Pacific Journal of Mathematics

ON THE UNIFORM CONVERGENCE OF A CERTAIN EIGENFUNCTION SERIES

LUNA MISHOE AND G. C. FORD

Vol. 6, No. 2 December 1956

ON THE UNIFORM CONVERGENCE OF A CERTAIN EIGENFUNCTION SERIES

L. I. MISHOE AND G. C. FORD

1. Introduction. In the attempt to solve certain problems in mathematical physics, such as diffraction of an arbitrary pulse by a wedge as considered by Irvin Kay [1], one encounters a hyperbolic differential equation of the type

$$u_{xx}-q(x)u=u_{xt}-p(x)u_t \tag{a}$$

where $u(x, t)$ must satisfy the boundary conditions $u(1, t)=u(0, t)=0$ and $u(x, 0)=F(x)$. In attempting to solve equation (a) by separation of variables, one is led to the consideration of expanding an arbitrary function $F(x)$ in terms of the eigenfunctions $u_n(x)$ of the equation

$$u''+q(x)u+\lambda(p(x)u-u')=0$$

satisfying the boundary conditions $u(0)=u(1)=0$.

In the previous paper [2] by B. Friedman and L. I. Mishoe, it was proved that a function $F(x)$ of bounded variation for $0\leq x\leq 1$ could be expanded in terms of the eigenfunctions $u_n(x)$ of the system $u''+qu+\lambda(pu-u')=0$, $u(0)=u(1)=0$, provided $F(0^+)+F(1^-)\exp\left(-\int_0^1 pdt\right)=0$. However, the question of uniform convergence of the series $\sum_{-\infty}^{\infty} a_n u_n(x)$ to $F(x)$ was not considered. In this paper we establish sufficient conditions for the series $\sum_{-\infty}^{\infty} a_n u_n(x)$ to converge uniformly to $F(x)$ for $0<x<1$.

The following theorem has already been proved [2]:

THEOREM 1. *Let $F(x)$ be a function of bounded variation for $0\leq x\leq 1$. Let $u_n(x)$ be the eigenfunctions of the system*

$$(A+\lambda B)u=0\ ; \qquad u(0)=u(1)=0\ , \tag{1}$$

where A is the operator $d^2/dx^2+q(x)$, and where B is the operator $-d/dx+p(x)$.

Let $q(x)$ be continuous and $p(x)$ have a continuous second derivative. Furthermore, let $v_n(x)$ be the eigenfunctions of the system adjoint to (1). If

Received February 2, 1955. The research in this document has been made possible through a grant to Morgan State College by the National Science Foundation, Washington, D. C. Research Grant NSF-G 1224, November 1954.

$$F(0^+)+F(1^-)\exp\left(-\int_0^1 p(t)dt\right)=0\ , \tag{2}$$

then the series

$$\sum_{-\infty}^{\infty} a_n u_n(x)\ , \tag{3}$$

where

$$a_n=\int_0^1 F(\xi)\left[\frac{p(\xi)v_n(\xi)+v_n'(\xi)}{C'(\lambda_n)}\right]d\xi\ , \tag{4}$$

and where the Wronskian $\omega(x)$ of the two independent solutions $u_1(x)$ and $u_2(x)$ has the form

$$\omega(x)=u_1u_2'-u_2u_1'=C(\lambda)e^{\lambda x} \tag{5}$$

with

$$C(\lambda)=\lambda^{-1}\exp\left(-\int_0^1 p(t)dt\right)-\exp\left(-\lambda+\int_0^1 p(t)dt\right)+O(\lambda^{-2})\ , \tag{6}$$

converges to $F(x)$ at every point where $F(x)$ is continuous in $0<x<1$. At all other points, the series converges to $\frac{1}{2}(F(x+0)+F(x-0))$. If $F(x)$ does not satisfy the boundary conditions (2), then the series (3) converges to

$$\frac{1}{2}\left[F(x+0)+F(x-0)-\left\{F(0^+)+F(1^-)\exp\left(-\int_0^1 p(t)dt\right)\right\}\exp\left(\int_0^x p(t)dt\right)\right]. \tag{7}$$

In this paper, we prove:

THEOREM 2. *If $F'(x)$ exists and is of bounded variation for $0\leqq x\leqq 1$, then a sufficient condition for the series $\sum_{-\infty}^{\infty} a_n u_n(x)$ to converge uniformly to $F(x)$ for $0<x<1$ is that $F(0)=F(1)=0$.*

2. An asymptotic form for $C'(\lambda_n)$. Using (5) and the boundary conditions $u(0)=u(1)=0$ and $u'(0)=u'(1)=1$, we have

$$C(\lambda)=e^{-\lambda}u_1(1,\lambda)\ . \tag{8}$$

Then it follows that

$$C'(\lambda)=\frac{d}{d\lambda}C(\lambda)=-C(\lambda)+\tfrac{1}{2}e^{-\lambda}u_1(1,\lambda)+e^{-\lambda/2}\frac{d}{d\lambda}w_1(x,\lambda)\qquad \text{at } x=1 \tag{9}$$

where

$$u_1=e^{\lambda x/2}w_1\ . \tag{10}$$

Now (10) transforms the equation $(A+\lambda B)u=0$ into $w_1''+\left(q+\lambda p-\frac{\lambda^2}{4}\right)w_1=0$. It can be verified [2] that w_1 satisfies the equation

$$w_1=\frac{\sinh R(0,x)}{[r(x)r(0)]^{1/2}}-\int_0^x\frac{\sinh R(\xi,t)}{[r(x)r(\xi)]^{1/2}}g(\xi)w_1(\xi)d\xi \tag{11}$$

where

$$g(x)=p^2+\frac{p''}{\lambda-2p}+\frac{3p'^2}{(\lambda-2p)^2}+q \tag{12}$$

and

$$r(x)=\frac{\lambda}{2}-p(x)\ ,\qquad R(\xi,x)=\int_\xi^x r(t)dt\ . \tag{13}$$

We note that $g(x)$ and $g'(x)=\frac{d}{d\lambda}g(x)$ are bounded for $|\lambda|$ sufficiently large. Also, if in (11) we make the substitution

$$w_1=\lambda^{-1}\exp\left(\tfrac{1}{2}|\sigma|x\right)Z_1(x) \tag{14}$$

where $\sigma=\mathscr{R}\lambda$, we note that $Z_1(x)$ is bounded [2] for $|\lambda|$ sufficiently large. Differentiating (11) with respect to λ, we obtain

$$\begin{aligned}w_1'=&\frac{x\cosh R(0,x)}{2[r(x)r(0)]^{1/2}}-\frac{[r(x)+r(0)]\sinh R(0,x)}{4[r(x)r(0)]^{3/2}}\\&-\int_0^x\frac{\sinh R(\xi,x)}{[r(x)r(\xi)]^{1/2}}g(\xi)w_1'(\xi)d\xi-\int_0^x\frac{\sinh R(\xi,x)}{[r(x)r(\xi)]^{1/2}}g'(\xi)w_1(\xi)d\xi\\&-\int_0^x(x-\xi)\frac{\cosh R(\xi,x)}{[r(x)r(\xi)]^{1/2}}g(\xi)w_1(\xi)d\xi\\&+\frac{1}{4}\int_0^x\frac{[r(x)+r(\xi)]\sinh R(\xi,x)}{[r(x)r(\xi)]^{3/2}}g(\xi)w_1(\xi)d\xi\ .\end{aligned} \tag{15}$$

If we substitute

$$w_1=\lambda^{-1}\rho(x)\exp\left(\tfrac{1}{2}|\sigma|x\right)\ , \tag{16}$$

we obtain that

$$\begin{aligned}\rho(x)=&\frac{\lambda x\exp\left(-\frac{1}{2}|\sigma|x\right)\cosh R(0,x)}{[r(x)r(0)]^{1/2}}\\&-\frac{\lambda[r(x)+r(0)]\exp\left(-\frac{1}{2}|\sigma|x\right)\sinh R(0,x)}{4[r(x)r(0)]^{3/2}}\end{aligned} \tag{17}$$

$$-\int_0^x \frac{\exp[-\frac{1}{2}|\sigma|(x-\xi)]\sinh R(\xi,x)}{[r(x)r(\xi)]^{1/2}} g(\xi)\rho(\xi)d\xi$$

$$-\int_0^x \frac{\exp[-\frac{1}{2}|\sigma|(x-\xi)]\sinh R(\xi,x)}{[r(x)r(\xi)]^{1/2}} g'(\xi)Z_1(\xi)d\xi$$

$$-\int_0^x \frac{1}{2}(x-\xi)\frac{\exp[-\frac{1}{2}|\sigma|(x-\xi)]\cosh R(\xi,x)}{[r(x)r(\xi)]^{1/2}} g(\xi)Z_1(\xi)d\xi$$

$$+\int_0^x \frac{[r(x)+r(\xi)]\exp[-\frac{1}{2}|\sigma|(x-\xi)]\sinh R(\xi,x)}{4[r(x)r(\xi)]^{3/2}} g(\xi)Z_1(\xi)d\xi$$

where $g'(\xi)=\frac{d}{d\lambda}g(\xi)$. Now $\lambda[r(x)r(\xi)]^{-1/2}$ and hence $\lambda[r(x)+r(\xi)][r(x)r(\xi)]^{-3/2}$ are both bounded by some constant C as $|\lambda|\to\infty$. Also, $\exp[-\frac{1}{2}|\sigma|(x-\xi)] \times\cosh R(\xi,x)$, and $\exp[-\frac{1}{2}|\sigma|(x-\xi)]\sinh R(\xi,x)$ are both bounded by some constant C as $|\lambda|\to\infty$ and $0\leqq\xi\leqq x$. Using these results we obtain from equation (17) that

$$(18)\qquad |\rho(x)|\leqq 2C^2+\frac{C^2}{|\lambda|}\int_0^x |g(\xi)\rho(\xi)|d\xi+\frac{C^2}{|\lambda|}\int_0^x |g'(\xi)Z_1(\xi)|d\xi + \frac{C^2}{|\lambda|}\int_0^x \frac{1}{2}|x-\xi||g(\xi)Z_1(\xi)|d\xi+\frac{C^2}{|\lambda|}\int_0^x |g(\xi)Z_1(\xi)|d\xi\,.$$

If we set $\mu(\lambda)$ equal to the maximum of $|\rho(x)|$ in $0\leqq x\leqq 1$, then we certainly have that

$$\mu\leqq\frac{2C^2}{1-\frac{C^2}{|\lambda|}\int_0^x |g(\xi)|d\xi}+\frac{\frac{C^2}{|\lambda|}\int_0^x (|g'(\xi)Z_1(\xi)|+\frac{1}{2}|x-\xi|\,|g(\xi)Z_1(\xi)|+|g(\xi)Z_1(\xi)|)d\xi}{1-\frac{C^2}{|\lambda|}\int_0^x |g(\xi)|d\xi}\,.$$

Therefore, μ, and consequently $\rho(x)$ are bounded as $|\lambda|\to\infty$. Rewrite equation (15) as follows:

$$(19)\quad w_1'(x,\lambda)=\frac{x\cosh R(0,x)}{2[r(x)r(0)]^{1/2}}-\frac{[r(x)+r(0)]\sinh R(0,x)}{4[r(x)r(0)]^{3/2}}$$

$$-\lambda^{-1}\exp(\tfrac{1}{2}|\sigma|x)\int_0^x \frac{\exp[-\frac{1}{2}|\sigma|(x-\xi)]\sinh R(\xi,x)}{[r(x)r(\xi)]^{1/2}} g(\xi)\rho(\xi)d\xi$$

$$-\lambda^{-1}\exp(\tfrac{1}{2}|\sigma|x)\int_0^x \frac{\exp[-\frac{1}{2}|\sigma|(x-\xi)]\sinh R(\xi,x)}{[r(x)r(\xi)]^{1/2}} g(\xi)Z_1(\xi)d\xi$$

$$-\lambda^{-1}\exp(\tfrac{1}{2}|\sigma|x)\int_0^x (x-\xi)\frac{\exp[-\frac{1}{2}|\sigma|(x-\xi)]\sinh R(\xi,x)}{2[r(x)r(\xi)]^{1/2}} g(\xi)Z_1(\xi)d\xi$$

$$+\lambda^{-1}\exp(\tfrac{1}{2}|\sigma|x)\int_0^x \frac{\exp[-\frac{1}{2}|\sigma|(x-\xi)][r(x)+r(\xi)]\sinh R(\xi,x)}{4[r(x)r(\xi)]^{3/2}} g(\xi)Z_1(\xi)d\xi\,.$$

The above four integrals are all at least $O(\lambda^{-2}\exp[\frac{1}{2}|\sigma|x])$. Also,

$$[r(x)r(0)]^{-1/2} \doteq \lambda^{-1}+O(\lambda^{-2}) ,$$

and

$$[r(x)]^{-1/2}[r(0)]^{-3/2}=\lambda^{-2}+O(\lambda^{-3}) .$$

So it follows that

$$(20) \quad w_1'(1, \lambda)=\lambda^{-1}\cosh\left(\int_0^1 r(t)dt\right)+O(\lambda^{-2}\exp[\tfrac{1}{2}|\sigma|])-\frac{\lambda^{-2}}{2}\sinh\left(\int_0^1 r(t)dt\right)$$
$$+O(\lambda^{-3}\exp[\tfrac{1}{2}|\sigma|]) .$$

Using this result for $w_1'(1, \lambda)$ in equation (9), we have, for $\mathscr{R}\lambda>0$,

$$(21) \quad C'(\lambda)=\lambda^{-1}\left[\exp\left(-\lambda+\int_0^1 p(t)dt\right)\right]$$
$$-\frac{1}{2}\lambda^{-2}\left[\exp\left(-\int_0^1 p(t)dt\right)-\exp\left(-\lambda+\int_0^1 p(t)dt\right)\right]+O(\lambda^{-3})+O(\lambda^{-2}) ,$$

and for $\mathscr{R}\lambda<0$,

$$(22) \quad C'(\lambda)=\lambda^{-1}\left[\exp\left(-\lambda+\int_0^1 p(t)dt\right)\right]$$
$$-\frac{1}{2}\lambda^{-2}\left[\exp\left(-\int_0^1 p(t)dt\right)-\exp\left(-\lambda+\int_0^1 p(t)dt\right)\right]+O(\lambda^{-3}e^{\lambda}) .$$

3. Distribution of the eigenvalues. Since by [2]

$$C(\lambda)=\lambda^{-1}\exp[-\lambda a]\left(\exp\left[-\int_a^b p(t)dt\right]-\exp\left[-\lambda(b-a)+\int_a^b p(t)dt\right]+O(\lambda^{-1})\right)$$
$$=\lambda^{-1}\exp[-\lambda a]C_1(\lambda) \qquad \text{for } \mathscr{R}\lambda\geqq 0,$$
$$C(\lambda)=\lambda^{-1}\exp[-\lambda b]\left(\exp\left[-\lambda(a-b)-\int_a^b p(t)dt\right]-\exp\left[\int_a^b p(t)dt\right]+O(\lambda^{-1})\right)$$
$$=\lambda^{-1}\exp[-\lambda b]C_2(\lambda) \qquad \text{for } \mathscr{R}\lambda\leqq 0,$$

and where a and b equal 0 and 1 respectively.
The condition that λ be an eigenvalue is that $C(\lambda)$ and hence either $C_1(\lambda)$ or $C_2(\lambda)$ be zero. Equating $C_1(\lambda)$ to zero we obtain

$$(23) \quad \exp\left[-\lambda(b-a)+\int_a^b p(t)dt\right]=\exp\left[-\int_a^b p(t)dt\right]+O(\lambda^{-1})$$
$$=\exp\left[-\int_a^b p(t)dt\right](1+O(\lambda^{-1})) .$$

By taking the logarithm of both sides of the above equation (23) and expanding the term $\log(1+O(\lambda^{-1}))$ we obtain that the large eigenvalues satisfy the equation

$$-\lambda_n=-2\int_0^1 p(t)dt+2n\pi i+O(\lambda_n^{-1}), \quad n=\pm N, \pm N+1, \cdots .$$

Hence the eigenvalues with positive real parts, if they exist, are given by

$$\lambda_n=2n\pi i+2\int_0^1 p(t)dt+O\left(\frac{1}{n}\right) . \tag{24}$$

The equation $C_2(\lambda)=0$ leads to the same result for those eigenvalues with negative real parts. Consequently, all the eigenvalues are represented by equation (24).

4. On the uniform convergence of series (3). Consider equation (4). In [2] it was shown that

$$B^* V_1(x)=\exp\left[-\lambda x+\int_0^x p(t)dt\right]+\Omega_1 , \tag{25}$$

where

$$\Omega_1=\begin{cases} O(\lambda^{-2})+O(\lambda^{-1}\exp[-\lambda x]) & \text{for } \mathscr{R}\lambda\geqq 0 \\ O(\lambda^{-1}\exp[-\lambda x]) & \text{for } \mathscr{R}\lambda\leqq 0. \end{cases}$$

Similarly,

$$B^* V_2(x)=\exp\left[-\lambda x+\int_1^x p(t)dt\right]+\Omega_2 , \tag{26}$$

where

$$\Omega_2=\begin{cases} O(\lambda^{-1}\exp[-\lambda x]) & \text{for } \mathscr{R}\lambda\geqq 0 \\ O(\lambda^{-1}\exp[-\lambda x])+O(\lambda^{-1}\exp[-\lambda b]) & \text{for } \mathscr{R}\lambda\leqq 0. \end{cases}$$

Also from [2], we have that

$$u_1(x)=\lambda^{-1}\left\{\exp\left[\lambda x-\int_0^x p(t)dt\right]-\exp\int_0^x p(t)dt\right]\right\}+O(\lambda^{-2}e^{\lambda x}) \tag{27}$$

and

$$u_2(x)=\lambda^{-1}\left\{\exp\left[\lambda(x-1)-\int_1^x p(t)dt\right]-\exp\int_1^x p(t)dt\right]\right\}+O(\lambda^{-2}) . \tag{28}$$

Using equations (26) and (27), for $\mathscr{R}\lambda>0$, we have

$$a_n u_1 = u_1 \int_0^1 F(\xi) \frac{B^* V_2(\xi)}{C'(\lambda_n)} d\xi$$

$$= O\left(\frac{e^{\lambda_n x}}{\lambda_n}\right) \int_0^1 \frac{F(\xi)\left[\exp\left(-\lambda_n \xi + \int_1^\xi p(t)dt\right)\right] d\xi}{C'(\lambda_n)}$$

$$+ O(\lambda_n^{-1} e^{\lambda_n x}) \int_0^1 \frac{F(\xi) O(\lambda_n^{-1} e^{-\lambda_n \xi})}{C'(\lambda_n)} d\xi$$

$$= \frac{A}{\lambda_n} \int_0^1 \frac{F(\xi) \exp\left(-\lambda_n \xi + \int_1^\xi p(t)dt\right)}{C'(\lambda_n)} d\xi + \frac{A}{\lambda_n} \int_0^1 \frac{F(\xi) O(\lambda_n^{-1} e^{-\lambda_n \xi})}{C'(\lambda_n)} d\xi$$

where A is bounded.

By equation (22), $C'(\lambda_n) = O(\lambda_n^{-1})$, therefore $\frac{1}{C'(\lambda_n)} = O(\lambda_n)$.

Hence

$$a_n u_1 = B_n \int_0^1 F(\xi) \exp\left(-\lambda_n \xi + \int_0^\xi p(t)dt\right) d\xi + B_n \int_0^1 F(\xi) O(\lambda_n^{-1} e^{-\lambda_n \xi}) d\xi \tag{29}$$

where $B_n = \lambda^{-1} O(\lambda_n) A$ is also bounded. Using equation (26) for $\mathscr{R}\lambda > 0$, and observing that $O(\lambda^{-1} \exp(-\lambda_n \xi))$ is the indefinite integral of a bounded function, it can be easily shown that

$$\int_0^1 F(\xi) O(\lambda_n^{-1} e^{-\lambda_n \xi}) d\xi = O(\lambda_n^{-2}) . \tag{30}$$

Consider now the first integral in equation (29). Setting $H(\xi) = F(\xi) \exp\left(\int_1^\xi p(t)dt\right)$ and integrating by parts, we obtain

$$\int_0^1 H(\xi) \exp(-\lambda_n \xi) d\xi = -\lambda_n^{-1} H(\xi) \exp(-\lambda_n \xi) \Big]_0^1 + \int_0^1 \lambda_n^{-1} H'(\xi) \exp(-\lambda_n \xi) d\xi . \tag{31}$$

Since $F(1) = F(0) = 0$, then $H(1) = H(0) = 0$, and the first term on the right hand side of equation (31) vanishes.

Now

$$H'(\xi) = p(\xi) F(\xi) \exp\left(\int_1^\xi p(t)dt\right) + F'(\xi) \exp\left(\int_1^\xi p(t)dt\right) . \tag{32}$$

$F'(\xi)$ is of bounded variation on $(0, 1)$ and $p'(\xi)$ is continuous on $(0, 1)$. Therefore, $H'(\xi)$ is of bounded variation on $(0, 1)$. Hence,

$$H'(\xi) = \varphi_1(\xi) - \varphi_2(\xi)$$

where $\varphi_1(\xi)$ and $\varphi_2(\xi)$ are two bounded, positive, monotone functions, either both nonincreasing or both nondecreasing. Now $\lambda_n^{-1}\exp(-\lambda_n\xi)$ is bounded and integrable for $0\leqq\xi\leqq 1$. Assume $\varphi_1(\xi)$ to be a monotone decreasing function, then

$$(33)\qquad \int_0^1 \lambda_n^{-1}H'(\xi)\exp(-\lambda_n\xi)d\xi$$

$$=\varphi_1(0)\int_0^{\xi_0}\lambda_n^{-1}\exp(-\lambda_n)d\xi-\varphi_2(0)\int_0^{\xi_1}\lambda_n^{-1}\exp(-\lambda_n\xi)d\xi=O(\lambda_n^{-2})$$

where ξ_0 and ξ_1 are on the interval (0,1).

Combining the results of (30) and (33) we have

$$\sum_{-\infty}^{\infty} a_n u_n(x)=\sum_{-\infty}^{-(N+1)} O(\lambda_n^{-2})+\sum_{-N}^{N-1} a_n u_n(x)+\sum_{N}^{\infty} O(\lambda_n^{-2})$$

where $\sum_{-N}^{N-1} a_n u_n(x)$ is finite for $0<x<1$. From (24) it is clear that $\lambda_n=O(n)$ for $n=\pm N,\ \pm N+1,\cdots$ Therefore

$$\sum_{-\infty}^{\infty} a_n u_n(x)=\sum_{-\infty}^{-(N+1)} \frac{O(1)}{n^2}+\sum_{-N}^{N-1} a_n u_n(x)+\sum_{N}^{\infty}\frac{O(1)}{n^2}$$

where $O(1)$ is a bounded function.

Since $\left|\frac{O(1)}{n^2}\right|\leqq\frac{M}{n^2}$, $M>0$ and the series $M\sum_{N}^{\infty}\frac{1}{n^2}$ converges, it is clear that $\sum_{-\infty}^{\infty} a_n u_n(x)$ converges uniformly to $F(x)$ for $0<x<1$. And our theorem is proved.

We note, however, that while Theorem 2 is sufficient, it is not a necessary condition for uniform convergence. For suppose $F(0)$ and $F(1)$ differ from zero, then by equations (31) and (33) we have $\sum_{-\infty}^{\infty} a_n u_n=\sum_{-\infty}^{\infty} O\left(\frac{1}{n}\right)$ which may or may not converge uniformly.

In fact, *a necessary and sufficient condition* for the uniform convergence of this series does not seem to be known.

The authors wish to thank Professor Bernard Friedman and Mr. Bertram Levy of New York University, and Miss Noel Cousins of Morgan State College for their assistance in preparing this paper for publication.

REFERENCES

1. Irvin Kay, *Diffraction of an arbitrary pulse by a wedge*, Mathematics Research Group, New York University, Report No. EM-43.

2. B. Friedman and L. I. Mishoe, *Eigenfunction expansions associated with a non-self-adjoint differential equation*, Pacific J. Math., **6** (1956), 249–270.

MORGAN STATE COLLEGE

PACIFIC JOURNAL OF MATHEMATICS

Mathematical papers intended for publication in the *Pacific Journal of Mathematics* should be typewritten (double spaced), and the author should keep a complete copy. Manuscripts may be sent to any of the editors. Manuscripts intended for the outgoing editors should be sent to their successors. All other communications to the editors should be addressed to the managing editor, Alfred Horn at the University of California, Los Angeles 24, California.

50 reprints of each article are furnished free of charge; additional copies may be obtained at cost in multiples of 50.

The *Pacific Journal of Mathematics* is published quarterly, in March, June, September, and December. The price per volume (4 numbers) is $12.00; single issues, $3.50. Back numbers are available. Special price to individual faculty members of supporting institutions and to individual members of the American Mathematical Society: $4.00 per volume; single issues, $1.25.

Subscriptions, orders for back numbers, and changes of address should be sent to Pacific Journal of Mathematics, c/o University of California Press, Berkeley 4, California.

Printed at Kokusai Bunken Insatsusha (International Academic Printing Co., Ltd.), No. 10, 1-chome, Fujimi-cho, Chiyoda-ku, Tokyo, Japan.

* During the absence of E. G. Straus.

PUBLISHED BY PACIFIC JOURNAL OF MATHEMATICS, A NON-PROFIT CORPORATION

Pacific Journal of Mathematics

Vol. 6, No. 2 December, 1956