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**BOUNDS FOR EIGENVALUES OF THE FREE AND FIXED
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CHAPTER I

Introduction. This paper gives explicit upper and lower bounds for the eigenvalues of both the free and fixed membrane problems in terms of the eigenvalues of analogous finite difference problems. For the fixed membrane we seek eigenvalues of the Laplace operator on a bounded region R of the Euclidean plane under the added condition that the solution must vanish on the boundary C . In the case of the free membrane it is the normal derivative which vanishes on C . If C is sufficiently smooth the difference between the upper and lower bounds is of the order of the grid width h .

Generally upper and lower bounds for the membrane eigenvalues, $\lambda_1 \leq \lambda_2 \leq \dots$, have been obtained separately and by distinctly different methods. Upper bounds for λ_k can always be obtained by the Rayleigh-Ritz process [53]. This depends on the minimum-maximum property of λ_k which was first discovered by Poincaré [46]. Following a suggestion of R. Courant [13] (also appearing implicitly in a paper of L. Collatz [8]) we can express this upper bound in terms of the eigenvalues of a related finite difference problem. Such upper bounds have been obtained by G. Pólya [47] and H. Weinberger [67] for particular finite difference analogues of the fixed membrane.

Lower bounds, in general, present a more formidable problem. In this case the maximum-minimum property [12] of λ_k is usually exploited. Perhaps the best known method is that of A. Weinstein [71] which gives arbitrarily close lower bounds for a symmetric linear elliptic differential operator A defined on a subspace V of a Hilbert space H . This method presupposes a knowledge of the eigenvalues and eigenvectors of an extension A' of A to a space V' which contains V . The method of A. Weinstein has been extended by N. Aronszajn [1, 2] and N. Bazley [3, 4, 5].

Certain methods for obtaining lower bounds have been pointed out by E. Trefftz [58], G. Temple [57], T. Kato [28], and others. Many of these and other methods are summarized in an excellent survey by J. B. Diaz [15]. Many of these have been unified into a single theory in

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a recent paper of H. Weinberger [69]. In order to apply any of the methods mentioned it is necessary to have already an exact lower bound for at least one of the higher eigenvalues of A . The finite difference methods presented here require no such additional knowledge.

The Poincaré inequality can also be used to obtain upper bounds for the eigenvalues, $\mu_1(h) \leq \mu_2(h) \leq \dots$, of a suitably defined finite difference problem. If the vectors in this inequality are appropriately defined in terms of the first k eigenfunctions of the membrane problem it is possible to relate this upper bound for $\mu_k(h)$ to λ_k . The resulting inequality is then solved to give a lower bound for λ_k in terms of $\mu_k(h)$. This technique has been applied by L. Collatz [8], G. Forsythe [19, 20], J. Hersch [24], and H. Weinberger [67]. The present paper relies heavily on the results of the latter author.

In Chapter II a finite difference eigenvalue problem is posed for the fixed membrane. Both upper and lower bounds for λ_k are obtained in terms of $\mu_k(h)$ and quantities involving the geometry of the boundary C . Most of the effort is expended in bounding a certain integral defined over a thin strip of R in the neighborhood of the boundary. Once this is accomplished the difference between the upper and lower bound is explicitly bounded by a term of the order of the mesh size, h . In the papers alluded to above, distinct finite difference problems were formulated to give upper and lower bounds. In most cases the difference between the upper and lower bounds was not explicitly bounded. It is known however, that the eigenvalues of such finite difference problems converge to the membrane eigenvalues as $h \rightarrow 0$. Clearly, by posing only one matrix eigenvalue problem we cut the actual computation involved in obtaining bounds in half.

In Chapter III bounds for the eigenvalues of the free membrane problem are obtained in a similar manner. Again the difference between the upper and lower bounds is bounded by a term of order h . Lower bounds for the free membrane eigenvalues by finite difference methods appear to have been left relatively unexplored. In this problem the eigenvalues do not have a monotone dependence on the region as is the case with fixed membrane eigenvalues. Hence it is difficult to obtain a lower bound for some higher eigenvalue; a requirement for most methods of finding lower bounds. Therefore, this result appears to mark a much greater advance than that obtained in the case of the fixed membrane.

The error bounds depend upon the choice of a vector field defined on $R + C$ subject to weak conditions with one possible method of choosing such a field illustrated in Appendix A. In Appendix B certain parameters depending on the geometry of C are bounded. Finally, the bounds of this paper are computed in two cases for which the solution is known explicitly; (a) the unit square, and (b) the unit circle. They are then

compared with the actual quantities being bounded.

I wish to express my gratitude to Professor H. F. Weinberger for his valuable assistance in directing my researches into this subject.

CHAPTER II

THE FIXED MEMBRANE PROBLEM

1. The problem and its finite difference analogue. Let R be a bounded, simply connected region in the $x - y$ plane with boundary C . The boundary itself is assumed to be composed of piecewise smooth arcs with interior angles at corners of C which are less than 180° .

Denote the eigenvalues of the fixed membrane problem

$$(2.1) \quad \begin{aligned} \Delta u + \lambda u &= 0 \quad \text{in } R, \\ u &= 0 \quad \text{on } C, \end{aligned}$$

by $\lambda_1 \leq \lambda_2 \leq \dots$. The eigenfunctions u_1, u_2, \dots are normalized by

$$(2.2) \quad \iint_R u_i^2 dx dy = 1.$$

Divide the $x - y$ plane into squares of width h by the two families of lines $x = mh$ and $y = nh$; $m, n = 1, 2, 3, \dots$. The dependent variables of the finite difference problem are defined at certain of the intersections of these lines, called "mesh points". Superposing a third family of lines of slope 1 through the mesh points divides the plane into isosceles right triangles so that each mesh point has six "associated" triangles of which

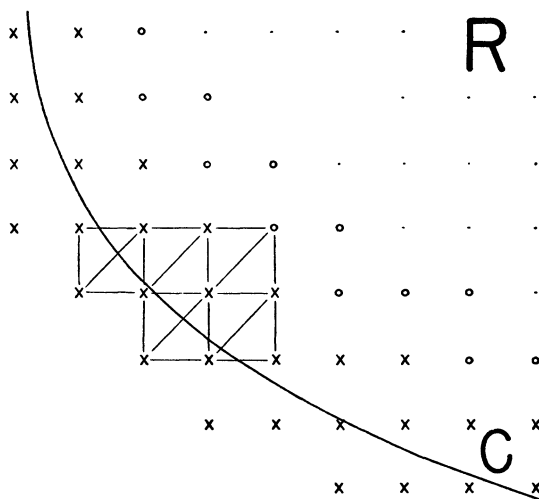


Fig. 1. Symbol Set to which point belongs
 \times C_h
 \circ B_h
 \cdot $R_h - C_h$

it is a vertex. C_h is defined to be the set of those mesh points which have at least one associated triangle intersecting C . R_h is the collection of mesh points of C_h . The set $R_h - C_h$ is composed of those points of R_h which do not belong to C_h , i.e., interior points whose triangles do not cut C , (see Fig. 1). In what follows the "mesh functions" are assumed to be defined only at points of R_h .

The eigenvalues of 2.1 are to be approximated by the eigenvalues $\mu_1(h) \leq \mu_2(h) \leq \dots$ of the finite difference problem

$$(2.3) \quad \begin{aligned} \Delta_h V + \mu(h) V &= 0 \quad \text{on } R_h - C_h, \\ V &= 0 \quad \text{on } C_h, \end{aligned}$$

where $V(m, n)$ is a mesh function evaluated at the point (mh, nh) of the plane. Here, Δ_h is the centered second order difference operator

$$(2.4) \quad \begin{aligned} \Delta_h V(m, n) &= h^{-2} [V(m+1, n) + V(m, n+1) + V(m-1, n) \\ &\quad + V(m, n-1) - 4V(m, n)]. \end{aligned}$$

Let V_1, V_2, \dots be the eigenvectors of 2.3 with the normalization

$$(2.5) \quad h^2 \sum_{R_h - C_h} V_i^2 = 1.$$

The eigenvalue $\mu_k(h)$ can be obtained as a solution of the minimum-maximum problem

$$(2.6) \quad \mu_k(h) = \min_{W_1 \dots W_k} \max_{a_1 \dots a_k} \frac{D_{R_h}^{(h)}(a_1 W_1 + \dots + a_k W_k)}{h^2 \sum_{R_h - C_h} (a_1 W_1 + \dots + a_k W_k)^2}.$$

The vectors W_1, \dots, W_k are linearly independent mesh functions which vanish at points of C_h and a_1, \dots, a_k are real numbers. The numerator of the above Rayleigh quotient is given as

$$(2.7) \quad \begin{aligned} D_{R_h}^{(h)}(W) &= \sum_{R_h} \{ [W(m+1, n) - W(m, n)]^2 \\ &\quad + [W(m, n+1) - W(m, n)]^2 \}, \end{aligned}$$

where the sum is over all differences of neighboring points of R_h .

The eigenvalues of the fixed membrane are defined by

$$(2.8) \quad \lambda_k = \min_{v_1 \dots v_k} \max_{a_1 \dots a_k} \frac{D(a_1 v_1 + \dots + a_k v_k)}{\iint_R (a_1 v_1 + \dots + a_k v_k)^2 dx dy},$$

where v_1, \dots, v_k are also linearly independent, piecewise continuously differentiable functions in R which vanish on C and $D(v)$ is the Dirichlet integral

$$(2.9) \quad D(v) = \iint_R (v_x^2 + v_y^2) dx dy .$$

The procedure for obtaining upper bounds for λ_k in terms of μ_k is to define k admissible functions for the inequality obtained from 2.8 in terms of the first k eigenvectors of the discrete problem. Inequalities are further developed for the numerator and denominator of 2.8 in terms of like quantities of 2.6. The indicated maximization with respect to the real numbers a_1, \dots, a_k is then effected to achieve the bound. An upper bound for $\mu_k(h)$ is found in a similar manner. This in turn can be solved by using the upper bound for λ_k to yield a lower bound for λ_k .

2. The upper bound. Starting with the mesh eigenfunctions V_1, \dots, V_k of the finite difference problem, we define functions v_1, \dots, v_k admissible in the continuous problem 2.8 as suggested by L. Collatz [8] and R. Courant [13], (see also G. Pólya [47]). Let $v_i(x, y)$ be the piecewise continuously differentiable function which is linear in each triangle and coincides with the eigenvector V_i at the vertices. Clearly the linear independence of v_1, \dots, v_k assures the linear independence of v_1, \dots, v_k and the vanishing of each V_i at points of C_h causes each v_i to be zero on C . Therefore these functions are admissible in the inequality

$$(2.10) \quad \lambda_k \leq \max_{a_1, \dots, a_k} \frac{D(v)}{\iint_R v^2 dx dy} ,$$

obtained from 2.8 where

$$(2.11) \quad \begin{aligned} v &= a_1 v_1 + \dots + a_k v_k , \\ V &= a_1 V_1 + \dots + a_k V_k . \end{aligned}$$

We next obtain inequalities for the numerator and denominator of 2.10 in terms of $D_{Rh}^{(h)}(V)$ and $h^2 \sum V^2$. It is easily shown (see [67], page 361) that

$$(2.12) \quad D(v) \leq D_{Rh}^{(h)}(V) ,$$

and further that

$$(2.13) \quad \begin{aligned} \iint_R v^2 dx dy &\geq h^2 \sum_{Rh-\sigma_h} V(m, n)^2 - \frac{h^2}{12} \sum_{Rh} \{ [V(m+1, n) - V(m, n)]^2 \\ &\quad + [V(m, n+1) - V(m, n)]^2 + [(V(m+1, n+1) - V(m, n))]^2 \} , \\ &\geq h^2 \sum_{Rh-\sigma_h} V^2 - \frac{h^2}{4} D_{Rh}^{(h)}(V) . \end{aligned}$$

Applying these inequalities to 2.10 yields the known upper bound (c.f. Weinberger [67])

$$\begin{aligned}
 (2.14) \quad \lambda_k &\leq \max_{a_1 \cdots a_k} \frac{D_{R_h}^{(h)}(V)}{h^2 \sum_{R_h - \mathcal{O}_h} V^2 - \frac{h^2}{4} D_{R_h}^{(h)}(V)}, \\
 &= \max_{a_1 \cdots a_k} \frac{\sum_{i=1}^k a_i^2 \mu_i(h)}{1 - \frac{h^2}{4} \sum_{i=1}^k a_i^2 \mu_i(h)} \leq \frac{\mu_k(h)}{1 - \frac{h^2}{4} \mu_k(h)}.
 \end{aligned}$$

We have chosen a finite difference problem which yields a particularly simple expression for the upper bound. This increases the labor involved in finding the lower bound as we shall see.

3. The lower bound. The following technique was used by Weinberger [67] to find lower bounds for λ_k in terms of the eigenvalues of a slightly different finite difference problem. The bound obtained in [67] is that of 2.27 with the integral 2.28 omitted.

Extend u_i as 0 outside R and let the mesh functions W_1, \dots, W_k be defined in terms of the eigenfunctions u_1, \dots, u_k as follows:

$$(2.15) \quad W_i(m, n) = \begin{cases} h^{-2} \iint_{S(m, n)} u_i dx dy, & (m, n) \in R_h - C_h \\ 0 & (m, n) \in C_h, \end{cases}$$

where $S(m, n)$ is the square with center (m, n) and sides of length h oriented in the directions of the $x - y$ axes. Let

$$\begin{aligned}
 (2.16) \quad W &= a_1 W_1 + \cdots + a_k W_k, \\
 u &= a_1 u_1 + \cdots + a_k u_k.
 \end{aligned}$$

Deferring for the moment the question of linear independence we see that in other respects the functions W_1, \dots, W_k are admissible in the inequality

$$(2.17) \quad \mu_k(h) \leq \max_{a_1 \cdots a_k} \frac{D_{R_h}^{(h)}(W)}{h^2 \sum_{R_h - \mathcal{O}_h} W^2}$$

which arises from 2.6.

We now seek inequalities relating the numerator and denominator of 2.17 to $D(u)$ and $\iint u^2 dx dy$. It is easily seen that if both (m, n) and $(m + 1, n) \in R_h$ then

$$(2.18) \quad W(m + 1, n) - W(m, n) = h^{-2} \int_0^{2h} d\xi \int_0^h d\gamma \psi(\xi) \frac{\partial u}{\partial x}(mh + \xi, nh + \gamma),$$

with a similar formula for y -differences. Here we have put

$$\psi(\xi) = \begin{cases} \xi & , 0 \leq \xi \leq h \\ 2h - \xi & , h \leq \xi \leq 2h \end{cases} .$$

In the mixed case, e.g. when $(m-1, n) \in C_h$, $(m, n) \in R_h - C_h$, an application of Schwarz's inequality yields

$$(2.19) \quad [W(m, n) - W(m-1, n)]^2 = W(m, n)^2 \leq h^{-2} \iint_{S(m, n)} u^2 dx dy .$$

We see that (m, n) might have points of C_h on as many as three sides with the result that $W(m, n)^2$ could appear with a factor of three in the Dirichlet sum, $D_{R_h}^{(h)}(W)$. If B_h is defined to be the set of those points of $R_h - C_h$ adjacent to points of C_h (see Fig. 1), then

$$(2.20) \quad D_{R_h}^{(h)}(W) \leq 3 \sum_{B_h} W^2 + D_{R_h - C_h}^{(h)}(W) .$$

A computation using the representation 2.18 shows that

$$(2.21) \quad D(u) - D_{R_h - C_h}^{(h)}(W) \geq h^{-1} \sum_{R_h - C_h} \int_0^{2h} d\xi \int_0^h d\gamma \psi(\xi) \cdot \left[\frac{\partial u}{\partial x}(mh + \xi, nh + \gamma) - h^{-1} \langle W(m+1, n) - W(m, n) \rangle \right]^2 + \left[\frac{\partial u}{\partial y}(mh + \gamma, nh + \xi) - h^{-1} \langle W(m, n+1) - W(m, n) \rangle \right]^2 \geq 0 .$$

Using 2.19 and 2.21 we rewrite 2.20 as

$$(2.22) \quad D_{R_h}^{(h)}(W) \leq 3h^{-2} \sum_{B_h} \iint_{S(m, n)} u^2 dx dy + D(u) ,$$

which is the desired inequality for the numerator of the Rayleigh quotient.

By definition 2.15 it is seen that for $(m, n) \in R_h - C_h$

$$(2.23) \quad \iint_{S(m, n)} [u(x, y) - W(m, n)] dx dy = 0 ,$$

so that the integrand is admissible in the variational problem for the second eigenvalue of the free membrane for the square $S(m, n)$. Consequently

$$(2.24) \quad \frac{\pi^2}{h^2} \leq \frac{\iint_{S(m, n)} |\text{grad } u|^2 dx dy}{\iint_{S(m, n)} [u(x, y) - W(m, n)]^2 dx dy} .$$

This can be written as

$$(2.25) \quad h^2 W(m, n)^2 \geq \iint_{S(m, n)} u^2 dx dy - \frac{h^2}{\pi^2} \iint_{S(m, n)} |\text{grad } u|^2 dx dy .$$

Then

$$\begin{aligned}
 (2.26) \quad h^2 \sum_{R_h - C_h} W^2 &\geq \sum_{R_h - C_h} \iint_{S(m,n)} \left[u^2 - \frac{h^2}{\pi^2} |\text{grad } u|^2 \right] dx dy, \\
 &\geq \iint_R u^2 dx dy - \iint_S u^2 dx dy - \frac{h^2}{\pi^2} D(u),
 \end{aligned}$$

where S is a strip of depth αh inside R with outer boundary C which completely covers the sets $S(m, n) \cap R$, for $(m, n) \in C_h + B_h$.

The parameter α which appears in the strip width, αh , depends to some extent on the region R and the manner in which the mesh is placed on R . In Appendix B it is shown that

$$\alpha \leq \frac{1}{2} \sqrt{34}$$

under the hypotheses of the next section. Substituting 2.22 and 2.26 into 2.17 we have the desired inequality

$$\begin{aligned}
 (2.27) \quad \mu_k(h) &\leq \max_{a_1, \dots, a_k} \frac{D(u) + 3h^{-2} \iint_R u^2 dx dy}{\iint_R u^2 dx dy - \frac{h^2}{\pi^2} D(u) - \iint_S u^2 dx dy}, \\
 &\leq \frac{\lambda_k + 3h^{-2} \max_{a_1, \dots, a_k} \iint_S u^2 dx dy}{1 - \frac{h^2}{\pi^2} \lambda_k - \max_{a_1, \dots, a_k} \iint_S u^2 dx dy}.
 \end{aligned}$$

Inequality 2.27 gives an explicit bound for λ_k in terms of $\mu_k(h)$ if one has a bound for

$$(2.28) \quad \iint_S u^2 dx dy.$$

The above lower bound without the strip integrals was obtained by Weinberger [67] for a finite difference problem on a set of mesh points which includes R_h . The admissible functions in his problem vanish at points outside of this set. The bound for $\lambda_k - \mu_k$, given by 2.27, is $O(h)$ as we shall show in the next paragraph through an explicit bound of order h^3 for the strip integral 2.28.

4. The strip integral. For simplicity of presentation we assume C to be a smooth arc whose curvature, K , is differentiable. Restrict h to be so small that the center of curvature corresponding to any point p on C does not lie in S near p ; i.e., choose h so that $\alpha h < \min_{\sigma} K^{-1}$. Let the parametric representation of C be

$$C: x_i = g_i(s) \qquad i = 1, 2,$$

where $x_1 = x$, $x_2 = y$ and s is the arc length along C . Inside the strip S we make the transformation to geodesic normal coordinates (s, n) in the following manner:

$$(2.29) \quad x_i = g_i(s) + n_i(s)n, \quad i = 1, 2, \dots$$

Here (n_1, n_2) is taken to be the unit inward normal on C . The Jacobian of the transformation is

$$(2.30) \quad \frac{\partial(x_1, x_2)}{\partial(s, n)} = (1 - K(s)n),$$

and the components of the metric tensor are

$$(2.31) \quad a^{11} = (1 - K(s)n)^{-2}, \quad a^{12} = a^{21} = 0, \quad a^{22} = 1.$$

The normal derivative of a function u on C has the property that

$$(2.32) \quad \frac{\partial u}{\partial \nu} = a^{ij}u_{,i}(-n_j) = -\frac{\partial u}{\partial n},$$

and hence we shall interpret $\partial u / \partial n$ in both senses. In terms of the new coordinates we have

$$(2.33) \quad |\text{grad } u|^2 = \left\{ (1 - K(s)n)^{-2} \left(\frac{\partial u}{\partial s} \right)^2 + \left(\frac{\partial u}{\partial n} \right)^2 \right\},$$

and the membrane equation 2.1 takes the form

$$(2.34) \quad (1 - Kn)^{-2} \frac{\partial^2 u}{\partial s^2} + \frac{\partial^2 u}{\partial n^2} + K'n(1 - Kn)^{-3} \frac{\partial u}{\partial s} - K(1 - Kn)^{-1} \frac{\partial u}{\partial n} + \lambda u = 0.$$

Rewriting our strip integral in terms of the new coordinates and noting that 2.16 implies $u = 0$ on C leads to

$$(2.35) \quad \begin{aligned} \iint_S u^2 dx dy &= \int_0^L ds \int_0^{\alpha h} u^2(s, n)(1 - K(s)n)dn \\ &= \int_0^L ds \int_0^{\alpha h} \left\{ u(s, 0) + \frac{\partial u}{\partial n}(s, 0)n + \int_0^n (n - \xi) \frac{\partial^2 u}{\partial \xi^2}(s, \xi) d\xi \right\}^2 (1 - K(s)n)dn, \\ &\leq 2 \int_0^L ds \int_0^{\alpha h} n^2 \left[\frac{\partial u}{\partial n}(s, 0) \right]^2 (1 - K(s)n)dn \\ &\quad + 2 \int_0^L ds \int_0^{\alpha h} \left[\int_0^n (n - \xi) \frac{\partial^2 u}{\partial \xi^2}(s, \xi) d\xi \right]^2 (1 - K(s)n)dn, \\ &\leq \frac{2}{3} (\alpha h)^3 (1 + \alpha h K(-)) \int_0^L \left(\frac{\partial u}{\partial \nu} \right)^2 ds \\ &\quad + \frac{(\alpha h)^4}{6} \left[\frac{1 + \alpha h K(-)}{1 - \alpha h K(+)} \right] \iint_S \left(\frac{\partial^2 u}{\partial n^2} \right) (1 - K(s)n) ds dn, \end{aligned}$$

where we have used the notation

$$(2.36) \quad \begin{aligned} K(-) &= \text{absolute value of maximum negative curvature} \\ K(+) &= \text{absolute value of maximum positive curvature.} \end{aligned}$$

The last lines of 2.35 display the desired $0(h^3)$ character of the strip integral.

We now seek explicit bounds for the integrals appearing on the right side of 2.35. It is true that

$$(2.37) \quad \left(\frac{\partial^2 u}{\partial n^2} \right)^2 \leq \left(\frac{\partial^2 u}{\partial x^2} \right)^2 + 2 \left(\frac{\partial^2 u}{\partial x \partial y} \right)^2 + \left(\frac{\partial^2 u}{\partial y^2} \right)^2,$$

where the right side is recognized to be the second order differential invariant. An application of the divergence theorem to this invariant, pointed out by Payne [41], is very useful here. Integrating 2.37 over S gives

$$(2.38) \quad \begin{aligned} \iint_S \left(\frac{\partial^2 u}{\partial n^2} \right)^2 dx dy &\leq \iint_R \left[\left(\frac{\partial^2 u}{\partial x^2} \right)^2 + 2 \left(\frac{\partial^2 u}{\partial x \partial y} \right)^2 + \left(\frac{\partial^2 u}{\partial y^2} \right)^2 \right] dx dy \\ &= D(u_x) + D(u_y) = - \iint_R [u_x \Delta u_x + u_y \Delta u_y] dx dy + \int_\sigma \left(u_x \frac{\partial u_x}{\partial \nu} + u_y \frac{\partial u_y}{\partial \nu} \right) ds. \end{aligned}$$

In view of 2.16 and the fact that $\partial u_i / \partial x$, $\partial u_i / \partial y$ each satisfy the membrane equation, 2.38 can be continued to

$$(2.39) \quad \begin{aligned} &\iint_S \left(\frac{\partial^2 u}{\partial n^2} \right)^2 dx dy \\ &\leq \iint_R \left\{ \left[\sum_1^k a_i \frac{\partial u_i}{\partial x} \right] \left[\sum_1^k a_i \lambda_i \left(\frac{\partial u_i}{\partial x} \right) \right. \right. \\ &\quad \left. \left. + \left[\sum_1^k a_i \frac{\partial u_i}{\partial y} \right] \left[\sum_1^k a_i \lambda_i \frac{\partial u_i}{\partial y} \right] \right\} dx dy + \frac{1}{2} \int_\sigma \frac{\partial}{\partial \nu} |\text{grad } u|^2 ds, \\ &\leq \sum_1^k a_i^2 \lambda_i^2 - \frac{1}{2} \int_\sigma \left\{ \lim_{n \rightarrow 0} \frac{\partial}{\partial n} \left[(1 - Kn)^{-2} \left(\frac{\partial u}{\partial s} \right)^2 + \left(\frac{\partial u}{\partial n} \right)^2 \right] \right\} ds. \end{aligned}$$

The latter inequality follows from 2.33 and the orthogonality of the eigenfunctions in Dirichlet norm, i.e.,

$$(2.40) \quad D(u_i, u_j) = \delta_{ij} \lambda_{(i)},$$

Continuing the inequality we have

$$(2.41) \quad \iint_S \left(\frac{\partial^2 u}{\partial n^2} \right)^2 dx dy \leq \lambda_k^2 - \int_\sigma \left(\frac{\partial u}{\partial n} \right) \left(\frac{\partial^2 u}{\partial n^2} \right) ds.$$

If we now assume that the differential equation 2.34 is satisfied on the boundary in the limiting sense 2.41 becomes

$$(2.42) \quad \int_s \left(\frac{\partial^2 u}{\partial n^2} \right)^2 dx dy \leq \lambda_k^2 + K(-) \int_o \left(\frac{\partial u}{\partial n} \right)^2 ds .$$

Estimates of the contour integral above can be made from an integral identity due to Payne and Weinberger [45]. This formula which was obtained for hyperbolic operators by L. Hörmander [27] is a generalization of an integral identity of Rellich [52]. Using the summation convention the identity in two dimensions becomes

$$(2.43) \quad \int_o \left\{ f^j \nu_j \left[\left(\frac{\partial u}{\partial s} \right)^2 - \left(\frac{\partial u}{\partial \nu} \right)^2 \right] - 2 f^j t_j \left(\frac{\partial u}{\partial s} \right) \left(\frac{\partial u}{\partial \nu} \right) \right\} ds \\ = -2 \int_R f^i u_{,i} \Delta u dx dy + \int_R \left[f^{,j}_j \delta^{ip} - 2 f^{i,p}_{,p} \right] u_{,i} u_{,p} dx dy .$$

In this formula (t_1, t_2) is the unit tangent vector to the curve C while (f^1, f^2) is an arbitrary piecewise continuously differentiable vector field defined on the closure of R . Since $u = 0$ on C , 2.43 reduces to

$$(2.44) \quad \int_C (f^i \nu_i) \left(\frac{\partial u}{\partial n} \right) ds = 2 \int_R f^i u_{,i} \Delta u dx dy \\ - \int_R \left[f^{,j}_j \delta^{ip} - 2 f^{i,p}_{,p} \right] u_{,i} u_{,p} dx dy .$$

Let us further assume that $f^i \nu_i > 0$ on C so that by Schwarz's inequality

$$(2.45) \quad \int_o \left(\frac{\partial u}{\partial n} \right)^2 ds \\ \leq \frac{1}{\min_C f^i \nu_i} \left\{ 2 \left[\max_R \sqrt{f^i f^i} \right] \left[\int_R (\Delta u)^2 dx dy \right]^{1/2} D(u)^{1/2} + D(u) \tau \right\} .$$

Here we have used the notation

$$(2.46) \quad \tau = \max_R \left| (f^1_{,1} - f^2_{,2})^2 + (f^1_{,2} + f^2_{,1})^2 \right|^{1/2} \\ = \max_R \left| f^{i,j}_j f^{j,i}_i + f^{i,j}_j f^{i,j}_j - (f^i_{,i})^2 \right|^{1/2} ,$$

which arises from the largest eigenvalue of the matrix

$$(2.47) \quad (f^{i,j}_j + f^{j,i}_i - \delta^i_j f^{k,k}_{,k}) .$$

From the definition 2.16 of u the inequality 2.45 gives rise to the result

$$(2.48) \quad \int_o \left(\frac{\partial u}{\partial n} \right)^2 ds \leq \frac{\lambda_k}{\min_C f^i \nu_i} \left\{ 2 \sqrt{\lambda_k} \max_R \sqrt{f^i f^i} + \tau \right\} .$$

The vector field (f^1, f^2) can be defined in many ways depending upon the region R . One possible method is described in detail in Appendix A. If R is star-shaped with respect to the origin, i.e., if every ray from

the origin cuts C in one point, we can define $f^i = x^i$. In this case 2.48 takes the form

$$(2.49) \quad \int_o \left(\frac{\partial u}{\partial n} \right)^2 ds \leq 2\lambda_k^{3/2} \frac{\max_R \sqrt{x^i x^i}}{\min_o x^i \nu_i}.$$

Finally, upon substituting 2.42 and 2.48 into 2.35, our original inequality 2.27 takes the form

$$(2.50) \quad \mu_k(h) \leq \frac{\lambda_k(1 + hB_k)}{(1 - h^2A_k)}.$$

Here A_k and B_k are given by

$$(2.51) \quad A_k = \frac{\bar{\lambda}_k}{\pi^2} + h\bar{\lambda}_k \left\{ \frac{\alpha^3(1 + \alpha hK(-))}{3\min_o f^i \nu_i} \right\} \left\{ 2 + \frac{\alpha hK(-)}{2(1 - \alpha hK(+))} \right\} \left\{ 2\sqrt{\bar{\lambda}_k} \max_R \sqrt{f^i f^i} + \tau \right\} + h^2 \bar{\lambda}_k^2 \frac{\alpha^4}{6} \left(\frac{1 + \alpha hK(-)}{1 - \alpha hK(+)} \right),$$

$$B_k = \frac{\alpha^3(1 + \alpha hK(-))}{\min_o f^i \nu_i} \left\{ 2 + \frac{\alpha hK(-)}{2(1 - \alpha hK(+))} \right\} \left\{ 2\sqrt{\bar{\lambda}_k} \max_R \sqrt{f^i f^i} + \tau \right\} + \frac{h\bar{\lambda}_k \alpha^4}{2} \left(\frac{1 + \alpha hK(-)}{1 - \alpha hK(+)} \right),$$

and $\bar{\lambda}_k$ is an upper bound for λ_k . In particular we can use the upper bound 2.14 as $\bar{\lambda}_k$. Solving 2.50 yields the lower bound

$$(2.52) \quad \lambda_k \geq \frac{\mu_k(h)(1 - h^2A_k)}{1 + hB_k}.$$

It is seen that the difference between the upper and lower bounds is the expression

$$(2.53) \quad \mu_k(h) \left\{ \left[1 - \frac{h^2}{4} \mu_k(h) \right]^{-1} - \left[1 - h^2 A_k \right] \left[1 + h B_k \right]^{-1} \right\},$$

which is indeed $O(h)$ in terms of easily computed quantities. These are $\mu_k(h)$, $K(-)$, $K(+)$, and the vector field (f^1, f^2) which in most instances will strongly reflect the geometry of C .

We note that by our choosing the mesh width h to be sufficiently small the inequality 2.26 will yield

$$(2.54) \quad h^2 \sum_{R_k - \mathcal{O}_h} W^2 > 0.$$

We assume h to be so chosen. This in turn assures the linear inde-

pendence of W_1, \dots, W_k which was assumed earlier.

At the beginning of this section we had restricted C to be a smooth arc for reasons of simplicity. However, discontinuities in the derivatives of the boundary C offer no essential difficulty in carrying through the previous development. As before we assume that the interior angles at corners are less than 180° . Difficulties arise at corners because of possible bad behavior of various derivatives of u . This could invalidate the inequalities 2.38, 2.39 and 2.41.

We construct the sequence of regions R_1, R_2, \dots with the properties:

- (1) $R_1 \subseteq R_2 \subseteq \dots \subseteq R$
- (2) $x \in R \Rightarrow x \in R_i$ for some i
- (3) the boundary C_i of R_i is continuously differentiable
- (4) at a corner of C the locus of centers of curvature of C_i lying with the strip S_i of depth αh has the property that the normal to C_i intersects the bisector of the angle before reaching its center of curvature.
- (5) the curvature of each C_i is positive in the vicinity of a corner.

The bounds of this section apply in the case of each R_i using the various parameters $h, \alpha, (f^1, f^2)$, which previously have been selected for R . At corners of R the locus of centers of curvature of R_i may enter the strip S . Condition (4) above, assures that we can change the upper limit of integration in 2.35 from αh to $K^{-1}(s)$ where such penetration occurs and still cover the domain of integration. Let S_i^η be the sector of S_i at the corner η such that the locus of centers of curvature of C_i^η lies in the strip S_i . At this corner 2.35 takes the form

$$\begin{aligned}
 (2.56) \quad & \iint_{S_i^\eta} u^2 dx dy \leq \int_0^{x_i^\eta} ds \int_0^{b_i(s)} u^2(s, n) (1 - K_i(s)n) dn \\
 & \leq \frac{2}{3} (\alpha h)^3 \int_{S_i^\eta} \left(\frac{\partial u}{\partial \nu} \right)^2 ds + \frac{(\alpha h)^4}{6} \iint_{S_i^\eta} \left(\frac{\partial^2 u}{\partial n^2} \right)^2 (1 - K_i(s)n) ds dn.
 \end{aligned}$$

Here $b_i(s) \leq K_i^{-1}(s)$, where $(s, b_i(s))$ is on the bisector of the angle at η . The bound 2.35 is then seen to be valid for R_i where $K_i(+)$ is the maximum value of the curvature of C_i at other than points of C_i^η . The inequalities following 2.35 are also seen to be valid under this interpretation of $K_i(+)$, and with λ_k replaced by $\lambda_k(R_i)$.

We assume the mesh to be placed on R so that no mesh points lie on C in the vicinity of a corner. Then for some N and all $i > N$, we have

$$(2.57) \quad \mu_k(h) = \mu_k(h, R_i).$$

Hence 2.50 becomes

$$(2.58) \quad \mu_k(h) \leq \frac{\lambda_k(R_i)(1 + h B_k(R_i))}{(1 - h^2 A_k(R_i))} .$$

The hypotheses of 2.55 are sufficient conditions for

$$(2.59) \quad \lim_{i \rightarrow \infty} \lambda_k(R_i) = \lambda_k ,$$

and also for

$$(2.60) \quad \begin{aligned} \lim_{i \rightarrow \infty} B_k(R_i) &= B_k , \\ \lim_{i \rightarrow \infty} A_k(R_i) &= A_k , \end{aligned}$$

as was stated by D. M. Eidus [16, 17]. Therefore the bound 2.50 also holds in the case of corners whose interior angles are less than 180° . There seems to be little hope of applying the procedures of this paper toward bounds in the case of angles greater than 180° since the factor $(1 + K(-)\alpha h)$ is present, and $K_i(-) \rightarrow \infty$.

CHAPTER III

THE FREE MEMBRANE PROBLEM

1. The finite difference problem. The free membrane problem is given by the equations

$$(3.1) \quad \begin{aligned} \Delta u + \lambda u &= 0 \quad \text{in } R \\ \frac{\partial u}{\partial \nu} &= 0 \quad \text{on } C \end{aligned}$$

with eigenvalues $0 = \lambda_1 \leq \lambda_2 \leq \dots$. The eigenfunctions are u_1, u_2, \dots with the normalization

$$(3.2) \quad \iint_R u_i^2 dx dy = 1 .$$

The approximating finite difference problem is chosen to be

$$(3.3) \quad \begin{aligned} \Delta_h V(m, n) + \mu V(m, n) &= 0, & (m, n) \in R_h - C_h \\ \nabla(\nu) V(m, n) &= 0, & (m, n) \in C_h \end{aligned}$$

where $\nabla(\nu) V(m, n) = 0$ is the condition that $V(m, n)$ be the average of the values of V at adjacent (vertical and horizontal) points of R_h only. For example in Fig. 2 the point (m, n) belongs to C_h so that

$$\begin{aligned} 0 = \nabla(\nu) V(m, n) &= V(m, n+1) + V(m+1, n) \\ &+ V(m, n-1) - 3V(m, n) . \end{aligned}$$

Now define C_h to be the set of grid points which have an associated triangle intersecting the boundary C . R_h is the set of mesh points interior to R augmented by points of C_h . The set B_h is defined to be those points in $R_h - C_h$ which have an associated triangle with a vertex in C_h . See Fig. 1.

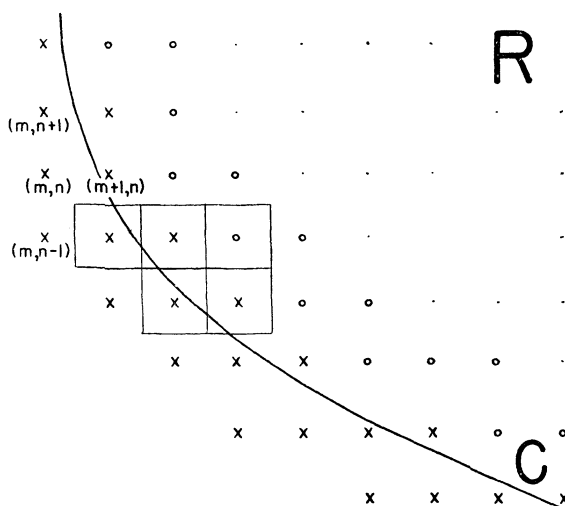


Fig. 2. Symbol Set to which point belongs
 \times C_h^*
 \circ B_h^*
 \cdot $R_h - C_h^*$

Let the eigenvalues of 3.3 be $\mu_1(h) \leq \mu_2(h) \leq \dots$ with corresponding mesh eigenfunctions V_1, V_2, \dots which have the normalization

$$(3.4) \quad h^2 \sum_{R_h - C_h} V_i^2 = 1.$$

The minimum-maximum principle 2.8 yields the eigenvalues of the freemembrane problem 3.1 where in this case the functions involved need not vanish on the boundary C to be admissible. In like manner the variational principle 2.6 gives the eigenvalues of our finite difference analogue 3.3 where any set of k linearly independent mesh functions is allowed. The procedure followed to obtain bounds for λ_k in terms of $\mu_k(h)$ is very similar to that of Chapter II for the fixed membrane problem.

To obtain an upper bound for λ_k we define a set of functions v_1, \dots, v_k which are admissible in the variational problem 2.8 in terms of the mesh eigenfunctions V_1, \dots, V_k . Let $v_j(x, y)$ be the piecewise continuously differentiable function which is linear in each triangle and which coincides at the vertices with $V_j(m, n)$. By considerations identical to those of Chapter II, § 2, we arrive at the upper bound

$$(3.5) \quad \lambda_k \leq \frac{\mu_k(h)}{1 - \frac{h^2}{4}\mu_k(h)} .$$

A lower bound for λ_k is achieved by finding an appropriate upper bound for $\mu_k(h)$ in terms of λ_k . As before $S(m, n)$ is defined to be the “associated” square with (mh, nh) as center and sides of length h having the same orientation as the grid lines. We can achieve somewhat better bounds by altering the set C_h slightly. Define C_h^* to be those points of C_h whose associated squares intersect the boundary. Let B_h^* be the set of interior mesh points of R whose squares share a common vertex with at least one of the squares of C_h^* , e.g. see Fig. 2. We define a linearly independent set of k mesh functions W_1, \dots, W_k in terms of the eigenfunctions u_1, \dots, u_k as follows:

$$(3.6) \quad W_i(m, n) = h^{-2} \iint_{S(m, n)} u_i(x, y) dx dy, \quad (m, n) \in R_h - C_h^*,$$

$$\nabla(\nu) W_i(m, n) = 0, \quad (m, n) \in C_h^* .$$

Note that $\nabla(\nu) W_i = 0$ represents as many linear equations as unknowns.

We further define

$$(3.7) \quad W = a_1 W_1 + \dots + a_k W_k ,$$

$$u = a_1 u_1 + \dots + a_k u_k .$$

By the same considerations used in proving 2.21 we see that

$$D_{R_h - C_h^*}^{(h)}(W) \leq D(u) .$$

The above sum is taken over all differences of neighboring points of $R_h - C_h^*$. It follows immediately that

$$(3.8) \quad D_{R_h}^{(h)}(W) \leq D(u) + \beta(W) D_{B_h^*}^{(h)}(W) ,$$

where

$$(3.9) \quad \beta(W) = \frac{D_{R_h}^{(h)}(W) - D_{R_h - C_h^*}^{(h)}(W)}{D_{B_h^*}^{(h)}(W)} .$$

The quantity $\beta(W)$ is minimized with respect to possible choices of W at points of C_h^* by the manner in which the functions W_1, \dots, W_k were defined at these points. In fact if $(m, n) \in C_h^*$ as in Fig. 2 we see that

$$(3.10) \quad \frac{d}{dW(m, n)} \beta(W) = 2\{3W(m, n) - W(m, n+1) - W(m+1, n) \\ - W(m, n-1)\} D_{B_h^*}^{(h)}(W)^{-1} \\ = 2 \nabla(\nu) W(m, n) D_{B_h^*}^{(h)}(W)^{-1} = 0$$

which is the desired result. The bound for $\beta(W)$ depends to some extent on the geometry of C and consequently must be determined anew for each region R . Assume that h is chosen small enough so that the points of C_h^* constitute at most a double "fence" around B_h^* . In such a case we can choose the values of W_i on C_h^* in terms of nearby points of B_h^* so that

$$(3.11) \quad \beta(W) \leq 4.$$

For a further discussion of β see Appendix B. Finally we can write 3.8 in the form

$$(3.12) \quad D_{R_h}^{(h)}(W) \leq D(u) + \beta(W) \iint_S |\text{grad } u|^2 dx dy,$$

where again S is a strip of suitable depth, $\bar{\alpha}h$, so that the squares associated with points of B_h^* lie in S . In Appendix B the parameter $\bar{\alpha}$ is discussed and it is shown that, under the hypotheses of this section, we can take

$$\bar{\alpha} \leq 2\sqrt{2}.$$

We observe, in addition, that 2.26 is valid for this same strip S so that we can write the lower bound in the form

$$(3.13) \quad \mu_k(h) \leq \frac{\lambda_k + \bar{\beta} \max_{a_1 \dots a_k} \iint_S |\text{grad } u|^2 dx dy}{1 - \frac{h^2}{\pi^2} \lambda_k - \max_{a_1 \dots a_k} \iint_S u^2 dx dy},$$

where $\bar{\beta}$ is a numerical bound for $\beta(W)$, i.e., $\beta(W) \leq \bar{\beta}$.

2. Bounds for the strip integrals. We are faced with the task of bounding each of two integrals in the free membrane problem over a strip of depth $\bar{\alpha}h$ adjacent to the boundary C . We desire the bounds themselves to be of order h . As in Chapter II, § 4, we assume C to be a smooth arc whose curvature K is differentiable. Restrict h to be so small that the center of curvature corresponding to any point p on C does not lie in S near p ; i.e., choose h so that $\alpha h < \min_o K^{-1}$. The geodesic normal coordinates (s, n) are introduced as before so that equations 2.29 through 2.34 remain valid. We see from 2.30 that this guarantees the single-valuedness of our transformation. An additional identity involving the second order differential invariant will be useful in the estimates which follow, and is included at this time. If $u|_i$ denotes covariant differentiation then

$$\begin{aligned}
(3.14) \quad u_{xx}^2 + 2u_{xy}^2 + u_{yy}^2 &= a^{ij}a^{pm}u_{|ip}u_{|jm} = (a^{11}u_{|11})^2 + 2a^{11}a^{22}(u_{|12})^2 \\
&+ (a^{22}u_{|22})^2 = \left\{ (1 - Kn)^{-2} \left[\frac{\partial^2 u}{\partial s^2} + \frac{K'n}{(1 - Kn)} \frac{\partial u}{\partial s} \right. \right. \\
&- \left. \left. K(1 - Kn) \frac{\partial u}{\partial n} \right] \right\}^2 + 2(1 - Kn)^{-2} \left[\frac{\partial^2 u}{\partial s \partial n} + \frac{K}{(1 - Kn)} \frac{\partial u}{\partial s} \right]^2 \\
&+ \left(\frac{\partial^2 u}{\partial n^2} \right)^2.
\end{aligned}$$

The strip integral appearing in the numerator of 3.13 can be bounded using 2.33 and 3.14 as follows:

$$\begin{aligned}
(3.15) \quad \iint_S |\text{grad } u|^2 dx dy &= \int_0^L \int_0^{\bar{\alpha}h} \left\{ (1 - Kn)^{-2} \left(\frac{\partial u}{\partial s} \right)^2 + \left(\frac{\partial u}{\partial n} \right)^2 \right\} (1 - Kn) dn ds \\
&= \int_0^L \int_0^{\bar{\alpha}h} \left\{ \left[\frac{\partial u}{\partial s}(s, 0) + \int_0^n \frac{\partial}{\partial t} \left\langle (1 - Kt)^{-1} \frac{\partial u}{\partial s}(s, t) \right\rangle dt \right]^2 \right. \\
&\quad \left. + \left[\int_0^n \frac{\partial^2 u}{\partial t^2}(s, t) dt \right]^2 \right\} (1 - Kn) dn ds \\
&\leq \int_0^L \int_0^{\bar{\alpha}h} \left\{ 2 \left[\frac{\partial u}{\partial s}(s, 0) \right]^2 \right. \\
&\quad \left. + n \int_0^n \left\langle (1 - Kt)^{-2} \left[\frac{\partial^2 u}{\partial s \partial t}(s, t) + \frac{K}{(1 - Kt)} \frac{\partial u}{\partial s}(s, t) \right]^2 \right. \right. \\
&\quad \left. \left. + \left[\frac{\partial^2 u(s, t)}{\partial t^2} \right]^2 \right\rangle dt \right\} (1 - Kn) dn ds \\
&\leq 2(\bar{\alpha}h)(1 + \bar{\alpha}hK(-)) \int_0^L \left(\frac{\partial u}{\partial s} \right)^2 ds \\
&+ \frac{(\bar{\alpha}h)^2}{2} \left(\frac{1 + \bar{\alpha}hK(-)}{1 - \bar{\alpha}hK(+)} \right) \iint_S a^{ij}a^{pm}u_{|ip}u_{|jm} dx dy.
\end{aligned}$$

The integral over the strip S of the second order differential invariant can be bounded as follows:

$$\begin{aligned}
(3.16) \quad \iint_S a^{ij}a^{pm}u_{|ip}u_{|jm} dx dy &= D(u_x) + D(u_y) \\
&\leq - \iint_R [u_x \Delta u_x + u_y \Delta u_y] dx dy + \frac{1}{2} \int_0^L \frac{\partial}{\partial \nu} |\text{grad } u|^2 ds \\
&\leq \sum_1^k \alpha_i^2 \lambda_i^2 - \frac{1}{2} \int_0^L \lim_{n \rightarrow 0} \frac{\partial}{\partial n} \left[(1 - Kn)^{-2} \left(\frac{\partial u}{\partial s} \right)^2 \right] ds \\
&\leq \lambda_k^2 + K(-) \int_0^L \left(\frac{\partial u}{\partial s} \right)^2 ds.
\end{aligned}$$

Using this bound, the inequality 3.15 can be expressed in the form

$$(3.17) \quad \iint_s |\text{grad } u|^2 dx dy \leq \frac{\lambda_k^2 (\bar{\alpha}h)^2 (1 + \bar{\alpha}hK(-))}{2(1 - \bar{\alpha}hK(+))} \\ + (2\bar{\alpha}h)(1 + \bar{\alpha}hK(-)) \left[1 + \left(\frac{\bar{\alpha}h}{2} \right)^2 \frac{K(-)}{1 - \bar{\alpha}hK(+)} \right] \iint_o \left(\frac{\partial u}{\partial s} \right)^2 ds .$$

Given a piecewise continuously differentiable vector field (f^1, f^2) defined in R and on C we can apply the integral identity 2.43 due to Payne and Weinberger to our function u defined by 3.7 to yield

$$(3.18) \quad \int_o f^k \nu_k \left(\frac{\partial u}{\partial s} \right)^2 ds = -2 \iint_R f^i u_{,i} \Delta u dx dy \\ + \iint_R \left[f^k_{,k} \delta^{ij} - 2f^i_{,j} \right] u_{,i} u_{,j} dx dy .$$

If we impose the further condition that $f^i \nu_i > 0$ on C , equation 3.18 yields the inequality

$$(3.19) \quad \int_o \left(\frac{\partial u}{\partial s} \right)^2 ds \leq \frac{\lambda_k}{\min_o f^i \nu_i} \left\{ 2 \sqrt{\lambda_k} \max_R \sqrt{f^i f^i} + \tau \right\} ,$$

where τ is given by 2.44. If R is star shaped with respect to the origin we can let $f^i = x^i$ and 3.19 takes the form

$$(3.20) \quad \int_o \left(\frac{\partial u}{\partial s} \right)^2 ds \leq 2\lambda_k^{3/2} \left[\frac{\max_R \sqrt{x^i x^i}}{\min_o x^i \nu^i} \right] .$$

Methods for prescribing the vector field (f^1, f^2) for more general regions are discussed in Appendix A.

We now return to the estimation of the strip integral appearing in the denominator of 3.13.

$$(3.21) \quad \iint_s u^2 dx dy = \int_0^L \int_0^{\bar{\alpha}h} \left[u(s, 0) + \int_0^n \frac{\partial u}{\partial t}(s, t) dt \right]^2 (1 - Kn) dn ds \\ \leq 2(\bar{\alpha}h)(1 + \bar{\alpha}hK(-)) \int_o u^2 ds + (\bar{\alpha}h)^2 \left(\frac{1 + \bar{\alpha}hK(-)}{1 - \bar{\alpha}hK(+)} \right) \iint_s |\text{grad } u|^2 dx dy .$$

An application of the divergence theorem, used in a similar manner by Payne and Weinberger (45), gives

$$(3.22) \quad \int_o g^j \nu_j u^2 ds = \iint_R u^2 g^j_{,j} dx dy + 2 \iint_R u u_{,j} g^j dx dy ,$$

where (g^1, g^2) is again a piecewise continuously differentiable vector field. Assuming further that $g^j \nu_j > 0$ on C we have

$$\begin{aligned}
(3.23) \quad \int_o u^2 ds &\leq \frac{1}{\min_o g^j \nu_j} \left\{ \max_R (g^i_{,j}) \int_R u^2 dx dy \right. \\
&\quad \left. + 2 \max_R \sqrt{g^i g^j} \sqrt{\int_R u^2 dx dy} \sqrt{D(u)} \right\} \\
&\leq \frac{1}{\min_o g^j \nu_j} \left[\max_R g^i_{,j} + 2 \sqrt{\bar{\lambda}_k} \max_R \sqrt{g^i g^j} \right].
\end{aligned}$$

For star shaped regions we let $g^j = x^j$ and this inequality becomes

$$(3.24) \quad \int_o u^2 ds \leq \frac{2}{\min_o g^j \nu_j} \left[1 + \sqrt{\bar{\lambda}_k} \max_R \sqrt{x^i x^j} \right].$$

Substituting 3.23 into 3.21 yields the inequality

$$\begin{aligned}
(3.25) \quad \iint_S u^2 dx dy &\leq \frac{2\bar{\alpha}h(1 + \bar{\alpha}hK(-))}{\min_o g^j \nu_j} \left[\max_R g^i_{,j} + 2 \sqrt{\bar{\lambda}_k} \max_R \sqrt{g^i g^j} \right] \\
&\quad + (\bar{\alpha}h)^2 \bar{\lambda}_k \left(\frac{1 + \bar{\alpha}hK(-)}{1 - \bar{\alpha}hK(+)} \right),
\end{aligned}$$

which has the desired $O(h)$ property.

By using 3.17, 3.19 and 3.25 in the inequality 3.13 we arrive at the result

$$(3.26) \quad \mu_k(h) \leq \frac{\lambda_k(1 + h B_k)}{(1 - h A_k)},$$

where

$$\begin{aligned}
(3.27) \quad A_k &= \frac{2\bar{\alpha}(1 + \bar{\alpha}hK(-))}{\min_o g^j \nu_j} \left[\max_R g^i_{,j} + 2 \sqrt{\bar{\lambda}_k} \max_R \sqrt{g^i g^j} \right] \\
&\quad + \bar{\alpha}^2 h \bar{\lambda}_k \left(\frac{1 + \bar{\alpha}hK(-)}{1 - \bar{\alpha}hK(+)} \right) + \frac{h \bar{\lambda}_k}{\pi^2}, \\
B_k &= \frac{2\bar{\alpha}\bar{\beta}(1 + \bar{\alpha}hK(-))}{\min_o f^i \nu_i} \left[1 + \left(\frac{\bar{\alpha}h}{2} \right)^2 \frac{K(-)}{(1 - \bar{\alpha}hK(+))} \right] \\
&\quad \cdot \left[2 \sqrt{\bar{\lambda}_k} \max_R \sqrt{f^i f^i} + \tau \right] + \frac{\bar{\beta} \bar{\lambda}_k \bar{\alpha}^2 h}{2} \left(\frac{1 + \bar{\alpha}hK(-)}{1 - \bar{\alpha}hK(+)} \right)
\end{aligned}$$

and $\bar{\lambda}_k$ is an upper bound for λ_k . In particular we can use the upper bound 3.5 for $\bar{\lambda}_k$. Solving 3.26 yields the lower bound

$$(3.28) \quad \lambda_k \geq \frac{\mu_k(h)(1 - h A_k)}{(1 + h B_k)}.$$

As is the case with the fixed membrane the presence of corners on

C causes difficulties in applying the previous development directly. This is because the function $u(x, y)$ may have singular behavior at the corner.

Following the approach of Chapter II we define a sequence of regions $R_1 \subseteq R_2 \subseteq \dots \subseteq R$ with the properties of 2.55. The bounds given by 3.17 and 3.25 can be shown to hold for each R_i where the parameters $\bar{\alpha}, \beta, h, (f^1, f^2), (g^1, g^2)$ are those chosen for R . Of course the locus of centers of curvature for R_i may now enter the strip S near a corner. Where this happens we obtain the bounds for the strip integrals at the corner by special considerations.

Condition (4) of 2.55 assures that we can change the upper limit of integration at the corners in the strip integrals appearing in 3.15 to 3.22 from αh to $K^{-1}(s)$. If C^η and S^η pertain to a typical corner, η , then at that corner 3.15 becomes

$$\begin{aligned}
 (3.29) \quad & \iint_{S_i^\eta} |\text{grad } u|^2 dx dy \\
 &= \int_0^{x_i^\eta} \int_0^{b_i^{(s)}} \left\{ (1 - Kn)^{-2} \left(\frac{\partial u}{\partial s} \right)^2 + \left(\frac{\partial u}{\partial n} \right)^2 \right\} (1 - Kn) dn ds, \\
 &\leq 2(\bar{\alpha}h) \int_{C_i^\eta} \left(\frac{\partial u}{\partial s} \right)^2 ds + \frac{(\bar{\alpha}h)^2}{2} \iint_{S_i^\eta} a^{ij} a^{pm} u|_{i,p} u|_{j,m} dx dy.
 \end{aligned}$$

This is the same as 3.15 with the factor $(1 + \bar{\alpha}hK(-))$ missing. By the same reasoning 3.21 becomes

$$(3.30) \quad \iint_{S_i^\eta} u^2 dx dy \leq 2(\bar{\alpha}h) \int_{C_i^\eta} u^2 ds + (\bar{\alpha}h)^2 \iint_{S_i^\eta} |\text{grad } u|^2 dx dy.$$

The inequalities 3.15–3.24 are true for R_i if we interpret $K(s)$ as $K(s, R_i)$, λ_k as $\lambda_k(R_i)$ and

$$(3.31) \quad K(+, R_i) = \max_{\sigma_i - \sigma_i^\eta} K(s, C_i).$$

We assume the mesh to be placed on R so that no mesh points lie on C in the vicinity of a corner. Then for some N and $i > N$ we have

$$(3.32) \quad \mu_k(h) = \mu_k(h, R_i).$$

Then 3.26 becomes

$$(3.33) \quad \mu_k(h) \leq \frac{\lambda_k(R_i)(1 + hB_k(R_i))}{(1 - hA_k(R_i))}.$$

By a result of D. M. Eidus [16, 17] the hypotheses of 2.55 imply that

$$\lim_{i \rightarrow \infty} \lambda_k(R_i) = \lambda_k.$$

Also

$$\lim_{i \rightarrow \infty} B_k(R_i) = B_k ,$$

$$\lim_{i \rightarrow \infty} A_k(R_i) = A_k ,$$

since B_k and $B_k(R_i)$ differ only in the value of the argument λ_k . Hence the bound 3.28 holds in the case of corners with interior angle less than 180° .

APPENDIX A

THE VECTOR FIELD (f^1, f^2)

One manner of choosing the vector field (f^1, f^2) for regions star-shaped with respect to the origin which generalizes readily to more arbitrary domains is the following:

Introduce polar coordinates (r, θ) given by

$$(4.1) \quad \begin{aligned} x^1 &= r \cos \theta , \\ x^2 &= r \sin \theta . \end{aligned}$$

Let the boundary C be a smooth curve with the polar representation

$$(4.2) \quad C: r = \rho(\theta) .$$

The vector field (f^1, f^2) is defined as

$$(4.3) \quad \begin{aligned} f^1(x^1, x^2) &= \frac{r}{\rho(\theta)} \nu^1(\theta) = \frac{r}{\rho(\theta)} t^2(\theta) , \\ f^2(x^1, x^2) &= \frac{r}{\rho(\theta)} \nu^2(\theta) = - \frac{r}{\rho(\theta)} t^1(\theta) . \end{aligned}$$

In 4.3, (t^1, t^2) , the unit tangent vector, is given by

$$(4.4) \quad \begin{aligned} t^1(\theta) &= \frac{dx^1}{ds} = \left(\frac{dx^1}{d\theta} \right) \left(\frac{d\theta}{ds} \right) = H^{-1} \frac{d}{d\theta} \left[\rho(\theta) \cos \theta \right] , \\ t^2(\theta) &= H^{-1} \frac{d}{d\theta} \left[\rho(\theta) \sin \theta \right] , \end{aligned}$$

with

$$(4.5) \quad H = \sqrt{\rho^2 + \dot{\rho}^2} .$$

A computation shows that

$$(4.6) \quad \begin{aligned} \min_{\sigma} f^i \nu_i &= 1 , \\ \max_R \sqrt{f^i f^i} &= 1 , \end{aligned}$$

$$\tau = \max_R \left\{ \frac{1}{H^4 \rho^4} \left[\rho^2 \dot{\rho}^2 (\rho^2 + 3\dot{\rho}^2) + \rho^3 \ddot{\rho} (\rho \ddot{\rho} - 2\dot{\rho}^2) + \dot{\rho}^6 \right] \right\}^{1/2}.$$

It is easily verified that (f^1, f^2) is piecewise continuously differentiable since, by our assumption, $\nu^i(s)$ is continuously differentiable on C . Indeed, a discontinuity in the derivative of $(\nu^1(\theta), \nu^2(\theta))$ at θ_0 will be propagated as a discontinuity in the derivative of (f^1, f^2) along the radius vector joining $(0, 0)$ to $(\rho(\theta_0), \theta_0)$. If $(\nu^1(\theta), \nu^2(\theta))$ is itself discontinuous at a point θ_0 , i.e., if the boundary C has a corner, then the quantity τ in 4.6 does not exist. In this case we replace $(\nu^1(\theta), \nu^2(\theta))$ in 4.3 by a continuously differentiable vector field $(g^1(\theta), g^2(\theta))$ which has the property

$$(4.7) \quad g^i(\theta) \nu_i(\theta) > 0.$$

The values of the quantities

$$(4.8) \quad f^i \nu_i, \sqrt{f^i f^i}, \tau,$$

are then computed as functions of (g^1, g^2) .

If R is not itself star-like we subdivide R into star shaped regions R_1, R_2, \dots, R_N whose bounding curves are given by C_1, C_2, \dots, C_N . Let

$$(4.9) \quad C_i = C'_i \cup C''_i, \quad i = 1, 2, \dots, N,$$

where C'_i is that portion of C_i which coincides with C , (see Fig. 3) We define the vector field (g^1_i, g^2_i) on C_i in terms of polar coordinates (R_i, θ_i)

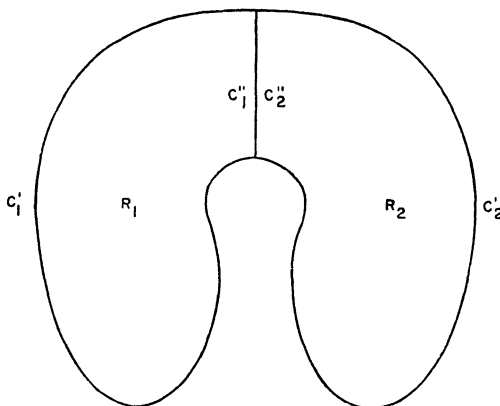


Fig. 3.

introduced in R_i . The functions $(g^1_i(\theta_i), g^2_i(\theta_i))$ are assumed continuous on C_i with the property

$$(4.10) \quad g^j_i(\theta_i) \nu_j > 0$$

on C'_i . We further assume that the functions (g^1, g^2) given by

$$(4.11) \quad g^j(x^1, x^2) = g_i^j(\theta_i), \quad (x^1, x^2) \in C_i,$$

form a piecewise continuously differentiable vector field defined on $C_1 UC_2 U \dots UC_N$. We define the vector field

$$(4.12) \quad f^j(x^1, x^2) = \frac{r_i}{\rho_i(\theta_i)} g_i^j(\theta_i), \quad j = 1, 2, (x^1, x^2) \in R_i,$$

in R . The quantities 4.8 are then found to be

$$(4.13) \quad \begin{aligned} \min_O f^j \nu_j &= \min_{i=1, \dots, N} \min_{O_i^1} f^j \nu_j, \\ \max_R \sqrt{f^j f^j} &= \max_{i=1, \dots, N} \max_{R_i} \sqrt{f^j f^j}, \\ \tau &= \max_{i=1, \dots, N} \max_{R_i} \left[(f_{,1}^1 - f_{,2}^2)^2 + (f_{,2}^1 + f_{,1}^2)^2 \right]^{1/2}. \end{aligned}$$

APPENDIX B

THE PARAMETERS α , $\bar{\alpha}$ AND β

1. **A bound for α .** For the parameter α we have the bound

$$(5.1) \quad \alpha \leq \frac{1}{2} \sqrt{34},$$

as will now be shown under the assumptions of Chapter II, §4. We recall that the strip S is the region swept out by a segment of the inward normal length αh as that normal moves along the curve C . The parameter α must be chosen large enough so that S will cover the squares associated with the points of B_h .

In Fig. 4 we have pictured a typical case. Here the points A, E, G , and H belong to C_h and B belongs to the set B_h . The point F belonging to the square associated with B is the center of the two circular arcs shown in the figure. Let J be a point on C with inward normal \vec{n} and tangent \vec{t} such that \vec{n} passes through F . For inequality 5.1 to hold we must show that J cannot lie outside the circle through G about F whose radius is $h/2\sqrt{34}$.

Assume that J does lie outside the arc through G , as shown in the figure, with the center of curvature of C at F . This assumption gives our boundary curve the maximum curvature at J allowable under the hypotheses. As the curve C continues onward it must intersect one of the triangles associated with the point A . At the same time the family of normals of C cannot intersect within the strip S , i.e., the locus of centers of curvature cannot move toward J . The curve C can most rapidly turn toward the triangle associated with A if the center of curvature remains

R

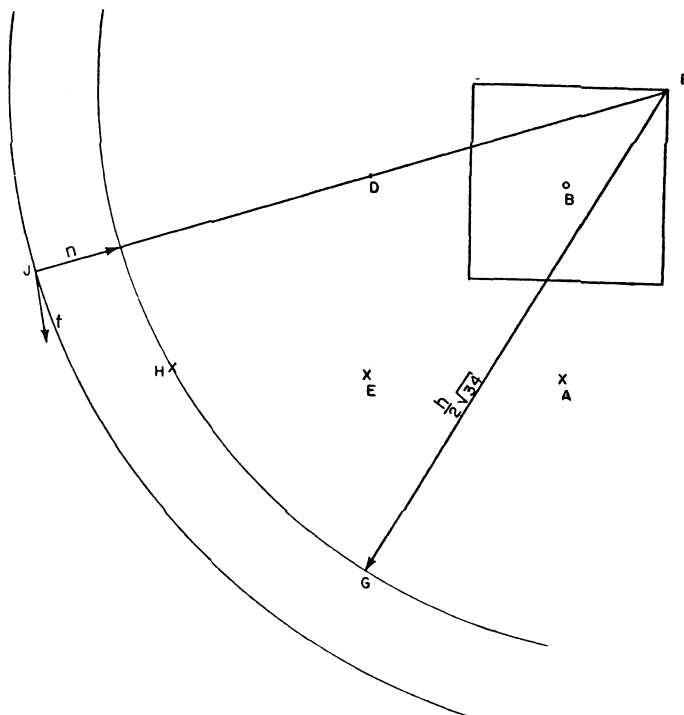


Fig. 4.

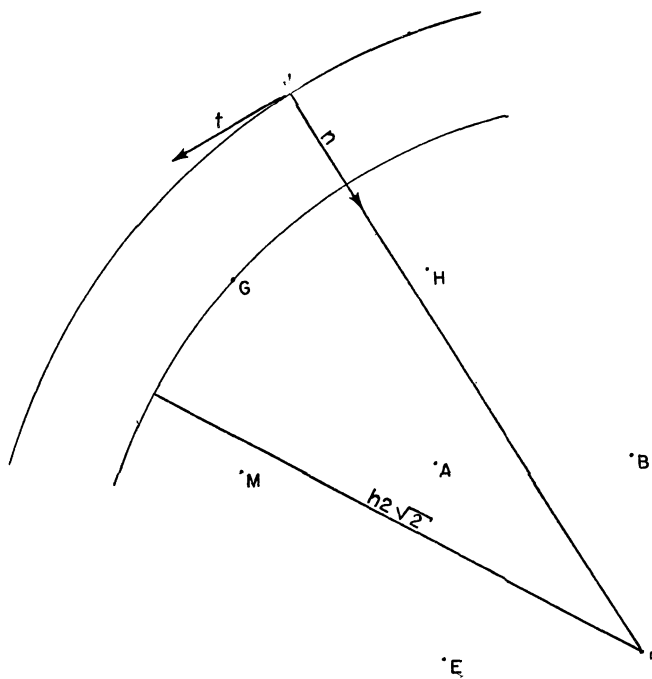


Fig. 5.

at F . Clearly such a circle with center at F will never intersect any triangle associated with A and we have arrived at a contradiction.

2. A bound for $\bar{\alpha}$. For the parameter $\bar{\alpha}$ we have the bound

$$(5.2) \quad \bar{\alpha} \leq 2\sqrt{2} ,$$

which is smaller than 5.1 since the strip S as defined in Chapter III § 1, is somewhat narrower. In fact S must cover only those squares which intersect C and their neighbors in R with whom they share a common side.

In Fig. 5 we have a typical situation where the square $FBAE$ shares a common vertex with $AHGM$ through which C is assumed to pass. The point F is taken as center of the two circles in the figure. Again we assume J to be a point on C whose inward normal \vec{n} passes through F . For inequality 5.2 to hold we must show that J cannot lie outside the circle through G about F whose radius is $2h\sqrt{2}$.

As before we assume that J does lie outside the arc through G , as shown in the figure, with the center of curvature of C at F . This assumption gives our boundary curve the maximum curvature at J allowed by the hypotheses. As the curve C continues onward, it must intersect the square $AHGM$. The curve C can most rapidly turn toward $AHGM$ under the constraints imposed if it is the circle with center at F . Hence the contradiction.

3. A bound for β . We shall first show that, under the assumption the points of C_h^* form a double fence of the type shown in Fig. 6, an upper bound $\bar{\beta}$ for β is

$$(5.3) \quad \beta \leq \bar{\beta} = 4 .$$

The actual selection of W_i minimizes $\beta(W)$ over C_h^* and hence any other assignment of values to W_i will provide an upper bound for β . First, we shall see how to assign values to the points of Fence 2 in terms of the values at points of Fence 1 so that

$$(5.4) \quad D^{(h)}_{\text{Fence 1} + \text{Fence 2}}(V) - D^{(h)}_{\text{Fence 1}}(V) \leq 2 D^{(h)}_{\text{Fence 1}}(V) .$$

In fact assume that the capital letters represent the values of V on Fence 1 in the figure, and consider the assignment made there at points of Fence 2. It is readily verified that property 5.4 holds for this scheme of assigning values to Fence 2. Therefore see that one can step inward from Fence n one fence at a time in assigning values with the result that

$$(5.5) \quad D^{(h)}_{\sum_{n=1}^N \text{Fence } n}(V) - D^{(h)}_{\text{Fence 1}}(V) \leq 2^{N-1} D^{(h)}_{\text{Fence 1}}(V) .$$

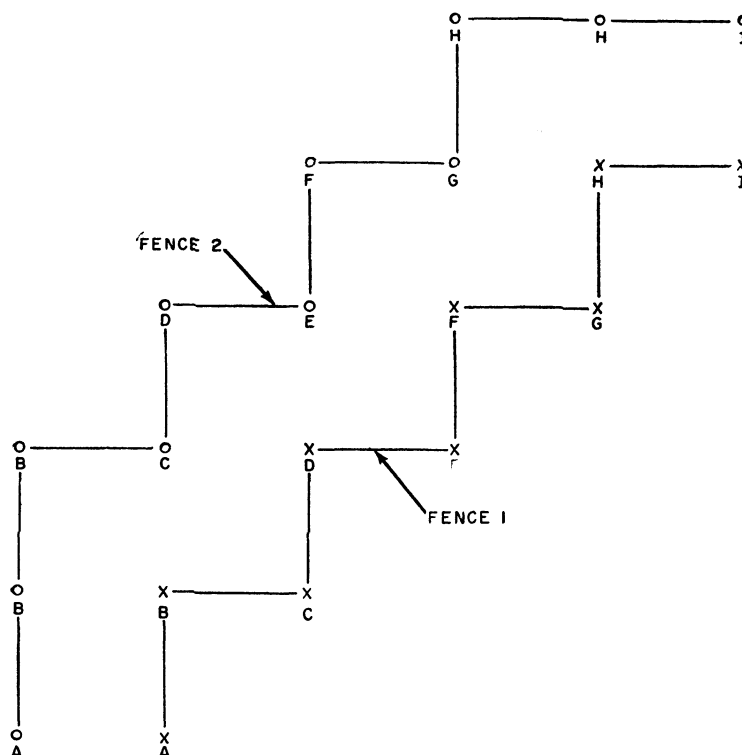


Fig. 6.

Applying 5.5 to 3.9 with $N = 3$ gives the result 5.3. The requirement on h that $(1 - \alpha h K) > 0$ isn't sufficient to allow us to take $N = 3$. In addition to a local condition of this nature a restriction "in the large" is also called for. In particular we must assure that the squares of $R_h - C_h^*$ form approximately the same geometrical figure as R . For example, if R is hour-glass shaped, the squares of $R_h - C_h^*$ might form doubly-connected regions. If R is snake-like the set $R_h - C_h^*$ might even be empty. Such an additional condition is as follows: No adjoining squares of C_h^* may represent remote sections of the boundary C .

APPENDIX C

SOME EXAMPLES

In this section we seek to gain some knowledge of the relative size of the quantities involved in the bounds for the fixed membrane. Computations have been made here for the first eigenvalue of two geometric configurations; (a) the square, and (b) the unit circle. In each case λ_1 and u_1 are known so the strip integral 2.28 can be evaluated explicitly

and compared with the bound obtained from § 2 of Chapter II.

1. The unit square. Let R be the unit square with vertices $(0, 0)$, $(0, 1)$, $(1, 1)$, and $(1, 0)$. The first eigenvalue and eigenfunction are

$$(6.1) \quad \begin{aligned} \lambda_1 &= 2\pi^2, \\ u_1 &= 2 \sin \pi x \sin \pi y. \end{aligned}$$

If we assume $h = \frac{1}{N}$, N an integer, then we can take

$$(6.2) \quad \alpha = \frac{3}{2}.$$

Let the vector field (f^1, f^2) be chosen as

$$(6.3) \quad \begin{aligned} f^1 &= x - \frac{1}{2}, \\ f^2 &= y - \frac{1}{2}, \end{aligned}$$

so that

$$(6.4) \quad \begin{aligned} \tau &= 0, \\ \min_{\sigma} f^i \nu_i &= \frac{1}{2}, \\ \max_R \sqrt{f^i f^i} &= \frac{\sqrt{2}}{2}. \end{aligned}$$

A direct computation shows that in 2.51

$$(6.5) \quad B_1 = 27\pi + O(h).$$

The quantity B_1 involves bounds on the strip integral. Define \bar{B}_1 to be the analogue of B_1 with the strip integral substituted for the bound. Then

$$(6.6) \quad \bar{B}_1 = \frac{2 \iint_S u_1^2 dx dy}{h^3 \lambda_1}.$$

We see that in this case

$$(6.7) \quad \bar{B}_1 = \frac{27}{2} + O(h^2).$$

The difference between the upper and lower bounds given by 2.53 be-

comes

$$(6.8) \quad \mu_1(h)27\pi h + O(h^2) .$$

For this difference to have the same magnitude as $\mu_1(h)$ would require our putting

$$(6.9) \quad h \approx \frac{7}{27\pi} .$$

4. The unit circle. Let R be the the unit circle. Its first eigenvalue and eigenfunction are

$$(6.10) \quad \begin{aligned} \lambda_1 &\approx 2.4^2 , \\ u_1 &= [\sqrt{\pi} J_1(\sqrt{\lambda_1})]^{-1} J_0(r \sqrt{\lambda_1}) . \end{aligned}$$

Let the vector field (f^1, f^2) be given by

$$(6.11) \quad \begin{aligned} f^1 &= x , \\ f^2 &= y . \end{aligned}$$

so that

$$(6.12) \quad \begin{aligned} \tau &= 0 , \\ \min_o f^i \nu^i &= 1 , \\ \max_R \sqrt{f^i f^i} &= 1 . \end{aligned}$$

Again a direct computation yields

$$(6.13) \quad \begin{aligned} B_1 &\approx 38.4 \alpha^3 + O(h) , \\ \bar{B}_1 &\approx 1.7 \alpha^3 + O(h) . \end{aligned}$$

It is clear from these two examples that bounds of this paper are not sufficiently sharp to be used in the actual computation of eigenvalues. In fact, the finite difference problem posed appears far from optimal. It is hoped, however, that the techniques employed here may be useful in obtaining bounds by other finite difference approximations.

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