Pacific Journal of Mathematics

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Vol. 96, No. 1 November 1981

PACIFIC JOURNAL OF MATHEMATICS
Vol. 96, No. 1, 1981

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The sufficiency in von Neumann algebras is discussed with some applications to classification of normal states. It is shown that the concept of sufficiency characterizes the KMS-states and the invariant states with respect to a modular automorphism group. The relations between the sufficiency and the relative entropy are established.

Introduction. Since the investigation of sufficient statistics in abstract measure theoretic terms was initiated by Halmos and Savage [10], the concept of sufficiency has been developed by many mathematical statisticians in terms of various relations given by comparison of experiments, risk functions within the framework of statistical decision problems and so on. A characterization of sufficiency was given in [12] through the measure of Kullback-Leibler information.

The concept of sufficiency was first generalized by Umegaki [22, 23] to the noncommutative case of semi-finite von Neumann algebras with some extension of the Kullback-Leibler information (usually called the relative entropy). Later the related discussions especially concerning the relative entropy for quantum systems have been made by several authors, e.g., Araki [2, 3], Gudder and Marchand [7], and Lindblad [13].

As defined precisely and explained in §§ 1 and 4 of this paper, the concept of sufficiency is more or less considered through the informativity of a certain subalgebra with respect to a given algebra for a dynamical system of interest. Namely, in the case that such a subalgebra is sufficient, the relative entropy on the subalgebra is equal to that on the given algebra. This fact may or may not be a reason why the concept of sufficiency has not been entered into analysis of physical systems, in which the change of entropy is thought of more relevant.

The Kubo-Martin-Schwinger (KMS) condition was introduced by these three authors [11, 14] as a boundary condition of the thermal Green function. Haag, Hugenholtz and Winnink [8] showed that in the operator algebraic framework this condition is a fundamental one describing thermal equilibrium of quantum systems. The KMS condition through the Tomita-Takesaki theory now becomes a core of studying von Neumann algebras.

Under the above historical basis, our main motivation of this

work is as follows: How useful for quantum systems is the concept of sufficiency? How much of the related topics of sufficiency, mostly done for the commutative case, can be generalized to the noncommutative case?

Having these questions in our mind, we discuss the sufficiency with some applications to classification of normal states on the basis of recent development of von Neumann algebras.

In § 1 of this paper, we establish definitions and notations used throughout and also give some simple facts.

In § 2, it is shown that the concept of sufficiency characterizes the invariant states and the KMS-states with respect to a modular automorphism group.

In § 3, we prove some formulas on the relative entropy using Araki's definition of relative entropy.

In § 4, combining several theorems obtained in the previous sections, we establish some results which indicate the relations between the sufficiency and the relative entropy.

As a whole, we like to claim that the concept of sufficiency might be very useful for analysing von Neumann algebras and hence some quantum systems.

1. Definition and preliminaries. Throughout this paper, let $\mathfrak{N}$ be a von Neumann algebra with unity I acting on a Hilbert space $\mathscr{H}$, and $\mathfrak{S}$ be the set of all normal states of $\mathfrak{N}$. A dynamical system of physically interest is described by a triple $(\mathfrak{N}, \mathfrak{S}, \alpha)$, where α_t, $t \in \boldsymbol{R}$, is a strongly continuous one-parameter automorphism group of $\mathfrak{N}$. A state $\varphi \in \mathfrak{S}$ is said to satisfy the *Kubo-Martin-Schwinger* (KMS) *boundary condition* at a certain constant $\beta > 0$ with respect to α_t if for every pair $A, B \in \mathfrak{N}$ there exists a bounded function $F_{A,B}(z)$ continuous on and holomorphic in the strip $0 \leqq \mathrm{Im}\, z \leqq \beta$ with boundary values:

$$F_{A,B}(t) = \varphi(\alpha_t(A)B) \quad \text{and} \quad F_{A,B}(t + i\beta) = \varphi(B\alpha_t(A)) .$$

If φ satisfies the KMS condition with respect to α_t, then φ is proved to be α_t-invariant, i.e., $\varphi \circ \alpha_t = \varphi$. Considering $\alpha_{\beta t}$, we may take $\beta = 1$ in the sequel discussions. Takesaki showed in [17] using Tomita's theory that to every faithful state $\varphi \in \mathfrak{S}$ there exists a unique one-parameter automorphism group (i.e., the so-called *modular automorphism group*) σ_t^{φ} with respect to which φ satisfies the KMS condition at $\beta = 1$.

In this paper, a subalgebra $\mathfrak{M}$ always means a von Neumann subalgebra of $\mathfrak{N}$ with I. For a subalgebra $\mathfrak{M}$ and a state $\varphi \in \mathfrak{S}$, let $E_{\varphi}(\cdot|\mathfrak{M})$ denote the conditional expectation with respect to $\mathfrak{M}$ and φ (if it exists), which is characterized as a norm one normal projec-

tion from $\mathfrak{N}$ onto $\mathfrak{M}$ satisfying $\varphi(A) = \varphi(E_\varphi(A|\mathfrak{M}))$ for all $A \in \mathfrak{N}$ (cf. [19, 21]). It was shown by Takesaki [18] that for a faithful state $\varphi \in \mathfrak{S}$ the conditional expectation $E_\varphi(\cdot|\mathfrak{M})$ exists if and only if $\mathfrak{M}$ is invariant under the modular automorphism group σ_t^φ.

According to [5], for any two faithful states $\varphi, \psi \in \mathfrak{S}$ there exists a strongly continuous function $t \mapsto u_t$ of $\boldsymbol{R}$ into the unitary group of $\mathfrak{N}$ which is a φ-cocycle, i.e.,

$$u_{s+t} = u_s \sigma_s^\varphi(u_t) , \quad s, t \in \boldsymbol{R} ,$$

and which satisfies

$$\sigma_t^\psi(A) = u_t \sigma_t^\varphi(A) u_t^* , \quad t \in \boldsymbol{R} , \quad A \in \mathfrak{N} .$$

This φ-cocycle u_t is denoted by $u_t = (D\psi : D\varphi)_t$ and is called the *Connes Radon-Nikodym derivative* of ψ with respect to φ. Some discussions are found in [4, 9] concerning Connes Radon-Nikodym derivatives and conditional expectations.

Let S be a subset of $\mathfrak{S}$. A subalgebra $\mathfrak{M}$ is said to be *sufficient* for S if $E_\varphi(\cdot|\mathfrak{M})$ exists for each $\varphi \in S$ and for every $A \in \mathfrak{N}$ there exists an $A_0 \in \mathfrak{M}$ such that

$$A_0 = E_\varphi(A|\mathfrak{M}) \quad \text{a.e. } [\varphi] , \quad \varphi \in S ,$$

where $A = B$ a.e. $[\varphi]$ means $\varphi(|A - B|) = 0$. This definition of sufficient subalgebras is somewhat weaker than that in [22]. Also we call $\mathfrak{M}$ to be *minimal sufficient* for S if $\mathfrak{M}$ is sufficient for S and any subalgebra being sufficient for S includes $\mathfrak{M}$.

For $\varphi, \psi \in \mathfrak{S}$, it is said that ψ is absolutely continuous with respect to φ (we write $\psi \ll \varphi$) if for each $A \in \mathfrak{N}$, $\varphi(A^*A) = 0$ implies $\psi(A^*A) = 0$; that is, $\psi \ll \varphi$ if and only if $s(\psi) \leqq s(\varphi)$ where $s(\varphi)$ is the support projection of φ. We give here the elementary facts of sufficiency which are readily seen from the definition.

(1°) Let $\varphi, \psi \in \mathfrak{S}$ with $\psi \ll \varphi$. Then a subalgebra $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$ if and only if $E_\varphi(\cdot|\mathfrak{M})$ exists and $\psi(A) = \psi(E_\varphi(A|\mathfrak{M}))$ for all $A \in \mathfrak{N}$.

(2°) If a subalgebra $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$, then $\varphi = \psi$ on $\mathfrak{N}$ if and only if $\varphi = \psi$ on $\mathfrak{M}$.

When $S(\subset \mathfrak{S})$ contains a faithful state φ, then:

(3°) A subalgebra $\mathfrak{M}$ is sufficient for S if and only if $\mathfrak{M}$ is sufficient for every pair $\{\varphi, \psi\}$ with $\psi \in S$.

(4°) If $\mathfrak{M}$ is sufficient for S, then any subalgebra $\mathfrak{M}_1$ including $\mathfrak{M}$ is sufficient for S whenever $E_\varphi(\cdot|\mathfrak{M}_1)$ exists.

2. Sufficiency and characterization of states. The following lemma is a restatement of [4, Lemma 1.6] in our terminology. We

give the proof for completeness.

LEMMA 2.1. *For each subalgebra* $\mathfrak{M}$ *and two faithful states* $\varphi, \psi \in \mathfrak{S}$, *the following conditions are equivalent*:

(i) $\mathfrak{M}$ *is sufficient for* $\{\varphi, \psi\}$;

(ii) $E_\varphi(\cdot|\mathfrak{M})$ *exists and* $(D\psi: D\varphi)_t \in \mathfrak{M}$ *for every* $t \in \boldsymbol{R}$.

Proof. Let $\hat{\varphi} = \varphi \restriction \mathfrak{M}$ and $\hat{\psi} = \psi \restriction \mathfrak{M}$. Assume that $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$. Then the conditional expectation $E_\varphi(\cdot|\mathfrak{M})$ exists and $\psi(A) = \psi(E_\varphi(A|\mathfrak{M}))$ for $A \in \mathfrak{N}$. By [5, Lemma 1.4.4], we have

$$(D\psi: D\varphi)_t = (D(\hat{\psi}\circ E): D(\hat{\varphi}\circ E))_t = (D\hat{\psi}: D\hat{\varphi})_t \in \mathfrak{M}$$

for every $t \in \boldsymbol{R}$, where $E(\cdot) = E_\varphi(\cdot|\mathfrak{M})$. Conversely assume that $E_\varphi(\cdot|\mathfrak{M})$ exists and $(D\psi: D\varphi)_t \in \mathfrak{M}$ for all $t \in \boldsymbol{R}$. Since $\sigma_t^{\hat{\varphi}} = \sigma_t^\varphi \restriction \mathfrak{M}$, it follows that $u_t = (D\psi: D\varphi)_t$ is a $\hat{\varphi}$-cocycle. By [5, Theorem 1.2.4], there exists a unique faithful normal semi-finite weight $\tilde{\psi}$ on $\mathfrak{M}$ such that $(D\tilde{\psi}: D\hat{\varphi})_t = u_t$. Define a faithful normal semi-finite weight ψ' on $\mathfrak{N}$ by $\psi'(A) = \tilde{\psi}(E_\varphi(A|\mathfrak{M}))$ for $A \in \mathfrak{N}$. Then it follows that

$$(D\psi': D\varphi)_t = (D\tilde{\psi}: D\hat{\varphi})_t = (D\psi: D\varphi)_t\,, \quad t \in \boldsymbol{R}\,.$$

Hence we have $\psi' = \psi$, so that $\psi(A) = \psi(E_\varphi(A|\mathfrak{M}))$ for every $A \in \mathfrak{N}$. This shows that $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$. □

In this section, let φ be a fixed faithful normal state of $\mathfrak{N}$ and σ_t^φ its modular automorphism group. Let Z_φ be the subalgebra consisting of all $A \in \mathfrak{N}$ such that $\varphi(AB) = \varphi(BA)$ for every $B \in \mathfrak{N}$. The subalgebra Z_φ is called the *centralizer* of φ and is exactly the fixed point algebra of σ_t^φ (cf. [17, Lemma 15.8]), i.e.,

$$Z_\varphi = \{A \in \mathfrak{N}: \sigma_t^\varphi(A) = A,\ t \in \boldsymbol{R}\}\,.$$

Let $\mathfrak{Z}$ be the center of $\mathfrak{N}$, i.e., $\mathfrak{Z} = \mathfrak{N} \cap \mathfrak{N}'$. Clearly $\mathfrak{Z} \subset Z_\varphi$. Let $I(\varphi)$ be the set of all σ_t^φ-invariant states in $\mathfrak{S}$, and $K(\varphi)$ be the set of all states in $\mathfrak{S}$ satisfying the KMS condition with respect to σ_t^φ at $\beta = 1$. Then we have:

THEOREM 2.2. (1) *For each* $\psi \in \mathfrak{S}$, $\psi \in I(\varphi)$ *if and only if* Z_φ *is sufficient for* $\{\varphi, \psi\}$.

(2) *The centralizer* Z_φ *is minimal sufficient for* $I(\varphi)$.

Proof. (1) Let $\psi \in \mathfrak{S}$ and take $\psi_1 = (\psi + \varphi)/2$. Then we easily see that $\psi \in I(\varphi)$ is equivalent to $\psi_1 \in I(\varphi)$, and the sufficiency of Z_φ for $\{\varphi, \psi\}$ is equivalent to that for $\{\varphi, \psi_1\}$. Therefore we can assume that ψ is faithful. Since Z_φ is elementwise invariant under σ_t^φ,

there exists the conditional expectation $E_\varphi(\cdot|Z_\varphi)$ from $\mathfrak{N}$ onto Z_φ. Hence, in view of Lemma 2.1, it suffices to show that $\psi \in I(\varphi)$ if and only if $(D\psi: D\varphi)_t \in Z_\varphi$ for every $t \in \boldsymbol{R}$. If $\psi \in I(\varphi)$, then by [5, Lemma 1.2.3] there exists a positive self-adjoint operator h affiliated with Z_φ such that $(D\psi: D\varphi)_t = h^{it} \in Z_\varphi$ for all $t \in \boldsymbol{R}$. Conversely suppose that $(D\psi: D\varphi)_t \in Z_\varphi$ for every $t \in \boldsymbol{R}$. Since

$$\sigma_t^\psi(A) = (D\psi: D\varphi)_t \sigma_t^\varphi(A)(D\psi: D\varphi)_t^* , \tag{2.1}$$

we have

$$\varphi(\sigma_t^\psi(A)) = \varphi(\sigma_t^\varphi(A)) = \varphi(A) , \quad A \in \mathfrak{N} .$$

Hence it follows that φ is σ_t^ψ-invariant, and thus ψ is σ_t^φ-invariant (cf. [17, Theorem 15.2]).

(2) It follows from (1) that Z_φ is sufficient for every pair $\{\varphi, \psi\}$ with $\psi \in I(\varphi)$. Hence Z_φ is sufficient for $I(\varphi)$. To show the minimality of Z_φ, let $\mathfrak{M}$ be any subalgebra which is sufficient for $I(\varphi)$. We now prove that $Z_\varphi \subset \mathfrak{M}$. Take any positive invertible operator $h \in Z_\varphi$ with $\varphi(h) = 1$, and define a faithful state $\psi \in \mathfrak{S}$ by $\psi(A) = \varphi(hA)$ for $A \in \mathfrak{N}$. Then we have $\psi \in I(\varphi)$ and $(D\psi: D\varphi)_t = h^{it}$. Since $(D\psi: D\varphi)_t \in \mathfrak{M}$ for every $t \in \boldsymbol{R}$ by Lemma 2.1, it follows that $h \in \mathfrak{M}$. Thus $Z_\varphi \subset \mathfrak{M}$. □

THEOREM 2.3. (1) *For each* $\psi \in \mathfrak{S}$, $\psi \in K(\varphi)$ *if and only if* $\mathfrak{Z}$ *is sufficient for* $\{\varphi, \psi\}$.

(2) *The center* $\mathfrak{Z}$ *is minimal sufficient for* $K(\varphi)$.

Proof. As in the proof of Theorem 2.2, we can assume that ψ is faithful. If $\psi \in K(\varphi)$, then by [15, Theorem 5.4] there exists a positive self-adjoint operator h affiliated with $\mathfrak{Z}$ such that $\psi(A) = \varphi(hA)$ for $A \in \mathfrak{N}$, so that $(D\psi: D\varphi)_t = h^{it} \in \mathfrak{Z}$ for every $t \in \boldsymbol{R}$. Conversely if $(D\psi: D\varphi)_t \in \mathfrak{Z}$ for every $t \in \boldsymbol{R}$, then by (2.1) we have $\sigma_t^\psi = \sigma_t^\varphi$ and hence $\psi \in K(\varphi)$. Thus (1) is proved. The proof of (2) is analogous to that of Theorem 2.2. □

3. Relative entropy. When $\mathfrak{N}$ is finite dimensional, for each φ and ψ in $\mathfrak{S}$ the relative entropy $S(\varphi|\psi)$ is defined by

$$S(\varphi|\psi) = \mathrm{tr}(\rho_\psi \log \rho_\psi - \rho_\psi \log \rho_\varphi) ,$$

where ρ_φ and ρ_ψ are density matrices for φ and ψ. Araki [2, 3] extended the relative entropy to the case for normal positive linear functionals of general von Neumann algebras, and studied its several properties such as joint convexity, lower semiconitinuity and monotonicity.

In this section, we assume as in [3] that $\mathfrak{N}$ has a cyclic and separating vector. Let V be a natural positive cone (cf. [1]) for $\mathfrak{N}$ and let φ and ψ be states in $\mathfrak{S}$. By [1, Theorem 6], there exist unique vector representatives Φ and Ψ of φ and ψ in V such that $\varphi(A) = \langle \Phi, A\Phi \rangle$ and $\psi(A) = \langle \Psi, A\Psi \rangle$ for all $A \in \mathfrak{N}$. The operator $S_{\Phi,\Psi}$ with the domain

$$D(S_{\Phi,\Psi}) = \mathfrak{N}\Psi + (I - s^{\mathfrak{N}'}(\Psi))\mathscr{H}$$

is defined by

$$S_{\Phi,\Psi}(A\Psi + \Omega) = s^{\mathfrak{N}}(\Psi)A^*\Phi\,, \quad A \in \mathfrak{N}\,, \quad s^{\mathfrak{N}'}(\Psi)\Omega = 0\,,$$

where $s^{\mathfrak{N}}(\Psi)(\in \mathfrak{N})$ denotes the $\mathfrak{N}$-support of the vector Ψ. Then $S_{\Phi,\Psi}$ is a closable conjugate-linear operator (cf. [3]) and the relative modular operator $\Delta_{\Phi,\Psi}$ is defined by

$$\Delta_{\Phi,\Psi} = (S_{\Phi,\Psi})^*\overline{S_{\Phi,\Psi}}\,.$$

The relative entropy $S(\varphi|\psi)$ is now given by

$$S(\varphi|\psi) = \begin{cases} -\langle \Psi, (\log \Delta_{\Phi,\Psi})\Psi \rangle & \text{if} \quad \psi \ll \varphi\,, \\ +\infty & \text{otherwise}\,. \end{cases}$$

Then $S(\varphi|\psi) \geqq 0$, and $S(\varphi|\psi) = 0$ if and only if $\varphi = \psi$. For each subalgebra $\mathfrak{M}$, let $S_{\mathfrak{M}}(\varphi|\psi)$ denote the relative entropy of the restrictions of φ and ψ to $\mathfrak{M}$. By the monotonicity of relative entropy generally proved in [20], it holds that

$$S_{\mathfrak{M}}(\varphi|\psi) \leqq S(\varphi|\psi) \tag{3.1}$$

for every subalgebra $\mathfrak{M}$ (also see [2, 3, 13, 23]).

THEOREM 3.1. *For each* $\varphi, \psi \in \mathfrak{S}$,

$$\|\varphi - \psi\| \leqq \{2S(\varphi|\psi)\}^{1/2}\,.$$

Proof. By [16, p.31], we can take two normal positive linear functionals φ_1 and φ_2 such that $\varphi - \psi = \varphi_1 - \varphi_2$, $\|\varphi - \psi\| = \|\varphi_1\| + \|\varphi_2\|$ and $s(\varphi_1) \perp s(\varphi_2)$. Let $e = s(\varphi_1)$. Then it follows that

$$\begin{aligned} \|\varphi - \psi\| &= (\varphi - \psi)(e) - (\varphi - \psi)(I - e) \\ &= 2(\varphi(e) - \psi(e))\,. \end{aligned}$$

Let $\mathfrak{M}$ be the subalgebra generated by e and $I - e$. By using the monotonicity, we have

$$\begin{aligned} S(\varphi|\psi) &\geqq S_{\mathfrak{M}}(\varphi|\psi) \\ &= \psi(e)\log\frac{\psi(e)}{\varphi(e)} + \psi(I - e)\log\frac{\psi(I - e)}{\varphi(I - e)}\,. \end{aligned}$$

It was shown in [6] that

$$2|\alpha-\beta| \leqq \left\{2\left(\beta\log\frac{\beta}{\alpha}+(1-\beta)\log\frac{1-\beta}{1-\alpha}\right)\right\}^{1/2}$$

for $0 \leqq \alpha, \beta \leqq 1$. Taking $\alpha = \varphi(e)$ and $\beta = \psi(e)$, we deduce the desired inequality. □

THEOREM 3.2. *Let $\varphi, \psi \in \mathfrak{S}$ be faithful and $\mathfrak{M}$ be a subalgebra such that $\mathfrak{M} \subset Z_\varphi$. Define $\psi' \in \mathfrak{S}$ by $\psi'(A) = \psi(E_\varphi(A|\mathfrak{M}))$ for $A \in \mathfrak{N}$. If $S_{\mathfrak{M}}(\varphi|\psi) < +\infty$, then*

$$S(\psi'|\psi) = S(\varphi|\psi) - S_{\mathfrak{M}}(\varphi|\psi)\ .$$

Proof. First note that ψ' is well defined from $\mathfrak{M} \subset Z_\varphi$. Since $\varphi \restriction \mathfrak{M}$ is a faithful normal trace, there exists a positive self-adjoint operator h affiliated with $\mathfrak{M}$ such that $\psi(A) = \varphi(hA)$ for all $A \in \mathfrak{M}$. Take the spectral decomposition $h = \int_0^\infty \lambda de(\lambda)$ and $h_n = \int_0^n \lambda\, de(\lambda)$. Since $h_n \in \mathfrak{M}$, we have for every $A \in \mathfrak{N}$

$$\begin{aligned}\psi'(A) &= \psi(E_\varphi(A|\mathfrak{M})) = \lim_{n\to\infty}\varphi(h_n E_\varphi(A|\mathfrak{M})) \\ &= \lim_{n\to\infty}\varphi(E_\varphi(h_n A|\mathfrak{M})) = \lim_{n\to\infty}\varphi(h_n A) = \varphi(hA)\ .\end{aligned}$$

Hence it follows (cf. [5, Lemma 1.2.3]) that $(D\psi': D\varphi)_t = h^{it}$ for all $t \in \boldsymbol{R}$. By the relations

$$\begin{aligned}(D\psi': D\psi)_t &= (D\psi': D\varphi)_t(D\varphi: D\psi)_t = h^{it}(D\varphi: D\psi)_t\ , \\ (D\psi': D\psi)_t &= (\Delta_{\Psi',\Psi})^{it}\Delta_\Psi^{-it}\ , \\ (D\varphi; D\psi)_t &= (\Delta_{\Phi,\Psi})^{it}\Delta_\Psi^{-it}\ ,\end{aligned}$$

where $\Delta_\Psi = \Delta_{\Psi,\Psi}$, we deduce that $(\Delta_{\Psi',\Psi})^{it} = h^{it}(\Delta_{\Phi,\Psi})^{it}$ for all $t \in \boldsymbol{R}$. Moreover since $h^{it} \in \mathfrak{M} \subset Z_\varphi$ and

$$\sigma_t^\varphi(A) = (\Delta_{\Phi,\Psi})^{it}A(\Delta_{\Phi,\Psi})^{-it}\ , \quad A \in \mathfrak{N}\ ,$$

it follows that h^{it} and $(\Delta_{\Phi,\Psi})^{it}$ commute. Now let $\hat{\Phi}$ and $\hat{\Psi}$ be vector representatives of $\varphi \restriction \mathfrak{M}$ and $\psi \restriction \mathfrak{M}$ in a natural positive cone $\hat{V}$ for $\mathfrak{M}$. Since $\varphi \restriction \mathfrak{M}$ is a trace, it follows that $\Delta_{\hat{\psi},\hat{\phi}} = h$. By $\log\Delta_{\hat{\phi},\hat{\psi}} = -J(\log\Delta_{\hat{\psi},\hat{\phi}})J$ where J is the modular conjugation operator associated with $\hat{V}$ (cf. [3, Remark 3.4]), we have

$$\begin{aligned}S_{\mathfrak{M}}(\varphi|\psi) &= -\langle\hat{\Psi}, (\log\Delta_{\hat{\phi},\hat{\psi}})\hat{\Psi}\rangle \\ &= \langle\hat{\Psi}, (\log h)\hat{\Psi}\rangle = \langle\Psi, (\log h)\Psi\rangle\ ,\end{aligned} \tag{3.2}$$

which is finite from the assumption. Therefore we obtain

$$S(\psi'|\psi) = -\langle\Psi, (\log\Delta_{\Psi',\Psi})\Psi\rangle$$

$$= -\langle \Psi, (\log \Delta_{\phi,\Psi})\Psi \rangle - \langle \Psi, (\log h)\Psi \rangle$$
$$= S(\varphi|\psi) - S_{\mathfrak{M}}(\varphi|\psi) . \qquad \square$$

THEOREM 3.3. *Let $\varphi \in \mathfrak{S}$ be faithful and $\mathfrak{M}$ be a subalgebra such that $\mathfrak{M} \subset Z_\varphi$. Let $\psi \in \mathfrak{S}$ and define $\psi' \in \mathfrak{S}$ by $\psi'(A) = \psi(E_\varphi(A|\mathfrak{M}))$ for $A \in \mathfrak{N}$. Suppose either* (a) *ψ is faithful or* (b) *$\psi \leqq \lambda\varphi$ for some $\lambda > 0$. If $S_{\mathfrak{M}}(\varphi|\psi) < +\infty$, then*

$$\|\psi' - \psi\| \leqq \{2(S(\varphi|\psi) - S_{\mathfrak{M}}(\varphi|\psi))\}^{1/2} .$$

Proof. For the case (a), the desired inequality is immediate from Theorems 3.1 and 3.2. Now suppose that $\psi \leqq \lambda\varphi$ for some $\lambda > 0$. For each $\varepsilon > 0$, let $\psi_\varepsilon = (1+\varepsilon)^{-1}(\psi + \varepsilon\varphi) \in \mathfrak{S}$ and define $\psi'_\varepsilon \in \mathfrak{S}$ by $\psi'_\varepsilon(A) = \psi_\varepsilon(E_\varphi(A|\mathfrak{M}))$. By the convexity of relative entropy (cf. [3, Theorem 3.8]), we have

$$S_{\mathfrak{M}}(\varphi|\psi_\varepsilon) \leqq (1+\varepsilon)^{-1} S_{\mathfrak{M}}(\varphi|\psi) < +\infty .$$

Hence it follows from the case (a) that

$$\|\psi'_\varepsilon - \psi_\varepsilon\| \leqq \{2(S(\varphi|\psi_\varepsilon) - S_{\mathfrak{M}}(\varphi|\psi_\varepsilon))\}^{1/2} . \tag{3.3}$$

Since $\psi_\varepsilon \leqq \lambda\varphi$ for each $\varepsilon > 0$, by [3, Theorem 3.7] we have

$$\lim_{\varepsilon \to +0} S(\varphi|\psi_\varepsilon) = S(\varphi|\psi) ,$$
$$\lim_{\varepsilon \to +0} S_{\mathfrak{M}}(\varphi|\psi_\varepsilon) = S_{\mathfrak{M}}(\varphi|\psi) .$$

Since $\psi'_\varepsilon = (1+\varepsilon)^{-1}(\psi' + \varepsilon\varphi)$, we obtain the desired inequality by letting $\varepsilon \to +0$ in (3.3). $\square$

Before closing this section, we have to note that Professor Araki gave us very important comments to some results of our first version of this paper, which make us enable to write them in the above form.

4. Sufficiency and relative entropy. In this section, let a faithful state $\varphi \in \mathfrak{S}$ be fixed as in § 2.

THEOREM 4.1. *For each subalgebra $\mathfrak{M} \subset Z_\varphi$ and each $\psi \in \mathfrak{S}$, the following statements hold:*

(1) *Suppose the condition* (a) *or* (b) *in Theorem* 3.3. *If $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$, then $S_{\mathfrak{M}}(\varphi|\psi) = S(\varphi|\psi)$, and conversely if $S_{\mathfrak{M}}(\varphi|\psi) = S(\varphi|\psi) < +\infty$, then $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$.*

(2) *If $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$, then $S_{\mathfrak{M}}(\varphi|(\psi+\varphi)/2) = S(\varphi|(\psi+\varphi)/2)$, and conversely if $S_{\mathfrak{M}}(\varphi|(\psi+\varphi)/2) = S(\varphi|(\psi+\varphi)/2) < +\infty$, then $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$.*

Proof. (1) We use the notations in the proof of Theorem 3.2. Let $\mathfrak{M}$ be sufficient for $\{\varphi, \psi\}$, and suppose the condition (a). There exists a positive self-adjoint operator h affiliated with $\mathfrak{M}$ such that $\psi(A) = \varphi(hA)$ for all $A \in \mathfrak{M}$. Then we have

$$\psi(A) = \psi(E_\varphi(A|\mathfrak{M})) = \varphi(hA) , \quad A \in \mathfrak{N} ,$$

as in the proof of Theorem 3.2. Hence it follows that $(D\psi: D\varphi)_t = h^{it}$, and we have

$$(\Delta_{\Phi,\Psi})^{it} = (D\varphi: D\psi)_t \Delta_\Psi^{it} = h^{-it}\Delta_\Psi^{it} .$$

Since h^{it} and $(\Delta_{\Phi,\Psi})^{it}$ commute, it follows that h^{-it} and Δ_Ψ^{it} commute. We thus have

$$\begin{aligned} S(\varphi|\psi) &= -\langle \Psi, (-\log h + \log \Delta_\Psi)\Psi \rangle \\ &= \langle \Psi, (\log h)\Psi \rangle , \end{aligned} \tag{4.1}$$

by $\Delta_\Psi \Psi = \Psi$. From (3.2) and (4.1), we obtain $S_{\mathfrak{M}}(\varphi|\psi) = S(\varphi|\psi)$. The case (b) is proved from the case (a) by taking limits as in the proof of Theorem 3.3.

Assume conversely that $S_{\mathfrak{M}}(\varphi|\psi) = S(\varphi|\psi) < +\infty$. Then it follows from Theorem 3.3 that $\psi' = \psi$, which implies that $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$.

(2) is immediate from (1), since $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$ if and only if $\mathfrak{M}$ is sufficient for $\{\varphi, (\psi + \varphi)/2\}$. □

The above fact (1) extends the result [23, Theorem 5] which was proved under some strong assumptions. Combining Theorem 4.1 with Theorems 2.2 and 2.3, we have the following:

COROLLARY 4.2. (1) *Suppose the condition* (a) *or* (b) *in Theorem* 3.3. *If* $\psi \in I(\varphi)$, *then* $S_{Z_\varphi}(\varphi|\psi) = S(\varphi|\psi)$, *and conversely if* $S_{Z_\varphi}(\varphi|\psi) = S(\varphi|\psi) < +\infty$, *then* $\psi \in I(\varphi)$.

(2) *If* $\psi \in I(\varphi)$, *then* $S_{Z_\varphi}(\varphi|(\psi + \varphi)/2) = S(\varphi|(\psi + \varphi)/2)$, *and conversely if* $S_{Z_\varphi}(\varphi|(\psi + \varphi)/2) = S(\varphi|(\psi + \varphi)/2) < +\infty$, *then* $\psi \in I(\varphi)$.

COROLLARY 4.3. (1) *Suppose the condition* (a) *or* (b) *in Theorem* 3.3. *If* $\psi \in K(\varphi)$, *then* $S_3(\varphi|\psi) = S(\varphi|\psi)$, *and conversely if* $S_3(\varphi|\psi) = S(\varphi|\psi) < +\infty$, *then* $\psi \in K(\varphi)$.

(2) *If* $\psi \in K(\varphi)$, *then* $S_3(\varphi|(\psi + \varphi)/2) = S(\varphi|(\psi + \varphi)/2)$, *and conversely if* $S_3(\varphi|(\psi + \varphi)/2) = S(\varphi|(\psi + \varphi)/2) < +\infty$, *then* $\psi \in K(\varphi)$.

The monotonicity (3.1) says that the restriction of measurement to a subalgebra $\mathfrak{M}$ usually makes it more difficult to discriminate between two states. From Theorem 4.1, the physical meaning of suf-

ficiency might be explained as follows: If a subalgebra $\mathfrak{M}$ is sufficient for $\{\varphi, \psi\}$, then we obtain from the measurement of $\mathfrak{M}$ as much information as from that of $\mathfrak{N}$ to discriminate between φ and ψ. In particular, to distinguish $\psi \in I(\varphi)$ (resp. $\psi \in K(\varphi)$) from φ, the measurement of Z_φ (resp. $\mathfrak{Z}$) gives as much information as $\mathfrak{N}$.

ACKNOWLEDGMENT. The authors wish to express their gratitude to Professor H. Umegaki for his valuable advice and constant encouragement. They thank Professor H. Araki for his critical reading and essential comments for this work. Professor J. Tomiyama and the referee kindly indicated the useful paper [4] to the authors.

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Received March 13, 1979 and in revised form January 4, 1980.

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NODA CITY, CHIBA 278, JAPAN

PACIFIC JOURNAL OF MATHEMATICS

Printed in Japan by International Academic Printing Co., Ltd., Tokyo, Japan

Pacific Journal of Mathematics
Vol. 96, No. 1 November, 1981