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**BOUNDARY BEHAVIOR OF RANDOM VALUED HEAT  
POLYNOMIAL EXPANSIONS**

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# BOUNDARY BEHAVIOR OF RANDOM VALUED HEAT POLYNOMIAL EXPANSIONS

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This paper is concerned with random series of the form  $\sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t)$  where the  $X_n$ 's are random variables, the  $a_n$ 's are real numbers, and the  $v_n$ 's are heat polynomials as introduced by P. C. Rosenbloom and D. V. Widder. The sequences  $\{a_n\}$  are assumed to satisfy  $\limsup_{n \rightarrow \infty} |a_n|^{2/n} (2n/e) = 1$  which implies  $\sum_{n=0}^{\infty} a_n v_n(x, t)$  has  $|t| < 1$  as its strip of convergence, i.e., it converges to a  $C^2$ -solution of the heat equation in  $|t| < 1$  and does not converge everywhere in any larger open strip. Associated with each sequence  $\{a_n\}$  is its classification number from  $[0, 1]$  which indicates how rapidly  $a_n$  tends to zero. Assumptions are placed on the random variables which imply that for almost every  $\omega$  the series  $\sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t)$  has  $|t| < 1$  as its strip of convergence.

The main results of the paper are two theorems. The first states that if  $\{a_n\}$  has its classification number in  $[0, 1/2]$ , then for almost every  $\omega$  the lines  $t = 1$  and  $t = -1$  form the natural boundary for  $\sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t)$ . The second is concerned with sequences having their classification numbers in  $(1/2, 1]$ . The conclusion implies that for almost every  $\omega$  no interval of either of the lines  $t = 1$  or  $t = -1$  is part of the natural boundary for  $\sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t)$ .

The present work had its original motivation in the study of the boundary behavior of random power series. These are series of the form  $\sum_{n=0}^{\infty} a_n(\omega) z^n$  where the  $a_n$ 's are complex valued random variables and  $z$  is a complex number. Reference [1] contains a history of results in this area. One of the early results which helped to motivate the first part of the proof of our Theorem 1 is due to Paley and Zygmund in a 1932 paper [see 6, p. 220]. In this theorem it is assumed that  $\sum_{n=0}^{\infty} a_n z^n$  is an ordinary power series with a finite radius of convergence. Letting  $\{\phi_n\}$  be the sequence of Rademacher functions, the conclusion is that for almost every  $\omega$  in  $[0, 1]$  the series  $\sum_{n=0}^{\infty} \phi_n(\omega) a_n z^n$  has its circle of convergence as its natural boundary.

More recently [see 3] V. L. Shapiro has considered series of the form  $\sum_{n=0}^{\infty} X_n(\omega) H_n(x)$  where the  $X_n$ 's are random variables and

$$\sum_{n=0}^{\infty} H_n(x)$$

is the spherical harmonic representation of a harmonic function in the unit ball. The harmonic continuability across the boundary of the unit ball of the functions  $\sum_{n=0}^{\infty} X_n(\omega) H_n(x)$  was investigated. This

work further motivated the first part of the proof of our Theorem 1 and influenced our choice of the class of random variables to be considered.

2. **Definitions and preliminary comments.** For a point  $(x_0, t_0)$  in the plane and a number  $\rho > 0$  we let

$$S(x_0, t_0; \rho) = \{(x, t): |x - x_0| < \rho \text{ and } |t - t_0| < \rho\}.$$

If  $u(x, t)$  is a  $C^2$ -solution to the heat equation in the strip  $|t| < \sigma$  we say the line  $t = -\sigma$  ( $t = \sigma$ ) is part of the natural boundary for  $u$  in case for every  $x_0$  and every  $\rho > 0$  there is no  $C^2$ -solution  $v(x, t)$  in  $S(x_0, -\sigma; \rho)$  ( $S(x_0, \sigma; \rho)$ ) which agrees with  $u(x, t)$  where  $u$  and  $v$  are both defined.

Let  $E_0$  be the set of all sequences  $\{a_n\}_{n=0}^\infty$  of real numbers. For  $r > 0$  let

$$E_r = \{\{a_n\} \in E_0: |a_n| (2n/e)^{n/2} = O(e^{-nr}) \text{ as } n \rightarrow \infty\}.$$

We call  $\sup \{r: \{a_n\} \in E_r\}$  the classification number of  $\{a_n\}$ .

Explicitly, from [2, p. 222]

$$(2.1) \quad v_n(x, t) = n! \sum_{k=0}^{\lfloor n/2 \rfloor} \frac{x^{n-2k}}{(n-2k)!} \frac{t^k}{k!}, \quad n = 0, 1, \dots.$$

In [2, Th. 5.3, p. 231] it was shown that the series  $\sum_{n=0}^\infty a_n v_n(x, t)$  converges to a  $C^2$ -solution of the heat equation in the strip  $|t| < \sigma$  where

$$(2.2) \quad \sigma = (\limsup |a_n|^{2/n} (2n/e))^{-1}$$

and that this strip is the largest open strip of convergence of the series. One easily shows that sequences  $\{a_n\}$  satisfying

$$\limsup |a_n|^{2/n} (2n/e) = 1$$

have their classification numbers in  $[0, 1]$ .

We will make repeated use of the following bounds which appear in [4] by S. Täcklind. Assume  $u(x, t)$  is continuous on the rectangle  $R = \{(x, t): |x| \leq \mathcal{L}, 0 \leq t \leq T\}$ , is a  $C^2$ -solution to the heat equation in the interior of  $R$ , and  $\mu$  is an upper bound for  $|u(x, t)|$  on  $R$ ; then  $u(x, t)$  is in class  $C^\infty$  on the interior of  $R$  and for  $n = 2, 3, \dots$ ,  $|x| < \mathcal{L}$ , and  $0 < t \leq T$

$$(2.3) \quad \left| \frac{\partial^n u}{\partial x^n}(x, t) \right| \leq \frac{\mu}{2\sqrt{\pi}} \frac{2^{(n+3)/2}}{t^{n/2}} \Gamma((n+1)/2) \\ + \frac{\mu}{\sqrt{\pi}} \left(\frac{\pi}{2}\right)^{5/2} \frac{2^{3n/2}}{(\mathcal{L} - |x|)^n} \Gamma(n+1).$$

3. THEOREM 1. Let  $\{X_n\}_{n=0}^\infty$  be a sequence of symmetric independent random variables defined on the complete probability space  $(\Omega, \mathcal{F}, P)$  and satisfying

(i) there exists a number  $M$  such that

$$\int_{\Omega} |X_n(\omega)|^2 dP(\omega) \leq M \text{ for } n = 0, 1, \dots, \text{ and}$$

(ii) there exists  $N > 0$  such that

$$N \leq \int_{\Omega} |X_n(\omega)| dP(\omega), n = 0, 1, \dots$$

Assume  $\{a_n\}$  satisfies  $\limsup |a_n|^{2/n} (2n/e) = 1$  and has its classification number in  $[0, 1/2)$ . Then for almost every  $\omega$  in  $\Omega$  the lines  $t = 1$  and  $t = -1$  form the natural boundary for

$$u_{\omega}(x, t) = \sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t).$$

*Proof.* Letting  $\Omega' = \{\omega \in \Omega : \sum_{n=0}^{\infty} X_n(\omega) a_n v_n(x, t) \text{ converges in the strip } |t| < 1\}$ , we will first show  $P(\Omega') = 1$ . Clearly

$$[\limsup |X_n|^{2/n} \leq 1] \supseteq \bigcup_{k=1}^{\infty} \bigcap_{n=k}^{\infty} [|X_n| \leq nM^{1/2}]$$

and by the Borel-Cantelli Lemma the last set has probability 1 since  $P[|X_n| > nM^{1/2}] \leq 1/n^2$  from (i). Hence

$$P\{\omega : \limsup |X_n(\omega) a_n|^{2/n} (2n/e) \leq 1\} = 1$$

which by (2.2) shows  $P(\Omega') = 1$ .

The following fact is essentially a merger of Lemma 1 from [3] and a special case of Lemma 2 from [3]. There exist numbers  $A$  in  $(0, 1)$  and  $B > 0$  with the following property: for  $E \in \mathcal{F}$  with  $P(E) > A$  there is a positive integer  $n_0$  such that for  $n \geq n_0$ , every sequence  $\{c_j\}_{j=0}^{\infty}$  of real numbers, and  $k \geq 1$  we have

$$(3.1) \quad \sum_{j=n}^{n+k} c_j^2 \leq B \int_E \left\{ \sum_{j=n}^{n+k} c_j X_j(\omega) \right\}^2 dP(\omega).$$

We will show that for almost every  $\omega$  the line  $t = -1$  is part of the natural boundary for  $u_{\omega}$  and will use this in the proof for the line  $t = 1$ .

Assume it is false that for a.e.  $\omega$  in  $\Omega$  the line  $t = -1$  is part of the natural boundary for  $u_{\omega}$ . The first part of the argument we give in order to obtain a contradiction is analogous to parts of the proof of Theorem 1 in [3] by V. L. Shapiro. We will employ (2.3), (3.1), and an asymptotic estimate for heat polynomials from [2] in

order to obtain conditions on the sequence  $\{a_n\}$  which contradict the fact that the classification number of  $\{a_n\}$  is in  $[0, 1/2)$ .

Let  $E = \{\omega \in \Omega': t = -1 \text{ is not part of the natural boundary for } u_\omega\}$ . Then either (i)  $E \notin \mathcal{F}$ , or (ii)  $E \in \mathcal{F}$  and  $P(E) > 0$ . Using the fact that the real line is separable and the countable additivity of the probability  $P$ , it follows that there exist a real number  $x_0$  and  $\rho_0 > 0$  such that  $E_1 = \{\omega \in E: \text{there is a } C^2\text{-solution to the heat equation in } S(x_0, -1; \rho_0) \text{ which agrees with } u_\omega \text{ where they are both defined}\}$  satisfies either (i)  $E_1 \notin \mathcal{F}$ , or (ii)  $E_1 \in \mathcal{F}$  and  $P(E_1) > 0$ . For  $i = 1, 2, \dots$  define

$$E_{2,i} = \left\{ \omega \in \Omega': \left| \frac{\partial^m u_\omega}{\partial x^m}(x, t) \right| \leq i^m m^m \text{ for } (x, t) \text{ in } S\left(x_0, -1; \frac{\rho_0}{2}\right), \right. \\ \left. |t| < 1, \text{ and } m = i, i+1, \dots \right\}$$

and let  $E_2 = \bigcup_{i=1}^{\infty} E_{2,i}$ .  $E_2$  is in the tail  $\sigma$ -field generated by the independent  $X_n$ 's. From (2.3) it follows that  $E_1 \subseteq E_2$ . By Kolmogorov's zero-one law  $P(E_2) = 1$ . Let  $A$  and  $B$  be as in (3.1). Take  $i_0$  sufficiently large that  $P(E_{2,i_0}) > A$  and let  $n_0$  correspond to  $E_{2,i_0}$  as in (3.1). Now let  $m \geq \max\{n_0, i_0\}$  and let  $(x, t)$  be in  $S(x_0, -1; \rho_0/2)$  with  $|t| < 1$ . Then by (3.1) for  $k = 1, 2, \dots$

$$\sum_{n=m}^{m+k} \left[ \frac{n!}{(n-m)!} a_n v_{n-m}(x, t) \right]^2 \\ \leq B \int_{E_{2,i_0}} \left[ \sum_{n=m}^{m+k} \frac{n!}{(n-m)!} a_n v_{n-m}(x, t) X_n(\omega) \right]^2 dP(\omega) .$$

Making use of the independence and symmetry of the random variables and of condition (i) we see that the integrand of the last integral is Cauchy in the variable  $k$  in  $L^1(\Omega)$  and thus in  $L^1(E_{2,i_0})$ . Hence

$$\sum_{n=m}^{\infty} \left[ \frac{n!}{(n-m)!} a_n v_{n-m}(x, t) \right]^2 \\ \leq B \int_{E_{2,i_0}} \left[ \sum_{n=m}^{\infty} \frac{n!}{(n-m)!} a_n v_{n-m}(x, t) X_n(\omega) \right]^2 dP(\omega) \\ = B \int_{E_{2,i_0}} \left| \frac{\partial^m u_\omega}{\partial x^m}(x, t) \right|^2 dP(\omega) \leq B i_0^{2m} m^{2m}$$

with the last inequality following from the definition of  $E_{2,i_0}$ . We conclude that for every  $m \geq \max\{n_0, i_0\}$ , every  $n \geq m$ , and every  $(x, t)$  in  $S(x_0, -1; \rho_0/2)$  with  $|t| < 1$ ; we have

$$(3.2) \quad \frac{n!}{(n-m)!} |a_n| |v_{n-m}(x, t)| \leq B^{1/2} i_0^m m^m .$$

It follows from Theorem 3.1 of [2] that there exist numbers  $A$  and  $l_0$  such that for  $n \geq l_0$

$$\sup_{|x-x_0| < \rho_0/2} |v_n(x, -1)| \geq A[2n/e]^{n/2}.$$

Thus from (3.2) we have for  $n > m + l_0 > m \geq \max\{n_0, i_0\}$

$$|a_n| \frac{n!}{(n-m)!} A[2(n-m)/e]^{(n-m)/2} \leq B^{1/2} i_0^m m^m.$$

Employing Stirling's theorem we see there is a number  $C$  such that for  $n > m + l_0 > m \geq \max\{n_0, i_0\}$

$$(3.3) \quad |a_n| (2n/e)^{n/2} \leq \left[ \frac{Cm}{\sqrt{n-m}} \right]^m \cdot ((n-m)/n)^{(n+1)/2}.$$

Let  $r$  be a number which is strictly greater than the classification number of  $\{a_n\}$  and strictly less than  $1/2$ . Let  $m$  be related to  $n$  by  $m = [4n^r] + 1$  where the brackets denote the greatest integer function. Then from (3.3), for sufficiently large  $n$ ,

$$(3.4) \quad |a_n| (2n/e)^{n/2} \leq (1 - 4/n^{1-r})^{(n^{1-r/4}) \cdot 2 \cdot n^r}.$$

For large enough  $n$ ,  $(1 - 4/n^{1-r})^{(n^{1-r/4}) \cdot 2} \leq 1/e$  and thus from (3.4) we have for such  $n$ ,  $|a_n| (2n/e)^{n/2} \leq 1/e^{n^r}$ . Hence  $\{a_n\} \in E_r$  which contradicts the fact that  $r$  is strictly greater than the classification number of  $\{a_n\}$  and concludes the proof for the line  $t = -1$ .

For the last part of the proof we find it convenient to introduce the probability space  $(R^\omega, \mathscr{A}', \mu')$  which we now describe.

$$R^\omega = \prod_{n=0}^{\infty} R_n$$

where each  $R_n$  is the set of real numbers. Let  $\mathscr{A}_0$  be the field of all subsets of  $R^\omega$  of the form  $B \times (\prod_{n=n_0+1}^{\infty} R_n)$  where  $n_0$  is a positive integer and  $B$  is a Borel set in  $\prod_{n=0}^{n_0} R_n$ . Let  $\mathscr{A}$  be the  $\sigma$ -field generated by  $\mathscr{A}_0$ . Let  $\mu$  be the probability on  $(R^\omega, \mathscr{A})$  which is induced by the  $X_n$ 's. Then  $(R^\omega, \mathscr{A}', \mu')$  is the completion of  $(R^\omega, \mathscr{A}, \mu)$ .

Let  $\{\eta_i\}_{i=0}^{\infty}$  be a sequence of  $\pm 1$ 's. Define  $T: R^\omega \rightarrow R^\omega$  by

$$T((\xi_0, \xi_1, \dots)) = (\eta_0 \xi_0, \eta_1 \xi_1, \dots).$$

Notice that

$$\begin{aligned} \mu \left( \prod_{n=0}^{n_0} (a_n, b_n] \times \prod_{n=n_0+1}^{\infty} R_n \right) &= \prod_{n=0}^{n_0} P[X_n \in (a_n, b_n]] \\ &= \prod_{n=0}^{n_0} P[X_n \in \eta_n(a_n, b_n)] = \mu \left( T \left( \prod_{n=0}^{n_0} (a_n, b_n] \times \prod_{n=n_0+1}^{\infty} R_n \right) \right) \end{aligned}$$

where we have used both the independence and symmetry of the  $X_n$ 's. From this it follows that for  $A \in \mathcal{N}'$ ,  $\mu'(A) = \mu'(T(A))$ . We will make use of this fact twice in the remainder of this proof.

To finish the proof it suffices to show that for a.e.  $p \in R^\omega$  the line  $t = 1$  is part of the natural boundary for

$$u_p(x, t) = \sum_{n=0}^{\infty} \pi_n(p) a_n v_n(x, t)$$

where the  $\pi_n$ 's are the projection random variables. Suppose this is false. From the first paragraph of the present proof we know  $R^{\omega'} = \{p \in R^\omega: \sum_{n=0}^{\infty} \pi_n(p) a_n v_n(x, t) \text{ converges in } |t| < 1\}$  has  $\mu'$ -measure 1. Now let  $F = \{p \in R^{\omega'}: t = 1 \text{ is not part of the natural boundary for } u_p\}$ . Then either (i)  $F \in \mathcal{N}'$ , or (ii)  $F \in \mathcal{N}'$  and  $\mu'(F) > 0$ . It follows that there exist numbers  $a, b, \rho$  with  $a < b$  and  $\rho > 0$  such that  $F_1 = \{p \in R^{\omega'}: \text{there is a function } v_p(x, t) \text{ which is continuous on } a \leq x \leq b, 0 \leq t \leq 1 + \rho; \text{ is a } C^2\text{-solution to the heat equation for } a < x < b, 0 < t < 1 + \rho; \text{ and agrees with } u_p(x, t) \text{ in } a \leq x \leq b, 0 \leq t < 1\}$  satisfies either (i)  $F_1 \in \mathcal{N}'$ , or (ii)  $F_1 \in \mathcal{N}'$  and  $\mu'(F_1) > 0$ . But  $F_1 = \{p \in R^{\omega'}: \lim_{t \uparrow 1} u_p(a, t) \text{ and } \lim_{t \uparrow 1} u_p(b, t) \text{ both exist}\}$ .  $F_1$  is in the tail  $\sigma$ -field generated by the independent  $\pi_n$ 's. From the zero-one law,  $\mu'(F_1) = 1$ .

Either  $a \neq 0$  or  $b \neq 0$  and for definiteness we assume  $a \neq 0$ . Then  $F_2 = \{p \in R^{\omega'}: \lim_{t \uparrow 1} u_p(a, t) \text{ exists}\}$  has  $\mu'(F_2) = 1$ . Let  $T: R^\omega \rightarrow R^\omega$  be defined by  $T((\xi_0, \xi_1, \dots)) = (\xi_0, -\xi_1, \xi_2, -\xi_3, \dots)$ . By our earlier comments concerning such mappings we have  $\mu'(F_2 \cap T(F_2)) = 1$ . For  $p \in R^{\omega'}$  and  $|t| < 1$  one checks that  $u_{T(p)}(-a, t) = u_p(a, t)$ . Hence for  $p \in F_2 \cap T(F_2)$ ,  $\lim_{t \uparrow 1} u_p(-a, t)$  and  $\lim_{t \uparrow 1} u_p(a, t)$  both exist. Thus for  $p \in F_2 \cap T(F_2)$  there is a function  $w_p(x, t)$  which is continuous in  $|x| \leq a, 0 \leq t \leq 2$ ; is a  $C^2$ -solution to the heat equation in  $|x| < a, 0 < t < 2$ ; and agrees with  $u_p$  in  $|x| \leq a, 0 \leq t < 1$ . For  $p \in F_2 \cap T(F_2)$  and  $0 \leq t \leq 2$  let  $\phi_p(t) = w_p(0, t)$  and  $\psi_p(t) = (\partial w_p / \partial x)(0, t)$ . Then, employing (2.3), we see that  $\phi_p$  and  $\psi_p$  are in class  $C\{(2n)!\}$  on  $[0, 2]$  (a function  $f$  is in class  $C\{(2n)!\}$  on an interval  $I$  if  $f$  is in class  $C^\infty$  on  $I$  and there exist constants  $\beta$  and  $B$  such that for every  $t$  in  $I$ ,  $|f^{(n)}(t)| \leq \beta B^n (2n)!, n = 0, 1, \dots$ ).

Now let  $T': R^\omega \rightarrow R^\omega$  be defined by

$$T'((\xi_0, \xi_1, \dots)) = (\xi_0, \xi_1, -\xi_2, -\xi_3, \xi_4, \xi_5, -\xi_6, -\xi_7, \dots).$$

Then for  $p \in R^{\omega'}$  and  $|t| < 1$ ,  $u_p(0, t) = u_{T'(p)}(0, -t)$  and

$$\frac{\partial u_p}{\partial x}(0, t) = \frac{\partial u_{T'(p)}}{\partial x}(0, -t).$$

For  $p$  in the almost sure set  $T'(F_2 \cap T(F_2))$  we have  $T'(p) \in F_2 \cap T(F_2)$  and we define  $\phi'_p$  and  $\psi'_p$  on  $[-2, 0]$  by  $\phi'_p(t) = \phi_{T'(p)}(-t)$  and

$$\psi'_p(t) = \psi_{T'(p)}(-t)$$

thereby obtaining class  $C\{(2n)!\}$  extensions of  $u_p(0, t)$  and  $(\partial u_p / \partial x)(0, t)$  on  $[-1, 0]$ . Thus for  $p \in T'(F_2 \cap T(F_2))$

$$u'_p(x, t) = \sum_{n=0}^{\infty} \frac{\phi'_p{}^{(n)}(t)x^{2n}}{(2n)!} + \sum_{n=0}^{\infty} \frac{\psi'_p{}^{(n)}(t)x^{2n+1}}{(2n+1)!}$$

provides a solution to the heat equation which is a  $C^2$ -extension of  $u_p$  into some rectangle  $|x| < r, -2 < t < 0$  which contradicts the first part of the proof.

**4. THEOREM 2.** *Let  $\{X_n\}$  be a sequence of independent random variables over a probability space  $(\Omega, \mathcal{F}, P)$  which satisfies (i) and (ii) of Theorem 1. Assume  $\{a_n\}$  satisfies  $\limsup |a_n|^{2/n}(2n/e) = 1$  and has its classification number in  $(1/2, 1]$ . Then for almost every  $\omega$  in  $\Omega$  the following holds:  $|t| < 1$  is the strip of convergence of  $\sum_{n=0}^{\infty} X_n(\omega)a_nv_n(x, t)$  which for every  $\mathcal{L} > 0$  can be extended as a  $C^2$ -solution of the heat equation into  $\{|t| < 1\} \cup \{|x| < \mathcal{L}\}$ .*

*Proof.* We will first show for almost every  $\omega$  in  $\Omega$  that  $|t| < 1$  is the strip of convergence of  $\sum_{n=0}^{\infty} X_n(\omega)a_nv_n(x, t)$ . By (2.2) we must show that almost surely  $\limsup |X_n(\omega)a_n|^{2/n}(2n/e) = 1$ . The argument given in the first part of the proof of Theorem 1 shows that almost surely the last limit superior does not exceed 1. Let  $\{n_j\}$  be a strictly increasing sequence of positive integers such that

$$\lim_{j \rightarrow \infty} |a_{n_j}|^{2/n_j}(2n_j/e) = 1.$$

Then  $\limsup |X_n(\omega)a_n|^{2/n}(2n/e) \geq \limsup_{j \rightarrow \infty} |X_{n_j}(\omega)a_{n_j}|^{2/n_j}(2n_j/e) \geq \limsup_{j \rightarrow \infty} |X_{n_j}(\omega)|^{2/n_j}$  which by the zero-one law is equal to some number  $c$  almost surely. Suppose  $c < 1$ . Then  $X_{n_j} \rightarrow 0$  almost surely. By (ii) for  $A > 0$  and  $j = 0, 1, \dots$

$$N \leq \int_{[|X_{n_j}| \leq A]} |X_{n_j}(\omega)| dP(\omega) + A^{-1} \int_{[|X_{n_j}| > A]} |X_{n_j}(\omega)|^2 dP(\omega).$$

By the Lebesgue dominated convergence theorem the next to the last integral tends to 0 as  $j$  tends to  $\infty$ . From (i) the last term is uniformly bounded by  $A^{-1}M$ . Thus for every  $A > 0, N \leq A^{-1}M$  which is a contradiction. We conclude that  $c \geq 1$ . Thus almost surely

$$\limsup |X_n(\omega)a_n|^{2/n}(2n/e) \geq 1$$

which concludes the proof that almost surely this limit superior is 1.

In order to establish Theorem 2 for the line  $t = 1$  we first construct a function which is  $C^\infty$  on the closed strip  $|t| \leq 1$  and has a heat polynomial expansion in  $|t| < 1$ . Let  $r$  be a number which is strictly greater than  $1/2$  and strictly less than the classification num-



ber of  $\{a_n\}$ . For  $n = 0, 1, \dots$  define  $\alpha_n = (2n)e^{-nr}$ . Define  $f$  on  $[-1, 1]$  by  $f(t) = \sum_{k=0}^{\infty} \alpha_k t^k$ . We will show this definition makes sense and obtain some bounds on the derivatives of  $f$ .

Let  $n$  be a nonnegative integer. Differentiating  $\sum_{k=0}^{\infty} \alpha_k t^k$  term by term  $n$  times yields  $\sum_{k=n}^{\infty} k!/(k-n)! \alpha_k t^{k-n}$ . For  $|t| \leq 1$  the  $k^{\text{th}}$  term of this series is dominated by  $2 k^{n+1} e^{-kr}$ . One checks that

$$g_n(x) = x^{n+1} e^{-x^r}$$

is increasing on  $(0, (n+1/r)^{1/r})$  and decreasing on  $((n+1/r)^{1/r}, \infty)$ . Hence

$$\sum_{k=n}^{\infty} k^{n+1} e^{-kr} \leq \int_n^{\infty} g_n(x) dx + g_n\left(\left(\frac{n+1}{r}\right)^{1/r}\right) \leq 3\Gamma((n+2)/r)/r.$$

We conclude that  $f$  is a  $C^\infty$ -function with  $|f^{(n)}(t)| \leq 6\Gamma((n+2)/r)/r$  for  $n = 0, 1, \dots$  and  $|t| \leq 1$ .

Now define

$$(4.1) \quad u(x, t) = \sum_{n=0}^{\infty} \frac{f^{(n)}(t)x^{2n}}{(2n)!} + \sum_{n=0}^{\infty} \frac{f^{(n+1)}(t)x^{2n+1}}{(2n+1)!}.$$

Because of the bounds obtained in the preceding paragraph it can be shown that the series of (4.1) can be differentiated term by term and that  $u(x, t)$  is a  $C^\infty$ -solution to the heat equation in the closed strip  $|t| \leq 1$ . Since both  $u(0, t)$  and  $\partial u/\partial x(0, t)$ , as functions of  $t$  on  $(-1, 1)$ , are given by their Maclaurin expansions,  $u$  has a heat polynomial expansion in  $|t| < 1$  (see [5]). Thus

$$(4.2) \quad \begin{aligned} u(x, t) &= \sum_{n=0}^{\infty} b_n v_n(x, t), \\ b_{2n} &= f^{(n)}(0)/(2n)!, \\ b_{2n+1} &= f^{(n+1)}(0)/(2n+1)!. \end{aligned}$$

According to the first paragraph of the proof of Theorem 1,  $\bigcap_{k=1}^{\infty} \bigcap_{n=k}^{\infty} [|X_n| \leq nM^{1/2}]$  has probability 1. Let  $\omega$  be in this almost sure set. Let  $k_0$  be a positive integer such that for  $n \geq k_0$ ,  $|X_n(\omega)| \leq nM^{1/2}$ . Since  $r$  is less than the classification number of  $\{a_n\}$ , there is a number  $K$  such that  $|a_n|(2n/e)^{n/2} \leq Ke^{-nr}$ ,  $n = 1, 2, \dots$ . Using Stirling's theorem we have for  $2n \geq k_0$

$$b_{2n}(4n/e)^n \geq |X_{2n}(\omega)a_{2n}| (4n/e)^n (1/2)^{3/2}/KM^{1/2}.$$

Similarly for  $2n+1 \geq k_0$

$$b_{2n+1}(2(2n+1)/e)^{(2n+1)/2} \geq |X_{2n+1}(\omega)a_{2n+1}| (2(n+1)/e)^{(2n+1)/2} e^{-1/2}/KM^{1/2}.$$

Letting  $K' = K(Me)^{1/2}$  we have

$$|X_n(\omega)a_n| \leq K'b_n \text{ for } n \geq k_0.$$

Let  $\mathcal{L} > 0$ . Then for  $0 < t < 1$  we have

$$\begin{aligned} \left| \frac{\partial}{\partial t} \sum_{n=k_0}^{\infty} X_n(\omega)a_n v_n(\pm \mathcal{L}, t) \right| &= K' \sum_{n=k_0}^{\infty} b_n n(n-1) |v_{n-2}(\pm \mathcal{L}, t)| \\ &\leq K' \sum_{n=k_0}^{\infty} b_n n(n-1) v_{n-2}(\mathcal{L}, t) \leq K' \frac{\partial u}{\partial t}(\mathcal{L}, 1) < \infty. \end{aligned}$$

Thus  $\lim_{t \uparrow 1} \sum_{n=k_0}^{\infty} X_n(\omega)a_n v_n(\pm \mathcal{L}, t)$  both exist as is easily seen from the mean value theorem and the Cauchy criterion. Hence we can obtain an extension of  $\sum_{n=0}^{\infty} X_n(\omega)a_n v_n(x, t)$  into

$$\{(x, t): |t| < 1\} \cup \{(x, t): |x| < \mathcal{L}, 0 < t\}$$

which is a  $C^2$ -solution of the heat equation. (Notice at this point that we can also obtain an extension which is a bounded  $C^2$ -solution in  $\{(x, t): |x| < \mathcal{L}, 0 \leq t\}$ .) Since  $\omega$  was from the almost sure set

$$\bigcap_{k=1}^{\infty} \bigcap_{n=k}^{\infty} [|X_n| \leq nM^{1/2}],$$

this establishes the result for the line  $t = 1$ .

We now turn to the line  $t = -1$ . Define  $\{Y_n\}_{n=0}^{\infty}$  on  $\Omega$  by  $Y_{2n} = (-1)^n X_{2n}$  and  $Y_{2n+1} = (-1)^n X_{2n+1}$ . Then, applying the first part of the proof, there is a set  $F$  in  $\mathcal{F}$  with  $P(F) = 1$  such that for  $\omega$  in  $F$  and  $\mathcal{L} > 0$  the solution  $v_{\omega}(x, t) = \sum_{n=0}^{\infty} Y_n(\omega)a_n v_n(x, t)$  can be extended into  $\{|t| < 1\} \cup \{|x| < \mathcal{L} \text{ and } 0 < t\}$  so as to be a bounded  $C^2$ -solution of the heat equation in  $\{(x, t): |x| < \mathcal{L} \text{ and } 0 < t\}$ . One easily checks that for  $\omega$  in  $F$ ,

$$\sum_{n=0}^{\infty} X_n(\omega)a_n v_n(0, t) = \sum_{n=0}^{\infty} Y_n(\omega)a_n v_n(0, -t)$$

and  $\sum_{n=1}^{\infty} X_n(\omega)a_n n v_{n-1}(0, t) = \sum_{n=1}^{\infty} Y_n(\omega)a_n n v_{n-1}(0, -t)$ . Using these facts and (2.3) we see that for  $\omega$  in  $F$  and  $\mathcal{L} > 0$  the functions  $\phi(t) = \sum_{n=0}^{\infty} X_n(\omega)a_n v_n(0, t)$  and  $\psi(t) = \sum_{n=1}^{\infty} X_n(\omega)a_n n v_{n-1}(0, t)$  on  $(-1, 1)$  possess sufficiently well behaved extensions  $\phi'$  and  $\psi'$  to  $(-\infty, 1)$  that

$$\sum_{n=0}^{\infty} \frac{\phi'^{(n)}(t)x^{2n}}{(2n)!} + \sum_{n=0}^{\infty} \frac{\psi'^{(n)}(t)x^{2n+1}}{(2n+1)!}$$

is an extension of  $\sum_{n=0}^{\infty} X_n(\omega)a_n v_n(x, t)$  in  $|t| < 1$  to

$$\{(x, t): |t| < 1\} \cup \{(x, t): |x| < \mathcal{L} \text{ and } -\infty < t < 1\}.$$

**5. Examples.** The first example will show that our two theorems are best possible with respect to the allowable values of the classification number.

EXAMPLE 1. We will take  $[0, 1]$  with Lebesgue measure as the probability space and the sequence of Rademacher functions,  $\{\phi_n\}_{n=0}^\infty$ , for the random variables.

For  $k = 0, 1, \dots$  define  $\alpha_k = e^{-\sqrt{k}}$ . Then, as in the proof of Theorem 2, defining  $f$  on  $[-1, 1]$  by  $f(t) = \sum_{k=0}^\infty \alpha_k t^k$  yields a  $C^\infty$ -function whose  $n^{\text{th}}$  derivative on  $[-1, 1]$  is bounded in absolute value by  $6\Gamma(2(2n+1))$ . In the strip  $|t| < 1$  define  $u(x, t) = \sum_{n=0}^\infty (f^{(n)}(t)x^{2n})/(2n)!$ . To see that this definition makes sense and that term by term partial differentiation is permitted, we note that for every closed interval  $I \subseteq (-1, 1)$ ,  $f$  is in class  $C\{n!\}$  on  $I$ . Because of the bounds on the derivatives of  $f$  we see from the defining series for  $u$  that  $u$  may be extended as a  $C^\infty$ -solution of the heat equation to

$$\{|t| < 1\} \cup \{(x, 1): |x| < 1\}.$$

Since  $u(0, t)$  and  $\partial u/\partial x(0, t)$  are both given by their Maclaurin expansions in  $|t| < 1$ ,  $u$  possesses a heat polynomial expansion in the strip  $|t| < 1$  (see [5]). Thus for  $|t| < 1$ ,  $u(x, t) = \sum_{n=0}^\infty a_n v_n(x, t)$ ;  $a_{2n} = (e^{-\sqrt{n}}n!)/(2n)!$ ,  $a_{2n+1} = 0$ . One checks that  $\limsup |a_n|^{2/n}(2n/e) = 1$ . Also it is easily seen that  $\lim |a_{2n}|(4n/e)^n e^{\sqrt{2n}} = \infty$  which implies  $\{a_n\} \notin E_{1/2}$  and thus the classification number of  $\{a_n\}$  is in  $[0, 1/2]$ . As in the proof of Theorem 2,  $\lim_{t \uparrow 1} u_\omega(\pm 1/2, t)$  both exist for every  $\omega$  in  $[0, 1]$ . Thus for every  $\omega \in [0, 1]$  the line  $t = 1$  is not part of the natural boundary for  $u_\omega(x, t)$ . Using Theorem 1, we conclude that the classification number of  $\{a_n\}$  is  $1/2$  and that in Theorem 1 we cannot replace  $[0, 1/2)$  by  $[0, 1/2]$  as the allowable range for the classification number.

We will next show that the conclusion of Theorem 2 does not hold for  $\sum_{n=0}^\infty \phi_n(\omega)a_nv_n(x, t)$ . Assume there is a set  $A$  in  $[0, 1]$  with  $m(A) = 1$  such that for each  $\omega$  in  $A$  no interval of the line  $t = 1$  is part of the natural boundary for  $u_\omega(x, t)$ . Thus for  $\omega$  in  $A$ ,  $g_\omega(x) = \lim_{t \uparrow 1} u_\omega(x, t)$  is well defined and is the restriction of an entire function to the real axis (this last assertion can be seen by employing (2.3)). Thus for  $\omega$  in  $A$ ,  $\limsup (|g_\omega^{(n)}(0)|/n!)^{1/n} = 0$ . For  $\omega$  in  $A$ ,  $|g_\omega^{(2n+1)}(0)| = 0$  and  $|g_\omega^{(2n)}(0)| = |\sum_{k=2n}^\infty \phi_k(\omega)\alpha_k(k!/(k-2n)!)v_{k-2n}(0, 1)| = |\sum_{k=n}^\infty \phi_{2k}(\omega)(k!/(k-n)!)e^{-\sqrt{k}}|$ . Thus for  $\omega$  in  $A$ ,

$$\limsup \left[ \frac{\left| \sum_{k=n}^\infty \phi_{2k}(\omega) \frac{k!}{(k-n)!} e^{-\sqrt{k}} \right|}{(2n)!} \right]^{1/n} = 0.$$

Let  $\delta > 0$ . For  $m = 0, 1, \dots$  let

$$\begin{aligned} F_m &= \left\{ \omega \in A : \left( \left| \sum_{k=n}^\infty \phi_{2k}(\omega) \frac{k!}{(k-n)!} e^{-\sqrt{k}} \right| / (2n)! \right)^{1/n} \right. \\ &\quad \left. \leq \delta \text{ for } n = m, m+1, \dots \right\} \end{aligned}$$

and note  $F_m \uparrow A$ . Let  $A$  and  $B$  be two numbers associated with the sequence  $\{\phi_{2n}\}_{n=0}^\infty$  as in (3.1). Let  $m_0$  be sufficiently large that  $m(F_{m_0}) > A$ . Let  $n_0$  be an integer larger than  $m_0$  with  $n_0$  corresponding to  $F_{m_0}$  as in (3.1). Thus for  $n \geq n_0$  and  $k \geq 1$

$$(5.1) \quad \sum_{j=n}^{n+k} \left[ \frac{j!}{(j-n)!} e^{-\sqrt{j}} \right]^2 \leq B \int_{F_{m_0}} \left( \sum_{j=n}^{n+k} \phi_{2j}(\omega) \frac{j!}{(j-n)!} e^{-\sqrt{j}} \right)^2 dm(\omega).$$

As in the proof of Theorem 1, letting  $k$  tend to  $\infty$  yields (5.1) with  $n+k$  replaced by  $\infty$ . Using the definition of  $F_{m_0}$ , we have

$$\sum_{j=n}^{\infty} \left[ \frac{j!}{(j-n)!} e^{-\sqrt{j}} \right]^2 \leq B((2n)! \delta^n)^2,$$

for  $n \geq n_0$ . From this we conclude that

$$\limsup \left[ \frac{\left[ \sum_{k=n}^{\infty} \left( \frac{k!}{(k-n)!} e^{-\sqrt{k}} \right)^2 \right]^{1/2}}{(2n)!} \right]^{1/n} = 0.$$

On the other hand, letting  $L$  denote this last limit superior, we have

$$L \geq \limsup \left[ \frac{\left[ \sum_{k=n}^{\infty} (k-n)^{2n} \exp(-2\sqrt{k-n}) \exp(-(2\sqrt{k}-2\sqrt{k-n})) \right]^{1/2}}{(2n)!} \right]^{1/n}.$$

But  $\exp(-(2\sqrt{k}-2\sqrt{k-n})) \geq e^{-2\sqrt{n}}$  for  $k \geq n$  and  $\lim (e^{-\sqrt{n}})^{1/n} = 1$ . Hence  $L \geq \limsup ((\sum_{k=0}^{\infty} k^{2n} e^{-2\sqrt{k}})^{1/2} / (2n)!)^{1/n}$ . Define  $h_n$  on  $(0, \infty)$  by  $h_n(x) = x^{2n} e^{-2\sqrt{x}}$ . One checks that  $h_n$  is increasing on  $(0, (2n)^2)$  and decreasing on  $((2n)^2, \infty)$ . Thus  $\sum_{k=0}^{\infty} k^{2n} e^{-2\sqrt{k}} \geq \int_0^{\infty} h_n(x) dx - h_n((2n)^2) = (\Gamma(4n+2) - 2(4n)^{4n} e^{-4n}) / (2 \cdot 4^{2n})$ . Thus

$$L \geq \frac{1}{4} \limsup \left[ \left( \frac{\Gamma(4n+2)}{(4n)!} - \frac{2(4n)^{4n} e^{-4n}}{(4n)!} \right) ((4n)! / ((2n)!)^2) \right]^{1/2n} > 0.$$

This is a contradiction. Hence in Theorem 2 we cannot replace  $(1/2, 1]$  by  $[1/2, 1]$  as the allowable range for the classification number.

The next example shows that in Theorem 1 we cannot omit the symmetry of the random variables.

EXAMPLE 2. Let  $k(x, t) = e^{-x^2/4t} / (4\pi t)^{1/2}$  for  $t > 0$  and define

$$u(x, t) = k(x, t+1)$$

in the strip  $|t| < 1$ . Then [2, Th. 4.2, p. 227]

$$u(x, t) = (4\pi)^{-1/2} \sum_{n=0}^{\infty} \frac{(-1)^n}{n! 4^n} v_{2n}(x, t).$$

Let  $\{a_n\}_{n=0}^{\infty}$  be defined by  $a_{2n} = (-1)^n/n! 4^n$  and  $a_{2n+1} = 0$ . One easily checks that  $\limsup |a_n|^{2/n} (2n/e) = 1$  and that the classification number of  $\{a_n\} = 0$ . Let  $X_n = 1, n = 0, 1, \dots$  on some complete probability space. Then for every  $\omega, u_\omega$  can be continued above the line  $t = 1$ .

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