Journal of Mechanics of Materials and Structures

THE INVERSE DETERMINATION OF THE VOLUME FRACTION OF FIBERS IN A UNIDIRECTIONALLY REINFORCED COMPOSITE FOR A GIVEN EFFECTIVE THERMAL CONDUCTIVITY

Magdalena Mierzwiczak and Jan Adam Kołodziej

Volume 7, No. 3

and the second second

March 2012





THE INVERSE DETERMINATION OF THE VOLUME FRACTION OF FIBERS IN A UNIDIRECTIONALLY REINFORCED COMPOSITE FOR A GIVEN EFFECTIVE THERMAL CONDUCTIVITY

MAGDALENA MIERZWICZAK AND JAN ADAM KOŁODZIEJ

We consider the problem of determining the volume fraction of fibers in a unidirectionally reinforced composite in order to provide the appropriate effective thermal conductivity. The problem is formulated in such a way as to be treated as an inverse heat transfer problem. The thermal conductivities of the constituents (fibers and matrix) and fiber arrangement are known. The calculations are carried out for a perfect thermal contact between the fibers and matrix.

1. Introduction

In the literature the following problems are considered to be classical inverse heat conduction problems:

- determination of heat sources [Yan et al. 2008],
- determination of the heat transfer coefficient [Hon and Wei 2004],
- the Cauchy problem [Marin 2005], and
- determination of the temperature dependent thermal conductivity [Chantasiriwan 2002].

These problems usually apply to homogeneous media. In the case of composite materials (nonhomogeneous media) other practically important issues might have to be considered. One of them is the inverse problem of determination of the volume fraction of constituents in order to obtain the appropriate effective thermal conductivity. Let's consider a unidirectional fibrous composite with regular arrangement of fibers (Figure 1, left). If the thermal conductivity coefficients of constituents and their volume fractions are known then the composite can be treated as a homogeneous region for which effective thermal conductivity coefficient is determined for a regular arrangement of fibers for given thermal conductivity of constituents and volume fraction of fibers (the direct problem). The method of determination is usually based on the solution of the heat transfer equation at a microstructure level in repeated elements of an array [Han and Cosner 1981]. But to our knowledge no paper has considered the inverse problem of determination of the volume fraction of fibers for a given effective thermal conductivity. Here we propose an analytic-numerical algorithm for determination of the volume fraction of fibers in order to obtain a given value of the transverse effective thermal conductivity λ_z (the inverse problem).

Keywords: effective thermal conductivity, inverse heat transfer problem, boundary collocation method, Newton's method.

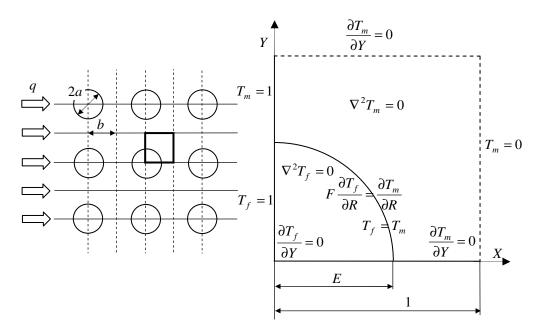


Figure 1. A unidirectional reinforced fibrous composite with fibers arranged in a square array. Left: general view, with marked repeated element. Right: formulation of the boundary value problem in the repeated element for nondimensional variables.

2. Direct problem: determination of the effective thermal conductivity coefficient of the composite material

Consider a unidirectional composite with fibers arranged in a matrix in a regular, square array with imperfect thermal contact between the fiber and matrix (Figure 1, left), where *a* is radius of the fibers, 2*b* is the distance between neighboring fibers, E = a/b, and $\varphi = \pi E^2/4$ is the volume fraction of the fibers. The ratio of the thermal conductivity of fibers λ_f to matrix λ_m is denoted as $F = \lambda_f/\lambda_m$, R = r/b is the dimensionless radius, X = x/b and Y = y/b are the dimensionless Cartesian coordinates, $T = (\hat{T} - \hat{T}_R)/(\hat{T}_R - \hat{T}_L)$ is the dimensionless temperature field, and \hat{T}_R and \hat{T}_L are the temperatures on the left and right boundaries of the repeated element, respectively. In order to solve the boundary value problem in the repeated element of the composite (Figure 1) the boundary collocation method is used [Kołodziej and Zieliński 2009]. The general solution of the Laplace equation in a polar coordinate system has the form

$$T = A_0 + A_1\theta + A_2\theta \ln R + A_3 \ln R + \sum_{k=1}^{\infty} (B_k R^k + C_k R^{-k}) \cos(k\theta) + (D_k R^k + E_k R^{-k}) \sin(k\theta), \quad (2-1)$$

where A_0, A_1, \ldots, E_k are integral constants.

Given the repetitive element of the square array $\Omega = \Omega_m + \Omega_f$ in the region of the fiber Ω_f and the matrix Ω_m , a solution is predicted with the form of (2-1). Some of the constants must be determined

strictly by the conditions at the bottom and on the left side of the repeated element:

$$\frac{\partial T_f}{\partial \theta} = \frac{\partial T_m}{\partial \theta} = 0 \quad \text{for} \quad \theta = 0, \qquad T_f = T_m = 1 \quad \text{for} \quad \theta = \frac{\pi}{2}, \tag{2-2}$$

and the contact conditions of the fiber-matrix:

$$F\frac{\partial T_f}{\partial R} = \frac{\partial T_m}{\partial R}$$
 for $R = E$, $T_f = T_m$ for $R = E$. (2-3)

After determining the constants from the boundary conditions (2-2) and from the contact conditions (2-3) of fiber-matrix, marking the remaining constants as w_k , and cutting off an infinite number of test functions to N expressions, we obtain a solution for the temperature field of the fiber and matrix:

$$T_f = 1 + \sum_{k=1}^{N} w_k R^{(2k-1)} \cos((2k-1)\theta), \qquad (2-4)$$

$$T_m = 1 + \sum_{k=1}^{N} \frac{w_k}{2} \bigg[(1+F)R^{(2k-1)} + (1-F)\frac{E^{2(2k-1)}}{R^{(2k-1)}} \bigg] \cos((2k-1)\theta).$$
(2-5)

The constants w_k are determined by fulfillment of the condition on the collocation points on the upper Γ_2 and on the right Γ_1 edges of the concerned region (Figure 2):

$$T_m = 0$$
 for $X = 1$, (2-6)

$$\frac{\partial T_m}{\partial Y} = 0 \quad \text{for} \quad Y = 1.$$
 (2-7)

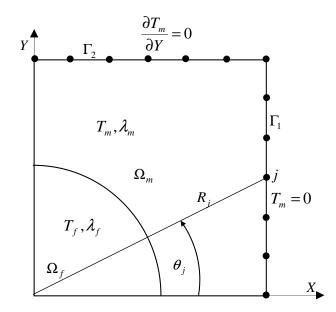


Figure 2. The collocation points at the upper and right boundaries of the matrix in a repeated element in which the boundary conditions are collocated.

The condition (2-7) can be written for polar coordinates:

$$\frac{\partial T_m}{\partial Y} = \frac{\partial T_m}{\partial R}\sin(\theta) + \frac{1}{R}\frac{\partial T_m}{\partial \theta}\cos(\theta).$$
(2-8)

Choosing N_1 points on the right boundary Γ_1 and N_2 points on the upper boundary Γ_2 and collocating the conditions (2-6) and (2-8) we obtain the system of $N_1 + N_2$ linear equations with N unknown coefficients $w_k, k = 1, ..., N$:

$$Aw = b, \tag{2-9}$$

$$\sum_{k=1}^{N} w_k \bigg[(1+F) R_j^{(2k-1)} + (1-F) \frac{E^{2(2k-1)}}{R_j^{(2k-1)}} \bigg] \cos((2k-1)\theta_j) = -2, \quad (R_j, \theta_j) \in \Gamma_1, \quad j = 1, \dots, N_1,$$

$$\sum_{k=1}^{N} w_k (2k-1) \bigg[(1+F) R_j^{(2k-1)} \sin((2k-1)\theta_j) + (1-F) \frac{E^{2(2k-1)}}{R_j^{2k}} \cos(2k\theta_j) \bigg] = 0,$$

$$(R_j, \theta_j) \in \Gamma_2, \quad j = N_1 + 1, \dots, N_1 + N_2.$$

The constants w_k obtained by the Gaussian elimination method provide estimates of the value of the global heat flux through the unit region of the considered element:

$$q = \frac{1}{b} \left[-\lambda_f \int_0^a \frac{\partial \hat{T}_f}{\partial x} \Big|_{x=0} dy + \lambda_m \int_a^b \frac{\partial \hat{T}_m}{\partial x} \Big|_{x=0} dy \right].$$
(2-10)

The transverse effective thermal conductivity is defined by the formula

$$\lambda_z = \frac{qb}{\Delta \hat{T}},\tag{2-11}$$

where *b* is the distance between the isothermal boundaries and $\Delta \hat{T} = \hat{T}_L - \hat{T}_R$ is the difference of the temperatures at the isothermal edges. After taking into consideration in formula (2-11) the definition of the nondimensional temperature and coordinates, the value of the effective thermal conductivity in relation to the thermal conductivity of the matrix can be calculated from the relationship:

$$\frac{\lambda_z}{\lambda_m} = -F \int_0^E \frac{1}{R} \frac{\partial T_f}{\partial \theta} \bigg|_{\theta = \frac{\pi}{2}} dR + \int_E^1 \frac{1}{R} \frac{\partial T_m}{\partial \theta} \bigg|_{\theta = \frac{\pi}{2}} dR, \qquad (2-12)$$

or

$$\frac{\lambda_z}{\lambda_m} = \sum_{k=1}^N \frac{w_k}{2} (-1)^k \left[(F+1) + (F-1)E^{2(2k-1)} \right].$$
(2-13)

3. The results of the numerical experiment

The results of the calculations of the effective thermal conductivity of the fibrous composite are shown in Figure 3. The value of the effective thermal conductivity in relation to the thermal conductivity of the matrix $\lambda_z/\lambda_m = (\lambda_z/\lambda_m)(\varphi)|_F$ is presented as a function of the volume fraction of fibers φ for the desired value of thermal conductivity ratio *F* of the fiber λ_f to the matrix λ_m , with $F \in \{0.5, 2, 10, 20\}$. In order to compare the results, for a flat layer composite consisting of two components with different

232

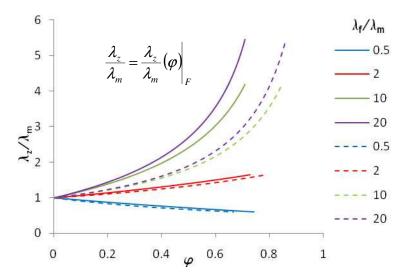


Figure 3. The effective thermal conductivity as a function of the volume fraction of the fibers in the matrix for different values of the ratio of thermal conductivity of the fibers to the matrix.

coefficients of thermal conductivity we calculate the effective thermal conductivity coefficient in relation to the thermal conductivity of the matrix for an ideal contact of the components from the formula

$$\frac{\lambda_z}{\lambda_m} = \left((1 - \varphi) + \frac{\varphi}{F} \right)^{-1}.$$

The comparative results for a flat layer of composite are shown in Figure 3 by the dotted line. The value of the effective thermal conductivity λ_z depends not only on the constants characterizing the composite, F and E, but also on the coefficients w_k involved in fulfilling the boundary conditions at the $N_1 + N_2$ collocation points. Table 1 shows the influence of the number of collocation points on the maximum error fulfilling the collocation boundary conditions calculated at the control points (between the collocation points). The analysis of the results shows that increase in the number of collocation points doesn't lead to an increase in the accuracy of the calculations. Increasing the number of collocation points entails a rise in the dimension of the matrix system of equations. In all four examples presented in Table 1, the smallest maximum error satisfying the boundary conditions was obtained for 7 collocation points on the right edge and for 6 points on the upper edge of the region considered.

4. Inverse problem: determination of the volume fraction of fibers in a composite for a given effective thermal conductivity

At times, when designing composites with specific properties of the fibers and matrix we must estimate the fraction of the volume of fibers to obtain effective thermal conductivity values. Assuming that E = a/b is unknown, we use the known value of the effective thermal conductivity in relation to the thermal conductivity of the matrix λ_z/λ_m . From the collocation of the boundary conditions in the $N_1 + N_2$ points on the right and upper edges of the considered region and from condition (2-11) we obtain a system of $N_1 + N_2 + 1$ nonlinear equations with the N + 1 unknowns w_k and E:

	E = 0.5, F = 10			E	E = 0.9, F = 10			
N_1	N_2	λ_z/λ_m	$\delta_{\max} _{T_m=0}$	$\delta_{\max} _{\partial T_m/\partial Y=0}$	λ_z/λ_m	$\delta_{\max} _{T_m=0}$	$\delta_{\max}\Big _{\partial T_m/\partial Y=0}$	
5	4	1.3829	$1.13 imes 10^{-4}$	1.33×10^{-3}	3.3401	0.004232	0.004027	
6	5	1.3829	2.42×10^{-5}	$3.83 imes 10^{-4}$	3.3408	2.61×10^{-3}	1.12×10^{-2}	
7	6	1.3829	6.20×10^{-7}	8.38×10^{-5}	3.3405	5.48×10^{-4}	1.63×10^{-2}	
8	7	1.3829	9.35×10^{-6}	$3.85 imes 10^{-4}$	3.3413	1.55×10^{-3}	8.48×10^{-2}	
9	8	1.3829	1.97×10^{-5}	$8.77 imes 10^{-4}$	3.3396	4.25×10^{-3}	2.02×10^{-1}	
10	9	1.3829	$4.96 imes 10^{-5}$	2.47×10^{-3}	3.3431	1.14×10^{-2}	0.578786	
11	10	1.3829	$7.29 imes 10^{-5}$	3.98×10^{-3}	3.3372	1.72×10^{-2}	0.943426	
12	11	1.3829	2.26×10^{-4}	0.013628	3.3502	0.053729	3.235684	
13	12	1.3829	$1.87 imes 10^{-4}$	0.012135	3.3328	4.46×10^{-2}	2.895936	
14	13	1.3831	1.98×10^{-3}	0.139725	3.4128	0.463741	32.72905	
15	14	1.3829	4.42×10^{-4}	0.033299	3.3242	$1.06 imes 10^{-1}$	7.967415	
	E = 0.5, F = 0.5			1	E = 0.9, F = 0.5			
N_1	N_2	λ_z/λ_m	$\delta_{\max} _{T_m=0}$	$\delta_{\max}\big _{\partial T_m/\partial Y=0}$	λ_z/λ_m	$\delta_{\max} _{T_m=0}$	$\delta_{\max} \Big _{\partial T_m / \partial Y = 0}$	
5	4	0.87713	$3.73 imes 10^{-5}$	0.000484	0.64835	4.10×10^{-4}	0.010038	
6	5	0.87714	$8.21 imes 10^{-6}$	$1.46 imes 10^{-4}$	0.64849	$2.28 imes 10^{-4}$	6.92×10^{-3}	
7	6	0.87714	$1.14 imes 10^{-7}$	$2.17 imes 10^{-5}$	0.64844	6.26×10^{-5}	2.18×10^{-3}	
8	7	0.87714	$2.96 imes 10^{-6}$	1.21×10^{-4}	0.64845	4.66×10^{-6}	2.22×10^{-4}	
9	8	0.87714	6.31×10^{-6}	$2.81 imes 10^{-4}$	0.64846	9.55×10^{-5}	4.29×10^{-3}	
10	9	0.87713	1.59×10^{-5}	0.000792	0.64842	2.96×10^{-4}	0.014795	
11	10	0.87714	2.34×10^{-5}	0.001275	0.64849	$4.58 imes 10^{-4}$	0.025030	
12	11	0.87713	$7.26 imes 10^{-5}$	0.00437	0.64834	1.44×10^{-3}	0.086817	
13	12	0.87714	6.00×10^{-5}	0.003891	0.64855	1.20×10^{-3}	0.077723	
14	13	0.87708	6.35×10^{-4}	0.044798	0.64758	0.012614	0.890239	
15	14	0.87715	1.42×10^{-4}	0.010678	0.64865	2.84×10^{-3}	0.213776	

Table 1. The impact of the number of collocation points on the value of the effective thermal conductivity of the composite and the maximum error of fulfilling the boundary conditions at control points.

$$f_{j} = 1 + \sum_{k=1}^{N} \frac{w_{k}}{2} \left[(1+F)R_{j}^{(2k-1)} + (1-F)\frac{E^{2(2k-1)}}{R_{j}^{(2k-1)}} + \right] \cos\left((2k-1)\theta_{j}\right) = 0,$$

$$(R_{j}, \theta_{j}) \in \Gamma_{1} \quad j = 1, \dots, N_{1},$$

$$f_{j} = \sum_{k=1}^{N} w_{k}(2k-1) \left[(1+F)R_{j}^{(2k-1)} \sin\left((2k-1)\theta_{j}\right) + (1-F)\frac{E^{2(2k-1)}}{R_{j}^{(2k-1)}} \cos(2k\theta_{j}) \right] = 0,$$

$$(R_{j}, \theta_{j}) \in \Gamma_{2} \quad j = N_{1} + 1, \dots, N_{1} + N_{2},$$

$$f_{N_1+N_2+1} = \sum_{k=1}^{N} \frac{w_k}{2} (-1)^k \left[(1+F) + (F-1)E^{2(2k-1)} \right] - \frac{\lambda_z}{\lambda_m} = 0.$$
(4-1)

The nonlinear system of $N_1 + N_2 + 1$ equations f with N + 1 unknowns $W = [w_1, \ldots, w_N, E]^T$ is solved by using Newton's iterative method:

$$\begin{bmatrix} w_{1} \\ \vdots \\ w_{N} \\ E \end{bmatrix}^{(i+1)} = \begin{bmatrix} w_{1} \\ \vdots \\ w_{N} \\ E \end{bmatrix}^{(i)} - \begin{bmatrix} f_{1}^{(i)} \\ \vdots \\ f_{N_{1}+N_{2}}^{(i)} \\ f_{N_{1}+N_{2}+1}^{(i)} \end{bmatrix} \begin{bmatrix} J_{1,1}^{(i)} & \cdots & J_{1,N+1}^{(i)} \\ \vdots & \ddots & \vdots \\ J_{N_{1}+N_{2},1}^{(i)} & \cdots & \vdots \\ J_{N_{1}+N_{2}+1,1}^{(i)} & \cdots & J_{N_{1}+N_{2}+1,N+1}^{(i)} \end{bmatrix}^{-1},$$

$$W^{(i+1)} = W^{(i)} - f(W^{(i)})J(W^{(i)})^{-1} \rightarrow W^{(i+1)} = W^{(i)} - Y(W^{(i)}),$$

$$Y(W^{(i)}) = f(W^{(i)})J(W^{(i)})^{-1} \rightarrow J(W^{(i)})^{-1}Y(W^{(i)}) = f(W^{(i)}).$$
(4-2)

The functions f_i are described by (4-1), while the Jacobi elements have the following form:

$$\begin{split} J_{j,k} &= \frac{\partial f_j}{\partial w_k} = \frac{1}{2} \bigg[(1+F) R_j^{(2k-1)} + (1-F) \frac{E^{2(2k-1)}}{R_j^{(2k-1)}} + \bigg] \cos((2k-1)\theta_j), \\ & j = 1, \dots, N_1, \quad k = 1, \dots, N, \end{split} \\ J_{j,N+1} &= \frac{\partial f_j}{\partial E} = \sum_{k=1}^N w_k (2k-1)(1-F) \frac{E^{(4k-3)}}{R_j^{(2k-1)}} \cos((2k-1)\theta_j), \quad j = 1, \dots, N_1, \\ J_{j,k} &= \frac{\partial f_j}{\partial w_k} = \frac{(2k-1)}{2} \bigg[(1+F) R_j^{2(k-1)} \sin(2(k-1)\theta_j) + (1-F) \frac{E^{2(2k-1)}}{R_j^{(2k)}} \sin(2k\theta_j) \bigg], \\ & j = N_1 + 1, \dots, N_1 + N_2, \quad k = 1, \dots, N, \end{split}$$

$$J_{j,N+1} &= \frac{\partial f_j}{\partial E} = \sum_{k=1}^N w_k (2k-1)^2 \bigg[(1-F) \frac{E^{(4k-3)}}{R_j^{(2k)}} \sin(2k\theta_j) \bigg], \quad j = N_1 + 1, \dots, N_1 + N_2, \\ J_{j,k} &= \frac{\partial f_j}{\partial w_k} = \frac{(-1)^k}{2} \bigg[(1+F) + (F-1)E^{2(2k-1)} \bigg], \quad j = N_1 + N_2 + 1, \quad k = 1, \dots, N, \end{aligned}$$

To start the Newton's iteration we need to know $W^{(0)} = [w_1^{(0)}, \ldots, w_N^{(0)}, E^{(0)}]^T$ as an initial condition. As an initial value of the constants $w_k^{(0)}$, $k = 1, \ldots, n$, the solution of the linear problem for E = 0.1 has been adopted. The condition for the end of the iteration was adopted at $\delta_{\text{Newton}} = ||W^{(i+1)} - W^{(i)}||_{\text{max}} \le 10^{-7}$, where $|| ||_{\text{max}}$ means the maximum norm.

5. The results of the numerical experiment

The results of the iterative calculation of the volume fraction of fibers for a composite are shown in Figure 4. The value of the volume fraction of fibers in a composite ϕ is presented as a function $\varphi = \varphi(\lambda_z/\lambda_m)|_F$ of the effective thermal conductivity in relation to the thermal conductivity of the matrix

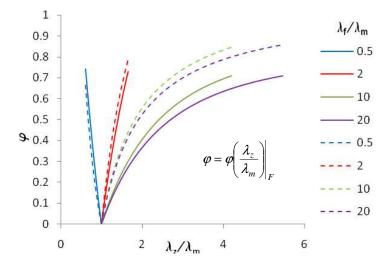


Figure 4. The volume fraction of fibers in the matrix as a function of the effective thermal conductivity of the composite for different relative values of thermal conductivity of the fiber and matrix.

 λ_z/λ_m for the assumed value of the thermal conductivity ratio of the fiber to the matrix, $F = \lambda_f/\lambda_m \in \{0.5, 2, 10, 20\}$. In order to compare the results for a flat composite layer consisting of two components with different thermal conductivities, the volume fraction of the fibers for a known value of the effective thermal conductivity for an ideal contact with the components can be calculated from

$$\varphi = \left(\left(\frac{\lambda_z}{\lambda_m} \right)^{-1} - 1 \right) \frac{F}{1 - F}.$$

			$\lambda_z/\lambda_m =$	= 1.4, F = 10)	λ	$\lambda_z/\lambda_m = 3$	3.35, F = 10	
N_1	N_2	φ	Ε	$\delta_{\max} _{T_m=0}$	$\delta_{\max}\Big _{\partial T_m/\partial Y=0}$	arphi	Ε	$\delta_{\max}\big _{T_m=0}$	$\delta_{\max}\Big _{\partial T_m/\partial Y=0}$
5	4	0.1904	0.4924	$5.82 \cdot 10^{-5}$	0.100879	0.5013	0.7989	$1.74 \cdot 10^{-3}$	2.688120
6	5	0.2036	0.5092	$4.99 \cdot 10^{-5}$	0.000256	0.6362	0.9	$2.05 \cdot 10^{-2}$	0.059686
7	6	0.2036	0.5092	$6.72 \cdot 10^{-7}$	$8.91 \cdot 10^{-5}$	0.6372	0.9007	$5.68 \cdot 10^{-4}$	0.016511
8	7	—	—	—	—	—		—	-
	$\lambda_z/\lambda_m = 0.88, F = 0.5$				$\lambda_z/\lambda_m = 0,65,F = 0.5$				
N_1	N_2	arphi	Ε	$\delta_{\max} _{T_m=0}$	$\delta_{\max}\big _{\partial T_m/\partial Y=0}$	arphi	Ε	$\delta_{\max}\big _{T_m=0}$	$\delta_{\max}\big _{\partial T_m/\partial Y=0}$
5	4	0.1790	0.4774	$1.49 \cdot 10^{-5}$	0.026341	0.7876	1.0014	0.00025	0.32083
6	5	0.1915	0.4937	$1.18 \cdot 10^{-5}$	$7.26 \cdot 10^{-5}$	0.6327	0.8976	0.00025	0.0005
7	6	0.1915	0.4938	$1.13 \cdot 10^{-7}$	$2.13 \cdot 10^{-5}$	0.6328	0.8976	$5.84 \cdot 10^{-5}$	0.00204
8	7	_	_	_	—	_	_	_	_

Table 2. The impact of the number of collocation points on the value of the volumetric fraction of the fibers in the composite and the maximal error fulfillment of the boundary conditions at checkpoints.

	F = 10							
λ_z/λ_m	$\lambda_z/\lambda_m = 1.4$		$\lambda_z/\lambda_m =$	3.35				
Iter.	$\delta_{ m Newton}$	Ε	$\delta_{ m Newton}$	Ε				
1	0.456063	0.5439	0.265565	0.7344				
2	0.032858	0.5111	0.179407	0.9138				
3	$1.89 imes 10^{-3}$	0.5092	0.013073	0.9008				
4	3.41×10^{-6}	0.5092	4.17×10^{-5}	0.9007				
5	1.13×10^{-11}	0.5092	2.71×10^{-9}	0.9007				
	F = 0.5							
λ_z/λ_m	$\lambda_z/\lambda_m =$	0.88	$\lambda_z/\lambda_m = 0.65$					
Iter.	$\delta_{ m Newton}$	Ε	$\delta_{ m Newton}$	E				
1	0.357932	0.6421	0.217974	0.9164				
2	1.31×10^{-1}	0.5113	0.018393	0.8980				
3	1.72×10^{-2}	0.4941	0.000336	0.8976				
4	$2.99 imes 10^{-4}$	0.4938	3.61×10^{-8}	0.8976				
5	$9.05 imes 10^{-8}$	0.4938	_	-				

Table 3. Convergence of Newton's method for the test examples, with $N_1 = 7$ and $N_2 = 8$.

The results for a flat composite layer are presented by the dotted lines in Figure 4. As with the problem of identification of λ_z/λ_m , so also for the iterative identification of the volume fraction of fibers φ in a composite; the number of collocation points $N_1 + N_2$ where the boundary condition is approximately fulfilled affects the accuracy of the calculations. Table 2 shows the impact of the number of collocation points on the value of the volumetric fraction of fibers φ in the composite and the maximum error of fulfillment of boundary conditions at the checkpoints (between collocation points). As in the case of the direct problem, we obtain the best results for 7 collocation points at the right edge Γ_1 of a large finite element and 6 points at the upper edge Γ_2 . In the case of the algorithm is lost. Table 3 presents the convergence of the used Newton's iterative method for four test examples. The method proves to be convergent very quickly, and just after five iterations we get the correct result of the iteration with an error of less than 10^{-7} .

6. Conclusions

The presented method of determining the volume fraction of fibers of a composite or the effective thermal conductivity except in the cases of maximal fiber density is easy to implement and efficient. It can be easily applied to other configurations of regular arrangement of fibers in the matrix, for example to a triangular or hexagonal mesh. This study compared the influence of the ratio of the thermal conductivity of fibers to the thermal conductivity of matrix F on the value of the volumetric fraction of the fibers and the value of the effective thermal conductivity of the composite. It was also shown that increasing the number of collocation points doesn't reduce the error of the approximation of the boundary conditions; it leads to the ill-conditioning of the system of equations.

References

- [Chantasiriwan 2002] S. Chantasiriwan, "Steady-state determination of temperature-dependent thermal conductivity", Int. Commun. Heat Mass 29:6 (2002), 811–819.
- [Han and Cosner 1981] L. S. Han and A. A. Cosner, "Effective thermal conductivities of fibrous composites", *J. Heat Transf.* (*ASME*) 103:2 (1981), 387–392.
- [Hon and Wei 2004] Y. C. Hon and T. Wei, "A fundamental solution method for inverse heat conduction problem", *Eng. Anal. Bound. Elem.* **28**:5 (2004), 489–495.
- [Kołodziej and Zieliński 2009] J. A. Kołodziej and A. P. Zieliński, *Boundary collocation techniques and their application in engineering*, WIT Press, Southampton, 2009.
- [Marin 2005] L. Marin, "Numerical solution of the Cauchy problem for steady-state heat transfer in two-dimensional functionally graded materials", *Int. J. Solids Struct.* **42**:15 (2005), 4338–4351.
- [Yan et al. 2008] L. Yan, C.-L. Fu, and F.-L. Yang, "The method of fundamental solutions for the inverse heat source problem", *Eng. Anal. Bound. Elem.* **32**:3 (2008), 216–222.

Received 29 Sep 2010. Revised 5 Jun 2011. Accepted 15 Dec 2011.

MAGDALENA MIERZWICZAK: magdalena.mierzwiczak@wp.pl Institute of Applied Mechanics, Poznan University of Technology, ul. Jana Pawla II 24, 60-965 Poznan, Poland

JAN ADAM KOŁODZIEJ: jan.kolodziej@put.poznan.pl Institute of Applied Mechanics, Poznan University of Technology, ul. Jana Pawla II 24, 60-965 Poznan, Poland

JOURNAL OF MECHANICS OF MATERIALS AND STRUCTURES

jomms.net

Founded by Charles R. Steele and Marie-Louise Steele

EDITORS

CHARLES R. STEELE	Stanford University, USA
DAVIDE BIGONI	University of Trento, Italy
Iwona Jasiuk	University of Illinois at Urbana-Champaign, USA
YASUHIDE SHINDO	Tohoku University, Japan

EDITORIAL BOARD

H. D. BUI	École Polytechnique, France
J. P. CARTER	University of Sydney, Australia
R. M. CHRISTENSEN	Stanford University, USA
G. M. L. GLADWELL	University of Waterloo, Canada
D. H. HODGES	Georgia Institute of Technology, USA
J. HUTCHINSON	Harvard University, USA
C. HWU	National Cheng Kung University, Taiwan
B. L. KARIHALOO	University of Wales, UK
Y. Y. Kim	Seoul National University, Republic of Korea
Z. Mroz	Academy of Science, Poland
D. PAMPLONA	Universidade Católica do Rio de Janeiro, Brazil
M. B. RUBIN	Technion, Haifa, Israel
A. N. Shupikov	Ukrainian Academy of Sciences, Ukraine
T. TARNAI	University Budapest, Hungary
F. Y. M. WAN	University of California, Irvine, USA
P. WRIGGERS	Universität Hannover, Germany
W. YANG	Tsinghua University, China
F. ZIEGLER	Technische Universität Wien, Austria
PRODUCTION	contact@msp.org
SILVIO LEVY	Scientific Editor

Cover design: Alex Scorpan

See http://jomms.net for submission guidelines.

JoMMS (ISSN 1559-3959) is published in 10 issues a year. The subscription price for 2012 is US \$555/year for the electronic version, and \$735/year (+\$60 shipping outside the US) for print and electronic. Subscriptions, requests for back issues, and changes of address should be sent to Mathematical Sciences Publishers, Department of Mathematics, University of California, Berkeley, CA 94720–3840.

JoMMS peer-review and production is managed by EditFLOW® from Mathematical Sciences Publishers.



Journal of Mechanics of Materials and Structures

Volume 7, No. 3 Ma

March 2012

Special issue				
Trends in Continuum Physics (TRECOP 2010)				
Preface BOGDAN T. MARUSZEWSKI, WOLFGANG MUSCHIK, JOSEPH N. GRIMA and KRZYSZTOF W. WOJCIECHOWSKI	225			
The inverse determination of the volume fraction of fibers in a unidirectionally reinforced composite for a given effective thermal conductivity MAGDALENA MIERZWICZAK and JAN ADAM KOŁODZIEJ	229			
Analytical-numerical solution of the inverse problem for the heat conduction				
equation MICHAŁ CIAŁKOWSKI, ANDRZEJ MAĆKIEWICZ,				
JAN ADAM KOŁODZIEJ, UWE GAMPE and ANDRZEJ FRĄCKOWIAK	239			
Analysis of stress-strain distribution within a spinal segment				
ZDENKA SANT, MARIJA CAUCHI and MICHELLE SPITERI	255			
A mesh-free numerical method for the estimation of the torsional stiffness of long				
bones ANITA USCILOWSKA and AGNIESZKA FRASKA	265			
Rayleigh-type wave propagation in an auxetic dielectric ANDRZEJ DRZEWIECKI	277			
Local gradient theory of dielectrics with polarization inertia and irreversibility of local mass displacement VASYL KONDRAT and OLHA HRYTSYNA	285			
Electromagnetoelastic waves in a vortex layer of a superconductor BOGDAN T. MARUSZEWSKI, ANDRZEJ DRZEWIECKI AND ROMAN STAROSTA	297			

