzelà–Ascoli theorem, and the third one is immediate from the boundedness of (p_{j_m}) in a finite dimensional space.

We can now make use of the stability result in [Fernández-Real and Ros-Oton 2024a, Proposition 2.2.36], and deduce that for $w_0 = d_D^{s-1}u_0$, it holds that

$$\|p_0\|_{L^\infty(D)} = 1, \quad \text{and} \quad \begin{cases} Lw_0 = p_0 & \text{in } D, \\ w_0 = 0 & \text{in } \mathbb{R}^n \setminus D, \\ w_0/d_D^{s-1} = 0 & \text{on } \partial D, \\ \|w_0/d_D^{s-1}\|_{L^\infty(D)} = 0. \end{cases}$$

Clearly, $w_0 = 0$ is not a solution to $Lw_0 = p_0$ in D , so we have obtained a contradiction, and conclude the proof of (4-11). □

5. Liouville theorem in the half-space

The proof of our main result (see Theorem 1.2) is based on a blow-up argument. A crucial ingredient in such proof is a suitable Liouville theorem in the half-space. In this section, we will establish such a result for nonlocal problems with local Neumann boundary conditions:

Theorem 5.1. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $k \in \mathbb{N}, \gamma \in (0, 1)$ with $\gamma \neq s$, and $K \in C^{k-1+\gamma-s+\delta}(\mathbb{S}^{n-1})$ for some $\delta > 0$. Let $u \in C(\mathbb{R}^n)$ be a viscosity solution to*

$$\begin{cases} L((x_n)_+^{s-1}u)^{k-1+\lceil\gamma-s\rceil} = 0 & \text{in } \{x_n > 0\}, \\ \partial_n u = p & \text{in } \{x_n = 0\}, \\ |u(x)| \leq C(1 + |x|)^{k+\gamma} & \text{for all } x \in \{x_n > 0\}, \end{cases}$$

for some $C > 0, p \in \mathcal{P}_{k-1}$. Then, there exist $a_\beta \in \mathbb{R}$ for any $\beta \in (\mathbb{N} \cup \{0\})^n$ with $|\beta| \leq k$ such that

$$u(x) = \sum_{|\beta| \leq k} a_\beta x_1^{\beta_1} \cdots x_n^{\beta_n} \quad \text{for all } x \in \{x_n > 0\}.$$

In order to prove Theorem 5.1, we first establish the following one-dimensional version, which can be proved by combination of the arguments in [Ros-Oton and Serra 2016a, Lemma 6.2; Abatangelo and Ros-Oton 2020, Lemma 3.3].

Lemma 5.2. *Let $k \in \mathbb{N}$, $\gamma \in (0, 1)$ with $\gamma \neq s$, and $u \in C(\mathbb{R})$ satisfying*

$$\begin{cases} (-\Delta)^s ((x_+)^{s-1} u)^{k-1+\lceil\gamma-s\rceil} \equiv 0 & \text{in } (0, \infty), \\ |u(x)| \leq C(1 + |x|)^{k+\gamma} & \text{for all } x > 0, \end{cases}$$

for some $C > 0$. Then, there exist $a_0, a_1, \dots, a_k \in \mathbb{R}$ such that

$$u(x) = \sum_{j=0}^k a_j x^j \quad \text{for all } x > 0.$$

Proof. In case $k = 1$ and $\gamma < s$, the proof is an application of [Ros-Oton and Serra 2016a, Lemma 6.2] with $u(x) := (x_+)^{s-1} u(x)$, and $\delta = s > 0$, $\beta = s + \gamma \in (0, 2s)$.

In case $k > 1$ or $\gamma > s$, we have $k - 1 + \lceil\gamma - s\rceil \geq 1$. Let us define $v(x) = (x_+)^{s-1} u(x)$, let $V : \mathbb{R} \times [0, \infty) \rightarrow \mathbb{R}$ be the harmonic extension of v in the sense of [Abatangelo and Ros-Oton 2020, Lemma 3.3], and finally define $\tilde{V}(x, y) = \int_{-\infty}^x V(z, y) dz$. Note that \tilde{V} satisfies (see [Abatangelo and Ros-Oton 2020, Lemma 3.3])

$$\begin{cases} \operatorname{div}(y^{1-2s} \nabla \tilde{V}(x, y)) = 0 & \text{in } \mathbb{R} \times (0, \infty), \\ \tilde{V}(x, y) = v(x) & \text{on } \mathbb{R} \times \{0\}, \\ |\tilde{V}(x, y)| \leq C(1 + |x|^{2(k-1+\lceil\gamma-s\rceil)+1+\gamma+s} + |y|^{(k-1+\lceil\gamma-s\rceil)+1+\gamma+s}) & \text{in } \mathbb{R} \times (0, \infty). \end{cases}$$

Next, by [Ros-Oton and Serra 2016a, Lemma 6.2] (see also [Fernández-Real and Ros-Oton 2024a, Theorem 1.10.16]), we have the representation formula

$$\tilde{V}(x, y) = \tilde{V}(r \cos \theta, r \sin \theta) = \sum_{j=0}^{\infty} a_j \Theta_j(\theta) r^{j+s}, \quad \text{for all } x \in \mathbb{R}, y \in [0, \infty),$$

where $a_j \in \mathbb{R}$, and $(\Theta_j)_j$ is a complete orthogonal system in the subspace of even functions in $L^2((0, \pi), (\sin \theta)^{1-2s} d\theta)$. By the Parseval identity, the bounds on $|\tilde{V}|$ imply

$$\sum_{j=0}^{\infty} a_j^2 R^{2+2j} = \int_{\partial B_R \cap \{y>0\}} \tilde{V}(x, y)^2 y^{1-2s} d\sigma \leq C R^{4(k-1+\lceil\gamma-s\rceil)+2+2\gamma+2} = C R^{4(k+\lceil\gamma-s\rceil)+2\gamma}.$$

Therefore, it must be $a_j = 0$ for any $j > j_0$, where $j_0 = \min\{j \in \mathbb{N} : 2 + 2j > 4(k + \lceil\gamma - s\rceil) + 2\gamma\}$, which implies

$$\tilde{V}(x, y) = \sum_{j=0}^{j_0} a_j \Theta_j(\theta) r^{j+s} \quad \text{for all } x \in \mathbb{R}, y \in [0, \infty).$$

Upon recalling the definition of V and \tilde{V} , this implies

$$v(x) = (x_+)^s \sum_{j=0}^{j_0-1} b_j x^j$$

for some $b_j \in \mathbb{R}$, and since $|v(x)| \leq C(1 + |x|)^{k+\gamma-1+s}$ and $\gamma \in (0, 1)$ by assumption, it must be $b_j = 0$ for any $j \geq k$. Recalling that by definition $u(x) = v(x)(x_+)^{1-s}$, we deduce that u must be a polynomial of degree at most k in $\{x > 0\}$, as desired. □

Moreover, we will need the following lemma (see also [Kukuljan 2021, Proposition 4.3]):

Lemma 5.3. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $k \in \mathbb{N}$, $K \in C^{k-2+\delta}(\mathbb{S}^{n-1})$ for some $\delta > 0$. Let $f \in \mathcal{P}_k$. Then,*

$$L((x_n)_+^{s-1} f) \stackrel{k-1}{=} 0 \quad \text{in } \{x_n > 0\}.$$

Proof. Let us first give a simple proof in case $k = 1$. Then, it suffices to prove that for any $i \in \{1, \dots, n\}$

$$L((x_n)_+^{s-1} x_i) = 0 \quad \text{in } \{x_n > 0\}.$$

First, by integrating $x \mapsto (x_n)_+^{s-1}$ in x_i , and using that $L((x_n)_+^{s-1}) = 0$, we deduce

$$L((x_n)_+^{s-1} x_i) \equiv c \quad \text{in } \{x_n > 0\}$$

for some constant $c \in \mathbb{R}$. Then, since $x \mapsto (x_n)_+^{s-1} x_i$ is homogeneous of degree s , we deduce that for any $\lambda > 0$ and $x \in \{x_n > 0\}$,

$$c = L((x_n)_+^{s-1} x_i)(\lambda x) = \lambda^{-2s} L((\lambda x_n)_+^{s-1} \lambda x_i)(x) = \lambda^{-s} L((x_n)_+^{s-1} x_i)(x) = \lambda^{-s} c.$$

This implies that $c = 0$, as desired.

For $k \geq 2$, we prove the result by induction. Assume that we know already

$$L((x_n)_+^{s-1} p) \stackrel{k-2}{=} 0 \quad \text{in } \{x_n > 0\} \tag{5-1}$$

for every $p \in \mathcal{P}_{k-1}$. Now, let $q \in \mathcal{P}_k$. Then, by integrating (5-1) with $p := \partial_i q$ for $i \in \{1, \dots, n\}$, by Lemma 2.12 we find that there exists a constant $c \in \mathbb{R}$ such that

$$L((x_n)_+^{s-1} q) \stackrel{k-1}{=} c \quad \text{in } \{x_n > 0\}.$$

Since $c \stackrel{k-1}{=} 0$ for any $k \geq 2$, we conclude the proof. □

Finally, we state a Hölder regularity estimate in the half-space, which follows from Corollary 4.4.

Corollary 5.4. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $k \in \mathbb{N} \cup \{0\}$, $\gamma > 0$, and $K \in C^{k-1+\delta}(\mathbb{S}^{n-1})$ for some $\delta > 0$. Let $f \in C(\{x_n > 0\} \cap B_2)$, $g \in C(\overline{\{x_n = 0\}} \cap B_2)$, and $u \in C(\{x_n \geq 0\})$ be a viscosity solution to*

$$\begin{cases} L((x_n)_+^{s-1} u) \stackrel{k}{=} f & \text{in } \{x_n > 0\} \cap B_2, \\ \partial_n u = g & \text{on } \{x_n = 0\} \cap B_2. \end{cases}$$

Then, there exists $\alpha_0 > 0$ such that if $(x_n)_+^{s+1-\alpha} f \in L^\infty(\{x_n > 0\} \cap B_2)$ for some $\alpha \in (0, \alpha_0]$, then the following holds true: if $k = 0$ and $(x_n)_+^{s-1} u \in L_{2s}^1(\mathbb{R}^n)$ it holds that $u \in C_{\text{loc}}^\alpha(\{x_n \geq 0\} \cap B_2)$, and

$$\begin{aligned} & \|u\|_{C^\alpha(\{x_n \geq 0\} \cap B_1)} \\ & \leq c \left(\|u\|_{L^\infty(\{x_n > 0\} \cap B_4)} + \|(x_n)_+^{s-1} u\|_{L_{2s}^1(\mathbb{R}^n \setminus B_4)} + \|(x_n)_+^{s+1-\alpha} f\|_{L^\infty(\{x_n > 0\} \cap B_2)} + \|g\|_{L^\infty(\{x_n = 0\} \cap B_2)} \right), \end{aligned}$$

and if $k \in \mathbb{N}$ and $(x_n)_+^{s-1} u \in L_{2s+(k-1+\delta)}^1(\mathbb{R}^n)$ it holds that $u \in C_{\text{loc}}^\alpha(\{x_n \geq 0\} \cap B_2)$, and

$$\begin{aligned} \|u\|_{C^\alpha(\{x_n \geq 0\} \cap B_1)} & \leq c \left(\|u\|_{L^\infty(\{x_n > 0\} \cap B_4)} + \|[(x_n)_+^{s-1} u] \cdot |^{-n-2s-(k-1+\delta)} \|_{L^1(\{x_n > 0\} \setminus B_4)} \right. \\ & \quad \left. + \|(x_n)_+^{s+1-\alpha} f\|_{L^\infty(\{x_n > 0\} \cap B_2)} + \|g\|_{L^\infty(\{x_n = 0\} \cap B_2)} \right), \end{aligned}$$

where $c > 0$ and α_0 depend only on $n, s, \lambda, \Lambda, \gamma, k, \delta$.

Proof. The result follows directly from Corollary 4.4 applied to some domain $\Omega \subset \mathbb{R}^n$ with $\partial\Omega \in C^{1,\gamma}$, which satisfies $\{x_n > 0\} \cap B_2 \subset \Omega \subset \{x_n > 0\} \cap B_4$. \square

With the help of the one-dimensional Liouville theorem in the half-space and the Hölder regularity estimate up to the boundary (see Corollary 5.4), the proof of Theorem 5.1 follows by a standard procedure, which is explained in detail for instance in [Abatangelo and Ros-Oton 2020, proof of Theorem 3.10].

Proof of Theorem 5.1. First, we observe that by scaling Corollary 5.4, we obtain that for any $R \geq 1$,

$$\begin{aligned}
 [u]_{C^\alpha(B_R)} &\leq cR^{-\alpha} \left[\|u\|_{L^\infty(B_{4R})} + R^{1+s} \|(x_n)_+^{s-1} u\| \cdot |^{-n-2s}\|_{L^1(\mathbb{R}^n \setminus B_{4R})} \mathbb{1}_{\{k=1 \text{ and } \gamma < s\}} \right. \\
 &\quad \left. + R^{s+k-1+\lceil\gamma-s\rceil+\eta} \|[(x_n)_+^{s-1} u]\| \cdot |^{-n-2s-(k-2+\lceil\gamma-s\rceil+\eta)}\|_{L^1(\{x_n>0\} \setminus B_{4R})} \mathbb{1}_{\{k \geq 2 \text{ or } \gamma > s\}} \right. \\
 &\quad \left. + R \|p\|_{L^\infty(\{x_n=0\} \cap B_{4R})} \right] \\
 &\leq cR^{k+\gamma-\alpha}, \tag{5-2}
 \end{aligned}$$

where we take $\eta = 1 + \gamma - s - \lceil\gamma - s\rceil + \delta$ and used in the last estimate the growth condition on u , the fact that $\|p\|_{L^\infty(\{x_n=0\} \cap B_R)} \leq cR^{k-1}$, and the following computation using polar coordinates with $y_n = r \cos \theta$ for some $\theta \in [0, 2\pi)$ (similar to the proof of Lemma 4.2), which is slightly different in case ($k = 1$ and $\gamma < s$) and ($k \geq 2$ or $\gamma > s$). In case $k = 1$ and $\gamma < s$, we obtain

$$\begin{aligned}
 R^{1+s} \|(x_n)_+^{s-1} u\| \cdot |^{-n-2s}\|_{L^1(\mathbb{R}^n \setminus B_{4R})} &\leq cR^{1+s} \int_{\mathbb{R}^n \setminus B_{4R}} (y_n)_+^{s-1} |y|^{-n-2s+1+\gamma} dy \\
 &\leq cR^{1+s} \int_0^{2\pi} \cos(\theta)_+^{s-1} \left(\int_{4R}^\infty r^{s-1} r^{-1-2s+1+\gamma} dr \right) d\theta \\
 &\leq cR^{1+s} R^{\gamma-s} \left(\int_0^{2\pi} \cos(\theta)_+^{s-1} d\theta \right) \leq cR^{1+\gamma}. \tag{5-3}
 \end{aligned}$$

In case ($k \geq 2$ or $\gamma > s$), we obtain, using that $\eta > 1 + \gamma - s - \lceil\gamma - s\rceil$,

$$\begin{aligned}
 R^{s+k-1+\lceil\gamma-s\rceil+\eta} \|[(x_n)_+^{s-1} u]\| \cdot |^{-n-2s-(k-2+\lceil\gamma-s\rceil+\eta)}\|_{L^1(\{x_n>0\} \setminus B_4)} \\
 &\leq cR^{s+k-1+\lceil\gamma-s\rceil+\eta} \int_{\mathbb{R}^n \setminus B_{4R}} (y_n)_+^{s-1} |y|^{-n-2s-(k-2+\lceil\gamma-s\rceil+\eta)+k+\gamma} dy \\
 &\leq cR^{s+k-1+\lceil\gamma-s\rceil+\eta} \int_0^{2\pi} \cos(\theta)_+^{s-1} \left(\int_{4R}^\infty r^{s-1} r^{-1-2s+2-\lceil\gamma-s\rceil+\gamma-\eta} dr \right) d\theta \\
 &\leq cR^{k+s} R^{\gamma-s} \left(\int_0^{2\pi} \cos(\theta)_+^{s-1} d\theta \right) \leq cR^{k+\gamma}.
 \end{aligned}$$

Next, let us take any $\tau \in \mathbb{S}^{n-1}$ such that $\tau_n = 0$ and $0 < h < R/2$. We consider the difference quotients

$$w_{1,\tau}(x) = \frac{u(x+h\tau) - u(x)}{h^\alpha}, \quad p_{1,\tau}(x) = \frac{p(x+h\tau) - p(x)}{h^\alpha}$$

and deduce from (5-2) (after applying the estimate to smaller balls of radius comparable to R inside B_R) that

$$\|w_{1,\tau}\|_{L^\infty(B_R)} \leq cR^{k+\gamma-\alpha} \quad \text{for all } R \geq 1.$$

Clearly, since $\tau_n = 0$, $w_{1,\tau}$ satisfies in the viscosity sense

$$\begin{cases} L((x_n)_+^{s-1} w_{1,\tau}) \stackrel{k-1+[\gamma-s]}{=} 0 & \text{in } \{x_n > 0\}, \\ \partial_n w_{1,\tau} = p_{1,\tau} & \text{on } \{x_n = 0\}. \end{cases} \tag{5-4}$$

Here, we are using that sums of viscosity solutions are again viscosity solutions by Lemma 2.14. Using (5-4) and also that $|p_{1,\tau}(x)| \leq c|x|^{k-1-\alpha}$ since $p \in \mathcal{P}_{k-1}$, we can apply the previous arguments to $w_{1,\tau}$. Eventually, this implies that $w_{2,\tau}(x) = (w_{1,\tau}(x+h\tau) - w_{1,\tau}(x))/h^\alpha$ satisfies $\|w_{2,\tau}\|_{L^\infty(B_R)} \leq cR^{k+\gamma-2\alpha}$. This way, we obtain higher order difference quotients $w_{j,\tau}$, $j \in \mathbb{N}$, and they satisfy

$$\begin{cases} L((x_n)_+^{s-1} w_{j,\tau}) \stackrel{k-1+[\gamma-s]}{=} 0 & \text{in } \{x_n > 0\}, \\ \partial_n w_{j,\tau} = p_{j,\tau} & \text{on } \{x_n = 0\}, \\ \|w_{j,\tau}\|_{L^\infty(B_R)} \leq cR^{k+\gamma-j\alpha} & \text{for all } R \geq 1, \\ \|p_{j,\tau}\|_{L^\infty(B_R)} \leq cR^{k-1-j\alpha} & \text{for all } R \geq 1. \end{cases}$$

Then, taking $j_0 \in \mathbb{N}$ as the smallest number such that $j_0\alpha > k + \gamma$, and upon taking the limit $R \rightarrow \infty$, we deduce that

$$\lim_{R \rightarrow \infty} \|w_{j_0,\tau}\|_{L^\infty(B_R)} \leq c \lim_{R \rightarrow \infty} R^{k+\gamma-j_0\alpha} = 0,$$

i.e., $w_{j_0,\tau} \equiv 0$ in \mathbb{R}^n . Thus, $w_{j_0-1,\tau}$ is a function that is constant in the τ -direction. Clearly, we can also take difference quotients of $w_{j_0-1,\tau}$ in other directions $\tau' \in \mathbb{S}^{n-1}$ with $\tau'_n = 0$, and the same arguments as before apply. Therefore, $w_{j_0-1,\tau}(x) = w_{j_0-1,\tau}(x_n)$ is one-dimensional for any $\tau \in \mathbb{S}^{n-1}$ with $\tau_n = 0$.

Unraveling the higher order difference quotients, we get that $w_{j_0-2,\tau}(x) = (V_1(x_n), x') + V_2(x_n)$ for some one-dimensional functions $V_1 : \mathbb{R} \rightarrow \mathbb{R}^{n-1}$ and $V_2 : \mathbb{R} \rightarrow \mathbb{R}$, and continuing this argument $j_0 - 1$ times, we deduce that u must be a polynomial in x' with coefficients that are one-dimensional functions from $\mathbb{R} \rightarrow \mathbb{R}$ in x_n .

Then, by the growth condition on u , for any multi-index $\beta \in (\mathbb{N} \cup \{0\})^{n-1}$ with $|\beta| \leq k$, we obtain functions A_β in x_n such that

$$u(x) = \sum_{|\beta| \leq k} (x')^\beta A_\beta(x_n).$$

In particular, this implies that $D_{x'}^\beta u(x) = c(\beta)A_\beta(x_n)$ for any $|\beta| = k$ and some constant $c(\beta) > 0$, where $D_{x'}^\beta$ denotes an incremental quotient approximating the partial derivative $\partial_{x'}^\beta$ in the x' -variables. Therefore, discretely differentiating the equation for u , we deduce

$$c(\beta)L((x_n)_+^{s-1} A_\beta)(x) = L((x_n)_+^{s-1} D_{x'}^\beta u)(x) = L(D_{x'}^\beta [(x_n)_+^{s-1} u])(x) \stackrel{k-1+[\gamma-s]}{=} 0 \quad \text{in } \{x_n > 0\}.$$

By the growth condition on u , it must be $|A_\beta(x_n)| \leq c(1+|x_n|)^{k-|\beta|+\gamma} = c(1+|x|)^\gamma$, and since A_β was also one-dimensional, i.e., $LA_\beta = (-\Delta)_{\mathbb{R}}^s A_\beta$, we can apply Lemma 5.2 to A_β , which yields $A_\beta(x_n) = p_\beta(x_n)$ for some polynomial $p_\beta \in \mathcal{P}_{k-|\beta|} = \mathcal{P}_0$. Next, we recall from Lemma 5.3,

$$L((x_n)_+^{s-1} (x')^\beta p_\beta(x_n)) \stackrel{k-1+[\gamma-s]}{=} 0 \quad \text{in } \{x_n > 0\}.$$

Thus, repeating the arguments from above, we deduce that, for every β with $|\beta| \leq k$ it holds that

$$\begin{cases} L((x_n)_+^{s-1} A_\beta)^{k-1+\lceil\gamma-s\rceil} 0 & \text{in } \{x_n > 0\}, \\ |A_\beta(x)| \leq C(1 + |x|)^{k-|\beta|+\gamma} & \text{for all } x \in \{x_n > 0\}, \end{cases}$$

and hence $A_\beta(x_n) = p_\beta(x_n)$ for some polynomial $p_\beta \in \mathcal{P}_{k-|\beta|}$. This implies $u(x) = p(x)$ for some polynomial p , and by the growth condition on u , it must be $p \in \mathcal{P}_k$, as desired. \square

6. Higher order boundary regularity

The goal of this section is to prove the desired higher order boundary regularity for nonlocal equations with local Neumann conditions (see Theorem 1.2). The proof goes by a blow-up argument and heavily uses the Liouville theorem in the half-space (see Theorem 5.1), as well as the boundary Hölder estimate (see Theorem 1.6).

Lemma 6.1. *Let $L \in \mathcal{L}_s^{\text{hom}}(\lambda, \Lambda)$. Let $k \in \mathbb{N}$ and $\Omega \subset \mathbb{R}^n$ be an open, bounded domain with $\partial\Omega \in C^{k+1,\gamma}$ for some $\gamma \in (0, 1)$ with $\gamma \neq s$, and $K \in C^{2k+2\gamma+3}(\mathbb{S}^{n-1})$. Let $v \in L^1_{2s}(\mathbb{R}^n)$ with $v/d^{s-1} \in C(\bar{\Omega})$ be a viscosity solution to*

$$\begin{cases} Lv = f & \text{in } \Omega \cap B_1, \\ v = 0 & \text{in } B_1 \setminus \Omega, \\ \partial_{v_n}(v/d^{s-1}) = g & \text{on } \partial\Omega \cap B_1. \end{cases}$$

(i) *If $k = 1$ and $\gamma < s$, $f \in C(\Omega \cap B_1)$ with $d^{s-\gamma} f \in L^\infty(\Omega \cap B_1)$, $g \in C^\gamma(\partial\Omega \cap B_1)$, then for any $x_0 \in \partial\Omega \cap B_{1/2}$ and $x \in \Omega \cap B_{1/2}$ it holds that*

$$\begin{aligned} & \left| \frac{v}{d^{s-1}}(x) - \left(\frac{v}{d^{s-1}}(x_0) - A(x_0) \cdot (x - x_0) \right) \right| \\ & \leq c \left(\left\| \frac{v}{d^{s-1}} \right\|_{L^\infty(\Omega)} + \|v\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + \|d^{s-\gamma} f\|_{L^\infty(\Omega \cap B_1)} + \|g\|_{C^\gamma(\partial\Omega \cap B_1)} \right) |x - x_0|^{1+\gamma} \end{aligned}$$

for some $c > 0$, which only depends on $n, s, \lambda, \Lambda, \gamma$, and the $C^{2,\gamma}$ radius of Ω . If in addition, $g \equiv 0$, then $A(x_0) \cdot v_{x_0} = 0$.

(ii) *If $k \geq 2$ or $\gamma > s$, $f \in C^{(k-1)-s+\gamma}(\Omega \cap B_1)$, $g \in C^{k-1+\gamma}(\partial\Omega \cap B_1)$, then for any $x_0 \in \partial\Omega \cap B_{1/2}$, there is $Q(\cdot; x_0) \in \mathcal{P}_k$ such that for any $x \in \Omega \cap B_{1/2}$ it holds that*

$$\begin{aligned} & \left| \frac{v}{d^{s-1}}(x) - Q(x; x_0) \right| \\ & \leq c \left(\left\| \frac{v}{d^{s-1}} \right\|_{L^\infty(\Omega)} + \|v\|_{L^\infty(\mathbb{R}^n \setminus \Omega)} + \|f\|_{C^{(k-1)-s+\gamma}(\Omega \cap B_1)} + \|g\|_{C^{k-1+\gamma}(\partial\Omega \cap B_1)} \right) |x - x_0|^{k+\gamma} \end{aligned}$$

for some $c > 0$, which only depends on $n, s, \lambda, \Lambda, \gamma, k$, and the $C^{k+1,\gamma}$ radius of Ω .

Proof. Let us assume without loss of generality that $x_0 = 0 \in \partial\Omega$ with $\partial_{v_0} = e_n$. We set $u := v/d^{s-1}$.

We will prove the desired result by a blow-up argument. To do so, we assume by contradiction that for any $j \in \mathbb{N}$ there exist $C^{k+1,\gamma}$ domains $\Omega_j \subset \mathbb{R}^n$, $f_j \in C^{k-1}(\Omega_j \cap B_1)$, $g_j \in C^{k-1+\gamma}(\partial\Omega_j \cap B_1)$, $r_j > 0$,

operators L_j with ellipticity constants λ, Λ , and $v_j \in C(\Omega_j) \cap L^1_{2s}(\mathbb{R}^n)$ viscosity solutions to

$$\begin{cases} L_j v_j = f_j & \text{in } \Omega_j \cap B_1, \\ v_j = 0 & \text{in } B_1 \setminus \Omega_j, \\ \partial_\nu(v_j/d_j^{s-1}) = g_j & \text{on } \partial\Omega_j \cap B_1, \end{cases}$$

such that

$$\begin{aligned} |\text{diam } \Omega_j| + \|u_j\|_{L^\infty(\Omega_j)} + \|v_j\|_{L^\infty(\mathbb{R}^n \setminus \Omega_j)} + \mathbb{1}_{\{k=1 \text{ and } \gamma < s\}} \|d_j^{s-\gamma} f_j\|_{L^\infty(\Omega_j \cap B_1)} \\ + \mathbb{1}_{\{k \geq 2 \text{ or } \gamma > s\}} \|f_j\|_{C^{(k-1)-s+\gamma}(\Omega_j \cap B_1)} + \|g_j\|_{C^{k-1+\gamma}(\Omega_j \cap B_1)} + \|d_j\|_{C^{k+1,\gamma}(\Omega_j \cap B_1)} \leq C \end{aligned}$$

for some $C > 0$, denoted $d_{\Omega_j} = d_j$, and used that $d_j \in C^{k+1,\gamma}$ by [Fernández-Real and Ros-Oton 2024a, Definition 2.7.5]. Finally, we assume by contradiction

$$\sup_{j \in \mathbb{N}} \sup_{r > 0} r^{-k-\gamma} \|u_j - Q\|_{L^\infty(\Omega_j \cap B_r)} = \infty \quad \text{for all } Q \in \mathcal{P}_k.$$

Observe that up to a rotation, $r_m^{-1}\Omega_{j_m} \cap B_{r_m^{-1}} \rightarrow \{x_n > 0\}$. Moreover, we will write

$$\tilde{d}_{j_m} \mathbb{1}_{r_m^{-1}\Omega_{j_m}} := \tilde{d}_{j_m} =: r_m^{-1}d_{j_m}(r_m \cdot)$$

for the (regularized) distance with respect to $r_m^{-1}\Omega_{j_m}$.

We consider the $L^2(\Omega_j \cap B_r)$ -projections of u_j over \mathcal{P}_k , and denote them by $Q_{j,r} \in \mathcal{P}_k$. They satisfy the following properties:

$$\begin{aligned} \|u_j - Q_{j,r}\|_{L^2(\Omega_j \cap B_r)} &\leq \|u_j - Q\|_{L^2(\Omega_j \cap B_r)} \quad \text{for all } Q \in \mathcal{P}_k, \\ \int_{\Omega_j \cap B_r} (u_j(x) - Q_{j,r}(x))Q(x) \, dx &= 0 \quad \text{for all } Q \in \mathcal{P}_k. \end{aligned}$$

Next, we introduce

$$\theta(r) := \sup_{j \in \mathbb{N}} \sup_{\rho \geq r} \rho^{-k-\gamma} \|u_j - Q_{j,\rho}\|_{L^\infty(\Omega_j \cap B_\rho)}. \tag{6-1}$$

Observe that $\theta(r) \nearrow \infty$, as $r \searrow 0$. This follows from [Abatangelo and Ros-Oton 2020, Lemma 4.3] applied with $s = 0$ (the proof remains exactly the same in this case).

As a consequence, there exist sequences $(r_m)_m$ and $(j_m)_m$ such that

$$\frac{\|u_{j_m} - Q_{j_m,r_m}\|_{L^\infty(\Omega_{j_m} \cap B_{r_m})}}{r_m^{k+\gamma} \theta(r_m)} \geq \frac{1}{2} \quad \text{for all } m \in \mathbb{N}. \tag{6-2}$$

Let us define for any $m \in \mathbb{N}$,

$$w_m(x) = \frac{u_{j_m}(r_m x) - Q_{j_m,r_m}(r_m x)}{r_m^{k+\gamma} \theta(r_m)}, \tag{6-3}$$

and observe that by construction, we have

$$\|w_m\|_{L^\infty(r_m^{-1}\Omega_{j_m} \cap B_1)} \geq \frac{1}{2}, \quad \int_{r_m^{-1}\Omega_{j_m} \cap B_1} w_m(x)Q(r_m x) \, dx = 0 \quad \text{for all } m \in \mathbb{N}, Q \in \mathcal{P}_k. \tag{6-4}$$

Next, we claim that

$$\|w_m\|_{L^\infty(r_m^{-1}\Omega_{j_m} \cap B_R)} \leq cR^{k+\gamma} \quad \text{for all } R \geq 1, m \in \mathbb{N}. \tag{6-5}$$

To see this, we estimate for any $R \geq 1$, using the definitions of $\theta(Rr_m)$ and w_m (see (6-1) and (6-3)),

$$\begin{aligned} \|w_m\|_{L^\infty(r_m^{-1}\Omega_{j_m} \cap B_R)} &\leq \frac{\|u_{j_m} - \mathcal{Q}_{j_m, Rr_m}\|_{L^\infty(\Omega_{j_m} \cap B_{Rr_m})}}{r_m^{k+\gamma}\theta(r_m)} + \frac{\|\mathcal{Q}_{j_m, Rr_m} - \mathcal{Q}_{j_m, r_m}\|_{L^\infty(\Omega_{j_m} \cap B_{Rr_m})}}{r_m^{k+\gamma}\theta(r_m)} \\ &\leq \frac{(Rr_m)^{k+\gamma}\theta(Rr_m)}{r_m^{k+\gamma}\theta(r_m)} + \frac{\|\mathcal{Q}_{j_m, Rr_m} - \mathcal{Q}_{j_m, r_m}\|_{L^\infty(\Omega_{j_m} \cap B_{Rr_m})}}{r_m^{k+\gamma}\theta(r_m)}. \end{aligned} \quad (6-6)$$

Moreover, it follows that for any $j \in \mathbb{N}$, $r > 0$, and $R \geq 1$,

$$\|\mathcal{Q}_{j, Rr} - \mathcal{Q}_{j, r}\|_{L^\infty(\Omega_j \cap B_{Rr})} \leq c\theta(r)(Rr)^{k+\gamma}. \quad (6-7)$$

Indeed, if we write

$$\mathcal{Q}_{j, r}(x) = \sum_{|\beta| \leq k} a_{j, r}^{(\beta)} x_1^{\beta_1} \cdots x_n^{\beta_n}, \quad \beta \in \mathbb{N}^n, \quad a_{j, r}^{(\beta)} \in \mathbb{R},$$

then by [Abatangelo and Ros-Oton 2020, Lemma A.10] we have for any $|\alpha| \leq k$

$$\begin{aligned} r^{|\beta|} |a_{j, r}^{(\beta)} - a_{j, 2r}^{(\beta)}| &\leq c \|\mathcal{Q}_{j, r} - \mathcal{Q}_{j, 2r}\|_{L^\infty(\Omega_j \cap B_r)} \\ &\leq c \|u_j - \mathcal{Q}_{j, r}\|_{L^\infty(\Omega_j \cap B_r)} + c \|u_j - \mathcal{Q}_{j, 2r}\|_{L^\infty(\Omega_j \cap B_{2r})} \\ &\leq c\theta(r)r^{k+\gamma} + c\theta(2r)(2r)^{k+\gamma} \leq c\theta(r)(2r)^{k+\gamma}. \end{aligned}$$

By iteration of this inequality, we obtain for any $l \in \mathbb{N}$

$$\begin{aligned} |a_{j, r}^{(\beta)} - a_{j, 2^l r}^{(\beta)}| &\leq \sum_{i=0}^{l-1} |a_{j, 2^i r}^{(\beta)} - a_{j, 2^{i+1} r}^{(\beta)}| \leq c \sum_{i=0}^{l-1} \theta(2^i r)(2^i r)^{k+\gamma-|\beta|} \\ &\leq c\theta(r)r^{k+\gamma-|\beta|} \sum_{i=0}^{l-1} \frac{\theta(2^i r)}{\theta(r)} 2^{i(k+\gamma-|\beta|)} \leq c\theta(r)(2^l r)^{k+\gamma-|\beta|}. \end{aligned}$$

This yields for any $R > 1$

$$|a_{j, r}^{(\beta)} - a_{j, Rr}^{(\beta)}| \leq c\theta(r)(Rr)^{k+\gamma-|\beta|},$$

which implies (6-7).

Thus, combining (6-6) and (6-7),

$$\|w_m\|_{L^\infty(\Omega_{j_m} \cap B_R)} \leq \frac{(Rr_m)^{k+\gamma}\theta(Rr_m)}{r_m^{k+\gamma}\theta(r_m)} + c \frac{(Rr_m)^{k+\gamma}\theta(r_m)}{r_m^{k+\gamma}\theta(r_m)} \leq cR^{k+\gamma},$$

where we used in the last step that $t \mapsto \theta(t)$ is monotone decreasing, proving (6-5).

Next, using (6-5), we will estimate the $L^{1}_{2s+(k+\lceil\gamma-s\rceil-1)}$ norm of w_m . We have the estimate

$$\begin{aligned} \int_{(\Omega_{j_m} \setminus B_{Rr_m}) \cap \{d_{j_m} \geq \kappa\}} d_{j_m}^{s-1}(y) |y|^{-n-2s-\lceil\gamma-s\rceil+1+\gamma} dy \\ \leq \kappa^{s-1} \int_{(\Omega_{j_m} \setminus B_{Rr_m})} |y|^{-n-s+\gamma-\lceil\gamma-s\rceil} |y|^{1-s} dy \\ \leq c\kappa^{s-1} \text{diam}(\Omega_{j_m})^{1-s} \int_{\mathbb{R}^n \setminus B_{Rr_m}} |y|^{-n-s+\gamma-\lceil\gamma-s\rceil} \leq c(Rr_m)^{\gamma-s-\lceil\gamma-s\rceil}, \end{aligned}$$

where we used that always $\gamma < s + \lceil\gamma - s\rceil < 0$. Moreover, by a similar computation as in Lemma 2.2

(with $\gamma := s - 1 < 2s + \lceil \gamma - s \rceil - 1 - \gamma =: \beta$), we have

$$\int_{(\Omega_{j_m} \setminus B_{Rr_m}) \cap \{d_{j_m} < \kappa\}} d_{j_m}^{s-1}(y) |y|^{-n-2s-\lceil \gamma - s \rceil + 1 + \gamma} dy \leq c(Rr_m)^{\gamma - s - \lceil \gamma - s \rceil}.$$

Thus, altogether, using (6-5) and $\gamma \in (0, 1)$ we obtain

$$\begin{aligned} & \|\tilde{d}_{j_m}^{s-1} w_m | \cdot |^{-n-2s-(k+\lceil \gamma - s \rceil)-1}\|_{L^1(\mathbb{R}^n \setminus B_R)} \\ & \leq c \int_{r_m^{-1} \Omega_{j_m} \setminus B_R} \tilde{d}_{j_m}^{s-1}(y) |y|^{-n-2s-\lceil \gamma - s \rceil + 1 + \gamma} dy \\ & \leq cr_m^{1-s} \int_{r_m^{-1} \Omega_{j_m} \setminus B_R} d_{j_m}^{s-1}(r_m y) |y|^{-n-2s-\lceil \gamma - s \rceil + 1 + \gamma} dy \\ & = cr_m^{s-\gamma+\lceil \gamma - s \rceil} \int_{\Omega_{j_m} \setminus B_{Rr_m}} d_{j_m}^{s-1}(y) |y|^{-n-2s-\lceil \gamma - s \rceil + 1 + \gamma} dy \\ & \leq cr_m^{s-\gamma+\lceil \gamma - s \rceil} \int_{(\Omega_{j_m} \setminus B_R) \cap \{d_{j_m} \geq \kappa\}} d_{j_m}^{s-1}(y) |y|^{-n-2s-\lceil \gamma - s \rceil + 1 + \gamma} dy \\ & \quad + cr_m^{s-\gamma+\lceil \gamma - s \rceil} \int_{(\Omega_{j_m} \setminus B_R) \cap \{d_{j_m} < \kappa\}} d_{j_m}^{s-1}(y) |y|^{-n-2s-\lceil \gamma - s \rceil + 1 + \gamma} dy \\ & \leq cr_m^{s-\gamma+\lceil \gamma - s \rceil} (Rr_m)^{\gamma - s - \lceil \gamma - s \rceil} \leq R^{\gamma - s - \lceil \gamma - s \rceil} \rightarrow 0 \quad \text{as } R \rightarrow \infty. \end{aligned} \quad (6-8)$$

Now, we investigate the equation that is satisfied by w_m . We claim that

$$\frac{|a_{j_m, r_m}^{(\beta)}|}{\theta(r_m)} \rightarrow 0, \quad \text{as } m \rightarrow \infty \quad \text{for all } |\beta| \leq k. \quad (6-9)$$

Indeed, from the considerations above, we deduce that for any $m, l \in \mathbb{N}$,

$$\frac{|a_{j_m, r_m}^{(\beta)} - a_{j_m, 2^l r_m}^{(\beta)}|}{\theta(r_m)} \leq c \sum_{i=1}^l \frac{\theta(2^{l-i} r_m)}{\theta(r_m)} (2^{l-i} r_m)^{k+\gamma-|\beta|}.$$

Hence, choosing $l \in \mathbb{N}$ such that $2^l r_m \in [1, 2)$, we deduce

$$\begin{aligned} \frac{|a_{j_m, r_m}^{(\beta)}|}{\theta(r_m)} & \leq \frac{|a_{j_m, 2^l r_m}^{(\beta)}|}{\theta(r_m)} + \frac{|a_{j_m, r_m}^{(\beta)} - a_{j_m, 2^l r_m}^{(\beta)}|}{\theta(r_m)} \\ & \leq c\theta(r_m)^{-1} \left(|a_{j_m, 2^l r_m}^{(\beta)}| + \sum_{i=1}^l \theta(2^{-i}) (2^{-i})^{k+\gamma-|\beta|} \right) \rightarrow 0 \quad \text{as } m \rightarrow \infty, \end{aligned}$$

which implies (6-9).

Let us now distinguish between the cases ($k = 1$ and $\gamma < s$) and ($k \geq 2$ or $\gamma > s$). In case $k = 1$ and $\gamma < s$, we find that it holds in the viscosity sense

$$\begin{aligned} \tilde{d}_{j_m}^{s-\gamma} L_{j_m}(\tilde{d}_{j_m}^{s-1} w_m) & = r^{\gamma-s} d_{j_m}^{s-\gamma}(r_m \cdot) r_m^{1-s} L_{j_m}(d_{j_m}^{s-1}(r_m \cdot) w_m) \\ & = d_{j_m}^{s-\gamma}(r_m \cdot) \frac{L_{j_m}(d_{j_m}^{s-1} u_{j_m}(r_m \cdot)) - L_{j_m}(d_{j_m}^{s-1} Q_{j_m, r_m}(r_m \cdot))}{r_m^{2s} \theta(r_m)} \\ & = d_{j_m}^{s-\gamma}(r_m \cdot) \frac{f_{j_m}(r_m \cdot) - L_{j_m}(d_{j_m}^{s-1} Q_{j_m, r_m})(r_m \cdot)}{\theta(r_m)} \quad \text{in } r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}}. \end{aligned} \quad (6-10)$$

Moreover, by Corollary 2.5(i),

$$\begin{aligned} \|d_{j_m}^{s-\gamma} L_{j_m}(d_{j_m}^{s-1} Q_{j_m, r_m})(r_m \cdot)\|_{L^\infty(r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}})} &= \|d_{j_m}^{s-\gamma} L_{j_m}(d_{j_m}^{s-1} Q_{j_m, r_m})\|_{L^\infty(\Omega_{j_m} \cap B_1)} \\ &\leq c \sum_{|\beta| \leq 1} |a_{j_m, r_m}^{(\beta)}|. \end{aligned} \quad (6-11)$$

Therefore, recalling $\|d_{j_m}^{s-\gamma} f_{j_m}\|_{L^\infty(\Omega_{j_m} \cap B_1)} \leq C$, and combining (6-10), (6-11), and (6-9), we obtain

$$\begin{aligned} \|\tilde{d}_{j_m}^{s-\gamma} L_{j_m}(\tilde{d}_{j_m}^{s-1} w_m)\|_{L^\infty(r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}})} \\ \leq c \frac{\|d_{j_m}^{s-\gamma} f_{j_m}\|_{L^\infty(\Omega_{j_m} \cap B_1)} + \sum_{|\beta| \leq 1} |a_{j_m, r_m}^{(\beta)}|}{\theta(r_m)} \rightarrow 0 \quad \text{as } m \rightarrow \infty. \end{aligned} \quad (6-12)$$

In case $k \geq 2$ or $\gamma > s$, we first deduce by an argument analogous to (6-10),

$$L_{j_m}(\tilde{d}_{j_m}^{s-1} w_m) = \frac{f_{j_m}(r_m \cdot) - L_{j_m}(d_{j_m}^{s-1} Q_{j_m, r_m})(r_m \cdot)}{r_m^{(k-1)-s+\gamma} \theta(r_m)} \quad \text{in } r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}}. \quad (6-13)$$

Next, using again Corollary 2.5(ii), we obtain

$$\begin{aligned} r^{-(k-1)+s-\gamma} [L_{j_m}(d_{j_m}^{s-1} Q_{j_m, r_m})(r_m \cdot)]_{C^{k-1-s+\gamma}(r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}})} &= [L_{j_m}(d_{j_m}^{s-1} Q_{j_m, r_m})]_{C^{k-1-s+\gamma}(\Omega_{j_m} \cap B_1)} \\ &\leq c \sum_{|\beta| \leq k} |a_{j_m, r_m}^{(\beta)}|, \end{aligned} \quad (6-14)$$

in analogy to (6-11). Finally, recalling

$$r^{-(k-1)+s-\gamma} [f_j(r_m \cdot)]_{C^{(k-1)-s+\gamma}(r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}})} = [f_j]_{C^{(k-1)-s+\gamma}(\Omega_{j_m} \cap B_1)} \leq C,$$

and combining (6-13), (6-14), and (6-9), we obtain

$$[L_{j_m}(\tilde{d}_{j_m}^{s-1} w_m)]_{C^{k-1-s+\gamma}(r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}})} \leq c \frac{[f_j]_{C^{(k-1)-s+\gamma}(\Omega_{j_m} \cap B_1)} + \sum_{|\beta| \leq 1} |a_{j_m, r_m}^{(\beta)}|}{\theta(r_m)} \rightarrow 0 \quad \text{as } m \rightarrow \infty.$$

Thus, there exists a polynomial $p_m \in \mathcal{P}_{k-2+[\gamma-s]}$ such that

$$|L_{j_m}(\tilde{d}_{j_m}^{s-1} w_m) - p_m| \rightarrow 0, \quad \text{as } m \rightarrow \infty \quad \text{in } L_{\text{loc}}^\infty(\{x_n > 0\}). \quad (6-15)$$

Next, considering again all values for γ, k at the same time, we treat the Neumann boundary condition:

$$\partial_\nu w_m = \frac{\partial_\nu u_{j_m}(r_m \cdot) - \partial_\nu(Q_{j_m, r_m})(r_m \cdot)}{r_m^{(k-1)+\gamma} \theta(r_m)} = \frac{g_{j_m}(r_m \cdot) - \partial_\nu(Q_{j_m, r_m})(r_m \cdot)}{r_m^{(k-1)+\gamma} \theta(r_m)} \quad \text{on } \partial r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}}.$$

We obtain

$$r_m^{-(k-1)-\gamma} [\partial_\nu Q_{j_m, r_m}(r_m \cdot)]_{C^{(k-1)+\gamma}(\partial r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}})} = [\partial_\nu Q_{j_m, r_m}]_{C^{(k-1)+\gamma}(\partial \Omega_{j_m} \cap B_1)} \leq c \sum_{|\beta| \leq k} |a_{j_m, r_m}^{(\beta)}|,$$

and using also that $g_{j_m} \in C^{k-1+\gamma}(\Omega_{j_m} \cap B_1)$ by the boundary condition, we deduce

$$[\partial_\nu w_m]_{C^{(k-1)+\gamma}(\partial r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}})} \leq c \frac{[g_{j_m}]_{C^{(k-1)+\gamma}(\partial \Omega_{j_m} \cap B_1)} + \sum_{|\beta| \leq k} |a_{j_m, r_m}^{(\beta)}|}{\theta(r_m)} \leq c \theta(r_m)^{-1} \rightarrow 0,$$

as $m \rightarrow \infty$. Consequently, for any $m \in \mathbb{N}$ there exists a polynomial $q_m \in \mathcal{P}_{k-1}$ such that

$$|\partial_\nu w_m(x) - q_m| \leq c \frac{|x|^\gamma}{\theta(r_m)} \rightarrow 0 \quad \text{for all } x \in \partial r_m^{-1} \Omega_{j_m} \cap B_{r_m^{-1}}. \tag{6-16}$$

Finally, we are in a position to apply the stability theorem (see Lemma 2.13) to w_m . The convergence results in (6-12), (6-15) and (6-16) establish the required convergence of the source terms and the Neumann boundary data.

Moreover, the operators L_{j_m} converge to an operator L with the same ellipticity constants. By the boundary Hölder regularity estimate for solutions to the nonlocal Neumann problem (see Corollary 4.4 applied with $k := k + \lceil \gamma - s \rceil$, $\delta := 0$, $\Omega := r_m^{-1} \Omega_{j_m}$, $v := \tilde{d}_{j_m}^{s-1} w_m$, $f := L_{j_m}(\tilde{d}_{j_m}^{s-1} w_m)$, and $g := \partial_\nu w_m$), together with the Arzelà–Ascoli theorem, the sequence $(w_m)_m$ converges in $L^\infty_{\text{loc}}(\{x_n \geq 0\})$ to some $w \in C(\{x_n \geq 0\})$. All the quantities on the right-hand side of the estimate in Corollary 4.4 will be bounded uniformly in k , due to (6-8), (6-12), (6-15), and (6-16). Thus, in particular $\tilde{d}_{j_m}^{s-1}(r_m \cdot) w_m \rightarrow (x_n)_+^{s-1} w$ locally uniformly in $\{x_n > 0\}$. Finally, in order to apply the stability result in Lemma 2.13, it remains to establish $\tilde{d}_{j_m}^{s-1}(r_m \cdot) w_m \rightarrow (x_n)_+^{s-1} w$ in $L^1_{2s+(k+\lceil \gamma - s \rceil - 1)}(\mathbb{R}^n)$. To see this, we also observe that by (6-5),

$$|w(x)| \leq C(1 + |x|)^{k+\gamma} \quad \text{for all } x \in \{x_n > 0\}. \tag{6-17}$$

Therefore, using also (6-8) and a computation based on polar coordinates (along the lines of (5-3)) we obtain, since $\gamma < 1$,

$$\begin{aligned} & \int_{\mathbb{R}^n \setminus B_R} |\tilde{d}_{j_m}^{s-1}(y) w_m(y) - (y_n)_+^{s-1} w(y)| |y|^{-n-2s-(k+\lceil \gamma - s \rceil - 1)} \, dy \\ & \leq C \int_{\mathbb{R}^n \setminus B_R} (y_n)_+^{s-1} |y|^{-n-2s-\lceil \gamma - s \rceil + 1 + \gamma} \, dy + C \int_{(r_m^{-1} \Omega_{j_m}) \setminus B_R} \tilde{d}_{j_m}^{s-1}(y) |y|^{-n-2s-\lceil \gamma - s \rceil + 1 + \gamma} \, dy \\ & \leq C R^{\gamma - s - \lceil \gamma - s \rceil} \rightarrow 0 \quad \text{as } R \rightarrow \infty. \end{aligned}$$

This implies $\tilde{d}_{j_m}^{s-1} w_m \rightarrow (x_n)_+^{s-1} w$ in $L^1_{2s+(k+\lceil \gamma - s \rceil - 1)}(\mathbb{R}^n)$, by combining it with the locally uniform convergence in $L^\infty_{\text{loc}}(\{x_n \geq 0\})$.

Thus, by stability of viscosity solutions (see Lemma 2.13), we deduce that in the viscosity sense

$$\begin{cases} L((x_n)_+^{s-1} w)^{k-1+\lceil \gamma - s \rceil} \equiv 0 & \text{in } \{x_n > 0\}, \\ \partial_n w = p & \text{on } \{x_n = 0\}, \end{cases}$$

where $p \in \mathcal{P}_{k-1}$ is a polynomial, and moreover, by (6-4), it must be that

$$\|w\|_{L^\infty(B_1 \cap \{x_n > 0\})} \geq \frac{1}{2}. \tag{6-18}$$

An application of the Liouville theorem (see Theorem 5.1, using (6-17)) yields now that $w \in \mathcal{P}_k$. Thus, we can choose $Q(x) = w(r_m^{-1}x)$ in (6-4). This implies that

$$0 = \lim_{m \rightarrow \infty} \int_{B_1 \cap r_m^{-1} \Omega_{j_m}} w_m(x) Q(r_m x) \, dx = \lim_{m \rightarrow \infty} \int_{B_1 \cap r_m^{-1} \Omega_{j_m}} w_m(x) w(x) \, dx = \int_{B_1 \cap \{x_n > 0\}} w^2(x) \, dx,$$

where we used in the last step $w_m \rightarrow w$ and $r_m^{-1} \Omega_{j_m} \rightarrow \{x_n > 0\}$. This yields $w = 0$, which however contradicts (6-18), and thus, we conclude the proof of (ii).

Finally, if $k = 1$ and $\gamma < s$, then by the Liouville theorem (see Theorem 5.1), there exist $a \in \mathbb{R}^n$ and $b \in \mathbb{R}$ such that

$$w(x) = (a, x) + b.$$

Moreover, if $g_j \equiv 0$, then also $g \equiv 0$. Thus, it must be that $\partial_n w = 0$ in $\{x_n = 0\}$, which implies $a_n = 0$. \square

We are now in a position to prove our main result.

Proof of Theorem 1.2. We define $u := v/d^{s-1}$. Let us assume that

$$\|u\|_{L^\infty(\mathbb{R}^n)} + \|d^{s-\gamma} f\|_{L^\infty(\Omega \cap B_2)} \mathbb{1}_{\{1+\gamma < 2s\}} + \|f\|_{C^{k-2s+\gamma}(\Omega \cap B_2)} \mathbb{1}_{\{1+\gamma > 2s\}} + \|g\|_{C^{k-1+\gamma}(\partial\Omega \cap B_2)} \leq 1.$$

First, we claim that for any $x_0 \in \Omega \cap B_{1/2}$ with $z \in \partial\Omega \cap B_{1/2}$ such that $|x_0 - z| = d(x_0) =: r \leq 1$, there exists a polynomial $Q \in \mathcal{P}_k$ of degree k such that

$$[u - Q]_{C^{k+\gamma}(B_{r/2}(x_0))} \leq c \tag{6-19}$$

for some constant $c > 0$, depending only on $n, s, \lambda, \Lambda, \gamma, \Omega, k$, where we assume without loss of generality that $v_z = e_n$. This estimate already yields the desired result since it implies

$$[u]_{C^{k+\gamma}(B_{r/2}(x_0))} \leq [u - Q]_{C^{k+\gamma}(B_{r/2}(x_0))} + [Q]_{C^{k+\gamma}(B_{r/2}(x_0))} \leq c.$$

From here, a covering argument (see [Fernández-Real and Ros-Oton 2024a, Lemma A.1.4]) together with Hölder interpolation (recall that $\|u\|_{L^\infty(\mathbb{R}^n)} \leq 1$) yields the desired regularity estimate in $\bar{\Omega} \cap B_1$. Improving the global L^∞ norm to the $L_{2s}^1(\mathbb{R}^n)$ norm, or the $L_{k+\gamma}^1(\mathbb{R}^n \setminus B_2)$ norm, respectively, in the estimate goes by the exact same arguments as in the proofs of Lemma 2.10 and Corollary 4.4.

To see (6-19), let us take $z \in \partial\Omega \cap B_{1/2}$ such that $|x_0 - z| = d(x_0) = r$, and apply Lemma 6.1 to see that there exists a polynomial $Q \in \mathcal{P}_k$ such that the function

$$u_r(x) := \frac{u(x_0 + rx) - Q(x_0 + rx)}{r^{k+\gamma}} \quad \text{satisfies} \quad \|u_r\|_{L^\infty(B_R)} \leq CR^{k+\gamma} \quad \text{for all } R \in [1, r^{-1}].$$

Moreover, since $\|u\|_{L^\infty(\mathbb{R}^n)} \leq 1$, and $Q \in \mathcal{P}_k$, we deduce

$$\|u_r\|_{L^\infty(B_R)} \leq Cr^{-k-\gamma}(1 + (rR)^k) \leq CR^{k+\gamma} \quad \text{for all } R \geq r^{-1}.$$

Together, this implies

$$\begin{aligned} \int_{\mathbb{R}^n \setminus B_{3/4}} \frac{d^{s-1}(x_0 + rx)|u_r(x)|}{|x|^{n+k+\gamma}} dx &\leq c \int_{\mathbb{R}^n \setminus B_{3/4}} \frac{d^{s-1}(x_0 + rx)}{|x|^n} dx = c \int_{\mathbb{R}^n \setminus B_{3/4}} \frac{d^{s-1}(x_0 + x)}{|x|^n} dx \\ &\leq c(1 + r^{s-1}), \end{aligned}$$

where we used Lemma 2.2 with $\gamma := s - 1 < 0 =: \beta$. Moreover, we have by the definition of r ,

$$\|d^{s-1}(x_0 + r \cdot)u_r\|_{L^\infty(B_{3/4})} \leq cr^{s-1}.$$

Now, we apply the interior regularity theory for nonlocal problems (see Lemma 2.10). To do so, we distinguish between the cases (i) $k = 1$ and $k + \gamma = 1 + \gamma \leq 2s$ and (ii) $k + \gamma > 2s$. In case (i), we apply

Lemma 2.10(i) with $\beta = 1 + \gamma$, observe that automatically $\gamma < s$, and obtain

$$\begin{aligned}
[d^{s-1}(u - Q)]_{C^{1+\gamma}(B_{r/2}(x_0))} &= [d^{s-1}(x_0 + r \cdot)u_r]_{C^{1+\gamma}(B_{1/2})} \\
&\leq c \|d^{s-1}(x_0 + r \cdot)u_r\|_{L^\infty(B_{3/4})} + c \left\| \frac{d^{s-1}(x_0 + r \cdot)u_r}{|x|^{n+1+\gamma}} \right\|_{L^1(\mathbb{R}^n \setminus B_{3/4})} \\
&\quad + c \|L(d^{s-1}(x_0 + r \cdot)u_r)\|_{L^\infty(B_{3/4})} \\
&\leq cr^{s-1} + cr^{2s-(1+\gamma)} \|Lv\|_{L^\infty(B_{3r/4}(x_0))} + cr^{2s-(1+\gamma)} \|L(d^{s-1}Q)\|_{L^\infty(B_{3r/4}(x_0))} \\
&\leq cr^{s-1} + cr^{s-1} \|d^{s-\gamma}f\|_{L^\infty(B_r(x_0))} + cr^{s-1} \leq cr^{s-1},
\end{aligned}$$

where we used Corollary 2.5(i) and that $d \geq r/4$ in $B_{3r/4}(x_0)$ by construction, and $r \leq 1$.

In case (ii), we apply Lemma 2.10(ii) with $\alpha := k + \gamma - 2s > 0$ and obtain

$$\begin{aligned}
[d^{s-1}(u - Q)]_{C^{k+\gamma}(B_{r/2}(x_0))} &= [d^{s-1}(x_0 + r \cdot)u_r]_{C^{k+\gamma}(B_{1/2})} \\
&\leq c \|d^{s-1}(x_0 + r \cdot)u_r\|_{L^\infty(B_{3/4})} + c \left\| \frac{d^{s-1}(x_0 + r \cdot)u_r}{|x|^{n+k+\gamma}} \right\|_{L^1(\mathbb{R}^n \setminus B_{3/4})} \\
&\quad + c [L(d^{s-1}(x_0 + r \cdot)u_r)]_{C^{k+\gamma-2s}(B_{3/4})} \\
&\leq cr^{s-1} + c [Lv]_{C^{k+\gamma-2s}(B_{3r/4}(x_0))} + c [L(d^{s-1}Q)]_{C^{k+\gamma-2s}(B_{3r/4}(x_0))} \\
&\leq cr^{s-1} + cr^{s-1} \|f\|_{C^{k+\gamma-2s}(B_r(x_0))} + cr^{s-1} \leq cr^{s-1},
\end{aligned}$$

where we used Corollary 2.5(iii) in the second to last step.

Moreover, using again the L^∞ estimate for u_r with $R = 1$, we have

$$\|d^{s-1}(u - Q)\|_{L^\infty(B_{r/2}(x_0))} \leq cr^{s-1} \|u - Q\|_{L^\infty(B_{r/2}(x_0))} \leq cr^{s+(k-1)+\gamma} \|u_r\|_{L^\infty(B_1)} \leq cr^{s+(k-1)+\gamma},$$

and hence by Hölder interpolation, we obtain that for any $\delta \in (0, k + \gamma)$ it holds that

$$[d^{s-1}(u - Q)]_{C^\delta(B_{r/2}(x_0))} \leq cr^{s-1+k+\gamma-\delta}.$$

Therefore, altogether by the product rule,

$$\begin{aligned}
[(u - Q)]_{C^{k+\gamma}(B_{r/2}(x_0))} &= [D^k(d^{1-s}d^{s-1}(u - Q))]_{C^\gamma(B_{r/2}(x_0))} \\
&\leq \sum_{|\beta|=k} \sum_{\alpha \leq \beta} [(\partial^\alpha d^{1-s})(\partial^{\beta-\alpha} d^{s-1}(u - Q))]_{C^\gamma(B_{r/2}(x_0))} \\
&\leq \sum_{|\beta|=k} \sum_{\alpha \leq \beta} (\|\partial^\alpha d^{1-s}\|_{L^\infty(B_{r/2}(x_0))} [\partial^{\beta-\alpha} d^{s-1}(u - Q)]_{C^\gamma(B_{r/2}(x_0))} \\
&\quad + \|\partial^{\beta-\alpha} d^{s-1}(u - Q)\|_{L^\infty(B_{r/2}(x_0))} [\partial^\alpha d^{1-s}]_{C^\gamma(B_{r/2}(x_0))}) \\
&\leq c \sum_{|\beta|=k} \sum_{\alpha \leq \beta} (r^{1-s-|\alpha|} r^{s-1+k+\gamma-(k-|\alpha|+\gamma)} + r^{s-1+k+\gamma-(k-|\alpha|)} r^{1-s-|\alpha|-\gamma}) \\
&\leq c \sum_{|\beta|=k} \sum_{\alpha \leq \beta} \leq c,
\end{aligned} \tag{6-20}$$

where we used that $r \leq 1$, and the following observation based on the fact that $d \in C^{k+1,\gamma}(\bar{\Omega})$ together with corresponding estimates $|D^j d| \leq c_j d^{1-j}$ (resp. $|D^j d^{1-s}| \leq c_j d^{1-s-j}$) in Ω for every $j \leq k$ (see [Fernández-Real and Ros-Oton 2024a, Lemma B.0.1]):

$$\begin{aligned} [\partial^\alpha d^{1-s}]_{C^\gamma(B_{r/2}(x_0))} &\leq \|D^{|\alpha|+1} d^{1-s}\|_{L^\infty(B_{r/2}(x_0))} \sup_{x,y \in B_{r/2}(x_0)} |x-y|^{1-\gamma} \\ &\leq cr^{1-s-|\alpha|-\gamma} \end{aligned} \quad \text{for all } |\alpha| \leq k. \quad (6-21)$$

This proves our claim (6-19). We can replace the L^∞ norm of u in $\mathbb{R}^n \setminus B_2$ by the $L^1_{2s}(\mathbb{R}^n)$ norm via a truncation argument, as in the proof of Corollary 4.4. We conclude the proof. \square

Finally, we explain how to prove Theorem 1.7.

Proof of Theorem 1.7. The result follows immediately from Theorem 1.2, however it remains to prove that the result only requires $K \in C^{k-2s+\gamma}(\mathbb{S}^{n-1})$ if $\Omega = \{x_n = 0\}$. First of all, we recall the Hölder estimate (see Corollary 5.4), which holds true without any regularity assumption on K if $k = 0$. For the Liouville theorem (see Theorem 5.1), we only require $K \in C^{k-1-s+\gamma+\delta}(\mathbb{S}^{n-1})$ for an arbitrarily small $\delta > 0$ and $k-1-s+\gamma < k-2s+\gamma$. In Lemma 6.1, additional regularity for K is assumed in order to apply Corollary 2.5. However, if $\Omega = \{x_n > 0\}$, we have

$$L(d^{s-1}Q) = L((x_n)_+^{s-1}Q) \stackrel{k-1}{=} 0 \quad \text{in } \{x_n > 0\}$$

for any $Q \in \mathcal{P}_k$. Hence, in case $k = 1$ and $\gamma < s$, the proof goes through exactly as before, without any restrictions on K . If $k \geq 2$ or $\gamma > s$, (6-13) needs to be interpreted as an equation up to a polynomial, but the rest of the proof remains the same. Moreover, we apply the Hölder estimate (see Corollary 5.4), which would force us to impose $K \in C^{k-1}(\mathbb{S}^{n-1})$ in case $k \geq 2$ or $\gamma > s$. Therefore, in this case, we need to proceed a little differently. Indeed, we replace the computation in (6-8) by the following estimate, based on polar coordinates (see also (5-3)) for $\eta = 1 + \gamma - s - \lceil \gamma - s \rceil + \delta$ for some very small $\delta > 0$:

$$\begin{aligned} \|(x_n)_+^{s-1} w_m \cdot |\cdot|^{-n-2s-(k-2+\lceil \gamma-s \rceil+\eta)}\|_{L^1(\mathbb{R}^n \setminus B_R)} &\leq c \int_{\mathbb{R}^n \setminus B_R} (x_n)_+^{s-1} |x|^{-n-2s-\lceil \gamma-s \rceil+2+\gamma+\eta} dx \\ &\leq c \int_0^{2\pi} \cos(\theta)_+^{s-1} \left(\int_R^\infty r^{s-1} r^{-1-2s-\lceil \gamma-s \rceil+2+\gamma+\eta} dr \right) d\theta \\ &= c \int_0^{2\pi} \cos(\theta)_+^{s-1} \left(\int_R^\infty r^{-1-\delta} dr \right) d\theta \leq cR^{-\delta} \rightarrow 0, \end{aligned}$$

as $R \rightarrow \infty$. Then, we can apply Corollary 5.4 with $k := k-1$ and $\delta := \eta$ and only need to assume that $K \in C^{k-2+\eta}(\mathbb{S}^{n-1})$, which is fine by the same reasoning as for the Liouville theorem above. Moreover, the stability theorem (see Lemma 2.13) can still be applied since $k-2+\eta \leq k-1$ if we choose $\delta < s-\eta$.

Finally, the proof of Theorem 1.2 relies on an application of the interior regularity result (see Lemma 2.10). In case $k = 1$ and $1 + \gamma \leq 2s$, we apply Lemma 2.10(i), so in this case, no regularity assumption on K is required, at all. In case $1 + \gamma > 2s$, we apply Lemma 2.10(ii) with $\alpha := k + \gamma - 2s$ (and interpret the equation up to a polynomial of degree $k-1$, which is possible due to Remark 2.11), so in this case, we need to assume only that $K \in C^{k-2s+\gamma}(\mathbb{S}^{n-1})$, as desired. \square

7. Nonlocal equations with local Dirichlet boundary conditions

Finally, we give the proof of the boundary regularity for nonlocal equations with Dirichlet boundary conditions (see Theorem 1.4).

Proof of Theorem 1.4. Let us first extend h in such a way that $h \in C^{k+\gamma}(\mathbb{R}^n)$. Then, we define $w := v - d^{s-1}h$ and observe that w solves

$$\begin{cases} Lw = \tilde{f} & \text{in } \Omega, \\ w = 0 & \text{in } \mathbb{R}^n \setminus \Omega, \\ w/d^{s-1} = 0 & \text{on } \partial\Omega, \end{cases}$$

where $\tilde{f} := f - L(d^{s-1}h)$. Moreover, for $x_0 \in \Omega$ an application of Corollary 2.5 yields $|\tilde{f}(x_0)| \leq c_1 d^{\gamma-s}(x_0)$ in case $k + \gamma < 1 + s$, as well as $[\tilde{f}]_{C^{k-1-s+\gamma}(\bar{\Omega})} \leq c_2$ in case $k + \gamma > 1 + s$, and also $[\tilde{f}]_{C^{k-2s+\gamma}(B_{d(x_0)/2}(x_0))} \leq c_3 d^{s-1}(x_0)$ in case $k + \gamma > 2s$. Since $w/d^{s-1} = 0$ on $\partial\Omega$, by the maximum principle (see Proposition 1.3) and a barrier argument (see for instance the proof of [Fernández-Real and Ros-Oton 2024a, Lemma 2.3.9], using the barrier from [Fernández-Real and Ros-Oton 2024a, Lemma 2.3.10] in case $k + \gamma > 1 + s$ and the barrier $\tilde{\psi}$ from the second claim in Lemma 2.7 in case $k + \gamma < 1 + s$) it holds that $w \in L^\infty(\Omega)$ and

$$\|w\|_{L^\infty(\Omega)} \leq C \|d^{s-\gamma} \tilde{f}\|_{L^\infty(\Omega)}. \quad (7-1)$$

Thus w is a solution in the setting of [Ros-Oton and Serra 2017; Abatangelo and Ros-Oton 2020]. We assume without loss of generality

$$\|w\|_{L^\infty(\Omega)} + \|d^{s-\gamma} \tilde{f}\|_{L^\infty(\bar{\Omega})} \mathbb{1}_{\{k+\gamma < 1+s\}} + \|\tilde{f}\|_{C^{k-1-s+\gamma}(\Omega)} \mathbb{1}_{\{k+\gamma > 1+s\}} \leq 1.$$

Then, [Ros-Oton and Serra 2017; Abatangelo and Ros-Oton 2020] imply that for any $z \in \partial\Omega$ there exists a polynomial $Q_z \in \mathcal{P}_{k-1}$ such that

$$|w(x) - Q_z(x)d^s| \leq c|x - z|^{k-1+\gamma+s} \leq c|x - z|^{k+\gamma} d^{s-1}(x) \quad \text{for all } x \in B_1(z).$$

By adjusting the proof of [Ros-Oton and Serra 2017, Proposition 3.2] in case $k + \gamma < 1 + s$, or the second part of the proof of [Abatangelo and Ros-Oton 2020, Proposition 4.1] in case $k + \gamma > 1 + s$, respectively, according to the slight modification of the upper bound in the previous estimate, we get that for any $x_0 \in \Omega \cap B_1(z)$, letting $r := d(x_0)$,

$$\|w - Q_z d^s\|_{L^\infty(B_{r/2}(x_0))} \leq cr^{k+\gamma+s-1}, \quad [w - Q_z d^s]_{C^{k+\gamma}(B_{r/2}(x_0))} \leq cr^{s-1}. \quad (7-2)$$

Indeed, while the first estimate is immediate from the expansion, the second result follows by letting

$$v_r(x) = r^{-k-\gamma} (u(x_0 + rx) - Q_z(x_0 + rx)d^s(x_0 + rx)),$$

and observing that by the previous estimate and the properties of \tilde{f} it holds that

$$\|v_r\|_{L^\infty(B_R)} \leq c(1 + r^{s-1}) \quad \text{for all } R > 0, \quad [\tilde{f}]_{C^{k+\gamma-2s}(B_{r/2}(x_0))} \mathbb{1}_{\{k+\gamma > 2s\}} \leq cr^{s-1}.$$

Plugging these findings into the remainder of [Ros-Oton and Serra 2017, proof of Theorem 1.2; Abatangelo and Ros-Oton 2020, proof of Theorem 1.4], we obtain (7-2). From there we can show, using Hölder interpolation, and also $d \in C^{k+1+\gamma}(\bar{\Omega})$, that for any $\delta \in (0, k + \gamma]$ it holds that

$$[w - Q_z d^s]_{C^\delta(B_{r/2}(x_0))} \leq cr^{k+\gamma+s-1-\delta}, \quad \|d^{1-s}\|_{L^\infty(B_{r/2}(x_0))} \leq cr^{1-s}, \quad [d^{1-s}]_{C^\delta(B_{r/2}(x_0))} \leq cr^{1-s-\delta}.$$

Thus, proceeding in a similar way as in the proof of Theorem 1.2, and using (7-2) as well as the previous estimate, we obtain

$$\begin{aligned} \left[\frac{w}{d^{s-1}} - Q_z d \right]_{C^{k+\gamma}(B_{r/2}(x_0))} &= [D^k(d^{1-s}(w - Q_z d^s))]_{C^\gamma(B_{r/2}(x_0))} \\ &\leq \sum_{|\beta|=k} \sum_{\alpha \leq \beta} [(\partial^\alpha d^{1-s})(\partial^{\beta-\alpha}(w - Q_z d^s))]_{C^\gamma(B_{r/2}(x_0))} \\ &\leq \sum_{|\beta|=k} \sum_{\alpha \leq \beta} (\|\partial^\alpha d^{1-s}\|_{L^\infty(B_{r/2}(x_0))} [\partial^{\beta-\alpha}(w - Q_z d^s)]_{C^\gamma(B_{r/2}(x_0))} \\ &\quad + [\partial^\alpha d^{1-s}]_{C^\gamma(B_{r/2}(x_0))} \|\partial^{\beta-\alpha}(w - Q_z d^s)\|_{L^\infty(B_{r/2}(x_0))}) \\ &\leq c \sum_{|\beta|=k} \sum_{\alpha \leq \beta} (r^{1-s-|\alpha|} r^{s-1+|\alpha|} + r^{1-s-|\alpha|-\gamma} r^{\gamma+s-1+|\alpha|}) \leq c. \end{aligned}$$

From here, by a covering argument, and using the continuity of the extension operator,

$$\begin{aligned} \left\| \frac{v}{d^{s-1}} \right\|_{C^{k,\gamma}(\bar{\Omega})} &\leq c(\|v - d^{s-1}h\|_{L^\infty(\mathbb{R}^n)} + \|f\|_{C^{k-1-s+\gamma}(\Omega)} + \|h\|_{C^{k+\gamma}(\partial\Omega)}) \\ &\leq c(\|f\|_{C^{k-1-s+\gamma}(\Omega)} + \|h\|_{C^{k+\gamma}(\partial\Omega)}). \end{aligned} \quad \square$$

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References

- [Abatangelo 2015] N. Abatangelo, “Large S -harmonic functions and boundary blow-up solutions for the fractional Laplacian”, *Discrete Contin. Dyn. Syst.* **35**:12 (2015), 5555–5607. MR
- [Abatangelo 2017] N. Abatangelo, “Very large solutions for the fractional Laplacian: towards a fractional Keller–Osserman condition”, *Adv. Nonlinear Anal.* **6**:4 (2017), 383–405. MR
- [Abatangelo and Ros-Oton 2020] N. Abatangelo and X. Ros-Oton, “Obstacle problems for integro-differential operators: higher regularity of free boundaries”, *Adv. Math.* **360** (2020), art. id. 106931. MR
- [Abatangelo et al. 2023] N. Abatangelo, D. Gómez-Castro, and J. L. Vázquez, “Singular boundary behaviour and large solutions for fractional elliptic equations”, *J. Lond. Math. Soc. (2)* **107**:2 (2023), 568–615. MR
- [Abels and Grubb 2023] H. Abels and G. Grubb, “Fractional-order operators on nonsmooth domains”, *J. Lond. Math. Soc. (2)* **107**:4 (2023), 1297–1350. MR

- [Alves and Torres Ledesma 2020] C. O. Alves and C. E. Torres Ledesma, “Fractional elliptic problem in exterior domains with nonlocal Neumann condition”, *Nonlinear Anal.* **195** (2020), art. id. 111732. MR
- [Audrito et al. 2023] A. Audrito, J.-C. Felipe-Navarro, and X. Ros-Oton, “The Neumann problem for the fractional Laplacian: regularity up to the boundary”, *Ann. Sc. Norm. Super. Pisa Cl. Sci. (5)* **24**:2 (2023), 1155–1222. MR
- [Barles et al. 2014a] G. Barles, E. Chasseigne, C. Georgelin, and E. R. Jakobsen, “On Neumann type problems for nonlocal equations set in a half space”, *Trans. Amer. Math. Soc.* **366**:9 (2014), 4873–4917. MR
- [Barles et al. 2014b] G. Barles, C. Georgelin, and E. R. Jakobsen, “On Neumann and oblique derivatives boundary conditions for nonlocal elliptic equations”, *J. Differential Equations* **256**:4 (2014), 1368–1394. MR
- [Barrios et al. 2014] B. Barrios, A. Figalli, and E. Valdinoci, “Bootstrap regularity for integro-differential operators and its application to nonlocal minimal surfaces”, *Ann. Sc. Norm. Super. Pisa Cl. Sci. (5)* **13**:3 (2014), 609–639. MR
- [Bass and Levin 2002] R. F. Bass and D. A. Levin, “Harnack inequalities for jump processes”, *Potential Anal.* **17**:4 (2002), 375–388. MR
- [Biswas and Jaroahs 2020] A. Biswas and S. Jaroahs, “On overdetermined problems for a general class of nonlocal operators”, *J. Differential Equations* **268**:5 (2020), 2368–2393. MR
- [Bogdan 1999] K. Bogdan, “Representation of α -harmonic functions in Lipschitz domains”, *Hiroshima Math. J.* **29**:2 (1999), 227–243. MR
- [Bogdan et al. 2003] K. Bogdan, K. Burdzy, and Z.-Q. Chen, “Censored stable processes”, *Probab. Theory Related Fields* **127**:1 (2003), 89–152. MR
- [Bogdan et al. 2009] K. Bogdan, T. Byczkowski, T. Kulczycki, M. Ryznar, R. Song, and Z. Vondraček, *Potential analysis of stable processes and its extensions*, edited by P. Graczyk and A. Stos, Lecture Notes in Mathematics **1980**, Springer, 2009. MR
- [Caffarelli and Silvestre 2009] L. Caffarelli and L. Silvestre, “Regularity theory for fully nonlinear integro-differential equations”, *Comm. Pure Appl. Math.* **62**:5 (2009), 597–638. MR
- [Caffarelli and Silvestre 2011a] L. Caffarelli and L. Silvestre, “The Evans–Krylov theorem for nonlocal fully nonlinear equations”, *Ann. of Math. (2)* **174**:2 (2011), 1163–1187. MR
- [Caffarelli and Silvestre 2011b] L. Caffarelli and L. Silvestre, “Regularity results for nonlocal equations by approximation”, *Arch. Ration. Mech. Anal.* **200**:1 (2011), 59–88. MR
- [Caffarelli et al. 2010] L. A. Caffarelli, J.-M. Roquejoffre, and Y. Sire, “Variational problems for free boundaries for the fractional Laplacian”, *J. Eur. Math. Soc. (JEMS)* **12**:5 (2010), 1151–1179. MR
- [Chan et al. 2021] H. Chan, D. Gómez-Castro, and J. L. Vázquez, “Blow-up phenomena in nonlocal eigenvalue problems: when theories of L^1 and L^2 meet”, *J. Funct. Anal.* **280**:7 (2021), art. id. 108845. MR
- [Chen and Kim 2002] Z.-Q. Chen and P. Kim, “Green function estimate for censored stable processes”, *Probab. Theory Related Fields* **124**:4 (2002), 595–610. MR
- [De Silva and Roquejoffre 2012] D. De Silva and J. M. Roquejoffre, “Regularity in a one-phase free boundary problem for the fractional Laplacian”, *Ann. Inst. H. Poincaré C Anal. Non Linéaire* **29**:3 (2012), 335–367. MR
- [De Silva and Savin 2012] D. De Silva and O. Savin, “ $C^{2,\alpha}$ regularity of flat free boundaries for the thin one-phase problem”, *J. Differential Equations* **253**:8 (2012), 2420–2459. MR
- [De Silva et al. 2014] D. De Silva, O. Savin, and Y. Sire, “A one-phase problem for the fractional Laplacian: regularity of flat free boundaries”, *Bull. Inst. Math. Acad. Sin. (N.S.)* **9**:1 (2014), 111–145. MR
- [Di Castro et al. 2014] A. Di Castro, T. Kuusi, and G. Palatucci, “Nonlocal Harnack inequalities”, *J. Funct. Anal.* **267**:6 (2014), 1807–1836. MR
- [Di Castro et al. 2016] A. Di Castro, T. Kuusi, and G. Palatucci, “Local behavior of fractional p -minimizers”, *Ann. Inst. H. Poincaré C Anal. Non Linéaire* **33**:5 (2016), 1279–1299. MR
- [Dipierro et al. 2017] S. Dipierro, X. Ros-Oton, and E. Valdinoci, “Nonlocal problems with Neumann boundary conditions”, *Rev. Mat. Iberoam.* **33**:2 (2017), 377–416. MR
- [Dipierro et al. 2019] S. Dipierro, O. Savin, and E. Valdinoci, “Definition of fractional Laplacian for functions with polynomial growth”, *Rev. Mat. Iberoam.* **35**:4 (2019), 1079–1122. MR

- [Dipierro et al. 2022] S. Dipierro, A. Dzhugan, and E. Valdinoci, “Integral operators defined “up to a polynomial””, *Fract. Calc. Appl. Anal.* **25**:1 (2022), 60–108. MR
- [Dipierro et al. 2023] S. Dipierro, G. Poggesi, J. Thompson, and E. Valdinoci, “Quantitative stability for the nonlocal overdetermined Serrin problem”, preprint, 2023. arXiv 2309.17119
- [Du et al. 2012] Q. Du, M. Gunzburger, R. B. Lehoucq, and K. Zhou, “Analysis and approximation of nonlocal diffusion problems with volume constraints”, *SIAM Rev.* **54**:4 (2012), 667–696. MR
- [Dyda 2012] B. Dyda, “Fractional calculus for power functions and eigenvalues of the fractional Laplacian”, *Fract. Calc. Appl. Anal.* **15**:4 (2012), 536–555. MR
- [Fall and Jarohs 2015] M. M. Fall and S. Jarohs, “Overdetermined problems with fractional Laplacian”, *ESAIM Control Optim. Calc. Var.* **21**:4 (2015), 924–938. MR
- [Fall and Ros-Oton 2022] M. M. Fall and X. Ros-Oton, “Global Schauder theory for minimizers of the $H^s(\Omega)$ energy”, *J. Funct. Anal.* **283**:3 (2022), art. id. 109523. MR
- [Felsinger et al. 2015] M. Felsinger, M. Kassmann, and P. Voigt, “The Dirichlet problem for nonlocal operators”, *Math. Z.* **279**:3-4 (2015), 779–809. MR
- [Fernández-Real and Ros-Oton 2024a] X. Fernández-Real and X. Ros-Oton, *Integro-differential elliptic equations*, Progress in Mathematics **350**, Springer, 2024. MR
- [Fernández-Real and Ros-Oton 2024b] X. Fernández-Real and X. Ros-Oton, “Stable cones in the thin one-phase problem”, *Amer. J. Math.* **146**:3 (2024), 631–685. MR
- [Feulefack and Jarohs 2023] P. A. Feulefack and S. Jarohs, “Nonlocal operators of small order”, *Ann. Mat. Pura Appl. (4)* **202**:4 (2023), 1501–1529. MR
- [Foghem and Kassmann 2024] G. Foghem and M. Kassmann, “A general framework for nonlocal Neumann problems”, *Commun. Math. Sci.* **22**:1 (2024), 15–66. MR
- [Gettoor 1961] R. K. Gettoor, “First passage times for symmetric stable processes in space”, *Trans. Amer. Math. Soc.* **101** (1961), 75–90. MR
- [Grubb 2014] G. Grubb, “Local and nonlocal boundary conditions for μ -transmission and fractional elliptic pseudodifferential operators”, *Anal. PDE* **7**:7 (2014), 1649–1682. MR
- [Grubb 2015] G. Grubb, “Fractional Laplacians on domains, a development of Hörmander’s theory of μ -transmission pseudodifferential operators”, *Adv. Math.* **268** (2015), 478–528. MR
- [Grubb 2018] G. Grubb, “Green’s formula and a Dirichlet-to-Neumann operator for fractional-order pseudodifferential operators”, *Comm. Partial Differential Equations* **43**:5 (2018), 750–789. MR
- [Grubb 2023] G. Grubb, “Resolvents for fractional-order operators with nonhomogeneous local boundary conditions”, *J. Funct. Anal.* **284**:7 (2023), art. id. 109815. MR
- [Grube and Hensiek 2023] F. Grube and T. Hensiek, “Maximum principle for stable operators”, *Math. Nachr.* **296**:12 (2023), 5684–5702. MR
- [Grube and Hensiek 2024] F. Grube and T. Hensiek, “Robust nonlocal trace spaces and Neumann problems”, *Nonlinear Anal.* **241** (2024), art. id. 113481. MR
- [Hmissi 1994] F. Hmissi, “Fonctions harmoniques pour les potentiels de Riesz sur la boule unité”, *Exposition. Math.* **12**:3 (1994), 281–288. MR
- [Jarohs and Weth 2019] S. Jarohs and T. Weth, “On the strong maximum principle for nonlocal operators”, *Math. Z.* **293**:1-2 (2019), 81–111. MR
- [Kassmann 2009] M. Kassmann, “A priori estimates for integro-differential operators with measurable kernels”, *Calc. Var. Partial Differential Equations* **34**:1 (2009), 1–21. MR
- [Korvenpää et al. 2016] J. Korvenpää, T. Kuusi, and G. Palatucci, “The obstacle problem for nonlinear integro-differential operators”, *Calc. Var. Partial Differential Equations* **55**:3 (2016), art. id. 63. MR
- [Kukuljan 2021] T. Kukuljan, “The fractional obstacle problem with drift: higher regularity of free boundaries”, *J. Funct. Anal.* **281**:8 (2021), art. id. 109114. MR

- [Landkof 1972] N. S. Landkof, *Foundations of modern potential theory*, Grundle. Math. Wissen. **180**, Springer, 1972. MR
- [Li and Liu 2023] C. Li and C. Liu, “Uniqueness and some related estimates for Dirichlet problem with fractional Laplacian”, *Bull. Lond. Math. Soc.* **55**:6 (2023), 2685–2704. MR
- [Lian and Zhang 2023] Y. Lian and K. Zhang, “Boundary pointwise regularity and applications to the regularity of free boundaries”, *Calc. Var. Partial Differential Equations* **62**:8 (2023), art. id. 230. MR
- [Liu and Zhuo 2025] C. Liu and R. Zhuo, “On the Dirichlet problem for fractional Laplace equation on a general domain”, *Commun. Contemp. Math.* **27**:6 (2025), art. id. 2450037. MR
- [Ros-Oton 2016] X. Ros-Oton, “Nonlocal elliptic equations in bounded domains: a survey”, *Publ. Mat.* **60**:1 (2016), 3–26. MR
- [Ros-Oton and Serra 2014] X. Ros-Oton and J. Serra, “The Dirichlet problem for the fractional Laplacian: regularity up to the boundary”, *J. Math. Pures Appl.* (9) **101**:3 (2014), 275–302. MR
- [Ros-Oton and Serra 2016a] X. Ros-Oton and J. Serra, “Boundary regularity for fully nonlinear integro-differential equations”, *Duke Math. J.* **165**:11 (2016), 2079–2154. MR
- [Ros-Oton and Serra 2016b] X. Ros-Oton and J. Serra, “Regularity theory for general stable operators”, *J. Differential Equations* **260**:12 (2016), 8675–8715. MR
- [Ros-Oton and Serra 2017] X. Ros-Oton and J. Serra, “Boundary regularity estimates for nonlocal elliptic equations in C^1 and $C^{1,\alpha}$ domains”, *Ann. Mat. Pura Appl.* (4) **196**:5 (2017), 1637–1668. MR
- [Ros-Oton and Weidner 2024] X. Ros-Oton and M. Weidner, “Optimal regularity for nonlocal elliptic equations and free boundary problems”, preprint, 2024. arXiv 2403.07793
- [Ros-Oton and Weidner 2025] X. Ros-Oton and M. Weidner, “Improvement of flatness for nonlocal free boundary problems”, *J. Eur. Math. Soc. (JEMS)* (online publication October 2025).
- [Ros-Oton et al. 2025] X. Ros-Oton, C. Torres-Latorre, and M. Weidner, “Semiconvexity estimates for nonlinear integro-differential equations”, *Comm. Pure Appl. Math.* **78**:3 (2025), 592–647. MR
- [Servadei and Valdinoci 2014] R. Servadei and E. Valdinoci, “Weak and viscosity solutions of the fractional Laplace equation”, *Publ. Mat.* **58**:1 (2014), 133–154. MR
- [Silvestre 2006] L. Silvestre, “Hölder estimates for solutions of integro-differential equations like the fractional Laplace”, *Indiana Univ. Math. J.* **55**:3 (2006), 1155–1174. MR
- [Silvestre 2007] L. Silvestre, “Regularity of the obstacle problem for a fractional power of the Laplace operator”, *Comm. Pure Appl. Math.* **60**:1 (2007), 67–112. MR
- [Soave and Valdinoci 2019] N. Soave and E. Valdinoci, “Overdetermined problems for the fractional Laplacian in exterior and annular sets”, *J. Anal. Math.* **137**:1 (2019), 101–134. MR
- [Vondraček 2021] Z. Vondraček, “A probabilistic approach to a non-local quadratic form and its connection to the Neumann boundary condition problem”, *Math. Nachr.* **294**:1 (2021), 177–194. MR

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