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## **THE SLOW STEADY MOTION OF LIQUID PAST A SEMI-ELLIPTICAL BOS**

G. POWER AND D. L. SCOTT-HUTTON

# THE SLOW STEADY MOTION OF LIQUID PAST A SEMI-ELLIPTICAL BOSS

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**1. Introduction.** In this problem of two-dimensional viscous flow, liquid is supposed to have a rigid boundary represented by  $ABCDE$  in Figure 1 and, apart from the disturbance caused by the presence of the elliptical boss  $BCD$ , is assumed to be in uniform shearing motion. The stream function is thus a biharmonic function vanishing together with its normal derivative at all points of the boundary, and proportional to  $y^2$  at a great distance from the boss. A series of functions is found, each of which satisfies all the boundary conditions save one. A linear combination of these functions will also satisfy the boundary conditions with this one exception, and by a particular choice of the arbitrary constants which it contains, the remaining condition can be satisfied at as many points as desired. Special cases are discussed, and a process of approximation is outlined which yields the most accurate results at  $C$ , and also gives a convenient function for determining at any point of the boundary the magnitude of the error in the unsatisfied boundary condition. A special case of this problem has previously been considered [1].

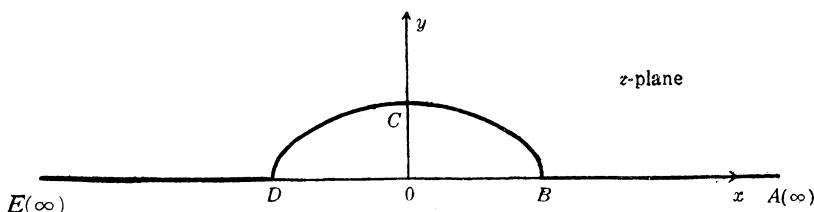


Figure 1.

**2. The stream function.** We take the equation of the boundary  $BCD$  to be  $x^2/a^2 + y^2/b^2 = 1$ , and note that the region occupied by the fluid, for which  $y$  is never negative, is transformed into the interior of the semi-circle of unit radius shown in Figure 2 by

$$(1) \quad -2z = (a-b)w + (a+b)/w.$$

The stream function  $\psi$  is biharmonic, that is to say it must satisfy  $\nabla^4\psi=0$ , and a satisfactory solution to the problem is

$$(2) \quad \psi = y^2 + U + yV,$$

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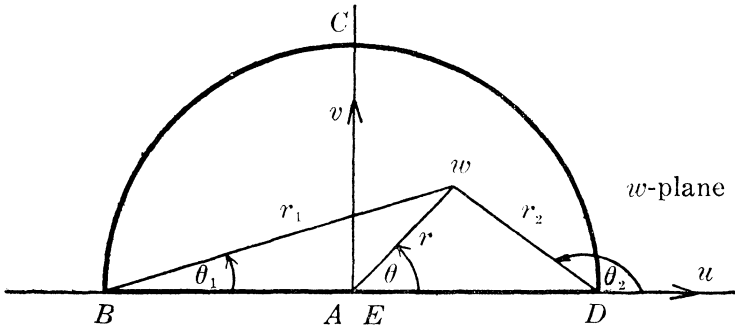


Figure 2.

provided  $U$  and  $V$  are harmonic functions which are chosen so that  $U+yV$  does not tend to infinity as  $z$  tends to infinity.

The boundary conditions to be satisfied are  $\psi=0$  and  $\partial\psi/\partial\nu=0$  along the boundary, where  $\delta\nu$  denotes an element of normal to the boundary. From (2), we see that four conditions are required, namely

$$(3) \begin{cases} (a) & U=0 & \text{along } AB \text{ and } DE, \text{ that is when } y=0, \\ (b) & V+\partial U/\partial y=0 & \text{along } AB \text{ and } DE, \text{ that is when } y=0, \\ (c) & y^2+U+yV=0 & \text{along } BCD, \text{ that is when } x^2/a^2+y^2/b^2=1, \\ (d) & (\partial/\partial\nu)(y^2+U+yV)=0 & \text{along } BCD, \text{ that is when } x^2/a^2+y^2/b^2=1. \end{cases}$$

Writing  $w=re^{i\theta}=u+iv$ , and using the transformation (1), we see that the boundary conditions (3) become, after a little reduction,

$$(4) \begin{cases} (a) & U=0, & \text{when } v=0, \\ (b) & (b-a)V' + \frac{2u^2}{(c+u^2)} \cdot \frac{\partial U}{\partial v} = 0, & \text{when } v=0, \\ (c) & U+b \sin \theta V' = 0, & \text{when } r=1, \\ (d) & b(a+b) \sin \theta = \frac{1}{\sin \theta} \frac{\partial U}{\partial r} + b \frac{\partial V'}{\partial r} - aV', & \text{when } r=1, \end{cases}$$

where  $V'=V+br \sin \theta$ , and  $c=(b+a)/(b-a)$ .

We will proceed to find pairs of harmonic functions  $U_{2n-1}, V'_{2n-1}$ , such that each pair will satisfy exactly the first three of the above boundary conditions. Any linear combination of these functions will also satisfy exactly these three conditions, and by giving special values to the arbitrary constants in this linear combination we can satisfy approximately the fourth equation. Physically, this means that in the fluid motion represented by our solution there will be a small velocity of slip along the boss  $BCD$ , which can be calculated from the error involved in the last boundary condition.

If we take

$$(5) \quad U = \sum_{n=1}^q a_{2n-1} U_{2n-1},$$

where

$$(6) \quad U_{2n-1} = \mathcal{S} \left\{ \frac{A_{2n-1}}{2n+1} w^{2n+1} + \frac{B_{2n-1}}{2n-1} w^{2n-1} + \frac{C_{2n-1}}{2n-3} w^{2n-3} \right\},$$

then  $U$  is harmonic, and 4 (a) is satisfied. Moreover, the consideration of symmetry shows that even powers of  $w$  are not required. Now we have

$$(7) \quad \left( \frac{\partial U_{2n-1}}{\partial v} \right)_{v=0} = (u^2 + c) \left( A_{2n-1} u^{2n-2} + \frac{C_{2n-1}}{c} u^{2n-4} \right),$$

provided  $cA_{2n-1} - B_{2n-1} + C_{2n-1}/c = 0$ ,

and

$$(8) \quad (U_{2n-1})_{r=1} = \sin \theta \left( \frac{2A_{2n-1}}{2n+1} \cos 2n\theta - \frac{2C_{2n-1}}{2n-3} \cos (2n-2)\theta \right),$$

provided  $\frac{A_{2n-1}}{2n+1} + \frac{B_{2n-1}}{2n-1} + \frac{C_{2n-1}}{2n-3} = 0$ .

From (7), (8), we see that we can take

$$(9) \quad \left\{ \begin{aligned} A_{2n-1} &= \frac{2(2n+1)}{b-a} \{2(n-1)b+a\} p_n, \\ B_{2n-1} &= \frac{4(2n-1)}{(b-a)^2} \{b^2+(2n-1)ab+a^2\} p_n, \\ C_{2n-1} &= -\frac{2(2n-3)(b+a)}{(b-a)^2} \{2nb+a\} p_n, \end{aligned} \right.$$

where the unknown  $p_n$  has yet to be determined. By setting

$$V_{2n-1} = W_{2n-1} + \frac{2C_{2n-1} r^{2n-2} \cos (2n-2)\theta}{(2n-3)b} - \frac{2A_{2n-1} r^{2n} \cos 2n\theta}{(2n+1)b}$$

we see that 4 (c) yields

$$(10) \quad W_{2n-1} = 0 \quad \text{when } r=1.$$

If further we take

$$(11) \quad p_n = \frac{b(b-a)^2}{4(2nb+a)\{2(n-1)b+a\}}$$

then condition 4 (b) gives

$$(12) \quad W_{2n-1} = u^{2n-2} - u^{2n} \quad \text{when } v=0.$$

Equations (9), (11) now give

$$(13) \quad \left\{ \begin{aligned} A_{2n-1} &= \frac{(2n+1)b(b-a)}{2(2nb+a)}, \\ B_{2n-1} &= \frac{(2n-1)b\{b^2 + (2n-1)ab + a^2\}}{(2nb+a)\{2(n-1)b+a\}}, \\ C_{2n-1} &= -\frac{(2n-3)b(b+a)}{2\{2(n-1)b+a\}}. \end{aligned} \right.$$

To find the function  $W_{2n-1}$  satisfying (10), (12), we consider  $\phi_{2n} = \mathcal{R}\{\chi_{2n}(w)\}$ , where

$$(14) \quad \left\{ \begin{aligned} \chi_{2n}(w) &= -w^{-2n} + \frac{i}{\pi} \left\{ w^{2n} + \frac{1}{w^{2n}} \right\} \{ \log(w+1) - \log(w-1) \} \\ &- \frac{2i}{\pi} \left\{ \left( w^{2n-1} + \frac{1}{w^{2n-1}} \right) + \frac{1}{3} \left( w^{2n-3} + \frac{1}{w^{2n-3}} \right) + \dots + \frac{1}{2n-1} \left( w + \frac{1}{w} \right) \right\}, \\ \chi_0(w) &= -1 + \frac{2i}{\pi} \{ \log(w+1) - \log(w-1) \}. \end{aligned} \right.$$

It is easy to verify that  $\mathcal{R}\{\chi_{2n}(w)\} = 0$  when  $v=1$ , and that  $\mathcal{R}\{\chi_{2n}(w)\} = u^{2n}$  when  $v=0$ , since from Figure 2

$$\log(w-1) - \log(w+1) = \log \frac{r_2}{r_1} + i(\theta_2 - \theta_1).$$

The function  $W_{2n-1}$  is thus given by

$$(15) \quad W_{2n-1} = \phi_{2n-2} - \phi_{2n}.$$

Finally, we see that the required stream function  $\psi$  is given by equation (2), where

$$(16) \quad \left\{ \begin{aligned} U &= \mathcal{I} \sum_{n=1}^q a_{2n-1} \left\{ \frac{A_{2n-1}}{2n+1} w^{2n+1} + \frac{B_{2n-1}}{2n-1} w^{2n-1} + \frac{C_{2n-1}}{2n-3} w^{2n-3} \right\}, \\ V &= \mathcal{R} \left\{ ibw + \sum_{n=1}^q a_{2n-1} \left\{ \frac{2C_{2n-1}}{(2n-3)b} w^{2n-2} - \frac{2A_{2n-1}}{(2n+1)b} w^{2n} \right. \right. \\ &\quad \left. \left. + \chi_{2n-1}(w) - \chi_{2n}(w) \right\} \right\}, \end{aligned} \right.$$

the constants  $A_{2n-1}$ ,  $B_{2n-1}$ ,  $C_{2n-1}$  being determined from (13),  $\chi_{2n}(w)$  from (14). It is quite easy to verify that  $\psi \rightarrow y^2$  as  $z \rightarrow \infty$ , that is as  $w \rightarrow 0$ , since the most significant terms in  $U$  and  $yV$  are respectively

$2a_1C_1y/(a+b)$  and  $a_1y(-2C_1/b+1)$ , the sum of these being clearly zero from (13).

**3. The fourth boundary condition.** It is now necessary to consider the boundary condition as yet unsatisfied, given by equation 4 (d), in the form

$$(17) \quad b(a+b) \sin \theta = \sum_{n=1}^q a_{2n-1} \lambda_{2n-1}(\theta) \quad \text{when } r=1,$$

where

$$\lambda_{2n-1}(\theta) = \frac{1}{\sin \theta} \frac{\partial U_{2n-1}}{\partial r} + b \frac{\partial V'_{2n-1}}{\partial r} - a V'_{2n-1}.$$

Theoretically, the constants  $a_{2n-1}$  must be chosen so that (17) is satisfied for all values of  $\theta$  lying between 0 and  $\pi$ , and this would require an infinite number of terms. Clearly, therefore, some form of approximation must be applied. Suppose that the constants  $a_{2n-1}$  are chosen so that

$$\sum_{n=1}^q a_{2n-1} \lambda_{2n-1} = b(a+b) \sin \theta + F(\theta),$$

then  $\sin \theta F(\theta)$  is the error involved in the boundary derivative  $(\partial\psi/\partial r)_{r=1}$ , and the actual velocity of slip on the boundary  $BCD$  in the  $z$ -plane is

$$\left\{ \left( \frac{\partial\psi}{\partial r} \right) \frac{dw}{dz} \right\}_{r=1} = \frac{\sin \theta F(\theta)}{(b^2 \cos^2 \theta + a^2 \sin^2 \theta)^{1/2}}.$$

This becomes infinite at  $\theta=\pi/2$  in the degenerate case  $a=0$ , unless  $F(\pi/2)$  happens to be zero. Therefore we must consider a method of approximation which gives no error at all when  $\theta=\pi/2$ . The coefficients  $a_{2n-1}$  will be chosen so that the expressions on each side of (17) have the same values and the same differential coefficients with regard to  $\theta$  when  $\theta=\pi/2$ , the number of differential coefficients that can be equated depending on the value of the integer  $q$ . From (16), we see that when  $r=1$ ,

$$(18) \quad \lambda_{2n-1}(\theta) = \frac{1}{\sin \theta} \{ A_{2n-1} \sin(2n+1)\theta + B_{2n-1} \sin(2n-1)\theta + C_{2n-1} \sin(2n-3)\theta \} \\ + 2 \left\{ \frac{2(n-1)b-a}{(2n-3)b} \right\} C_{2n-1} \cos(2n-2)\theta - 2 \left\{ \frac{2nb-a}{(2n+1)b} \right\} A_{2n-1} \cos 2n\theta \\ + b(f_{2n-2}(\theta) - f_{2n}(\theta)),$$

where

$$\begin{aligned}
 f_{2n}(\theta) &= \left( \frac{\partial \phi_{2n}}{\partial r} \right)_{r=1} = \mathcal{R} \left\{ w \frac{d\chi_{2n}(w)}{dw} \right\}_{r=1} \\
 &= \frac{8n}{\pi} \left\{ \sin (2n-1)\theta + \frac{1}{3} \sin (2n-3)\theta + \dots + \frac{1}{2n-1} \sin \theta \right\} \\
 &\quad + 2n \cos 2n\theta + \frac{4n}{\pi} \sin 2n\theta \log \tan \frac{\theta}{2} - \frac{2}{\pi} \csc \theta, \\
 f_0(\theta) &= -\frac{2}{\pi} \csc \theta.
 \end{aligned}$$

It is to be noted, that although  $f_{2n}(\theta)$  is infinite at  $\theta=0, \pi$ , the expression  $f_{2n-2}(\theta) - f_{2n}(\theta)$ , which occurs in  $\lambda_{2n-1}(\theta)$  is finite at these points. Equation (17) is satisfied exactly when  $\theta=0, \pi$ , and putting  $\theta=\pi/2$ , we have

$$(19) \quad b(a+b) = \sum_{n=1}^q a_{2n-1} \lambda_{2n-1}(\pi/2),$$

and by differentiation we are led to

$$(19') \quad \begin{cases} -b(a+b) = \sum_{n=1}^q a_{2n-1} \lambda'_{2n-1}(\pi/2), \\ b(a+b) = \sum_{n=1}^q a_{2n-1} \lambda''_{2n-1}(\pi/2), \end{cases}$$

and so forth.

It is from this set of equations that the constants  $a_{2n-1}$  are to be calculated.

**4. Special cases.** The two special cases of the semi-circular boss and projection will now be discussed.

( $\alpha$ ) *semi-circular boss.*

Setting  $a=b=1$ , equation (16) yields

$$\begin{aligned}
 U &= \sum_{n=1}^q \frac{a_{2n-1}}{2n-1} \{ r^{2n-1} \sin (2n-1)\theta - r^{2n-3} \sin (2n-3)\theta \}, \\
 V &= -r \sin \theta + \sum_{n=1}^q a_{2n-1} (\phi_{2n-2} - \phi_{2n}) - \sum_{n=1}^q \frac{2a_{2n-1}}{2n-1} r^{2n-2} \cos (2n-2)\theta,
 \end{aligned}$$

and from (18), we get when  $r=1$

$$\lambda_{2n-1}(\theta) = \frac{1}{\sin \theta} \left\{ \frac{2}{2n-1} \sin (2n-1)\theta \right\} + f_{2n-2}(\theta) - f_{2n}(\theta).$$

As an example, let us take  $q=3$ , so that from (19) we are required to solve the equations

$$\begin{aligned}
 2 &= a_1 \lambda_1(\pi/2) + a_3 \lambda_3(\pi/2) + a_5 \lambda_5(\pi/2) , \\
 -2 &= a_1 \lambda_1''(\pi/2) + a_3 \lambda_3''(\pi/2) + a_5 \lambda_5''(\pi/2) , \\
 2 &= a_1 \lambda_1^{IV}(\pi/2) + a_3 \lambda_3^{IV}(\pi/2) + a_5 \lambda_5^{IV}(\pi/2) .
 \end{aligned}$$

By substitution and straightforward calculation we obtain the following table :-

	$n=1$	$n=2$	$n=3$
$\lambda_{2n-1}(\pi/2)$	+1.45352	-0.72488	+0.38385
$\lambda_{2n-1}''(\pi/2)$	-0.36056	+5.18309	-9.65706
$\lambda_{2n-2}^{IV}(\pi/2)$	-1.10423	-22.24210	+191.53800

which leads directly to  $a_1 = +1.21058$ ,  $a_3 = -0.34379$ ,  $a_5 = -0.02299$ .

A more accurate result can be obtained by taking more terms in the linear expression for  $\psi$ , and it is found that the coefficients  $a_{2n-1}$  decrease rapidly in magnitude, but the numerical work involved soon becomes exceedingly heavy.

This choice of approximation method is seen to advantage if calculating the error function

$$F(\theta) = \sum_{n=1}^3 a_{2n-1} \lambda_{2n-1}(\theta) - 2 \sin \theta$$

at any point by means of the Taylor expansion about  $\theta = \pi/2$ , several of the significant differential coefficients being zero by definition. The following table gives the value of  $F(\theta)$  for various values of  $\theta$ , and Figure 3 shows the graph of  $F(\theta)$  plotted against values of  $\theta$  lying between 0 and  $\pi/2$ . The graph for  $\pi/2 \leq \theta \leq \pi$  will, of course, be similar, since  $F(\theta)$  is symmetrical about  $\theta = \pi/2$ .

$\theta$	0	$\pi/45$	$\pi/15$	$\pi/8$	$\pi/4$	$3\pi/8$	$\pi/2$
$F(\theta)$	0	0.05918	0.08225	0.05838	0.00915	0.00041	0

It will be noticed that the results are most accurate in the vicinity of  $\theta = \pi/2$ , as might be expected from the method of approximation. Although the value of  $F(\theta)$  becomes greater than 0.08 for a certain (small) values of  $\theta$ , the velocity of slip, given by  $\sin \theta \cdot F(\theta)$ , is really very small at these points.

( $\beta$ ) *degenerate boss.*

If  $a=0$ , the semi-elliptical boss degenerates to a projection into



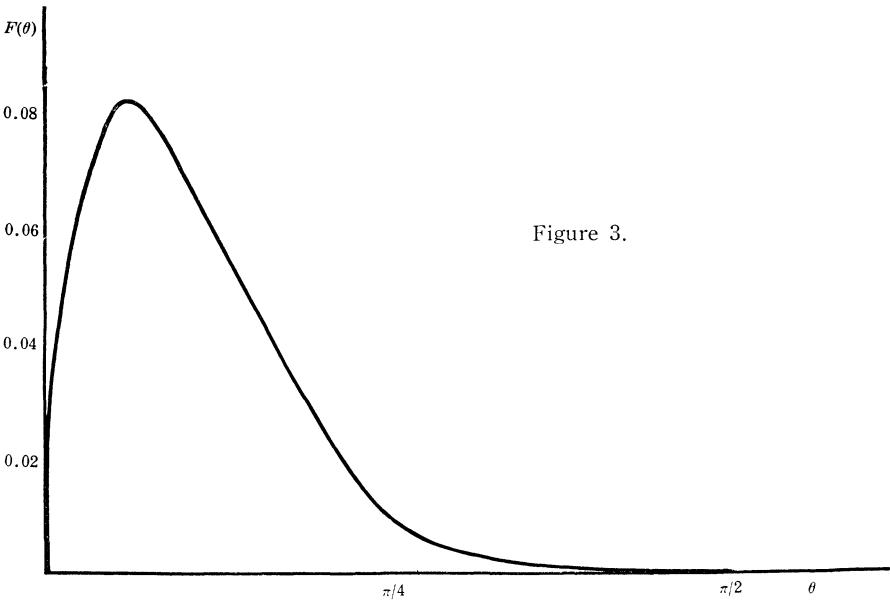


Figure 3.

the moving liquid, and the formulae (13) become

$$A_{2n-1} = \frac{(2n+1)b}{4n}, \quad B_{2n-1} = \frac{(2n-1)b}{4n(n-1)}, \quad C_{2n-1} = -\frac{(2n-3)b}{4(n-1)}.$$

These values will hold except for  $n=1$ . For this case we will follow W. R. Dean [1], and will find a pair of solutions  $U_1, V_1'$  which satisfy equations 4 (a), (b), (c). The procedure outlined in § 2 is again followed, and omitting details we are led to the two pairs of solutions

$$U_1 = \mathcal{S}bw, \quad V_1' = -\mathcal{R}\left\{ \frac{2w}{i+w} \right\},$$

and

$$\mathcal{S} \frac{b}{4} \{w^3 + 3w\}, \quad U_1 = \frac{w^3}{4} + w, \quad V_1' = \mathcal{R} \left\{ -1 - \frac{1}{2}w^2 + \chi_0(w) - \chi_2(w) \right\}.$$

The final solution is thus

$$(20) \left\{ \begin{aligned} U &= \mathcal{S} \left\{ a_0bw + \frac{a_1b}{4}(w^3 + 3w) + \sum_{n=2}^q a_{2n-1}b \left( \frac{w^{2n+1}}{4n} + \frac{w^{2n-1}}{4n(n-1)} - \frac{w^{2n-3}}{4(n-1)} \right) \right\}, \\ V &= \mathcal{R} \left\{ ibw + a_1 \left( -1 - \frac{1}{2}w^2 \right) + \sum_{n=2}^q a_{2n-1} \left( \frac{w^{2n-2}}{2(n-1)} + \frac{w^{2n}}{2n} \right) \right\} \\ &\quad + \mathcal{R} \left\{ -\frac{2a_0w}{i+w} + \sum_{n=1}^q a_{2n-1} (\chi_{2n-2}(w) - \chi_{2n}(w)) \right\}. \end{aligned} \right.$$

Again we note that  $\psi \rightarrow y^2$  as  $z \rightarrow \infty$ .

**5. The pressure equation.** The pressure  $p$  is determined from the equations of motion in the form

$$\frac{\partial p}{\partial x} = -\mu \frac{\partial}{\partial y} \nabla^2 \psi, \quad \frac{\partial p}{\partial y} = \mu \frac{\partial}{\partial x} \nabla^2 \psi,$$

where  $\mu$  is the coefficient of viscosity. Now  $\psi = y^2 + U + yV$ , where  $U$  and  $V$  are harmonic, so that  $\nabla^2 \psi = 2 + 2(\partial V / \partial y)$ , and hence

$$\frac{\partial p}{\partial x} = -2\mu \frac{\partial^2 V}{\partial y^2} = 2\mu \frac{\partial^2 V}{\partial x^2}, \quad \frac{\partial p}{\partial y} = 2\mu \frac{\partial^2 V}{\partial x \partial y}.$$

Ignoring an arbitrary constant, we have therefore

$$p = 2\mu \frac{\partial V}{\partial x} = 2\mu \mathcal{R} \left\{ \frac{dV(w)}{dz} \right\} = 2\mu \mathcal{R} \left\{ \frac{dw}{dz} \cdot \frac{dV(w)}{dw} \right\},$$

where  $V = \mathcal{R} \{ V(w) \}$ . From equation (16) we see that

$$V(w) = ibw + \sum_{n=1}^q a_{2n-1} \left\{ \frac{2C_{2n-1}w^{2n-2}}{(2n-3)b} - \frac{2A_{2n-1}w^{2n}}{(2n+1)b} + \chi_{2n-2}(w) - \chi_{2n}(w) \right\},$$

and so

$$(21) \quad p = 2\mu \mathcal{R} \left\{ \left[ \frac{2w^2}{(b+a) + (b-a)w^2} \right] \left[ ib + \sum_{n=1}^q a_{2n-1} \left\{ \frac{4(n-1)C_{2n-1}w^{2n-3}}{(2n-3)b} - \frac{4nA_{2n-1}w^{2n-1}}{(2n+1)b} + \frac{d\chi_{2n-2}(w)}{dw} - \frac{d\chi_{2n}(w)}{dw} \right\} \right] \right\}.$$

Equation (21) gives the pressure distribution, and on the plane boundary, where  $v=0$ , this becomes

$$(22) \quad p = \left\{ \frac{4\mu u^2}{(b+a) + (b-a)u^2} \right\} \left\{ \sum_{n=1}^q a_{2n-1} \left\{ \frac{4(n-1)C_{2n-1}u^{2n-3}}{(2n-3)b} - \frac{4nA_{2n-1}u^{2n-1}}{(2n+1)b} + 2(n-1)u^{2n-3} - 2nu^{2n-1} \right\} \right\}.$$

In particular the pressure at  $B$  exceeds that at  $D$  by

$$(23) \quad p_{\text{diff}} = \frac{8\mu}{b} \sum_{n=1}^q a_{2n-1} \left\{ -\frac{2(n-1)C_{2n-1}}{(2n-3)b} + \frac{2nA_{2n-1}}{(2n+1)b} + 1 \right\}.$$

For the special case  $a=b=1$  discussed in § 4 ( $\alpha$ ), this expression is

$$8\mu \sum_{n=1}^q \frac{(4n-3)a_{2n-1}}{2n-1}$$

and with the values for  $a_1$ ,  $a_3$ ,  $a_5$  substituted we obtain a difference of  $4.78 \mu$ .

For the degenerate case  $a=0$ ,  $b=1$ , Dean [1] obtains a difference of approximately  $5.80 \mu$  between the pressures at  $B$  and  $D$ .

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*Added in Proof:* The equations governing the slow steady flow of a viscous incompressible fluid are the same as those characterizing an equilibrium state of an incompressible elastic solid, if one simply replaces velocity and coefficient of viscosity by displacement and shear modulus. Thus the results here obtained can be used as the solution for the tension of a semi-infinite plane whose edge is indented and traction free.

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