

Solutions of boundary value problems at resonance with periodic and antiperiodic boundary conditions

Aldo E. Garcia and Jeffrey T. Neugebauer





# Solutions of boundary value problems at resonance with periodic and antiperiodic boundary conditions

# Aldo E. Garcia and Jeffrey T. Neugebauer

(Communicated by Johnny Henderson)

We study the existence of solutions of the second-order boundary value problem at resonance u'' = f(t, u, u') satisfying the boundary conditions u(0) + u(1) = 0, u'(0) - u'(1) = 0, or u(0) - u(1) = 0, u'(0) + u'(1) = 0. We employ a shift method, making a substitution for the nonlinear term in the differential equation so that these problems are no longer at resonance. Existence of solutions of equivalent boundary value problems is obtained, and these solutions give the existence of solutions of the original boundary value problems.

### 1. Introduction

Consider the second-order boundary value problem

$$u'' = f(t, u, u'), \quad t \in (0, 1),$$
 (1-1)

satisfying a combination of antiperiodic and periodic boundary conditions; either

$$u(0) + u(1) = 0, \quad u'(0) - u'(1) = 0.$$
 (1-2)

or

$$u(0) - u(1) = 0, \quad u'(0) + u'(1) = 0.$$
 (1-3)

Here we assume  $f(t, x, y): [0, 1] \times \mathbb{R} \times \mathbb{R} \to \mathbb{R}$  is continuous in each of its variables. Since the boundary value problem u'' = 0, (1-2) has the nontrivial solution  $u(t) = t - \frac{1}{2}$ , the problem (1-1), (1-2) is said to be at resonance. Similarly, since u'' = 0, (1-3) has the nontrivial solution  $u(t) \equiv 1$ , the problem (1-1), (1-3) is also at resonance. Hence, standard methods employing Green's functions cannot be used to show the existence of solutions of these boundary value problems directly. Thus, we consider a shifted boundary value problem so that Green's functions can be employed.

MSC2010: primary 34B15; secondary 34B27.

Keywords: boundary value problem, resonance, shift.

Han [2007] employed a shift argument when studying a three-point boundary value problem

$$x''(t) = f(t, x(t)), \quad t \in (0, 1),$$
  
 $x'(0) = 0, \quad x(\eta) = x(1).$ 

Here it was assumed  $g(t, x) = f(t, x) + \beta^2 x$  and the equivalent boundary value problem

$$x''(t) + \beta^2 x(t) = g(t, x(t)), \quad x'(0) = 0, \ x(\eta) = x(1),$$

was studied using the Krasnosel'skii–Guo fixed point theorem [Krasnosel'skii 1964]. Infante, Pietramala, and Tojo [Infante et al. 2016] also employed a shift argument when studying Neumann boundary value problems at resonance

$$u''(t) + h(t, u(t)) = 0, \quad t \in (0, 1),$$
  
 $u'(0) = u'(1) = 0.$ 

They assumed  $f(t, u) = h(t, u) + \omega^2 u$  and considered the equivalent boundary value problem

$$-u''(t) + \omega^2 u(t) = f(t, u(t)), \quad u'(0) = u'(1) = 0.$$

The Krasnosel'skii-Guo fixed point theorem was also used in their analysis.

More recently, Almansour and Eloe [2015] and Al Mosa and Eloe [2016] studied two-point boundary value problems

$$y''(t) = f(t, y(t)), \quad t \in [0, 1],$$
$$y'(0) = 0, \quad y'(1) = 0,$$

and

$$y''(t) = f(t, y(t), y'(t)), \quad t \in [0, 1],$$
  
$$y(0) = 0, \quad y'(0) = y'(1),$$

using shift arguments and the Krasnosel'skii–Guo fixed point theorem, the Schauder fixed point theorem, the Leray–Schauder nonlinear alternative [Zeidler 1990] in the former, and monotone methods coupled with upper and lower solutions in the latter.

When considering the first boundary value problem, they assumed  $g(t, y) = f(t, y) + \beta^2 y$  and studied the equivalent boundary value problem

$$y''(t) + \beta^2 y(t) = g(t, y(t)), \quad y'(0) = y'(1) = 0,$$

and when considering second, they assumed  $g(t, x, y) = f(t, x, y) + \beta y$  and studied the equivalent boundary value problem

$$y''(t) + \beta y'(t) = g(t, y(t), y'(t)), \quad y(0) = 0, \ y'(0) = y'(1).$$

Here, we make use of two substitutions, one of which has not been used previously in the literature. In Section 2, we study solutions of (1-1), (1-2) by employing the substitution  $g(t, x, y) := f(t, x, y) + \beta y$ . The shifted boundary value problem is no longer at resonance, and so a Green's function can be constructed. An appropriate integral operator is defined and fixed point methods are used to show the existence of solutions. In Section 3, we study solutions of (1-1), (1-3). The substitutions mentioned above do not help because in both cases the shifted boundary value problem is still at resonance. Thus, we use the substitution  $k(t, x, y) = f(t, x, y) + 2\alpha y + (\alpha^2 + \beta^2)x$ . This substitution has not been used in the prior literature. A similar approach to that in Section 2 is then used to show existence of solutions. The construction of the two Green's functions and the shift employed in Section 3 can both lead to more research in this area.

### 2. Solutions of (1-1), (1-2)

Notice that for  $\beta > 0$ ,  $\beta \neq n\pi$ ,  $n \in \mathbb{N}$ , the boundary value problem  $u'' + \beta^2 u = 0$ , (1-2) is at resonance, since  $u(t) = \cos \beta t - ((1 + \cos \beta)/\sin \beta) \sin \beta t$  is a nontrivial solution. If  $\beta = n\pi$ ,  $n \in \mathbb{N}$ , then  $u(t) = \sin \beta t$  is a nontrivial solution of the boundary value problem. Thus the substitution  $g(t, x, y) = f(t, x, y) + \beta^2 x$  cannot be applied.

Let  $\beta > 0$  be a constant and assume  $g(t, x, y) := f(t, x, y) + \beta y$ . We study the shifted differential equation

$$u'' + \beta u' = g(t, u, u'), \quad t \in (0, 1), \tag{2-1}$$

satisfying boundary conditions (1-2). The boundary value problem (2-1), (1-2) is not at resonance, since the unique solution of  $u'' + \beta u' = 0$ , (1-2), is  $u \equiv 0$ . Notice if u(t) is a solution of (2-1), (1-2), then

$$u''(t) = g(t, u(t), u'(t)) - \beta u'(t) = f(t, u(t), u'(t)),$$

implying u is a solution of (1-1), (1-2).

We first construct the Green's function associated with  $u'' + \beta u' = 0$ , (1-2).

**Lemma 2.1.** Let h(t) be a continuous function. Then u(t) is the unique solution of the boundary value problem

$$u'' + \beta u' = h(t), \quad t \in (0, 1),$$
 (2-2)

satisfying boundary conditions (1-2) if and only if

$$u(t) = \int_0^1 G(t, s) h(s) ds,$$

where

$$G(t,s) = \frac{1}{2\beta(1-e^{-\beta})} \begin{cases} 2e^{-\beta(1-s)} - 2e^{-\beta}e^{-\beta(t-s)} + e^{-\beta} - 1, & 0 \le t \le s \le 1, \\ 2e^{-\beta(1-s)} - 2e^{-\beta(t-s)} - e^{-\beta} + 1, & 0 \le s \le t \le 1. \end{cases}$$
(2-3)

*Proof.* Using Laplace transforms, one can show the general solution of (2-2) is given by

$$u(t) = c_1 + c_2 e^{-\beta t} + \frac{1}{\beta} \int_0^t (1 - e^{-\beta(t-s)}) h(s) ds.$$

Since u'(0) - u'(1) = 0, we have

$$-c_2\beta + c_2\beta e^{-\beta} - \int_0^1 e^{-\beta(1-s)} h(s) ds = 0.$$

Solving for  $c_2$  gives

$$c_2 = -\frac{1}{\beta(1 - e^{-\beta})} \int_0^1 e^{-\beta(1 - s)} h(s) ds.$$

The boundary condition u(0) + u(1) = 0 gives

$$c_1 + c_2 + c_1 + c_2 e^{-\beta} + \frac{1}{\beta} \int_0^1 (1 - e^{-\beta(t-s)}) h(s) ds = 0.$$

By substituting  $c_2$  from above, solving for  $c_1$ , and simplifying, we have

$$c_1 = \frac{1}{2\beta(1 - e^{-\beta})} \int_0^1 (-1 + e^{-\beta} + 2e^{-\beta(1-s)}) h(s) ds.$$

Thus

$$u(t) = \frac{1}{2\beta(1 - e^{-\beta})} \int_0^1 (-1 + e^{-\beta} + 2e^{-\beta(1 - s)}) h(s) ds$$
$$- \frac{e^{-\beta t}}{\beta(1 - e^{-\beta})} \int_0^1 e^{-\beta(1 - s)} h(s) ds + \frac{1}{\beta} \int_0^t (1 - e^{-\beta(t - s)}) h(s) ds$$
$$= \int_0^1 G(t, s) h(s) ds,$$

where

$$G(t,s) = \begin{cases} \frac{-1 + e^{-\beta} + 2e^{-\beta(1-s)}}{2\beta(1 - e^{-\beta})} - \frac{e^{-\beta t}e^{-\beta(1-s)}}{\beta(1 - e^{-\beta})}, & 0 \le t \le s \le 1, \\ \frac{-(1 - e^{-\beta}) + 2e^{-\beta(t-s)}}{2\beta(1 - e^{-\beta})} - \frac{e^{-\beta t}e^{-\beta(1-s)}}{\beta(1 - e^{-\beta})} + \frac{1 - e^{-\beta(t-s)}}{\beta}, & 0 \le s \le t \le 1. \end{cases}$$

Simplifying G(t, s) gives (2-3).

The reverse direction of the proof can be shown by direct computation.  $\Box$ 

Notice that

$$\frac{\partial}{\partial t}G(t,s) = \frac{1}{1 - e^{-\beta}} \begin{cases} e^{-\beta}e^{-\beta(t-s)}, & 0 \le t \le s \le 1, \\ e^{-\beta(t-s)}, & 0 \le s \le t \le 1. \end{cases}$$
(2-4)

We point out several properties of the Green's function.

**Lemma 2.2.** G(t, s) satisfies the following properties:

(1) 
$$G \in C([0, 1] \times [0, 1])$$
.

(2) 
$$G(0, s) = -\frac{1}{2\beta} < 0$$
 for all  $s \in [0, 1]$ .

(3) 
$$G(1, s) = \frac{1}{2\beta} > 0$$
 for all  $s \in [0, 1]$ .

(4) 
$$\frac{\partial}{\partial t}G(t,s) > 0$$
 for all  $(t,s) \in [0,1] \times [0,1]$ .

(5) 
$$\max_{t \in [0,1]} |G(t,s)| = \frac{1}{2\beta} \text{ for all } s \in [0,1].$$

(6) 
$$\max_{t \in [0,1]} \frac{\partial}{\partial t} G(t,s) \le \frac{1}{1 - e^{-\beta}} \text{ for all } s \in [0,1].$$

$$(7) \ \max_{t \in [0,1]} \int_0^1 |G(t,s)| \, ds \le \frac{(4+\beta)e^\beta + \beta - 4}{2\beta^2(e^\beta - 1)}.$$

(8) 
$$\max_{t \in [0,1]} \int_0^1 \frac{\partial}{\partial t} G(t,s) \, ds = \frac{1}{\beta}.$$

All of these properties can be shown directly, so a proof is not given. We point out that property (8) is obtained by making all the terms in G(t, s) positive, integrating, and finding an upper bound when  $t \in [0, 1]$ .

We employ Schauder's fixed point theorem in our analysis. Because of the fact that G(t, s) changes sign, many fixed point theorems using cones cannot be used.

**Theorem 2.3** (Schauder fixed point theorem [Hale and Verduyn Lunel 1993]). *If*  $\mathcal{M}$  is a closed, bounded, convex subset of a Banach space  $\mathcal{B}$  and  $T: \mathcal{M} \to \mathcal{M}$  is completely continuous, then T has a fixed point in  $\mathcal{M}$ .

Let  $\mathcal{B} = C^{(1)}[0, 1]$  be the Banach space of functions whose first derivatives are continuous endowed with the norm

$$||u|| = \max\{|u|_0, |u'|_0\},\$$

where  $|u|_0 = \max_{t \in [0,1]} |u(t)|$ . Let M > 0. Define  $\mathcal{M} = \{u \in \mathcal{B} : ||u|| \le M\}$ . Notice that  $\mathcal{M}$  is a closed, bounded, convex subset of  $\mathcal{B}$ .

Define the operator  $T: \mathcal{B} \to \mathcal{B}$  by

$$Tu(t) = \int_0^1 G(t, s) g(s, u(s), u'(s)) ds.$$

Thus if u is a fixed point of T, then u is a solution of (2-1), (1-2). A standard application of the Arzelà–Ascoli theorem gives us that T is completely continuous.

Define

$$\max_{t \in [0,1]} \int_0^1 |G(t,s)| \, ds := \overline{G} \quad \text{and} \quad \max_{t \in [0,1]} \int_0^1 \frac{\partial}{\partial t} G(t,s) \, ds := \overline{G}'.$$

**Theorem 2.4.** Assume f(t, x, y) is continuous in  $[0, 1] \times \mathbb{R} \times \mathbb{R}$  with

$$|f(t, x, y) + \beta y| \le \min \left\{ \frac{M}{\overline{G}}, \frac{M}{\overline{G}'} \right\}$$

for all  $(t, x, y) \in [0, 1] \times [-M, M] \times [-M, M]$ . Then (1-1), (1-2) has a solution  $u^* \in \mathcal{M}$ .

*Proof.* Since  $g(t, x, y) = f(t, x, y) + \beta y$ ,

$$|g(t, x, y)| \le \min \left\{ \frac{M}{\overline{G}}, \frac{M}{\overline{G}'} \right\}$$

for all  $(t, x, y) \in [0, 1] \times [-M, M] \times [-M, M]$ .

Now, for  $u \in \mathcal{M}$ ,

$$|Tu(t)| \le \int_0^1 |G(t,s)| |g(s,u(s),u'(s))| ds \le \frac{M}{\overline{G}} \int_0^1 |G(t,s)| ds = M,$$

$$|(Tu)'(t)| \le \int_0^1 \frac{\partial}{\partial t} G(t,s) |g(s,u(s),u'(s))| ds \le \beta M \int_0^1 \frac{\partial}{\partial t} G(t,s) ds = M.$$

So  $||Tu|| \le M$ , and  $T : \mathcal{M} \to \mathcal{M}$ . Thus T has a fixed point  $u^* \in \mathcal{M}$  which is a solution of (2-1), (1-2). Therefore,  $u^*$  is a solution of (1-1), (1-2).

## Example 2.5. Define

$$f(t, x, y) = \frac{5x^2t^2}{v^2 + 2} - 5y.$$

Let  $\beta = 5$ . Then from Lemma 2.2

$$\min\left\{\frac{M}{\overline{G}}, \frac{M}{\overline{G}'}\right\} \leq \min\left\{\frac{2\beta^2(e^{\beta}-1)}{(4+\beta)e^{\beta}+\beta-4}M, \beta M\right\} = 5M.$$

So

$$|f(t, x, y) + 5y| = \frac{5x^2t^2}{y^2 + 2} \le 5M^2 \le 5M$$

if  $M \le 1$ . So the boundary value problem

$$u'' = \frac{5u^2t^2}{(u')^2 + 2} - 5u', \quad t \in (0, 1),$$
  
$$u(0) + u(1) = 0, \quad u'(0) - u'(1) = 0,$$

has a solution  $u^*$  with  $||u^*|| \le 1$ .

### 3. Solutions of (1-1), (1-3)

For  $\beta > 0$ , the boundary value problem  $u'' + \beta^2 u = 0$ , (1-3) is at resonance, since

$$u(t) = \cos \beta t - \left(\frac{1 - \cos \beta}{\sin \beta}\right) \sin \beta t$$

gives a nontrivial solution. If  $\beta = n\pi$ ,  $n \in \mathbb{N}$ , then  $u(t) = \cos \beta t$  is a nontrivial solution of the boundary value problem. Thus the substitution  $k(t, x, y) = f(t, x, y) + \beta^2 x$  cannot be applied. Also, the boundary value problem  $u'' + \beta u' = 0$ , (1-3) is at resonance, since  $u(t) \equiv 1$  gives a nontrivial solution. This implies the substitution  $k(t, x, y) = f(t, x, y) + \beta^2 y$  cannot be used. Thus, neither substitution used in previous literature can be employed.

Let  $\alpha > 0$ ,  $\beta \in (0, \frac{\pi}{2})$  and define

$$k(t, x, y) = f(t, x, y) + 2\alpha y + (\alpha^2 + \beta^2)x.$$

Here we consider the equivalent boundary value problem

$$u'' + 2\alpha u' + (\alpha^2 + \beta^2)u = k(t, u, u'), \quad t \in (0, 1), \tag{3-1}$$

satisfying boundary conditions (1-3), which is not at resonance, since the unique solution of  $u'' + 2\alpha u' + (\alpha^2 + \beta^2)u = 0$ , (1-3) is  $u \equiv 0$ . If u is a solution of (3-1), (1-3), then u is a solution of (1-1), (1-3).

Again, we construct a corresponding Green's function.

# **Lemma 3.1.** The unique solution of

$$u'' + 2\alpha u' + (\alpha^2 + \beta^2)u = h(t), \quad t \in (0, 1), \tag{3-2}$$

satisfying the boundary conditions (1-3) is given by

$$u(t) = \int_0^1 H(t, s) h(s) ds,$$

where

$$H(t,s) = \frac{1}{2\beta(\beta \sinh \alpha - \alpha \sin \beta)} \Psi(t,s), \tag{3-3}$$

with

$$\Psi(t,s) = \begin{cases} e^{-\alpha(t-s)} \left[ -\beta e^{-\alpha} \sin(\beta(s-t)) + 2\alpha \sin(\beta(1-s)) \sin(\beta t) \\ -\beta \sin(\beta t) \cos(\beta(1-s)) + \beta \cos(\beta t) \sin(\beta(1-s)) \right], & 0 \le t \le s \le 1, \\ e^{-\alpha(t-s)} \left[ \beta e^{\alpha} \sin(\beta(t-s)) + 2\alpha \sin(\beta s) \sin(\beta(1-t)) \\ -\beta \sin(\beta s) \cos(\beta(1-t)) + \beta \cos(\beta s) \sin(\beta(1-t)) \right], & 0 \le s \le t \le 1. \end{cases}$$

*Proof.* If u satisfies (3-2), then, using Laplace transforms,

$$u(t) = e^{-\alpha t} (c_1 \cos(\beta t) + c_2 \sin(\beta t)) + \frac{1}{\beta} \int_0^t (e^{-\alpha(t-s)} \sin(\beta(t-s))) h(s) ds.$$

Solving the system u(0) - u(1) = 0, u'(0) + u'(1) = 0 gives

$$c_{1} = -\frac{1}{2\alpha e^{-\alpha}\sin(\beta) + \beta e^{-2\alpha} - \beta}$$

$$\times \left[ \int_{0}^{1} [e^{-\alpha(1-s)}\sin(\beta(1-s)) - e^{-\alpha(2-s)}\sin(\beta s)]h(s) ds \right],$$

$$c_{2} = -\frac{1}{\beta e^{-\alpha}\sin(\beta)(2\alpha e^{-\alpha}\sin(\beta) + \beta e^{-2\alpha} - \beta)}$$

$$\times \left[ \int_{0}^{1} \left[ \beta e^{-\alpha(3-s)}[\cos(\beta)\sin(\beta s) + \sin(\beta(1-s))] - e^{-\alpha(2-s)}[\beta\cos(\beta)\sin(\beta(1-s))] + \beta\sin(\beta s) - 2\alpha\sin(\beta)\sin(\beta(1-s)) \right] \right] h(s) ds \right]$$

The Green's function given in (3-3) can then be obtained.

Notice

$$\frac{\partial}{\partial t}H(t,s) = \frac{1}{2\beta(\beta \sinh \alpha - \alpha \sin \beta)}\Phi(t,s), \tag{3-4}$$

where

where 
$$\Phi(t,s) = \begin{cases} e^{-\alpha(t-s)} \left[ e^{-\alpha} \beta^2 \cos(\beta(s-t)) + 2\alpha\beta \sin(\beta(1-s)) \cos(\beta t) \\ -\beta^2 \sin(\beta(1-s)) \sin(\beta t) - \beta^2 \cos(\beta(1-s)) \cos(\beta t) \right] \\ -\alpha e^{-\alpha(t-s)} \left[ 2\alpha \sin(\beta(1-s)) \sin(\beta t) - e^{-\alpha} \beta \sin(\beta(s-t)) \\ +\beta \sin(\beta(1-s)) \cos(\beta t) - \beta \cos(\beta(1-s)) \sin(\beta t) \right], & 0 \le t \le s \le 1, \end{cases}$$

$$e^{-\alpha(t-s)} \left[ e^{\alpha} \beta^2 \cos(\beta(t-s)) - 2\alpha\beta \sin(\beta s) \cos(\beta(1-t)) \\ -\beta^2 \sin(\beta s) \sin(\beta(1-t)) - \beta^2 \cos(\beta s) \cos(\beta(1-t)) \right] \\ -\alpha e^{-\alpha(t-s)} \left[ 2\alpha \sin(\beta s) \sin(\beta(1-t)) + e^{\alpha} \beta \sin(\beta(t-s)) \\ -\beta \sin(\beta s) \cos(\beta(1-t)) + \beta \cos(\beta s) \sin(\beta(1-t)) \right], & 0 \le s \le t \le 1. \end{cases}$$

We point out several properties of the Green's function.

**Lemma 3.2.** H(t, s) satisfies the following properties:

(1)  $H \in C([0, 1] \times [0, 1])$ .

(2) 
$$H(0,s) = H(1,s) = \frac{e^{\alpha s} (\beta \sin(\beta(1-s)) - e^{-\alpha} \beta \sin(\beta s))}{2\beta(\beta \sinh \alpha - \alpha \sin \beta)}$$
 for all  $s \in [0, 1]$ .

$$(3) \max_{t \in [0,1]} |H(t,s)| \leq \frac{\beta e^{\alpha} + 2\alpha + 2\beta}{2\beta (\beta \sinh \alpha - \alpha \sin \beta)} \text{ for all } s \in [0,1].$$

$$(4) \max_{t \in [0,1]} \left| \frac{\partial}{\partial t} H(t,s) \right| \leq \frac{\alpha \beta e^{\alpha} + 2\alpha^2 + 2\beta^2 + 2\alpha \beta + \beta^2 e^{\alpha}}{2\beta (\beta \sinh \alpha - \alpha \sin \beta)} \text{ for all } s \in [0,1].$$

$$(5) \max_{t \in [0,1]} \int_0^1 |H(t,s)| \, ds \le \frac{\beta + \beta \sinh \alpha}{(\alpha^2 + \beta^2)(\beta \sinh \alpha - \alpha \sin \beta)}.$$

$$(6) \max_{t \in [0,1]} \int_0^1 \left| \frac{\partial}{\partial t} H(t,s) \right| ds \le \frac{\alpha^2 e^{\alpha} + \alpha^2 + \beta^2 + \beta^2 e^{\alpha} + \alpha e^{\alpha} + \alpha \beta + 3\beta}{(\alpha^2 + \beta^2)(\beta \sinh \alpha - \alpha \sin \beta)}.$$

Again, a proof is not given, since all these properties can be verified directly. Properties (5) and (6) are obtained by making all the terms in H(t, s) and  $(\partial/\partial t)H(t, s)$ , respectfully, positive, integrating, and finding an upper bound when  $t \in [0, 1]$ .

Define the operator  $T: \mathcal{B} \to \mathcal{B}$  by

$$Tu(t) = \int_0^1 H(t, s) \, k(s, u(s), u'(s)) \, ds.$$

Thus if u is a fixed point of T, then u is a solution of (3-1), (1-3). A standard application of the Arzelà–Ascoli theorem gives us that T is completely continuous.

Define

$$\max_{t \in [0,1]} \int_0^1 |H(t,s)| \, ds := \overline{H} \quad \text{and} \quad \max_{t \in [0,1]} \int_0^1 \left| \frac{\partial}{\partial t} H(t,s) \right| \, ds := \overline{H}'.$$

**Theorem 3.3.** Assume f(t, x, y) is continuous in  $[0, 1] \times \mathbb{R} \times \mathbb{R}$  with

$$\left| f(t, x, y) + 2\alpha y + (\alpha^2 + \beta^2) x \right| \le \min \left\{ \frac{M}{\overline{H}}, \frac{M}{\overline{H'}} \right\}$$

for all  $(t, x, y) \in [0, 1] \times [-M, M] \times [-M, M]$ . Then (1-1), (1-3) has a solution  $u^* \in \mathcal{M}$ .

The proof is similar to the proof of Theorem 2.4 and is therefore omitted.

### References

[Al Mosa and Eloe 2016] S. Al Mosa and P. Eloe, "Upper and lower solution method for boundary value problems at resonance", *Electron. J. Qual. Theory Differ. Equ.* **2016** (2016), art. id. 40. MR Zbl

[Almansour and Eloe 2015] A. Almansour and P. Eloe, "Fixed points and solutions of boundary value problems at resonance", *Ann. Polon. Math.* **115**:3 (2015), 263–274. MR Zbl

[Hale and Verduyn Lunel 1993] J. K. Hale and S. M. Verduyn Lunel, *Introduction to functional-differential equations*, Applied Mathematical Sciences **99**, Springer, 1993. MR Zbl

[Han 2007] X. Han, "Positive solutions for a three-point boundary value problem at resonance", J. Math. Anal. Appl. 336:1 (2007), 556–568. MR Zbl

[Infante et al. 2016] G. Infante, P. Pietramala, and F. A. F. Tojo, "Non-trivial solutions of local and non-local Neumann boundary-value problems", *Proc. Roy. Soc. Edinburgh Sect. A* **146**:2 (2016), 337–369. MR Zbl

[Krasnosel'skii 1964] M. A. Krasnosel'skii, *Topological methods in the theory of nonlinear integral equations*, Int. Series of Monographs in Pure Appl. Math. **45**, Macmillan, New York, 1964. MR Zbl

[Zeidler 1990] E. Zeidler, Nonlinear functional analysis and its applications, II/A: Linear monotone operators, Springer, 1990. MR Zbl

Received: 2018-01-24 Revised: 2018-02-13 Accepted: 2018-02-14

aldo\_garciaguinto@mymail.eku.edu

Department of Mathematics and Statistics,

Eastern Kentucky University, Richmond, KY, United States

Eastern Kentucky University, Richmond, KY, United States



#### INVOLVE YOUR STUDENTS IN RESEARCH

*Involve* showcases and encourages high-quality mathematical research involving students from all academic levels. The editorial board consists of mathematical scientists committed to nurturing student participation in research. Bridging the gap between the extremes of purely undergraduate research journals and mainstream research journals, *Involve* provides a venue to mathematicians wishing to encourage the creative involvement of students.

#### MANAGING EDITOR

Kenneth S. Berenhaut Wake Forest University, USA

### **BOARD OF EDITORS**

Colin Adams	Williams College, USA	Suzanne Lenhart	University of Tennessee, USA	
John V. Baxley	Wake Forest University, NC, USA	Chi-Kwong Li	College of William and Mary, USA	
Arthur T. Benjamin	Harvey Mudd College, USA	Robert B. Lund	Clemson University, USA	
Martin Bohner	Missouri U of Science and Technology,	USA Gaven J. Martin	Massey University, New Zealand	
Nigel Boston	University of Wisconsin, USA	Mary Meyer	Colorado State University, USA	
Amarjit S. Budhiraja	U of North Carolina, Chapel Hill, USA	Emil Minchev	Ruse, Bulgaria	
Pietro Cerone	La Trobe University, Australia	Frank Morgan	Williams College, USA	
Scott Chapman	Sam Houston State University, USA	Mohammad Sal Moslehian	Ferdowsi University of Mashhad, Iran	
Joshua N. Cooper	University of South Carolina, USA	Zuhair Nashed	University of Central Florida, USA	
Jem N. Corcoran	University of Colorado, USA	Ken Ono	Emory University, USA	
Toka Diagana	Howard University, USA	Timothy E. O'Brien	Loyola University Chicago, USA	
Michael Dorff	Brigham Young University, USA	Joseph O'Rourke	Smith College, USA	
Sever S. Dragomir	Victoria University, Australia	Yuval Peres	Microsoft Research, USA	
Behrouz Emamizadeh	The Petroleum Institute, UAE	YF. S. Pétermann	Université de Genève, Switzerland	
Joel Foisy	SUNY Potsdam, USA	Robert J. Plemmons	Wake Forest University, USA	
Errin W. Fulp	Wake Forest University, USA	Carl B. Pomerance	Dartmouth College, USA	
Joseph Gallian	University of Minnesota Duluth, USA	Vadim Ponomarenko	San Diego State University, USA	
Stephan R. Garcia	Pomona College, USA	Bjorn Poonen	UC Berkeley, USA	
Anant Godbole	East Tennessee State University, USA	James Propp	U Mass Lowell, USA	
Ron Gould	Emory University, USA	Józeph H. Przytycki	George Washington University, USA	
Andrew Granville	Université Montréal, Canada	Richard Rebarber	University of Nebraska, USA	
Jerrold Griggs	University of South Carolina, USA	Robert W. Robinson	University of Georgia, USA	
Sat Gupta	U of North Carolina, Greensboro, USA	Filip Saidak	U of North Carolina, Greensboro, USA	
Jim Haglund	University of Pennsylvania, USA	James A. Sellers	Penn State University, USA	
Johnny Henderson	Baylor University, USA	Andrew J. Sterge	Honorary Editor	
Jim Hoste	Pitzer College, USA	Ann Trenk	Wellesley College, USA	
Natalia Hritonenko	Prairie View A&M University, USA	Ravi Vakil	Stanford University, USA	
Glenn H. Hurlbert	Arizona State University, USA	Antonia Vecchio	Consiglio Nazionale delle Ricerche, Italy	
Charles R. Johnson	College of William and Mary, USA	Ram U. Verma	University of Toledo, USA	
K. B. Kulasekera	Clemson University, USA	John C. Wierman	Johns Hopkins University, USA	
Gerry Ladas	University of Rhode Island, USA	Michael E. Zieve	University of Michigan, USA	
PRODUCTION				
Silvio Levy, Scientific Editor				

Silvio Levy, Scientific Editor

Cover: Alex Scorpan

See inside back cover or msp.org/involve for submission instructions. The subscription price for 2019 is US \$/year for the electronic version, and \$/year (+\$, if shipping outside the US) for print and electronic. Subscriptions, requests for back issues and changes of subscriber address should be sent to MSP.

Involve (ISSN 1944-4184 electronic, 1944-4176 printed) at Mathematical Sciences Publishers, 798 Evans Hall #3840, c/o University of California, Berkeley, CA 94720-3840, is published continuously online. Periodical rate postage paid at Berkeley, CA 94704, and additional mailing offices.

Involve peer review and production are managed by EditFLOW® from Mathematical Sciences Publishers.



mathematical sciences publishers

nonprofit scientific publishing

http://msp.org/

© 2019 Mathematical Sciences Publishers

Optimal transportation with constant constraint	1	
WYATT BOYER, BRYAN BROWN, ALYSSA LOVING AND SARAH TAMMEN		
Fair choice sequences	13	
William J. Keith and Sean Grindatti		
Intersecting geodesics and centrality in graphs	31	
Emily Carter, Bryan Ek, Danielle Gonzalez, Rigoberto Flórez		
and Darren A. Narayan		
The length spectrum of the sub-Riemannian three-sphere	45	
DAVID KLAPHECK AND MICHAEL VANVALKENBURGH		
Statistics for fixed points of the self-power map	63	
MATTHEW FRIEDRICHSEN AND JOSHUA HOLDEN		
Analytical solution of a one-dimensional thermistor problem with Robin boundary	79	
condition		
VOLODYMYR HRYNKIV AND ALICE TURCHANINOVA		
On the covering number of $S_{14}$	89	
RYAN OPPENHEIM AND ERIC SWARTZ		
Upper and lower bounds on the speed of a one-dimensional excited random walk		
Erin Madden, Brian Kidd, Owen Levin, Jonathon Peterson,		
JACOB SMITH AND KEVIN M. STANGL		
Classifying linear operators over the octonions	117	
ALEX PUTNAM AND TEVIAN DRAY		
Spectrum of the Kohn Laplacian on the Rossi sphere	125	
TAWFIK ABBAS, MADELYNE M. BROWN, RAVIKUMAR RAMASAMI AND		
Yunus E. Zeytuncu		
On the complexity of detecting positive eigenvectors of nonlinear cone maps	141	
BAS LEMMENS AND LEWIS WHITE		
Antiderivatives and linear differential equations using matrices	151	
YOTSANAN MEEMARK AND SONGPON SRIWONGSA		
Patterns in colored circular permutations	157	
DANIEL GRAY, CHARLES LANNING AND HUA WANG		
Solutions of boundary value problems at resonance with periodic and antiperiodic	171	
boundary conditions		
ALDO E. GARCIA AND JEFFREY T. NEUGEBAUER		