

Pacific Journal of Mathematics

TRANSFORMATIONS OF SYSTEMS OF RELATIVISTIC PARTICLE MECHANICS

HERMAN RUBIN AND PATRICK COLONEL SUPPES

TRANSFORMATIONS OF SYSTEMS OF RELATIVISTIC PARTICLE MECHANICS

HERMAN RUBIN AND PATRICK SUPPES

1. Introduction. In [7] the axiomatic foundations of classical particle mechanics were investigated; and in [8] the transformations which carry systems of classical particle mechanics into systems of classical particle mechanics were determined. The purpose of the present paper is a similar investigation of relativistic particle mechanics (in the sense of the special theory of relativity). Some remarks on the general orientation of these studies are to be found in [7, § 1] and in [9].

In regard to our axiomatization of relativistic particle mechanics, we want to emphasize that we have in no sense attempted to use primitive notions which are logically or epistemologically simple. Investigations with these latter aims are to be found in [11], [12], [13], and [14]; but these studies are incomplete in the sense that they do not give axioms adequate for relativistic particle mechanics as it is ordinarily conceived by physicists. We have attempted to present such a complete set of axioms in a mathematically clear way.

The main result of the present paper is the determination under a certain weak hypothesis of the set of transformations which always carry systems of relativistic particle mechanics into systems of relativistic particle mechanics. Although this set of transformations is not a group (under the usual operation) we are able to show that it is essentially a Brandt groupoid. It is difficult precisely to compare our results with those in [6], but our results seem to represent an improvement in three respects: (i) we work within an explicit axiomatic framework; (ii) we consider transformations of the units of mass and force as well as position and time; (iii) we consider transformations from one value for the velocity of light to another.

We briefly summarize the mathematical notations we use, most of which are standard. We denote the ordered n -tuple whose first member is a_1 , whose second member is a_2 , and so on, by

Received May 7, 1953. The authors are grateful to Professor J.C.C. McKinsey for a large number of helpful suggestions and criticisms. This work was supported in part by a grant from the Office of Naval Research.

Pacific J. Math. 4 (1954), 563-601

$$\langle a_1, \dots, a_n \rangle.$$

By an n -dimensional vector we mean an ordered n -tuple of real numbers. Operations on vectors are defined in the usual way. We use the symbol "0" to denote the real number zero, the n -dimensional vector all of whose components are zero, and the matrix all of whose elements are zero. If $A = \langle a_1, \dots, a_n \rangle$ is any vector, the length $|A|$ of A is defined by

$$|A| = \sqrt{a_1^2 + \dots + a_n^2};$$

and by $[A]_{i_1, i_2, \dots, i_r}$ we mean the r -dimensional vector $\langle a_{i_1}, a_{i_2}, \dots, a_{i_r} \rangle$. Thus if $A = \langle 4, 7, 5 \rangle$, then $[A]_{2,3} = \langle 7, 5 \rangle$. If A is a vector, we sometimes write " A^2 " for " $|A|^2$." If \mathcal{A} is a matrix, we denote the transpose of \mathcal{A} by " \mathcal{A}^* ," and the determinant of \mathcal{A} by " $|\mathcal{A}|$." We denote the identity matrix by " \mathcal{I} ." Although we treat vectors as one-rowed matrices, if A is a vector we always mean by $|A|$ the length of A and not the determinant of A : the meaning should be clear from the context. We use both matrix notation and usual vector notation for the inner product of two vectors A and B . Thus we sometimes write: AB^* , and sometimes: $A \cdot B$, whichever is more convenient.

We use Menger's notation for derivatives (see [10]). If f is a function, then $D(f)$ is the derivative of f . Thus, for example,

$$D(\sin) = \cos, [D(\sin)](x) = \cos x, \text{ and } [D^2(\sin)](x) = -\sin x.$$

In this connection, we use the standard notation for sums, products, quotients, square roots, and so on, of functions. Thus, for example, if f and g are functions of a real variable, by $f + g$ we mean the function h such that for every real number x

$$h(x) = f(x) + g(x).$$

If f is a one-to-one function, f^{-1} is the inverse function of f . It is also convenient to introduce a special symbol for the *composition* of two functions: if f and g are functions of a real variable, by $g \circ f$ we mean the function h such that for every real number x

$$h(x) = g(f(x)).$$

To make some of our equations involving derivatives more perspicuous in relation to the notation ordinarily used in physics, we introduce formally the

following two symbols: if f and g are functions of a real variable, then the function df/dg is defined by the following equation (for all real numbers x)

$$\frac{df}{dg}(x) = \left(\frac{Df}{Dg} \right)(x),$$

and the function d^2f/dg^2 by the equation

$$\frac{d^2f}{dg^2}(x) = \left(\frac{D[Df/Dg]}{Dg} \right)(x).$$

Finally, we also use the following notation: I is the set of all positive integers, R is the set of all real numbers, R^+ is the set of all positive real numbers, and E_n is the set of all n -dimensional vectors. We sometimes use geometrical language, referring to vectors in E_n as points in n -dimensional Euclidean space, and so on.

2. Primitive notions. Our axioms for relativistic particle mechanics are based on six primitive notions: P , \mathfrak{V} , m , s , f , and c . P is a set, \mathfrak{V} and m are unary functions, s is a binary function, f is a ternary function, and c is a constant.

The intended physical interpretation of P is as the set of particles. For every p in P , $\mathfrak{V}(p)$ is to be interpreted physically as a set of real numbers measuring elapsed times (in terms of some unit of time and measured from some origin of time). There is a good physical reason for assigning (possibly) different sets of real numbers to different particles, instead of having one set of elapsed times for the whole system, as in [7]: two particles which have a simultaneous “life-span” with respect to one inertial frame of reference may have life-spans which do not even overlap with respect to another inertial frame.

For every p in P , $m(p)$ is to be interpreted physically as the numerical value of the rest mass of p . For every p in P and t in $\mathfrak{V}(p)$, $s(p, t)$ is a vector, to be thought of physically as giving the position of p at time t . Thus the primitive s fixes the choice of a coordinate system. It is also possible to take as a primitive the set of all admissible (that is, inertial) coordinate systems; this procedure is followed in [3]. We remark that for a fixed p in P , it is usually convenient to use in place of s the function s_p , which is defined on $\mathfrak{V}(p)$ and is such that, for every t in $\mathfrak{V}(p)$,

$$s_p(t) = s(p, t).$$

For every p in P and t in $\mathfrak{I}(p)$, and for i any positive integer, $f(p, t, i)$ is a vector giving the components (parallel to the axes of the coordinate system) of the i th force acting on p at time t . For further discussion of this primitive, applicable to relativistic as well as classical particle mechanics, see [7].

Our primitive constant c is to be interpreted as the numerical value of the velocity of light.

3. Axioms. Using the six primitive notions just described, we now give our axioms for relativistic particle mechanics.

An ordered sextuple $\Gamma = \langle P, \mathfrak{I}, m, s, f, c \rangle$ which satisfies the following Axioms A1-A7 is called an *n-dimensional system of relativistic particle mechanics* (or sometimes, simply a *system of relativistic particle mechanics*, for abbreviation, S.R.P.M.):

KINEMATICAL AXIOMS

A1. P is a nonempty, finite set.

A2. If $p \in P$, then $\mathfrak{I}(p)$ is an interval of real numbers.

A3. If $p \in P$ and $t \in \mathfrak{I}(p)$, then $s_p(t)$ is an n -dimensional vector; and, moreover, the second derivative of s_p exists throughout the interval $\mathfrak{I}(p)$.

A4. The constant c is a positive real number such that for every p in P and t in $\mathfrak{I}(p)$,

$$|(Ds_p)(t)| < c.$$

DYNAMICAL AXIOMS

A5. If $p \in P$, then $m(p)$ is a positive real number.

A6. If $p \in P$ and $t \in \mathfrak{I}(p)$, then $f(p, t, 1), f(p, t, 2), \dots$ are n -dimensional vectors such that the series

$$\sum_{i=1}^{\infty} f(p, t, i)$$

is absolutely convergent.

A7. If $p \in P$ and $t \in \mathfrak{I}(p)$, then

$$m(p) \left[D \frac{Ds_p}{\sqrt{1 - |Ds_p|^2/c^2}} \right] (t) = \sqrt{1 - \frac{|(Ds_p)(t)|^2}{c^2}} \cdot \sum_{i=1}^{\infty} f(p, t, i).$$

Since this set of axioms is similar in many ways to that given for classical mechanics in [7], a large number of remarks to be found in § 3 of that paper are also applicable here and will not be repeated. From Axiom A7 it is clear that the force concept we are using is that of Minkowski. In the solution of special problems this concept is not always the most useful one, but the relative simplicity of its transformation properties more than justifies its use here. Some readers may feel that there are good physical grounds for taking the notion of relativistic mass as primitive instead of that of rest mass; however, it is easy to define the notion of relativistic mass in terms of the notion of rest mass and our other primitives, and the use of the notion of rest mass as a primitive emphasizes the considerable formal similarity between our axioms for relativistic mechanics and the axioms in [7] for classical mechanics.

For p in P , $\mathfrak{I}(p)$ is a time interval for the particle p (with respect to the frame of reference fixed by our choice of primitives). It may seem that it would have been simpler to take $\mathfrak{I}(p)$ as the interval of proper time of the particle p . However, this approach would complicate the treatment of *systems* of particles. In the main, the notion of proper time is most convenient in discussions restricted to the consideration of a single particle. From the remark in the previous section it is clear that it is not reasonable to require that the intervals $\mathfrak{I}(p)$ be overlapping. A second argument against such an assumption is the prominence in modern physics of elementary particles with very short life-spans.¹ We note, however, that in studying certain special problems, such as that of defining a reasonable notion of center of mass of a S.R.P.M., it is desirable to restrict the discussion to systems in which $\mathfrak{I}(p) = (-\infty, +\infty)$ for every p in P .

If (i) “ c ” is replaced by “ $1/k$ ” in the inequality of Axiom A4 and the equation of Axiom A7, (ii) k is treated as a primitive replacing c , and (iii) Axiom A4 is modified to read: “The constant k is a nonnegative real number such that...,” then, by adding appropriate further axioms, we can get either classical or relativistic particle mechanics. Thus an additional axiom asserting that $k = 0$ gives us classical mechanics; and the assertion that $k > 0$ gives us relativistic mechanics.

¹Paulette Destouches-Février [2, pp.5-6] advocates the use of a three-valued logic to describe the creation and annihilation of elementary particles. Actually, the situation is easily handled by the simple device of introducing the function \mathfrak{I} defined on P instead of a fixed interval T for the whole system. Indeed, to our mind, her drastic proposal cannot be taken seriously until we know a great deal more about the mathematics which goes with a multi-valued logic. Even if such a body of mathematics existed (as it does not—we do not have even the general outlines of elementary set theory in three-valued logic), it would be reasonable to adopt such a proposal only after every feasible alternative in standard mathematics had been explored.

We close this section with a number of definitions which will be useful later.

For p in P and t in $\mathfrak{V}(p)$, we set

$$v_p(t) = (Ds_p)(t);$$

$v_p(t)$ is, of course, the *velocity* of p at time t . With respect to a fixed element t_0 in $\mathfrak{V}(p)$, we define the function τ_{t_0} (for p in P and t in $\mathfrak{V}(p)$) as follows:

$$\tau_{t_0}(p, t) = \int_{t_0}^t \sqrt{1 - \frac{|v_p(t)|^2}{c^2}} dt;$$

$\tau_{t_0}(p, t)$ is the *proper time* of p . Since we are interested only in the derivative of this function with respect to t , and since the derivative is independent of t_0 , we shall usually drop the subscript.

For p in P and t in $\mathfrak{V}(p)$, we define the function q as follows:

$$q(p, t) = \langle s(p, t), t \rangle.$$

It is natural to call q the *space-time* function.

For p in P , t in $\mathfrak{V}(p)$, and i any positive integer, we define what we call the *relativistic force* function f^{rel} as follows:

$$f^{\text{rel}}(p, t, i) = \left\langle f(p, t, i), \frac{f(p, t, i) \cdot v_p(t)}{c^2} \right\rangle$$

Although it is not usual to adopt a special name for this function, the function itself is used frequently in textbook treatments of relativity.

By a *c-particle path* (for any positive number c) we mean a set \mathfrak{P} of points (that is, vectors) in E_{n+1} for which there exists a S.R.P.M. $\langle \{1\}, \mathfrak{V}, m, s, f, c \rangle$ such that for every point X of E_{n+1} , X is in \mathfrak{P} if and only if there exists a t in $\mathfrak{V}(1)$ such that $X = \langle s(1, t), t \rangle$.² It is obvious that if g is any twice-differentiable function defined on an interval T of real numbers and taking vectors in E_n as values, then the set of vectors $\langle g(t), t \rangle$ for t in T is a *c-particle path*, provided that $|(Dg)(t)| < c$ for all t in T .

By the *slope* of a line α in E_{n+1} , whose projection on the $(n+1)$ st-axis

²The intuitive interpretation of E_{n+1} is as the space-time manifold of special relativity with the $(n+1)$ st coordinate representing the time coordinate. Thus, if $\langle Z, x \rangle$ is a point of E_{n+1} , then under the intended interpretation, the n -dimensional vector Z gives the spatial coordinates of the point and x its time coordinate.

is a nondegenerate segment, we mean the n -dimensional vector W such that for any two distinct points $\langle Z_1, x_1 \rangle$ and $\langle Z_2, x_2 \rangle$ of α ,

$$\frac{Z_1 - Z_2}{x_1 - x_2} = W.$$

By the speed of α we mean the nonnegative number $|W|$. By a c -inertial path we mean a line in E_{n+1} whose speed is less than c . We note that every segment of a c -inertial path is a c -particle path, but is not necessarily a c -inertial path (since a c -inertial path must be a whole line). By a c -line we mean a line in E_{n+1} whose speed is equal to c . The notion of a c -line corresponds to the intuitive notion of a light line.

If we want to refer to a S.R.P.M. Γ with numerical constant c , we shall write: S.R.P.M. Γ_c .

4. Transformation theorems. We begin by defining the notion of a *generalized Lorentz matrix*. An intuitive discussion of such matrices follows Theorem 1.

DEFINITION 1. Let c , c' , and λ be positive real numbers. Then a matrix \mathcal{Q} of order $n+1$ is said to be a *generalized Lorentz matrix with respect to $\langle c, c', \lambda \rangle$* if and only if there exist numbers δ and β , an n -dimensional vector U , and an orthogonal matrix \mathcal{E} of order n , such that

$$\delta^2 = 1, \quad \beta^2 \left(1 - \frac{U^2}{c'^2} \right) = 1,$$

and

$$(i) \quad \mathcal{Q} = \lambda \begin{pmatrix} \mathfrak{A} & 0 \\ 0 & \frac{c}{c'} \end{pmatrix} \begin{pmatrix} \mathcal{E} & 0 \\ 0 & \delta \end{pmatrix} \begin{pmatrix} \mathfrak{A} + \frac{\beta - 1}{U^2} U^* U & -\frac{\beta U^*}{c'^2} \\ -\beta U & \beta \end{pmatrix}.$$

The following two lemmas simplify the statement and proof of Theorem 1.

LEMMA 1. Let $\langle \{1\}, \mathfrak{V}, m, s, f, c \rangle$ be a S.R.P.M., let c' and λ be positive real numbers, and let \mathcal{Q} be a generalized Lorentz matrix with respect to $\langle c, c', \lambda \rangle$. Let the function h be defined by the equation (for every t in $\mathfrak{V}(1)$):

$$h(t) = [\langle s_1(t), t \rangle \mathcal{Q}]_{n+1}.$$

Then the function Dh exists; its values are either always positive or always negative; and the function h is one-to-one.

Proof. From Definition 1 and the hypothesis of the lemma we see that there are numbers δ and β , an n -dimensional vector U , and an orthogonal matrix \mathcal{E} , such that

$$\delta^2 = 1, \quad \beta^2 \left(1 - \frac{U^2}{c'^2} \right) = 1,$$

and

$$\mathcal{Q} = \begin{pmatrix} \lambda \left(\mathcal{E} + \frac{(\beta - 1) \mathcal{E} U^* U}{U^2} \right) - \frac{\lambda \beta \mathcal{E} U^*}{c'^2} & \\ - \frac{\lambda c \delta \beta U}{c'} & \frac{\lambda c \delta \beta}{c'} \end{pmatrix}.$$

Thus

$$h(t) = \frac{\lambda c \delta \beta t}{c'} + \frac{\lambda \beta s_1(t) \mathcal{E} U^*}{c'^2} = \left(\frac{\lambda c \delta \beta}{c'} \right) \left(t - \frac{\delta s_1(t) \mathcal{E} U^*}{c c'} \right).$$

Hence

$$\frac{c'(Dh)(t)}{\delta \lambda \beta c} = 1 - \frac{v_1(t) \delta \mathcal{E} U^*}{c c'} \geq 1 - \frac{|v_1(t) \mathcal{E}| |U|}{c c'}.$$

Using Axiom A4 and the fact that \mathcal{E} is orthogonal, we have

$$\frac{c'(Dh)(t)}{\delta \lambda \beta c} \geq 1 - \frac{|U|}{c'}.$$

Since $|U| < c'$, the function Dh is bounded away from zero, and it thus follows from Rolle's theorem that h is one-to-one.

The following lemma is a theorem of matrix theory.

LEMMA 2. Let c , c' , and λ be positive real numbers. Then a matrix \mathcal{Q} of order $n + 1$ is a generalized Lorentz matrix with respect to $\langle c, c', \lambda \rangle$ if and only if

$$(i) \quad \mathcal{Q} \begin{pmatrix} \mathfrak{L} & 0 \\ 0 & -c'^2 \end{pmatrix} \mathcal{Q}^* = \lambda^2 \begin{pmatrix} \mathfrak{L} & 0 \\ 0 & -c^2 \end{pmatrix}.$$

Proof. The proof of necessity is obtained by direct application of Definition 1.

For the proof of sufficiency, let

$$\mathcal{Q} = \begin{pmatrix} \mathfrak{K} & K^* \\ L & m \end{pmatrix},$$

where \mathfrak{K} is a matrix of order n , K and L are n -dimensional vectors, and m is a real number. From (i) we obtain at once:

$$(1) \quad \mathfrak{K}\mathfrak{K}^* - c'^2 K^*K = \lambda^2 \mathfrak{L},$$

$$(2) \quad \mathfrak{K}L^* - c'^2 mK^* = 0,$$

$$(3) \quad LL^* - c'^2 m = -\lambda^2 c^2.$$

From (3) it follows that

$$(4) \quad m \neq 0.$$

We define:

$$(5) \quad \beta = \frac{c'|m|}{c\lambda},$$

$$(6) \quad \delta = \frac{m}{|m|},$$

$$(7) \quad U = -\frac{L}{m},$$

$$(8) \quad \mathcal{E} = \frac{1}{\lambda} \left(\mathfrak{K} + \frac{(\beta - 1)c'^2 K^*U}{\beta U^2} \right).$$

Since the right member of equation (i) of Definition 1 can be written

$$\begin{pmatrix} \lambda \left(\mathcal{E} + \frac{(\beta - 1) \mathcal{E} U^* U}{U^2} \right) & - \frac{\lambda \beta \mathcal{E} U^*}{c'^2} \\ - \frac{\lambda c \delta \beta U}{c'} & \frac{\lambda c \delta \beta}{c'} \end{pmatrix},$$

in order to complete the proof it suffices to show that

$$(I) \quad \lambda \left(\mathcal{E} + \frac{(\beta - 1) \mathcal{E} U^* U}{U^2} \right) = \mathcal{H},$$

$$(II) \quad - \frac{\lambda \beta \mathcal{E} U^*}{c'^2} = K^*,$$

$$(III) \quad - \frac{\lambda c \delta \beta U}{c'} = L,$$

$$(IV) \quad \frac{\lambda c \delta \beta}{c'} = m,$$

$$(V) \quad \delta^2 = 1,$$

$$(VI) \quad \beta^2 \left(1 - \frac{U^2}{c'^2} \right) = 1,$$

$$(VII) \quad \mathcal{E} \mathcal{E}^* = \mathfrak{A}.$$

Equation (III) follows immediately from (5), (6), and (7), equation (IV) from (5) and (6), equation (V) from (6), and equation (VI) from (3), (5), and (7).

From (2) and (7) we get

$$(9) \quad \mathcal{H} U^* = - c'^2 K^*,$$

and then from (8) and (9) we have

$$\begin{aligned} (10) \quad \mathcal{E} \mathcal{E}^* &= \frac{1}{\lambda^2} \left(\mathcal{H} + \frac{(\beta - 1) c'^2 K^* U}{\beta U^2} \right) \left(\mathcal{H}^* + \frac{(\beta - 1) c'^2 U^* K}{\beta U^2} \right) \\ &= \frac{1}{\lambda^2} \left\{ \mathcal{H} \mathcal{H}^* + \frac{(\beta - 1) c'^2}{\beta^2 U^2} [\beta (-c'^2 K^*) K + \beta K^* (-c'^2 K) \right. \\ &\quad \left. + (\beta - 1) c'^2 K^* K] \right\} \end{aligned}$$

$$\begin{aligned}
 &= \frac{1}{\lambda^2} \left[\mathcal{H} \mathcal{H}^* + \frac{(\beta - 1) c'^4}{\beta^2 U^2} K^* K (-2\beta + \beta - 1) \right] \\
 &= \frac{1}{\lambda^2} \left[\mathcal{H} \mathcal{H}^* - \frac{(\beta^2 - 1) c'^4}{\beta^2 U^2} K^* K \right].
 \end{aligned}$$

From (VI), (1), and (10) we conclude that

$$\mathcal{E} \mathcal{E}^* = \mathcal{I},$$

which establishes equation (VII). Multiplying both sides of (8) on the right by $-\lambda \beta U^*/c'^2$, and using (9), we get equation (II). Equation (I) follows from (8) and (II), completing the proof of the lemma.

The following theorem is a generalization of the well-known result that the relativistic equation of motion is covariant under a Lorentz transformation.

THEOREM 1. *Let $\langle P, \mathfrak{V}, m, s, f, c \rangle$ be an n -dimensional S.R.P.M. Let c' , γ , and λ be positive real numbers, let B be an $(n+1)$ -dimensional vector, and let \mathcal{Q} be a generalized Lorentz matrix with respect to $\langle c, c', \lambda \rangle$. For each p in P let the function h_p be defined as follows (for all t in $\mathfrak{V}(p)$):*

$$h_p(t) = [\langle s_p(t), t \rangle \mathcal{Q} + B]_{n+1}.$$

(By Lemma 1 the inverse function h_p^{-1} exists.) Let the function \mathfrak{V}' be defined as follows: for p in P , $\mathfrak{V}'(p)$ is the range of the function h_p ; and let the functions m' , s' , and f' be defined by the following equations (for p in P , t' in $\mathfrak{V}'(p)$ and i in I):

$$m'(p) = \gamma m(p),$$

$$s'(p, t') = [\langle s(p, h_p^{-1}(t')), h_p^{-1}(t') \rangle \mathcal{Q} + B]_{1, \dots, n},$$

$$f'(p, t', i) = \frac{\gamma c'^2}{\lambda^2 c^2} \left[\left\langle f(p, h_p^{-1}(t'), i), \frac{f(p, h_p^{-1}(t'), i) \cdot v_p(h_p^{-1}(t'))}{c^2} \right\rangle \mathcal{Q} \right]_{1, \dots, n}.$$

Then $\Gamma' = \langle P, \mathfrak{V}', m', s', f', c' \rangle$ is an n -dimensional S.R.P.M.

Proof. It will suffice to show that Γ' satisfies Axioms A4 and A7, since the proof for the other axioms is trivial. Let

$$\mathcal{Q} = \begin{pmatrix} \mathbb{Q} & E^* \\ F & g \end{pmatrix}.$$

It is easy to show that for p in P , and t' in $\mathfrak{V}'(p)$,

$$(1) \quad v_p'(t') = \frac{\langle v_p(h_p^{-1}(t'), 1) \begin{pmatrix} \mathbb{Q} \\ F \end{pmatrix} \rangle}{\langle v_p(h_p^{-1}(t'), 1) \begin{pmatrix} E^* \\ g \end{pmatrix} \rangle},$$

with the denominator of the right member of (1) always unequal to zero. (Since in this proof we always consider a fixed particle p , we drop the subscript “ p ” from this point on.)

We have, from Axiom A4,

$$(2) \quad \lambda^2(|v(h^{-1}(t'))|^2 - c^2) < 0;$$

but

$$\lambda^2(|v(h^{-1}(t'))|^2 - c^2) = \lambda^2 \langle v(h^{-1}(t')), 1 \rangle \begin{pmatrix} \mathbb{Q} & 0 \\ 0 & -c'^2 \end{pmatrix} \langle v(h^{-1}(t')), 1 \rangle^*.$$

Then by Lemma 2 we have

$$(3) \quad \lambda^2(|v(h^{-1}(t'))|^2 - c^2) = \langle v(h^{-1}(t')), 1 \rangle \mathcal{Q} \begin{pmatrix} \mathbb{Q} & 0 \\ 0 & -c'^2 \end{pmatrix} \mathcal{Q}^* \langle v(h^{-1}(t')), 1 \rangle^*.$$

The right member of (3) is equal to

$$(4) \quad \langle v(h^{-1}(t')), 1 \rangle \begin{pmatrix} \mathbb{Q} \\ F \end{pmatrix} (\mathbb{Q}^* F^*) \langle v(h^{-1}(t')), 1 \rangle^* \\ - c'^2 \langle v(h^{-1}(t')), 1 \rangle \begin{pmatrix} E^* \\ g \end{pmatrix} (Eg) \langle v(h^{-1}(t')), 1 \rangle^*,$$

and using (1) we see that (4) is equal to

$$(5) \quad \left(\langle v(h^{-1}(t')), 1 \rangle \begin{pmatrix} E^* \\ g \end{pmatrix} v'(t') \right)^2 - c'^2 \left(\langle v(h^{-1}(t')), 1 \rangle \begin{pmatrix} E^* \\ g \end{pmatrix} \right)^2.$$

From (2), (3), (4), and (5) we conclude that

$$|v'(t')|^2 - c'^2 < 0,$$

which verifies Axiom A4 for Γ' .

It is not difficult to show that from Axiom A7 we have

$$(6) \quad \frac{m(p)}{(1 - |v(t)|^2/c^2)^{1/2}} D \left[\frac{Dq}{(1 - |v|^2/c^2)^{1/2}} \right] (t) = \sum_{i=1}^{\infty} f^{\text{rel}}(p, t, i).$$

Setting $q'(t') = \langle s'(t'), t' \rangle$ for all t' in $\mathfrak{D}'(p)$, we conclude from the hypothesis of our theorem that

$$q'(h(t)) = q(t) \mathcal{Q} + B,$$

and thus

$$(7) \quad ((Dq') \circ h)(t)(Dh)(t) = (Dq)(t) \mathcal{Q}.$$

Directly from the definition of q and q' we obtain

$$(8) \quad (Dq)(t) \begin{pmatrix} \lambda & 0 \\ 0 & -c^2 \end{pmatrix} ((Dq)(t))^* = |v(t)|^2 - c^2,$$

and

$$(9) \quad ((Dq') \circ h)(t) \begin{pmatrix} \lambda & 0 \\ 0 & -c'^2 \end{pmatrix} (((Dq') \circ h)(t))^* = |(v' \circ h)(t)|^2 - c'^2.$$

Using Lemma 1, Lemma 2, and (7) we obtain, from (8) and (9),

$$|(v' \circ h)(t)|^2 - c'^2 = \frac{\lambda^2}{((Dh)(t))^2} (|v(t)|^2 - c^2),$$

and thus

$$(10) \quad 1 - \frac{|(v' \circ h)(t)|^2}{c'^2} = \frac{\lambda^2 c^2}{c'^2 ((Dh)(t))^2} \left(1 - \frac{|v(t)|^2}{c^2} \right).$$

By Lemma 1, $(Dh)(t)$ is either always positive or always negative; the remainder of our proof is analogous in the two cases, so that we shall only consider the case where it is always positive. We then have, from (7) and (10),

$$\frac{[(Dq') \circ h](t)}{(1 - |(v' \circ h)(t)|^2/c'^2)^{1/2}} = \frac{c'(Dq)(t)\mathcal{Q}}{\lambda c(1 - |v(t)|^2/c^2)^{1/2}},$$

and hence

$$(11) \quad \left(\left[\frac{Dq'}{(1 - |v'|^2/c'^2)^{1/2}} \right] \circ h \right) (t) = \frac{c'(Dq)(t)\mathcal{Q}}{\lambda c(1 - |v(t)|^2/c^2)^{1/2}}.$$

Differentiating both sides of (11), and using (6), we obtain

$$(12) \quad (Dh)(t) \left(\left[D \left(\frac{Dq'}{(1 - |v'|^2/c'^2)^{1/2}} \right) \right] \circ h \right) (t) = \frac{c'}{\lambda c} D \left(\frac{Dq}{(1 - |v|^2/c^2)^{1/2}} \right) (t) \mathcal{Q} \\ = \left(\sqrt{1 - \frac{|v(t)|^2}{c^2}} \right) \left(\frac{c'}{\lambda cm(p)} \right) \left(\sum_{i=1}^{\infty} f^{\text{rel}}(p, t, i) \mathcal{Q} \right).$$

From (10), (12), and the hypothesis of our theorem, we infer that

$$(13) \quad \gamma m(p) \left(\left[D \left(\frac{Dq'}{(1 - |v'|^2/c'^2)^{1/2}} \right) \right] \circ h \right) (t) \\ = \sqrt{1 - \frac{|(v' \circ h)(t)|^2}{c'^2}} \sum_{i=1}^{\infty} f^{\text{rel}'}(p, h(t), i),$$

and from (13) we conclude immediately that Axiom A7 holds for Γ' .³

REMARK 1. All the transformations mentioned in Definition 1 and Theorem 1 have a clear intuitive interpretation if we consider $\langle P, \mathfrak{I}, m, s, f, c \rangle$ as a physical system whose mechanical properties are observed and measured with respect to some (inertial) frame of reference and some set of units of measurement, and $\langle P, \mathfrak{I}', m', s', f', c' \rangle$ as the same physical system observed and measured with respect to some other (inertial) frame of reference and some other set of units of measurement. Thus, c is the old and c' the new velocity of light. The introduction of the number γ amounts to changing the unit of mass by an amount $1/\gamma$, and the vector B corresponds to shifting the origin of the spatial frame of reference by $-[B]_1, \dots, n$, and the origin of time by an amount

³Readers familiar with the standard treatments of relativistic mechanics will note that (in the interests of rigor and explicitness) we have replaced " t " by " $h_p(t)$."

$-[B]_{n+1}$. The number λ represents a uniform stretch of space and time. When $\delta = -1$, we have a reversal of the direction of time. The matrix \mathcal{E} represents (for $n \leq 3$) a rotation of the spatial coordinates—or a rotation followed by a reflection. The vector U represents the relative velocity of the two inertial frames of reference, and the number β , which is determined by U and c' , is the well-known Lorentz contraction factor. Finally, it is easy to check that the last matrix in the factorization of the matrix \mathcal{Q} yields the ordinary Lorentz transformations. We note that the rather complicated transformation of the forces is the velocity-dependent transformation to be expected in relativistic mechanics.

REMARK 2. Theorem 2, our main result, is a sort of converse of Theorem 1: roughly speaking, we show that the transformations described in Theorem 1 are the only transformations which always take systems of relativistic particle mechanics into systems of relativistic particle mechanics. To facilitate the formulation and proof of Theorem 2, an additional lemma and some definitions will be useful.

LEMMA 3. Let $X_1 = \langle Z_1, x_1 \rangle$, $X_2 = \langle Z_2, x_2 \rangle$, and $X_3 = \langle Z_3, x_3 \rangle$ be any three points in E_{n+1} such that (i) $x_1 < x_2 < x_3$, (ii) there is a c -inertial path through X_1 and X_2 , and (iii) there is a c -inertial path through X_2 and X_3 . Then there is a c -particle path through X_1 , X_2 , and X_3 .

Proof. In view of the remark near the end of § 3, it will suffice to construct a function g which: (a) is defined on the closed interval $[x_1, x_3]$; (b) takes vectors in E_n as values; (c) is twice differentiable; (d) is such that for every t in $[x_1, x_3]$, $|(Dg)(t)| < c$; and (e) is such that

$$g(x_1) = Z_1, \quad g(x_2) = Z_2, \quad \text{and} \quad g(x_3) = Z_3.$$

Let

$$a = x_2 - x_1, \quad b = x_3 - x_2,$$

$$V = \frac{Z_2 - Z_1}{a}, \quad W = \frac{Z_3 - Z_2}{b},$$

$$\gamma = \frac{2c \log 2}{(c - \max(|V|, |W|)) \min(a, b)}, \quad A = \frac{aV \log \cosh \gamma b + bW \log \cosh \gamma a}{a \log \cosh \gamma b + b \log \cosh \gamma a},$$

$$B = \frac{ab\gamma(W - V)}{a \log \cosh \gamma b + b \log \cosh \gamma a}.$$

The reader may verify that the function g defined by the following equation (for t in $[x_1, x_3]$) has properties (a) – (e):

$$g(t) = Z_2 + (t - x_2)A + \frac{B}{\gamma} [\log \cosh \gamma(t - x_2)].$$

DEFINITION 2. Let ϕ_1 be a function mapping R^+ into R^+ ; let ϕ_2 be a function which is a one-to-one mapping of E_{n+1} into itself; and let ϕ_3 be a function mapping E_{2n} into E_n . Then we call the ordered triple $\langle \phi_1, \phi_2, \phi_3 \rangle$ an *eligible transformation*.

DEFINITION 3. Let $\Phi = \langle \phi_1, \phi_2, \phi_3 \rangle$ be an eligible transformation; let $\Gamma = \langle P, \mathfrak{I}, m, s, f, c \rangle$ be a S.R.P.M.; and for each p in P let the function H_p be defined as follows (for every t in $\mathfrak{I}(p)$):

$$H_p(t) = [\phi_2(s(p, t), t)]_{n+1}.$$

Then by the Φ -transform of Γ (which we also write: $\Phi(\Gamma)$), we mean the ordered quintuple $\langle P, \mathfrak{I}', m', s', f' \rangle$, where for p in P :

$$m'(p) = \phi_1(m(p));$$

$\mathfrak{I}'(p)$ is the range of the function H_p ; and s' and f' are defined by the following equations for t' in $\mathfrak{I}'(p)$, if the pre-image $H_p^{-1}(t')$ of t' under H_p is unique, and otherwise they are undefined:

$$s'(p, t') = [\phi_2(s(p, H_p^{-1}(t')), H_p^{-1}(t'))]_{1, \dots, n}$$

$$f'(p, t', i) = \phi_3(f(p, H_p^{-1}(t'), i) v(p, H_p^{-1}(t'))),$$

for $i \geq 1$.

We are now in a position to state and prove the main result of this paper.

THEOREM 2.⁴ Let $\Phi = \langle \phi_1, \phi_2, \phi_3 \rangle$ be an eligible transformation, and let

⁴The statement of Theorem 2 would be made more symmetrical to Theorem 1 if ϕ_2 were replaced by two functions ϕ' and ϕ'' such that

$$\phi'(Z, x) = [\phi_2(Z, x)]_{n+1} \quad \text{and} \quad \phi''(Z, x) = [\phi_2(Z, x)]_{1, \dots, n}.$$

This procedure was followed in [8] for classical mechanics; but in relativistic mechanics, it is natural to introduce the single transformation ϕ_2 for the space-time manifold.

c and c' be positive real numbers such that (i) for every n -dimensional system of relativistic particle mechanics Γ_c , $\langle \Phi(\Gamma_c), c' \rangle$ is a system of relativistic particle mechanics, and (ii) ϕ_2 carries no c -line into a c' -particle path. Then there exist positive real numbers γ and λ , an $(n+1)$ -dimensional vector B , and a generalized Lorentz matrix \mathcal{Q} with respect to $\langle c, c', \lambda \rangle$, such that, for any vectors Z_1 and Z_2 in E_n with $|Z_2| < c$, every x in R , and y in R^+ ,

$$\phi_1(\gamma) = \gamma\gamma,$$

$$\phi_2(Z_1, x) = \langle Z_1, x \rangle \mathcal{Q} + B,$$

$$\phi_3(Z_1, Z_2) = \frac{\gamma c'^2}{\lambda^2 c^2} \left[\left\langle Z_1, \frac{Z_1 \cdot Z_2}{c^2} \right\rangle \mathcal{Q} \right]_{1, \dots, n}.$$

Proof. We first want to show that if Z is any vector in E_n such that $|Z| < c$, then

$$\phi_3(0, Z) = 0.$$

Setting $P = \{1\}$, $\mathfrak{V}(1) = (-\infty, \infty)$, $m(1) = 1$, and, for t in $\mathfrak{V}(1)$,

$$s(1, t) = Zt,$$

$$f(1, t, i) = 0 \quad \text{for } i \geq 1,$$

we see that $\langle P, \mathfrak{V}, m, s, f, c \rangle$ is a S.R.P.M. Since for every t in $\mathfrak{V}(1)$, $Z = v(1, t)$, we conclude from the hypothesis of our theorem, Definition 2, and Axiom A6 that the series

$$\phi_3(0, Z) + \phi_3(0, Z) + \dots$$

is absolutely convergent. Hence

$$(1) \quad \phi_3(0, Z) = 0.$$

For every segment \mathfrak{A} of a c -inertial path there exists a one-particle S.R.P.M. $\langle \{1\}, \mathfrak{V}, m, s, f, c \rangle$ such that, for every t in $\mathfrak{V}(1)$,

$$f(p, t, i) = 0 \quad \text{for } i \geq 1,$$

and for every vector X in E_{n+1} , X is in \mathfrak{A} if and only if there is a t in $\mathfrak{V}(1)$ such that

$$X = \langle s(1, t), t \rangle.$$

Hence it follows immediately from (1) and the hypothesis of our theorem that:

(2) ϕ_2 carries segments of c -inertial paths into segments of c' -inertial paths.

Let $\Gamma_c = \langle P, \mathfrak{I}, m, s, f, c' \rangle$ be any S.R.P.M. with constant c . By hypothesis, $\langle \Phi(\Gamma_c), c \rangle$ is a S.R.P.M. For any p in P , if t_1 and t_2 are in $\mathfrak{I}(p)$ and $t_1 \neq t_2$, then

$$\phi_2(s(p, t_1), t_1) \neq \phi_2(s(p, t_2), t_2),$$

since ϕ_2 is one-to-one. Suppose now that

$$[\phi_2(s(p, t_1), t_1)]_{n+1} = [\phi_2(s(p, t_2), t_2)]_{n+1}.$$

Then we must have

$$[\phi_2(s(p, t_1), t_1)]_{1, \dots, n} \neq [\phi_2(s(p, t_2), t_2)]_{1, \dots, n};$$

but then $\langle \Phi(\Gamma_c), c' \rangle$ is not a S.R.P.M., for p is required to be in two places at the same time, which violates Axiom A3. We thus conclude:

(3) ϕ_2 is one-to-one in the last coordinate along the space-time path of any particle of a S.R.P.M. also, Γ_c , and thus the pre-image under ϕ_2 of any point t' in $\mathfrak{I}'(p)$, is unique.

Furthermore, since by hypothesis ϕ_2 takes the interval $\mathfrak{I}(p)$ into an interval $\mathfrak{I}'(p)$, we have:

(4) ϕ_2 is continuous in the last coordinate along the space-time path of any particle of a S.R.P.M.

From (4) and the fact that any two points $\langle Z, x \rangle$ and $\langle Z, y \rangle$ lie on a c -inertial path, we obtain:

(5) For any point $\langle Z, x \rangle$ and any $\epsilon > 0$, there exists a $\delta > 0$ such that for any point $\langle Z, y \rangle$ if $|x - y| < \delta$, then $|x' - y'| < \epsilon$, where

$$x' = [\phi_2(Z, x)]_{n+1} \quad \text{and} \quad y' = [\phi_2(Z, y)]_{n+1}.$$

We next show that

(I) ϕ_2 is continuous.

Let $\langle Z_1, x_1 \rangle$ be any point of E_{n+1} , and let ϵ be any positive number. Let $\epsilon^* = \epsilon/[2(1 + c')]$. Using (5), let δ^* be a positive number such that if $|x_1 - y| < \delta^*$ then $|x'_1 - y'| < \epsilon^*$, where

$$x'_1 = [\phi_2(Z_1, x_1)]_{n+1} \quad \text{and} \quad y' = [\phi_2(Z_1, y)]_{n+1};$$

and let $\delta = c\delta^*/(3c + 2)$. We shall show that if $\langle Z_2, x_2 \rangle$ is any point of E_{n+1} such that

$$(6) \quad |\langle Z_1, x_1 \rangle - \langle Z_2, x_2 \rangle| < \delta,$$

then

$$|\phi_2(Z_1, x_1) - \phi_2(Z_2, x_2)| < \epsilon.$$

Suppose for definiteness that

$$(7) \quad x_1 \geq x_2.$$

We may choose x_0 and x_3 so that

$$(8) \quad x_2 - \frac{|Z_2 - Z_1|}{c} - \delta < x_0 < x_2 - \frac{|Z_2 - Z_1|}{c}$$

and

$$(9) \quad x_1 + \frac{|Z_2 - Z_1|}{c} < x_3 < x_1 + \frac{|Z_2 - Z_1|}{c} + \delta$$

From (7), (8), and (9), we obtain

$$(10) \quad |x_3 - x_0| < |x_2 - x_1| + 2 \frac{|Z_2 - Z_1|}{c} + 2\delta;$$

and from (6) and (10) we then infer that

$$(11) \quad |x_3 - x_0| < \frac{(3c + 2)\delta}{c} = \delta^*.$$

Since from (7), (8), and (9) we have

$$(12) \quad x_0 < x_2 \leq x_1 < x_3,$$

we obtain from (11)

$$|x_3 - x_1| < \delta^*,$$

$$|x_1 - x_0| < \delta^*.$$

Consequently, by (5),

$$|x'_3 - x'_1| + |x'_1 - x'_0| < 2\epsilon^*,$$

and thus, by the triangle inequality,

$$(13) \quad |x'_3 - x'_0| < 2\epsilon^*,$$

where

$$x'_0 = [\phi_2(Z_1, x_0)]_{n+1},$$

$$x'_3 = [\phi_2(Z_1, x_3)]_{n+1}.$$

From the second part of (8) it follows that there is a c -inertial path through $\langle Z_1, x_0 \rangle$ and $\langle Z_2, x_2 \rangle$; and from (7) and the first part of (9) it follows that there is a c -inertial path through $\langle Z_2, x_2 \rangle$ and $\langle Z_1, x_3 \rangle$. We thus conclude from Lemma 3 that there exists a c -particle path through $\langle Z_1, x_0 \rangle$, $\langle Z_2, x_2 \rangle$, and $\langle Z_1, x_3 \rangle$. As before, for abbreviation, we set

$$x'_2 = [\phi_2(Z_2, x_2)]_{n+1},$$

$$Z'_2 = [\phi_2(Z_2, x_2)]_{1, \dots, n},$$

$$Z'_{1i} = [\phi_2(Z_1, x_i)]_{1, \dots, n} \quad (i = 0, 1, 3).$$

Since ϕ_2 is one-to-one and continuous in the last coordinate along any c -particle path, it is monotone in the last coordinate along any c -particle path, and we thus have: either

$$(14) \quad \left\{ \begin{array}{l} x'_0 < x'_1 < x'_3, \\ x'_0 < x'_2 < x'_3; \\ x'_3 < x'_1 < x'_0, \\ x'_3 < x'_2 < x'_0. \end{array} \right. \text{ or }$$

Also, since segments of c -inertial paths are carried by ϕ_2 into segments of c' -inertial paths, we have:

$$(15) \quad |Z'_{11} - Z'_2| \leq |Z'_{11} - Z'_{13}| + |Z'_{13} - Z'_2| < c' |x'_1 - x'_3| + c' |x'_3 - x'_2|$$

and

$$(16) \quad |Z'_{11} - Z'_2| \leq |Z'_{11} - Z'_{10}| + |Z'_{10} - Z'_2| < c' |x'_1 - x'_0| + c' |x'_0 - x'_2|.$$

We obtain from (14), (15), and (16):

$$(17) \quad 2|Z'_{11} - Z'_2| < c' [|x'_3 - x'_1| + |x'_1 - x'_0| + |x'_3 - x'_2| + |x'_2 - x'_0|] \\ < 2c' |x'_3 - x'_0|.$$

Thus from (13) and (17) we conclude that

$$|Z'_{11} - Z'_2| < 2c'\epsilon^*,$$

and from (13) and (14) that

$$|x'_1 - x'_2| < 2\epsilon^*;$$

and since $\epsilon^* = \epsilon/[2(1 + c')]$, we infer that

$$|\phi_2(Z_1, x_1) - \phi_2(Z_2, x_2)| < \epsilon,$$

which establishes (I).

We now establish:

(II) ϕ_2 carries parallel segments of c -inertial paths into parallel segments of c' -inertial paths.

It is clearly sufficient to show that ϕ_2 carries parallel c -inertial paths into parallel segments of c' -inertial paths. Let η_1 and η_2 be two parallel c -inertial paths, and let η_3 be a c -inertial path which intersects η_1 and η_2 in the points A_1 and A_2 , respectively (obviously such a c -inertial path η_3 exists). (See Figure 1, on following page.) As previously, we use a prime to designate the image under ϕ_2 of a point, line, and so on. We may construct a fourth c -inertial path which intersects η_3 between A_1 and A_2 and which intersects η_1 and η_2 at

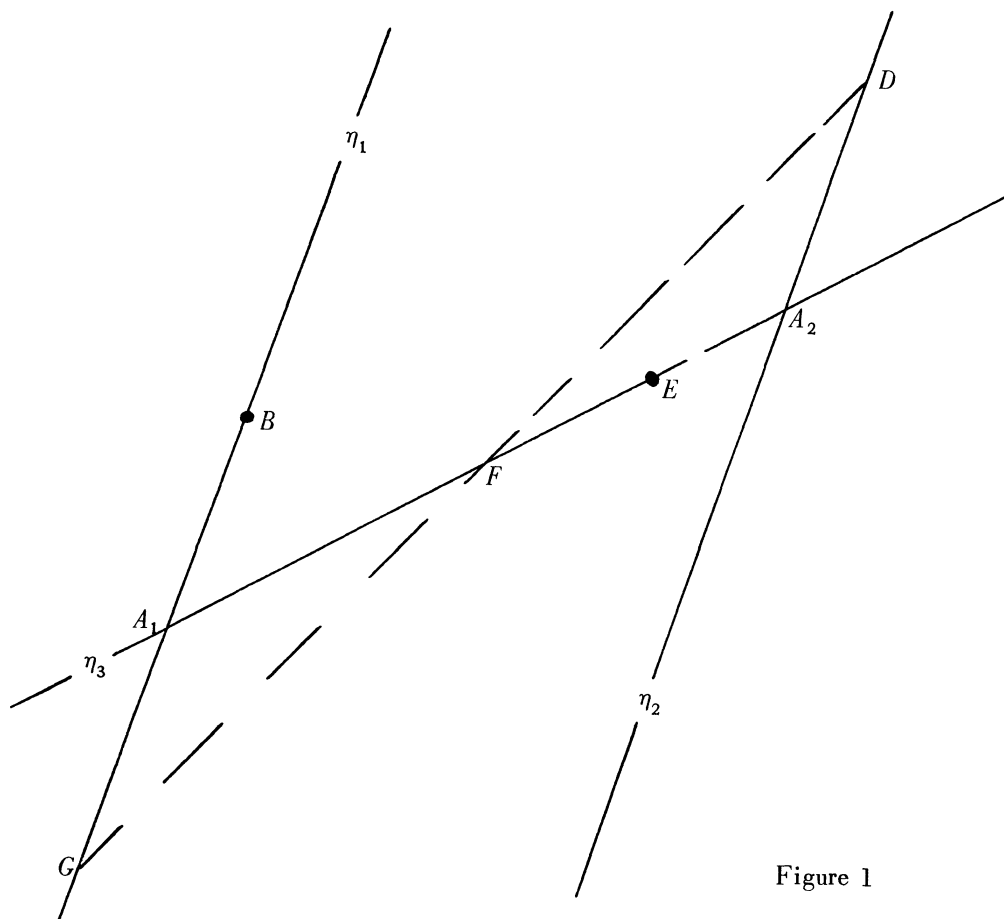


Figure 1

points distinct from A_1 and A_2 . Consequently, we infer from (2) that the segments η'_1 and η'_2 lie in the same plane in the image space of ϕ_2 . (See Figure 2 on following page.) Suppose now that η'_1 and η'_2 are not parallel. We extend (if necessary) η'_1 and η'_2 to their point of intersection, say J' . We next select B' on η'_1 between J' and A'_1 (we use "between" in such a way that B' must be distinct from J' and A'_1); similarly, we select D' on η'_2 between J' and A'_2 . We now consider the pre-images, B and D , of B' and D' . Since ϕ_2 is one-to-one and continuous, it is clear that B and D must be on the same side of η_3 ; that is, the segment BD does not intersect η_3 . Let E be a point on η_3 between A_1 and A_2 . Then, since η_3 is a c -inertial path, one of the numbers $[A_2]_{n+1} - [E]_{n+1}$ and $[A_1]_{n+1} - [E]_{n+1}$ is positive, and the other is negative. Since η_1 and η_2 are parallel, $[B]_{n+1} - [A_1]_{n+1}$ and $[D]_{n+1} - [A_2]_{n+1}$ have the same sign. We then construct a line through D' parallel to η'_1 or through B' parallel

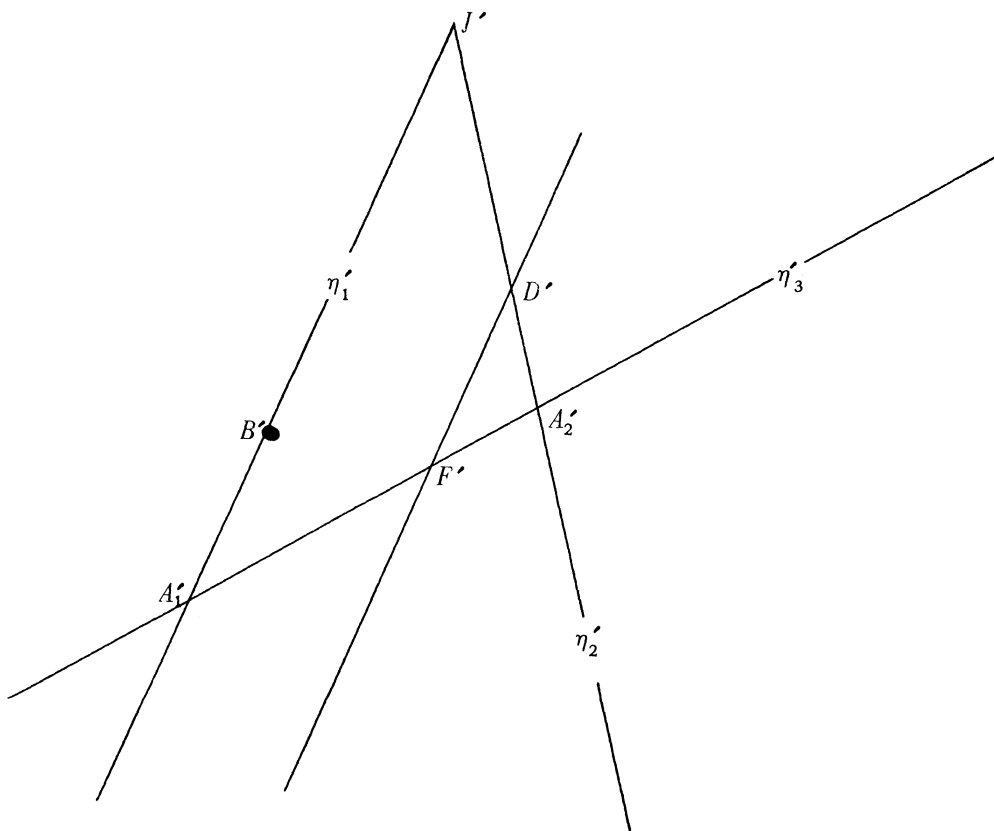


Figure 2

to η'_2 , according to whether $[A_2]_{n+1} - [E]_{n+1}$ or $[A_1]_{n+1} - [E]_{n+1}$ agrees in sign with $[B]_{n+1} - [A]_{n+1}$. Suppose, for definiteness, (see Figure 1) that $[A_2]_{n+1} - [E]_{n+1}$ agrees in sign and that this sign is positive. Let F' be the point of intersection of η'_3 with the line through D' parallel to η'_1 . By construction F' is between A'_1 and A'_2 , and thus F is between A_1 and A_2 .

We then have:

$$\begin{aligned} |[D]_{1,\dots,n} - [F]_{1,\dots,n}| &\leq |[D]_{1,\dots,n} - [A_2]_{1,\dots,n}| + |[A_2]_{1,\dots,n} - [F]_{1,\dots,n}| \\ &< c([D]_{n+1} - [A_2]_{n+1}) + c([A_2]_{n+1} - [F]_{n+1}) \\ &< c([D]_{n+1} - [F]_{n+1}). \end{aligned}$$

Hence the line through D and F is a c -inertial path. This line intersects η_1 at a

point, say G , and, furthermore, by construction DFG is a segment of a c -inertial path, and hence the image $D'F'G'$ is a segment of a c' -inertial path. But $D'F'$ is parallel to η'_1 , and the image of G does not lie on the extension of $D'F'$, which is a contradiction. Thus η_1 and η_2 are parallel, and the proof of (II) is complete.

We next show that

(III) ϕ_2 carries the midpoint of any finite segment α of a c -inertial path into the midpoint of α' .

We consider a fixed plane containing α and a line parallel to the t -axis (the $(n+1)$ st-coordinate axis). In this plane we construct, with α as a diagonal, a parallelogram whose sides and other diagonal are segments of c -inertial paths. Let the speed of the c -inertial path containing α be k . It is clear that through any point of our fixed plane there are exactly two lines with speed l , for every positive number l . Obviously, we may construct a parallelogram \mathcal{P} with α as one diagonal, with the other diagonal a segment of a c -inertial path with speed $(1/4)(3k+c)$, and with one side a segment of a c -inertial path with speed $(1/2)(k+c)$. The other side of the parallelogram \mathcal{P} is then a segment of a c -inertial path with speed $(1/6)(5k+c)$. We conclude from (II) that \mathcal{P} is carried by ϕ_2 into a parallelogram \mathcal{P}' , and the diagonals of \mathcal{P} are carried into the diagonals of \mathcal{P}' . Hence the midpoint of α is carried into the midpoint of α' and (III) is established.

We next show that

(IV) ϕ_2 carries arbitrary lines into lines.

Let α be an arbitrary line in E_{n+1} , and let $\langle Z_1, x_1 \rangle$ and $\langle Z_2, x_2 \rangle$ be any two points on α . We now construct an "inertial" parallelogram through these two points. For definiteness, we assume:

$$x_1 \geq x_2.$$

We set

$$Z_0 = \frac{Z_1 + Z_2}{2},$$

and we choose x_0 and x_3 so that:

$$x_0 < x_2 - \frac{|Z_1 - Z_2|}{2c}, \quad x_3 > x_1 + \frac{|Z_1 - Z_2|}{2c},$$

$$|\langle Z_1, x_1 \rangle - \langle Z_0, x_3 \rangle| = |\langle Z_0, x_0 \rangle - \langle Z_2, x_2 \rangle|,$$

$$|\langle Z_1, x_1 \rangle - \langle Z_0, x_0 \rangle| = |\langle Z_0, x_3 \rangle - \langle Z_2, x_2 \rangle|.$$

Let (see Figure 3)

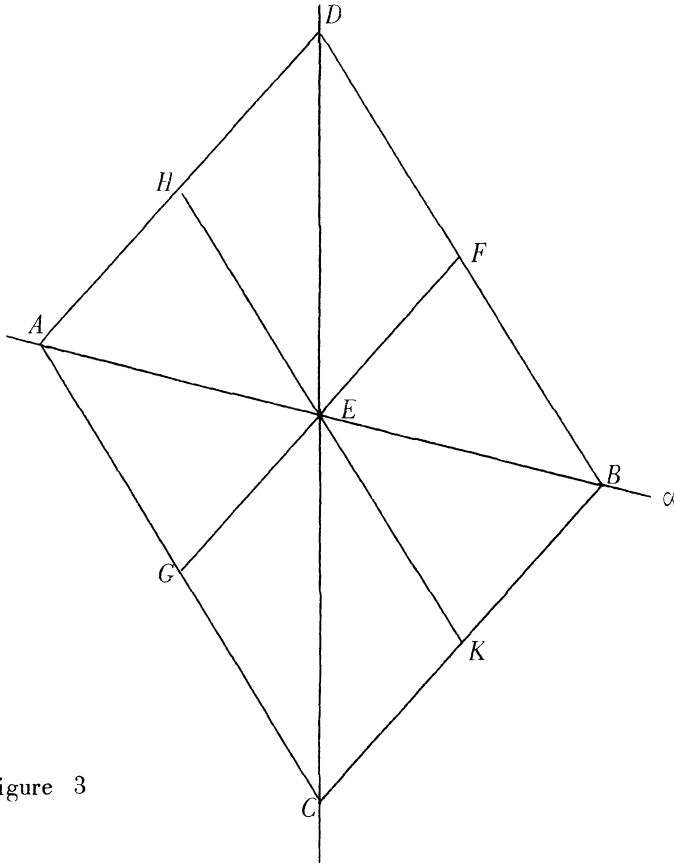


Figure 3

$$A = \langle Z_1, x_1 \rangle, \quad B = \langle Z_2, x_2 \rangle, \quad C = \langle Z_0, x_0 \rangle, \quad D = \langle Z_0, x_3 \rangle,$$

$$E = \frac{A + B}{2}, \quad F = \frac{B + D}{2}, \quad G = \frac{A + C}{2}, \quad H = \frac{A + D}{2}, \quad K = \frac{B + C}{2}.$$

Since the sides of the parallelogram $ACBD$ are by construction segments of c -inertial paths, we conclude from (II) that $A'C'B'D'$ is a parallelogram, where $A' = \phi_2(A)$, and so on, and that the sides of $A'C'B'D'$ are segments of c' -inertial paths. Moreover, it is clear that by construction CED , FEG , and HEK are segments of c -inertial paths, and consequently $C'E'D'$, $F'E'G'$, and

$H'E'K'$ are segments of c' -inertial paths. Hence, by (III), F' , G' , H' , and K' are the midpoints of the respective sides of $A'C'B'D'$. Thus E' , the point of intersection of the segments $F'G'$ and $H'K'$, is the point of intersection of the diagonals of $A'C'B'D'$. Consequently, E' is the midpoint of the segment $A'B'$. Since midpoints of finite segments are carried into midpoints of finite segments, and ϕ_2 is continuous, the proof of (IV) is complete.

From (IV) and the fact that ϕ_2 is one-to-one and continuous, we immediately infer that ϕ_2 is a projective transformation, and since it takes no finite point into a point at infinity, we conclude that

(18) ϕ_2 is a nonsingular affine transformation; that is, for every point $\langle Z, x \rangle$ in E_{n+1} ,

$$(19) \quad \phi_2(Z, x) = \langle Z, x \rangle \mathcal{Q} + B,$$

where \mathcal{Q} is a nonsingular matrix of order $n+1$ and B is an $(n+1)$ -dimensional vector.

Now let

$$(20) \quad \mathcal{Q} = \begin{pmatrix} \mathcal{Q} & E^* \\ F & g \end{pmatrix} \quad \text{and} \quad B = \langle B_1, b \rangle,$$

where \mathcal{Q} is a matrix of order n ; B_1 , E , and F are n -dimensional vectors; and b and g are real numbers. Then

$$(21) \quad \phi_2(Z, x) = \langle Z \mathcal{Q} + xF + B_1, ZE^* + gx + b \rangle.$$

Let α be a c -line such that, for any two distinct points $\langle Z_1, x_1 \rangle$ and $\langle Z_2, x_2 \rangle$ of α ,

$$\frac{Z_1 - Z_2}{x_1 - x_2} = W.$$

Obviously, $|W| = c$. Now α is carried by ϕ_2 into a line α' . We want to show that α is carried into a c' -line. From (21) it follows that the slope W' of α' is given by

$$(22) \quad W' = \frac{W\mathcal{Q} + F}{WE^* + g}.$$

By the hypothesis of our theorem,

$$(23) \quad |W'| \geq c'.$$

Consider now a sequence of c -inertial lines $\alpha_1, \alpha_2, \dots$, whose slopes W_1, W_2, \dots are such that

$$\lim_{i \rightarrow \infty} W_i = W.$$

From (21) and the hypothesis of our theorem we have

$$|W'_i| = \left| \frac{W_i \mathfrak{D} + F}{WE^* + g} \right| < c'.$$

Hence, if $WE^* + g \neq 0$, then

$$(24) \quad |W'| = \left| \frac{W \mathfrak{D} + F}{WE^* + g} \right| = \left| \lim_{i \rightarrow \infty} \frac{W_i \mathfrak{D} + F}{W_i E^* + g} \right| \leq c'.$$

Suppose now that $WE^* + g = 0$. Then

$$\lim_{i \rightarrow \infty} (W_i E^* + g) = 0,$$

and therefore

$$\lim_{i \rightarrow \infty} (W_i \mathfrak{D} + F) = 0.$$

Hence

$$W \mathfrak{D} + F = 0,$$

but then

$$\langle W, 1 \rangle \mathfrak{Q} = 0,$$

which is impossible, since \mathfrak{Q} is nonsingular. Thus we have

$$(25) \quad |W'| = c'.$$

For subsequent use we observe that for any S.R.P.M., $\langle P, \mathfrak{V}, m, s, f, c \rangle$, and any p in P and t in $\mathfrak{V}(p)$,

$$(26) \quad v(p, t)E^* + g \neq 0.$$

For $v(p, t) \neq 0$, the argument is the same as above; in case $v(p, t) = 0$ for some t , on the supposition that $v(p, t)E^* + g = 0$, we must have $g = 0$ and $F = 0$, which again contradicts the nonsingularity of \mathcal{Q} .

From (22) and (25) we get

$$\frac{W\mathbb{D}\mathbb{D}^*W^* + 2W\mathbb{D}F^* + |F|^2}{(WE^* + g)^2} = c'^2,$$

and hence

$$(27) \quad W(\mathbb{D}\mathbb{D}^* - c'^2E^*E)W^* + 2W(\mathbb{D}F^* - c'^2E^*g) + |F|^2 - c'^2g^2 = 0.$$

Since (27) holds for an arbitrary c -line, we may replace W by $-W$, and thus conclude that

$$W(\mathbb{D}F^* - c'^2E^*g) = 0.$$

Therefore, since the direction of W is arbitrary,

$$(28) \quad \mathbb{D}F^* = c'^2E^*g.$$

In view of the fact that (26) holds for $v(p, t) = 0$, we have

$$g \neq 0,$$

and we may then obtain, from (28),

$$(29) \quad E^* = \frac{\mathbb{D}F^*}{c'^2g}.$$

Using (20) and (29), we obtain

$$|\mathcal{Q}| = g \left| \mathbb{D} \left(\mathbb{D} - \frac{F^*F}{c'^2g^2} \right) \right|;$$

and since \mathcal{Q} is nonsingular, we have

$$(30) \quad |F|^2 - c'^2g^2 \neq 0.$$

From (27), (28), and (30) it follows that

$$\mathbb{D}\mathbb{D}^* - c'^2 E^* E \neq 0.$$

Thus, from (27), we have

$$W(\mathbb{D}\mathbb{D}^* - c'^2 E^* E)W^* = \frac{c'^2 g^2 - |F|^2}{c^2} |W|^2.$$

Using again the fact that the direction of W is arbitrary, we infer that

$$(31) \quad \mathbb{D}\mathbb{D}^* - c'^2 E^* E = \left(\frac{c'^2 g^2 - |F|^2}{c^2} \right) \mathbb{D} = \mu \mathbb{D},$$

where $\mu = [c'^2 g^2 - |F|^2]/c^2$. From (28) and (31) we obtain

$$(32) \quad \mathcal{Q} \begin{pmatrix} \mathbb{D} & 0 \\ 0 & -c'^2 \end{pmatrix} \mathcal{Q}^* = \begin{pmatrix} \mathbb{D}\mathbb{D}^* - c'^2 E^* E & \mathbb{D}F^* - c'^2 E^* g \\ (\mathbb{D}F^* - c'^2 E^* g)^* & FF^* - c'^2 g^2 \end{pmatrix} \\ = \begin{pmatrix} \mu \mathbb{D} & 0 \\ 0 & -\mu c^2 \end{pmatrix} = \mu \begin{pmatrix} \mathbb{D} & 0 \\ 0 & -c^2 \end{pmatrix}.$$

We next want to show that μ is positive. Let $\langle Z_1, x_1 \rangle$ and $\langle Z_2, x_2 \rangle$ be two points in E_{n+1} such that

$$\frac{|Z_1 - Z_2|}{|x_1 - x_2|} < c,$$

and let

$$V = \langle Z_1, x_1 \rangle - \langle Z_2, x_2 \rangle$$

and

$$V' = V\mathcal{Q}.$$

From (32) we obtain

$$V\mathcal{Q} \begin{pmatrix} \mathbb{D} & 0 \\ 0 & -c'^2 \end{pmatrix} \mathcal{Q}^* V^* = \mu V \begin{pmatrix} \mathbb{D} & 0 \\ 0 & -c^2 \end{pmatrix} V^*.$$

Hence

$$\mu = \frac{V' \begin{pmatrix} \mathfrak{J} & 0 \\ 0 & -c'^2 \end{pmatrix} (V')^*}{V \begin{pmatrix} \mathfrak{J} & 0 \\ 0 & -c^2 \end{pmatrix} V^*}.$$

By the hypothesis on V ,

$$V \begin{pmatrix} \mathfrak{J} & 0 \\ 0 & -c^2 \end{pmatrix} V^* < 0,$$

and from the fact that c -inertial paths are carried into c' -inertial paths, we have

$$V' \begin{pmatrix} \mathfrak{J} & 0 \\ 0 & -c'^2 \end{pmatrix} (V')^* < 0.$$

Thus μ is positive since it is the ratio of two negative numbers. We set

$$(33) \quad \lambda = \sqrt{\mu}.$$

We then conclude from (32), (33), Definition 1, and Lemma 2 that

$$(34) \quad \mathfrak{U} \text{ is a generalized Lorentz matrix with respect to } \langle c, c', \lambda \rangle.$$

We now turn to the function ϕ_3 which transforms the forces. In deducing the form of ϕ_3 it will be convenient to make use of the functions τ , q , and f^{rel} defined in § 3 (in the course of the present proof we obtain their transformation properties). It is also useful to introduce the function H defined by the following equation for every p in P and t in $\mathfrak{P}(p)$,

$$H(p, t) = [\phi_2(s(p, t), t)]_{n+1}.$$

We thus have that, for t' the element in $\mathfrak{P}'(p)$ corresponding to t in $\mathfrak{P}(p)$,

$$H_p(t) = t'.$$

We obtain, from (21),

$$(35) \quad D(H_p)(t) = v_p(t)E^* + g.$$

For any S.R.P.M. $\Gamma = \langle P, \mathfrak{P}, m, s, f, c \rangle$, the following equation is a direct consequence of Axiom A7 and the appropriate definitions (for any p in P and t in $\mathfrak{P}(p)$):

$$(36) \quad m(p) \frac{d^2 q_p}{d\tau_p^2}(t) = \sum_{i=1}^{\infty} f^{\text{rel}}(p, t, i);$$

and also, under the hypothesis of our theorem,

$$(37) \quad \phi_1(m(p)) \frac{d^2 q_p'}{d\tau_p'^2}(H_p(t)) = \sum_{i=1}^{\infty} f^{\text{rel}'}(p, H_p(t), i).$$

We now obtain the relationship between

$$\frac{d^2 q_p'}{d\tau_p'^2}(H_p(t)) \quad \text{and} \quad \frac{d^2 q_p}{d\tau_p^2}(t).$$

Using (35), we obtain

$$(38) \quad \begin{aligned} \frac{d(\tau_p' \circ H_p)}{d\tau_p}(t) &= \left(\frac{D(\tau_p' \circ H_p)}{D\tau_p} \right)(t) \\ &= \frac{[(D\tau_p')(H_p(t))][DH_p(t)]}{(D\tau_p)(t)} \\ &= \frac{\sqrt{1 - |v_p'(H_p(t))|^2/c'^2} (v_p(t)E^* + g)}{\sqrt{1 - |v_p(t)|^2/c^2}}. \end{aligned}$$

It is easy to show that

$$(39) \quad v_p'(H_p(t)) = \frac{v_p(t)\mathbb{J} + F}{v_p(t)E^* + g};$$

hence, using (39) and squaring (38), we get

$$(40) \quad \left[\left(\frac{D(\tau_p' \circ H_p)}{D\tau_p} \right)(t) \right]^2 = \frac{c^2}{c'^2} \left[\frac{c'^2(v_p(t)E^* + g)^2 - (v_p(t)\mathbb{J} + F)^2}{c^2 - |v_p(t)|^2} \right].$$

Using (27) to give us the expansion of the right member of (40) and then using (32) to simplify the result, we obtain

$$(41) \quad \left[\frac{D(\tau_p' \circ H_p)}{D\tau_p} (t) \right]^2 = \frac{c^2}{c'^2} \left[\frac{-v_p(t)\lambda^2 \Im(v_p(t))^* + \lambda^2 c^2}{c^2 - |v_p(t)|^2} \right] = \frac{c^2}{c'^2} \lambda^2 ;$$

hence

$$(42) \quad D(\tau_p' \circ H_p)(t) = \frac{\delta c \lambda}{c'} (D\tau_p)(t),$$

where $\delta^2 = 1$. We have, from (21) and Definition 3,

$$q_p'(H_p(t)) = q_p(t) \mathcal{U} + B,$$

and thus

$$(43) \quad [D(q_p' \circ H_p)(t)] = (Dq_p)(t) \mathcal{U}.$$

Since

$$\left[\frac{D(q_p' \circ H_p)}{D(\tau_p' \circ H_p)} \right] (t) = \frac{(Dq_p')(H_p(t))(DH_p)(t)}{(D\tau_p')(H_p(t))(DH_p)(t)} = \left(\frac{Dq_p'}{D\tau_p'} \right) (H_p(t)),$$

it is easily shown that

$$(44) \quad \frac{d^2 q_p'}{d\tau_p'^2} (H_p(t)) = \frac{D[D(q_p' \circ H_p)/D(\tau_p' \circ H_p)]}{D(\tau_p' \circ H_p)} (t).$$

From (42), (43), and (44) we infer that

$$(45) \quad \frac{d^2 q_p'}{d\tau_p'^2} (H_p(t)) = \frac{D[(Dq_p) \mathcal{U}]}{(\delta c \lambda / c')(D\tau_p)} (t) = \frac{c'^2}{c^2 \lambda^2} \frac{d^2 q_p}{d\tau_p^2} (t) \mathcal{U}.$$

Now let X and Y be any two vectors in E_n with $X \neq 0$ and $|Y| < c$. Then we set:

$$P = \{1\}, \quad m(1) = 1, \quad Z = \left(1 - \frac{Y^2}{c^2}\right) \left(X - \frac{(X \cdot Y)Y}{c^2}\right),$$

$$\mathfrak{U}(1) = \left(-\frac{c - |Y|}{|Z|}, \frac{c - |Y|}{|Z|}\right),$$

and for all t in $\mathfrak{V}(1)$,

$$s_1(t) = tY + \frac{1}{2} t^2 Z,$$

$$f(1, t, 1) = \frac{Z}{1 - |v_1(t)|^2/c^2} + \frac{v_1(t)(Z \cdot v_1(t))}{c^2(1 - |v_1(t)|^2/c^2)^2},$$

$$f(1, t, i) = 0 \quad \text{for } i > 1.$$

It is easy to verify that $\Gamma_{XY} = \langle P, \mathfrak{V}, m, s, f, c \rangle$ is a S.R.P.M., and consequently so is $\langle \Phi(\Gamma_{XY}), c' \rangle$. Thus there is a positive number γ such that

$$\phi_1(m(1)) = \gamma.$$

We note next that at $t = 0$:

$$s_1(0) = 0, \quad (Ds_1)(0) = Y, \quad (D^2s_1)(0) = Z, \quad \text{and } f(1, 0, 1) = X.$$

We thus have from (37), for $t = 0$,

$$\gamma \left[\frac{d^2 q_1'}{d\tau_1'^2} (H_1(0)) \right] = \phi_3(X, Y),$$

and thus, from (45),

$$\phi_3(X, Y) = \frac{\gamma c'^2}{c^2 \lambda^2} \left[\frac{d^2 q_1}{d\tau_1^2} (0) \hat{G} \right]_{1, \dots, n};$$

hence, from (36),

$$(46) \quad \phi_3(X, Y) = \frac{\gamma c'^2}{c^2 \lambda^2} \left[\left\langle X, \frac{X \cdot Y}{c^2} \right\rangle \hat{G} \right]_{1, \dots, n}.$$

In view of (1), (46) also holds for $X = 0$.

Now let x be any positive real number. Then we set:

$$\delta = \langle 1, \dots, 1 \rangle, \quad P = \{1\}, \quad m(1) = x, \quad \mathfrak{V}(1) = (-c, c);$$

for t in $\mathfrak{V}(1)$,

$$s_1(t) = \frac{1}{2} t^2 \delta,$$

$$f(1, t, 1) = \frac{x\delta}{1 - |v_1(t)|^2/c^2} + \frac{xv_1(t)(\delta \cdot v_1(t))}{c^2(1 - |v_1(t)|^2/c^2)^2},$$

$$f(1, t, i) = 0 \text{ for } i > 1$$

We easily verify that $\Gamma_x = \langle P, \mathfrak{I}, m, s, f, c \rangle$ is a S.R.P.M. such that for all t in $\mathfrak{I}(1)$,

$$\sum_{i=1}^{\infty} f^{\text{rel}}(1, t, i) \neq 0.$$

Furthermore, we infer from (36), (37), and (45) that, for every t in $\mathfrak{I}(1)$,

$$(47) \quad \sum_{i=1}^{\infty} f^{\text{rel}'}(1, H_1(t), i) = \frac{\phi_1(x)}{x} \frac{c'^2}{c^2 \lambda^2} \sum_{i=1}^{\infty} f^{\text{rel}}(1, t, i) \mathcal{Q};$$

hence, from (46),

$$(48) \quad \phi_1(x) = \gamma x.$$

Our theorem now follows from (19), (33), (34), (46), and (48).

REMARK 3. We want to emphasize the physically reasonable nature of the hypothesis of the theorem just proved. We have assumed that systems of relativistic mechanics are carried by our transformations into systems of relativistic mechanics and that light lines are not carried into particle paths. No assumptions concerning the continuity of either ϕ_1 , ϕ_2 , or ϕ_3 have been made. Our assumption that ϕ_2 is one-to-one may be justified physically by the argument that any two space-time positions of a particle distinct with respect to one observer must be distinct with respect to every observer.

The standard presentations of the special theory of relativity vary a good deal in their "derivations" of the Lorentz transformations. Almost without exception, however, the assumptions underlying these derivations are not clearly and completely stated. For the physicist who wants to begin with a set of axioms for relativistic particle mechanics with respect to a fixed coordinate system, our Theorem 2 provides a rigorous approach to the derivation of the

Lorentz transformations. The transformations we obtain in Theorem 2 are, of course, more general than the Lorentz transformations, but it is obvious how the hypothesis of Theorem 2 may be strengthened so as to obtain just the ordinary Lorentz transformations.

Theorem 2 is also pertinent to discussions of the relativity of size, (see, for example, [4]), since the determination of ϕ_1 , ϕ_2 , and ϕ_3 tells us exactly how the system of units of measurement may be changed in passing from one inertial frame of reference to another.

It is interesting to note that the set of transformations admissible (that is, satisfying the hypothesis of Theorem 2) in relativistic particle mechanics differs sharply from the set of those admissible (see the hypothesis of Theorem 3 of [8]) in classical particle mechanics: in the latter case, but not in the former, admissible transformations can change the unit of distance differently along different coordinates (with correspondingly different changes in the unit of force). Thus, although classical mechanics can in a certain sense be regarded as a limiting case of relativistic mechanics, the set of transformations admissible in classical mechanics is in no sense a limit of the set of transformations admissible in relativistic mechanics.

5. Algebraic structure of the set of admissible transformations. Let Φ be an eligible transformation which satisfies the hypothesis of Theorem 2 with respect to the positive real numbers c and c' . We then call the ordered triple $\langle \Phi, c, c' \rangle$ an *admissible triple*; and, corresponding to the informal usage at the end of the previous section, we call an eligible transformation an *admissible transformation* if it is the first element of some admissible triple. Since the set of admissible transformations is not a group under the obvious operation of composition, it is natural to ask what is its algebraic structure. We shall show that the structure of the set of admissible triples is that of a Brandt groupoid (formally defined below). Roughly speaking, the main difference between Brandt groupoids and groups is that a Brandt groupoid is not assumed to be closed under the binary operation corresponding to the group operation. Consequently, a Brandt groupoid may contain many identity elements, that is, many elements e such that $x * e = x = e * x$ whenever x , $x * e$, and $e * x$ are in the groupoid. If there is an e in the groupoid such that, for all x in the groupoid, $e * x = x = x * e$, then the groupoid is also a group. For this reason, we introduced the notion of an admissible triple: the admissible transformation which carries every S.R.P.M. into itself is an identity element whose composition with every admissible transformation is defined; consequently, the set of admissible transformations is neither a group nor a Brandt groupoid.

The notion of a Brandt groupoid was first defined in [1]; we use the formal definition given in [5].

DEFINITION 4. An algebraic system $\mathcal{G} = \langle G, *, J, {}^{-1} \rangle$ (where $*$ is an operation on a subset of $U \times U$ to U , J is a subset of U and ${}^{-1}$ is an operation on U to U) is called a *Brandt groupoid* if and only if the following conditions are satisfied:

- (i) For x, y, z in G , if $x * y \in G$ and $y * z \in G$, then $(x * y) * z \in G$ and $(x * y) * z = x * (y * z)$.
- (ii) For x, y, z in G , if $x * y \in G$ and $x * y = x * z$, then $y = z$.
- (iii) For x, y, z in G , if $x * z \in G$ and $x * z = y * z$, then $x = y$.
- (iv) For x in J , $x * x = x$.
- (v) For x in G , $x^{-1} * x \in J$ and $x * x^{-1} \in J$.
- (vi) For x, z in J , there exists a y in G such that $x * y \in G$ and $y * z \in G$.

Rather than deal directly with admissible triples, we find it somewhat simpler to use the following representation. From Theorem 2 we conclude that to each admissible triple there corresponds a unique ordered sextuple $\langle \mathcal{Q}, B, \gamma, \lambda, c, c' \rangle$; where B is an $(n+1)$ -dimensional vector, γ, λ, c , and c' are positive real numbers, and \mathcal{Q} is a generalized Lorentz matrix (of order $n+1$) with respect to $\langle c, c', \lambda \rangle$. Such an ordered sextuple $\langle \mathcal{Q}, B, \gamma, \lambda, c, c' \rangle$ we shall call a *carrier*. From Theorem 1, together with Theorem 2, it then follows that there is a one-to-one correspondence between the set of carriers and the set of admissible triples.

We say that the carrier $\langle \mathcal{Q}', B', \gamma', \lambda', c_1, c_2 \rangle$ is *left-conformable* to the carrier $\langle \mathcal{Q}, B, \gamma, \lambda, c_3, c_4 \rangle$ if and only if $c_1 = c_4$. By the *conformable subset* \mathcal{D} of $K \times K$ we mean the set of ordered pairs of elements of K such that the first element is left-conformable to the second.

We now define what we call *the carrier system*.

DEFINITION 5. By *the carrier system* we mean the ordered quadruple $\mathcal{K} = \langle K, *, J, {}^{-1} \rangle$, where:

- (i) K is the set of all carriers;
- (ii) $*$ is the operation on \mathcal{D} to K such that if the carrier $\langle \mathcal{Q}', B', \gamma', \lambda', c_1, c_2 \rangle$ is left-conformable to the carrier $\langle \mathcal{Q}, B, \gamma, \lambda, c_3, c_4 \rangle$ then

$$\langle \mathcal{Q}', B', \gamma', \lambda', c_1, c_2 \rangle * \langle \mathcal{Q}, B, \gamma, \lambda, c_3, c_4 \rangle = \langle \mathcal{Q}\mathcal{Q}', B\mathcal{Q}' + B', \gamma\gamma', \lambda\lambda', c_3, c_2 \rangle ;$$

(iii) J is the set of carriers of the form $\langle \mathfrak{I}, 0, 1, 1, c, c \rangle$, where \mathfrak{I} is the identity matrix of order $n + 1$; and

(iv) $^{-1}$ is the operation on K to K such that if $\langle \mathcal{Q}, B, \gamma, \lambda, c, c' \rangle \in K$ then

$$\langle \mathcal{Q}, B, \gamma, \lambda, c, c' \rangle^{-1} = \langle \mathcal{Q}^{-1}, -B\mathcal{Q}^{-1}, 1/\gamma, 1/\lambda, c', c \rangle.$$

We have then the following theorem, the proof of which we omit.

THEOREM 3. *The carrier system is a Brandt groupoid.*

We remark first that the operation $*$ of the carrier system corresponds to the composition of admissible triples; that is, if $\langle \Phi, c, c' \rangle$ corresponds to $\langle \mathcal{Q}, B, \gamma, \lambda, c, c' \rangle$, and $\langle \Psi, c', c'' \rangle$ corresponds to $\langle \mathcal{Q}', B', \gamma', \lambda', c', c'' \rangle$, then $\langle \mathcal{Q}\mathcal{Q}', B\mathcal{Q}' + B', \gamma\gamma', \lambda\lambda', c, c'' \rangle$ corresponds to $\langle \theta, c, c'' \rangle$, where $\langle \theta, c, c'' \rangle$ is the admissible triple such that, for any S.R.P.M. Γ_c ,

$$\langle \Psi(\langle \Phi(\Gamma_c), c' \rangle), c'' \rangle = \langle \theta(\Gamma_c), c'' \rangle.$$

Similarly, the inverse operation $^{-1}$ of the carrier system corresponds to the natural inverse operation on admissible triples; that is, if $\langle \Phi, c, c' \rangle$ corresponds to $\langle \mathcal{Q}, B, \gamma, \lambda, c, c' \rangle$, and $\langle \Psi, c', c \rangle$ corresponds to $\langle \mathcal{Q}^{-1}, -B\mathcal{Q}^{-1}, 1/\gamma, 1/\lambda, c', c \rangle$, then, for any S.R.P.M. Γ_c ,

$$\langle \Psi(\langle \Phi(\Gamma_c), c' \rangle), c \rangle = \Gamma_c.$$

It thus follows as a corollary to Theorem 3 that the set of admissible triples is a Brandt groupoid under the natural operations of composition and formation of inverses.

It is natural to ask how the hypothesis of Theorem 2 may be strengthened so that the set of eligible transformations satisfying it form a group. We state without proof some results concerning this question.

THEOREM 4. *Let $\Phi = \langle \phi_1, \phi_2, \phi_3 \rangle$ be an eligible transformation which carries every system of relativistic particle mechanics into a system of relativistic particle mechanics. Then there exist positive real numbers δ , γ , λ , and ρ , an $(n + 1)$ -dimensional vector B , an orthogonal matrix \mathfrak{E} of order n , and a matrix \mathcal{Q} of order $n + 1$, such that*

$$\delta^2 = 1, \quad \mathcal{Q} = \lambda \begin{pmatrix} \mathfrak{I} & 0 \\ 0 & \rho \end{pmatrix} \begin{pmatrix} \mathfrak{E} & 0 \\ 0 & \delta \end{pmatrix};$$

and for any vectors Z_1 and Z_2 in E_n , any x in R , and y in R^+ ,

$$\phi_1(y) = \gamma y,$$

$$\phi_2(Z_1, x) = \langle Z_1, x \rangle \hat{Q} + B,$$

$$\phi_3(Z_1, Z_2) = \frac{\gamma \rho^2 Z_1 \mathcal{E}}{\lambda}.$$

The interpretation of δ , γ , λ , B , and \mathcal{E} is the same as that stated in Remark 1. The number ρ is the ratio c/c' of the absolute values of the old and new velocities of light. The matrix \hat{Q} is a generalized Lorentz matrix with $U = 0$, which intuitively means that the old and new spatial frames of reference are at rest with respect to each other. The fact that the hypothesis of Theorem 4 thus excludes the possibility of transforming from one inertial frame of reference to another moving with respect to it is sufficient reason to regard this hypothesis as unnecessarily strong from the point of view of our intended physical interpretation. On the other hand, it is, of course, clear that the set of transformations satisfying this hypothesis constitute a group under the obvious operations.

REFERENCES

1. H. Brandt, *Über eine Verallgemeinerung des Gruppenbegriffes*, Math. Ann. 96 (1926-1927), 360-366.
2. P. Destouches-Février, *La structure des théories physiques*, Paris, 1951.
3. H. Hermes, *Eine Axiomatisierung des allgemeinen Mechanik*, Heft 3, Forschungen zur Logik und zur Grundlegung der Exakten Wissenschaften, Neue Folge, Leipzig, 1938.
4. B. Hoffmann, *The relativity of size*, Physical Rev. 89 (1953), 49-52.
5. B. Jónsson and A. Tarski, *Boolean algebras with operators, Part II*, Amer. J. Math. 74 (1952), 127-162.
6. L. A. MacColl, *Geometric aspects of relativistic dynamics*, Trans. Amer. Math. Soc. 46 (1939), 328-347.
7. J. C. C. McKinsey, A. C. Sugar, and P. Suppes, *Axiomatic foundations of classical particle mechanics*, J. Rational Mech. Anal. 2 (1953), 253-272.
8. J. C. C. McKinsey and P. Suppes, *Transformations of systems of classical particle mechanics*, J. Rational Mech. Anal. 2 (1953), 273-289.
9. ———, *Philosophy and the axiomatic foundations of physics*, Proceedings of the XIth International Congress of Philosophy, vol. VI (1953), 49-54.
10. K. Menger, *Are variables necessary in calculus*, Amer. Math. Monthly 56 (1949), 609-620.

11. H. Reichenbach, *Axiomatik der relativistischen Raum-Zeit-Lehre*, Braunschweig, 1924.
12. A. A. Robb, *Geometry of time and space*, Cambridge, 1936.
13. K. Schnell, *Eine Topologie der Zeit in logistischer Darstellung*, Dissertation, Münster, 1937.
14. A. G. Walker, *Foundations of relativity* I, II, Proc. Royal Society of Edinburgh, Section A, 62 (1948), 319-335.

STANFORD UNIVERSITY

PACIFIC JOURNAL OF MATHEMATICS

EDITORS

M.M. SCHIFFER*

Stanford University
Stanford, California

R.P. DILWORTH

California Institute of Technology
Pasadena 4, California

E. HEWITT

University of Washington
Seattle 5, Washington

E.F. BECKENBACH**

University of California
Los Angeles 24, California

ASSOCIATE EDITORS

H. BUSEMANN

P.R. HALMOS

BØRGE JESSEN

J.J. STOKER

HERBERT FEDERER

HEINZ HOPF

PAUL LÉVY

E.G. STRAUS

MARSHALL HALL

R.D. JAMES

GEORGE PÓLYA

KÔSAKU YOSIDA

SPONSORS

UNIVERSITY OF BRITISH COLUMBIA

CALIFORNIA INSTITUTE OF TECHNOLOGY

UNIVERSITY OF CALIFORNIA, BERKELEY

UNIVERSITY OF CALIFORNIA, DAVIS

UNIVERSITY OF CALIFORNIA, LOS ANGELES

UNIVERSITY OF CALIFORNIA, SANTA BARBARA

UNIVERSITY OF NEVADA

OREGON STATE COLLEGE

UNIVERSITY OF OREGON

UNIVERSITY OF SOUTHERN CALIFORNIA

STANFORD RESEARCH INSTITUTE

STANFORD UNIVERSITY

WASHINGTON STATE COLLEGE

UNIVERSITY OF WASHINGTON

* * *

AMERICAN MATHEMATICAL SOCIETY

HUGHES AIRCRAFT COMPANY

Mathematical papers intended for publication in the *Pacific Journal of Mathematics* should be typewritten (double spaced), and the author should keep a complete copy. Manuscripts may be sent to any of the editors. Manuscripts intended for the outgoing editors should be sent to their successors. All other communications to the editors should be addressed to the managing editor, E.G. Straus, at the University of California Los Angeles 24, California.

50 reprints of each article are furnished free of charge; additional copies may be obtained at cost in multiples of 50.

The *Pacific Journal of Mathematics* is published quarterly, in March, June, September, and December. The price per volume (4 numbers) is \$12.00; single issues, \$3.50; back numbers (Volumes 1, 2, 3) are available at \$2.50 per copy. Special price to individual faculty members of supporting institutions and to individual members of the American Mathematical Society: \$4.00 per volume; single issues, \$1.25.

Subscriptions, orders for back numbers, and changes of address should be sent to the publishers, University of California Press, Berkeley 4, California.

Printed at Ann Arbor, Michigan. Entered as second class matter at the Post Office, Berkeley, California.

* To be succeeded in 1955, by H.L. Royden, Stanford University, Stanford, California.

** To be succeeded in 1955, by E.G. Straus, University of California, Los Angeles 24, Calif.

UNIVERSITY OF CALIFORNIA PRESS • BERKELEY AND LOS ANGELES

COPYRIGHT 1954 BY PACIFIC JOURNAL OF MATHEMATICS

Pacific Journal of Mathematics

Vol. 4, No. 4

August, 1954

Paul Civin, <i>Orthonormal cyclic groups</i>	481
Kenneth Lloyd Cooke, <i>The rate of increase of real continuous solutions of algebraic differential-difference equations of the first order</i>	483
Philip J. Davis, <i>Linear functional equations and interpolation series</i>	503
F. Herzog and G. Piranian, <i>Sets of radial continuity of analytic functions</i> . . .	533
P. C. Rosenbloom, <i>Comments on the preceding paper by Herzog and Piranian</i>	539
Donald G. Higman, <i>Remarks on splitting extensions</i>	545
Margaret Jackson, <i>Transformations of series of the type ${}_3\Psi_3$</i>	557
Herman Rubin and Patrick Colonel Suppes, <i>Transformations of systems of relativistic particle mechanics</i>	563
A. Seidenberg, <i>On the dimension theory of rings. II</i>	603
Bertram Yood, <i>Difference algebras of linear transformations on a Banach space</i>	615