

*Pacific
Journal of
Mathematics*

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BRADLEY CURREY

Volume 170 No. 2

October 1995

SMOOTH DECOMPOSITION OF FINITE MULTIPLICITY MONOMIAL REPRESENTATIONS FOR A CLASS OF COMPLETELY SOLVABLE HOMOGENEOUS SPACES.

BRADLEY N. CURREY

Given a finite multiplicity monomial representation τ of a completely solvable Lie group G , a smooth decomposition of τ is a concrete direct integral decomposition into irreducibles parametrized by a manifold Σ , with the property that compactly-supported elements of \mathcal{H}_τ^∞ are mapped to smooth sections on Σ by the intertwining operator. A natural way of constructing such a decomposition is by means of the distribution-theoretic Plancherel formula for τ and a cross-section Σ for coadjoint orbits. However, for irreducible representations π_l , $l \in \Sigma$, the determination of appropriate distributions $\beta_l \in \mathcal{H}_l^{-\infty}$ is problematic. For the case where τ is induced from a "Levi" component, we overcome these problems and give an explicit and natural construction for a smooth decomposition. In the process we show that in this situation the nilradical of G must be two-step.

0. Introduction.

We are interested in the decomposition of the representation τ of a solvable Type I group G induced from a unitary character of closed, connected subgroup H . In one sense, to decompose τ means to describe the spectrum of τ , the multiplicities, and the equivalence class of the Plancherel measure, in terms of the coadjoint orbit picture. But there is also a stronger sense of what it means to decompose τ : one would like to give a construction for a direct integral, a unitary intertwining map, and a distribution-theoretic Plancherel formula. The goal here is that the construction be as explicit as possible, but at the same time natural: all objects should be naturally and uniquely determined up to the choice of a certain Jordan-Holder basis for the Lie algebra. The base space for the direct integral should be a smooth manifold which naturally parametrizes an explicitly determined set of coadjoint orbit data, and compactly supported smooth vectors for τ should be mapped under the intertwining map to smooth functions on this manifold. Under these circumstances we shall say that we have a *smooth* decomposition of τ . The present paper carries out this program for a particular class of completely solvable homogeneous spaces.

Let G and H be as above but also completely solvable. Then G is exponential, hence G is of Type I and one has the canonical bijection between the unitary dual \hat{G} and the space of coadjoint orbits [3]. For each $\lambda \in \hat{G}$, let O_λ denote the corresponding coadjoint orbit. Lie algebras of designated Lie groups will be denoted by corresponding German gothic letters. Given a unitary representation π of G acting in a Hilbert space \mathcal{H} , denote by \mathcal{H}^∞ the Frechet space of smooth vectors for π , and by $\mathcal{H}^{-\infty}$ the space of continuous anti-linear functionals on \mathcal{H}^∞ ; recall that $\mathcal{H}^\infty \subset \mathcal{H} \subset \mathcal{H}^{-\infty}$. Elements of $\mathcal{H}^{-\infty}$ will be called generalized vectors for π .

Let $f \in \mathfrak{g}^*$ have the property that \mathfrak{h} is subordinate to f (that is, $[\mathfrak{h}, \mathfrak{h}] \subset \ker(f)$), let χ be the corresponding unitary character of H , and let $\tau = \tau_f$ be the representation of G induced from χ . Denote by \mathfrak{h}^\perp the space of linear functionals that vanish on \mathfrak{h} . It was proved first by Corwin, Greenleaf, and Grelaud for G nilpotent [6], and by Lipsman for G completely solvable [14, 15], that the spectrum of τ consists of those $\lambda \in \hat{G}$ for which $O_\lambda \cap (f + \mathfrak{h}^\perp) \neq \emptyset$, and that the multiplicity m_λ of λ is given by the number of H -orbits contained in $O_\lambda \cap (f + \mathfrak{h}^\perp)$. Either $m_\lambda = \infty$ for a.e. λ , or there is $M > 0$ such that for a.e. λ , $m_\lambda < M$. One has the finite (indeed bounded) multiplicity case if and only if $\dim(Gl) = 2 \dim(Hl)$ for a.e. $l \in f + \mathfrak{h}^\perp$, whence each H -orbit in O_λ is a connected component of $O_\lambda \cap (f + \mathfrak{h}^\perp)$. Of course the preceding references contain more information than is conveyed here. We also remark that the spectral decomposition of τ for the more general class of exponential G was obtained by Fujiwara [10] and is similar to the above, though some questions surrounding the finite multiplicity case are still unresolved.

When $\dim(Gl) = 2 \dim(Hl)$ holds for generic $l \in f + \mathfrak{h}^\perp$, there is a one to one correspondence between H -orbits in O_λ and “appearances” of λ in decomposition of τ (for generic λ). The spectral decomposition formula of [14] is

$$(0.1) \quad [\tau] = \int_{(f+\mathfrak{h}^\perp)/H} \lambda_\theta d[\nu](\theta)$$

Here $[\tau]$ stands for the equivalence class of τ , and θ is an H -orbit contained in $O_{\lambda_\theta} \cap (f + \mathfrak{h}^\perp)$. The equivalence class $[\nu]$ is that of pushforwards of finite measures on $f + \mathfrak{h}^\perp$ equivalent to Lebesgue measure. Note that in contrast with the usual spectral decomposition with base space \hat{G} , the multiplicities are “spread out” within the H -orbit picture. Besides being elegant, the formula (0.1) calls our attention to the possibility of a smooth decomposition of τ over an H -orbit cross-section, wherein different realizations for the multiple “appearances” of each λ are allowed. The derivation (and application) of this kind of decomposition of τ first appears in [5], where G is nilpotent,

but the approach and methods there differ greatly from those of the present work.

Here the smooth decomposition is just a consequence of an explicit Plancherel formula. The distribution-theoretic Plancherel formula (the Penney-Fujiwara Plancherel formula) which is analogous to (0.1) is

$$(0.2) \quad \langle \tau(\omega) \alpha_\tau, \alpha_\tau \rangle = \int_{(f+\mathfrak{h}^\perp)/H} \langle \pi_\theta(\omega) \beta_\theta, \beta_\theta \rangle d\nu(\theta).$$

where α_θ is the canonical cyclic generalized vector for τ , π_θ is a realization of λ_θ , and β_θ is an (appropriately H -covariant) generalized vector for π_θ . The choice of ν depends on choices of various Haar measures. In the case that G is nilpotent, (0.2) was obtained by Fujiwara (in a different form) [9], and derives from the fundamental work of Penney [19]. Groundbreaking work on extending results of [9] to other classes of homogeneous spaces has been done by Fujiwara and Yamagami [11] and Lipsman [16, 17, 18]. However, beyond the nilpotent case, the technical difficulties involved in (0.2) are considerable. One constructs the model π_θ and the generalized vector β_θ for generic θ by first choosing $l \in \theta$ and a polarization $\mathfrak{b} = \mathfrak{b}(l)$ at l (satisfying the Pukanzsky condition), then β_θ is obtained by integrating $f \in (\mathcal{H}_{\pi_\theta})^\infty$ over $H \cap B \backslash H$ with respect to a certain appropriate measure. At issue is convergence of the integral, as well as the fact that β_θ must be appropriately H -covariant. One must make “good” choices for l and $\mathfrak{b}(l)$: examples show that not all choices will produce β_θ with the required properties. What exactly are the generic θ , and whether good choices for $\mathfrak{b}(l)$ actually exist are questions which are not settled in general. In any case, one would like to have a natural procedure for determining generic θ and making these choices, in the process describing $(f + \mathfrak{h}^\perp) / H$ by a smooth orbital cross-section Σ , and the measure ν as an explicit measure on Σ . In this paper we consider a special case in which the difficulties surrounding the construction of the β_θ are nevertheless very much present. Our main task is to overcome these difficulties and obtain an explicit version of (0.2) by the procedure outlined above.

The class of homogeneous spaces $H \backslash G$ with which we are concerned is that for which G is the semi-direct product $G = NH$, where N is nilpotent and normal in G , and H is abelian and acts semi-simply on N with real eigenvalues. In the context of algebraic groups H is sometimes called a Levi component [13] (but we need not assume algebraic here). The orbital spectrum formula for the quasi-regular representation τ_0 on this class of homogeneous spaces was known before the more general results of [10, 14, 15]: in [13] the spectrum of τ_0 is computed using the Mackey machine, and it is shown that τ_0 has uniform multiplicity (either a power of 2, or $+\infty$), which is in turn the number of H -orbits in a generic orbital intersection $O_\lambda \cap \mathfrak{h}^\perp$.

Now since H is co-normal, H -orbits in $f + \mathfrak{h}^\perp$ are just translates of H -orbits in \mathfrak{h}^\perp , so that if τ_f has finite multiplicity for some f , then it does for all f . If this holds we simply say that $H \backslash G$ is finite multiplicity, and we assume that this is the case for the present paper. Our first main result - one which was at first surprising to the present author - is that if $H \backslash G$ is finite multiplicity, then N must be two-step (in fact, a particular type of two-step group which includes the Heisenberg groups). It was already well-known that if $H \backslash G$ as above is symmetric, then the quasi-regular representation τ_0 is multiplicity free and N must be abelian. In our context we deduce that each τ_f has uniform multiplicity 2^u , where $u = \dim(\text{cent}(N) \sim N) / 2$. The reduction to two-step nilpotent groups here parallels a similar reduction that occurs in the situation where N is as above, but H is compact (whence G is no longer necessarily solvable). The result there is that if (H, N) is a Gelfand pair, then N is two-step [1].

A key to our method is a precise definition of what it means to be a generic element of $f + \mathfrak{h}^\perp$ by means of "jump sets" of indices, and this is the first instance known to this author where such techniques have been used to derive canonical structural information about the group itself. For generic l in $f + \mathfrak{h}^\perp$ the so-called Vergne polarizations $\mathfrak{b}(l)$ vary rationally with l and have central intersection with \mathfrak{h} . On the other hand, we show that the set of generic H -orbits admits a natural, smooth, algebraic cross-section Σ . Choosing $l \in \Sigma$, and Vergne polarizations $\mathfrak{b}(l)$, we obtain our models π_θ . The main result of Section 2 is that an appropriate integral formula for β_θ converges absolutely for every π_θ -smooth vector; thus the natural, smoothly varying $\mathfrak{b}(l)$ are in fact "good" polarization choices. In Section 3 we derive the Plancherel formula in terms of Σ . For this class of homogeneous spaces, the choices for Haar measures are natural and the resulting Plancherel measure on Σ is seen to be rational.

Thanks go to Professor Ronald Lipsman for many helpful conversations from which this work benefitted greatly.

1. Algebraic structure of \mathfrak{g} .

Let $\mathfrak{g} = \mathfrak{n} + \mathfrak{h}$ where \mathfrak{n} is nilpotent, $\mathfrak{n} \supset [\mathfrak{g}, \mathfrak{g}]$, and where \mathfrak{h} is an abelian subalgebra of \mathfrak{g} such that $\text{ad}(\mathfrak{h})$ consists of semisimple endomorphisms with real eigenvalues. For each $A \in \mathfrak{h}$, $a \in \mathbb{R}$, let $R(A, a) = \{X \in \mathfrak{n} : [A, X] = aX\}$. For any real numbers a and b , and for $A \in \mathfrak{h}$, we have the usual inclusion $[R(A, a), R(A, b)] \subset R(A, a + b)$. We fix once and for all a basis $\{Z_1, Z_2, \dots, Z_n\}$ for \mathfrak{n} with the properties that

- (i) $\text{span}\{Z_1, Z_2, \dots, Z_i\} = \mathfrak{n}_i$ is an ideal in \mathfrak{g} , and
- (ii) for each $A \in \mathfrak{h}$, Z_i is an eigenvector for $\text{ad}(A)$, $1 \leq i \leq n$.

Let $\lambda_1, \lambda_2, \dots, \lambda_n$ be the linear functional on \mathfrak{h} such that $[A, Z_i] = \lambda_i(A) Z_i$, $A \in \mathfrak{h}$, $1 \leq i \leq n$, and for each i let Λ_i be the corresponding positive character of G : $\Lambda_i(\exp Z) = e^{\lambda_i(Z)}$. We select a subset $\lambda_{i_1}, \lambda_{i_2}, \dots, \lambda_{i_n}$ as follows: $i_1 = \min\{1 \leq i \leq n : \lambda_i \neq 0\}$, $i_2 = \min\{1 \leq i \leq n : \lambda_i \text{ is not a multiple of } \lambda_{i_1}\}$, $i_3 = \min\{1 \leq i \leq n : \lambda_i \notin \text{span}\{\lambda_{i_1}, \lambda_{i_2}\}\}$, and so on. We thus obtain a minimal spanning set $\{\lambda_{i_1}, \lambda_{i_2}, \dots, \lambda_{i_d}\}$ for the root system $\{\lambda_1, \lambda_2, \dots, \lambda_n\}$. The minimality of our selection with respect to the ordering of the root system (as well as to cardinality) is crucial. Set $\Psi = \{i_1 < i_2 < \dots < i_d\}$. For each k , $1 \leq k \leq u$, let $A_k \in \mathfrak{h}$ be chosen such that $\lambda_{i_r}(A_s) = \delta_{rs}$, $1 \leq r, s \leq d$.

Lemma 1.1. *For each k , $1 \leq k \leq d$, we have*

$$[A_k, \mathfrak{n}_{i_k-1}] = [Z_{i_k}, \mathfrak{n}_{i_k-1}] = (0).$$

Proof. By minimality of the selection of Ψ , we have $[A_k, \mathfrak{n}_{i_k-1}] = (0)$. But since $\lambda_{i_k}(A_k) = 1$, $[Z_{i_k}, \mathfrak{n}_{i_k-1}] \subset R(A_k, 1) \cap \mathfrak{n}_{i_k-1} = (0)$. □

We have $\mathfrak{h} \cap \text{cent}(\mathfrak{g}) = \bigcap \{\ker \lambda_{i_k} : 1 \leq k \leq d\}$; choose any basis A_{d+1}, \dots, A_u for $\mathfrak{h} \cap \text{cent}(\mathfrak{g})$. This determines a Jordan-Holder sequence $\mathfrak{g} = \mathfrak{g}_m \supset \mathfrak{g}_{m-1} \supset \dots$ for \mathfrak{g} : namely

$$\mathfrak{g}_j = \mathfrak{n} + \text{span}\{A_u, A_{u-1}, \dots, A_{m-j+1}\}, \quad n < j \leq m,$$

and $\mathfrak{g}_j = \mathfrak{n}_j$, $1 \leq j \leq n$. The corresponding basis elements are $\{Z_1, Z_2, \dots, Z_m\}$, where $Z_{n+1} = A_u, Z_{n+2} = A_{u-1}, \dots, Z_m = A_1$.

Now that a (conveniently chosen) Jordan-Holder basis is in place, we can employ the ‘‘layering’’ construction of [7]. As is well-known, each $l \in \mathfrak{g}^*$ determines a degenerate alternating bilinear form

$$B_l(Z, W) = l([Z, W]), \quad Z, W \in \mathfrak{g}.$$

For any subset \mathfrak{s} of \mathfrak{g} , let \mathfrak{s}^l denote the orthogonal complement of \mathfrak{s} , and set $\mathfrak{r}(l, \mathfrak{s}) = \mathfrak{s}^l \cap \mathfrak{s}$ = the ‘‘radical’’ of the restriction of B_l to \mathfrak{s} . We also denote $\mathfrak{r}(l, \mathfrak{g})$ by $\mathfrak{g}(l)$. It is well-known that for each l ,

$$\mathfrak{b}(l) = \sum_j \mathfrak{r}(l, \mathfrak{g}_j)$$

is a subalgebra of \mathfrak{g} which is totally isotropic for B_l . Note that for any $f \in \mathfrak{n}^\perp$, $\mathfrak{b}(l + f) = \mathfrak{b}(l)$. Also associated to l and to the above Jordan-Holder sequence for \mathfrak{g} , we have the index pair $\alpha(l) = (i(l), j(l))$ (cf. [7] or [8]). Here

$$i(l) \cup j(l) = e(l) = \{1 \leq j \leq m : \mathfrak{g}_j + \mathfrak{g}(l) \neq \mathfrak{g}_{j-1} + \mathfrak{g}(l)\},$$

and

$$j(l) = \{1 \leq j \leq m : \mathfrak{g}_j + \mathfrak{b}(l) \neq \mathfrak{g}_{j-1} + \mathfrak{b}(l)\}.$$

Let $p : \mathfrak{g}^* \rightarrow \mathfrak{n}^*$ denote the restriction mapping. Obviously $p|_{\mathfrak{h}^\perp}$ is one-to-one, onto, and G -equivariant. In this way we can identify \mathfrak{h}^\perp and \mathfrak{n}^* . In a similar way we identify \mathfrak{n}^\perp and \mathfrak{h}^* (though in this case the identification is not G -equivariant).

For each index pair α , the corresponding “layer” in \mathfrak{g}^* is the set $\Omega_\alpha = \{l : \alpha(l) = \alpha\}$. Each layer is a real algebraic subset of \mathfrak{g}^* , determined by the polynomials which depend only on $p(l)$. We totally order the set of all non-empty layers as in [8, Prop. 1.2] and let Ω_0 denote the minimal layer, with $\alpha^0 = (i^0, j^0)$ its sequence pair. Ω_0 is Zariski open in \mathfrak{g}^* and consists of G -orbits having maximal dimension. Since the condition $l \in \Omega_0$ depends only on $p(l)$, we have that for $f \in \mathfrak{h}^*$, $\Omega_0 \cap (f + \mathfrak{h}^\perp) \neq \emptyset$ and hence $\Omega_0 \cap (f + \mathfrak{h}^\perp)$ is an H -invariant Zariski open subset of $F + \mathfrak{h}^\perp$. It is necessary that our notion of “generic” H -orbits require that they be contained in $\Omega_0 \cap (f + \mathfrak{h}^\perp)$, but this is not sufficient.

For any l , set $\mathfrak{h}(l) = \mathfrak{g}(l) \cap \mathfrak{h}$, let

$$\Omega_1 = \{l \in \mathfrak{g}^* : l(Z_{i_k}) \neq 0, 1 \leq k \leq d\}.$$

It is clear that for each $l \in \Omega_1$, $\mathfrak{h}(l) = \mathfrak{h} \cap \text{cent}(\mathfrak{g})$, that Ω_1 is H -invariant and consists of H -orbit of maximal dimension, and that $\Omega_1 \cap (f + \mathfrak{h}^\perp)$ is Zariski open in $f + \mathfrak{h}^\perp$. Note however that Ω_1 is not necessarily G -invariant.

Lemma 1.2. *For every $l \in \Omega_1$, $i(l) \supset \Psi$ and $\mathfrak{b}(l) \cap \mathfrak{h} = \text{cent}(\mathfrak{g}) \cap \mathfrak{h}$.*

Proof. For $l \in \Omega_1$, and for $i_k \in \Psi$, we have $A_k \in \mathfrak{g}_{i_k-1}^l \sim \mathfrak{g}_{i_k}^l$, from which it follows that $i_k \in e(l)$. But by Lemma 1.1, $Z_{i_k} \in \mathfrak{r}(l, \mathfrak{g}_{i_k}) \subset \mathfrak{b}(l)$. Thus $i_k \in e(l) \sim j(l) = i(l)$.

Now let $A \in \mathfrak{h} \sim \text{cent}(\mathfrak{g})$. For some $i \in \Psi$, $\lambda_i(A) \neq 0$, hence $l([A, Z_i]) \neq 0$. But $Z_i \in \mathfrak{b}(l)$, so $A \notin \mathfrak{b}(l)$. □

Set $\Omega = \Omega_0 \cap \Omega_1$; the functionals in Ω will be the “generic” ones: given any $f \in \mathfrak{h}^*$, the irreducible representations which correspond to G -orbits Gl , $l \in \Omega \cap (f + \mathfrak{h}^\perp)$, are sufficient to decompose τ_f . Containment in Ω_0 will insure that the subalgebras $\mathfrak{b}(l)$ vary smoothly with l , while containment in Ω_1 insures a nice cross-section for the H -orbits in Ω . The abelian group $\text{cent}(G) \cap H$ is a direct factor of G contained in H , and thus will have no effect on the analysis of τ_f . The point of the preceding lemma is therefore that for the remainder of this paper, we can assume that $\mathfrak{b}(l) \cap \mathfrak{h} = (0)$ for every $l \in \Omega$.

Now we introduce the finite multiplicity assumption, and derive from it additional algebraic information about \mathfrak{g} . In our context here we have

$H(f+l) = f + Hl, f \in \mathfrak{h}^*, l \in \mathfrak{n}^*$, so that $H \setminus G$ is a finite multiplicity homogeneous space simply means that for each $l \in \Omega$, $\dim(Gl) = 2 \dim(Hl)$. In the case of finite multiplicity, the above results mean the following.

Corollary 1.3. *Assume that $H \setminus G$ is a finite multiplicity homogeneous space. Then $i^0 = \Psi$, and for every $l \in \Omega$, $\mathfrak{b}(l) = \mathfrak{g}(l) \oplus \mathfrak{k} = \mathfrak{k}^l$, where $\mathfrak{k} = \text{span}\{Z_{i_1}, Z_{i_2}, \dots, Z_{i_d}\}$ is an abelian subalgebra of \mathfrak{g} .*

Proof. We have $\#(\Psi) = \dim(Hl) = \dim(Gl)/2 = \#(i^0)$, so $i^0 = \Psi$. In the proof of Lemma 1.2, we saw that $\mathfrak{k} \subset \mathfrak{b}(l)$ and so by definition of i^0 , $\mathfrak{b}(l) = \mathfrak{g}(l) \oplus \mathfrak{k}$. \square

To derive more algebraic information about \mathfrak{g} , we use the pairing between elements of i^0 and j^0 , and between their corresponding basis elements, established in [7]. There the set j^0 is written as a (not necessarily increasing) sequence $\{j_1, j_2, \dots, j_d\}$, and subalgebras $\mathfrak{b}_k(l), 1 \leq k \leq d$, are defined, according to the following inductive scheme: for $l \in \Omega$, set $\mathfrak{b}_0(l) = \mathfrak{g}$, define $\mathfrak{b}_k(l) = \mathfrak{b}_{k-1}(l) \cap (\mathfrak{g}_{i_k} \cap \mathfrak{b}_{k-1}(l))^l$, and

$$j_k = \min\{1 \leq j \leq m : \mathfrak{g}_j \cap \mathfrak{b}_{k-1}(l) \not\subset \mathfrak{b}_k(l)\},$$

$k = 1, 2, \dots, d$. (The sequence $i^0 = \Psi = \{i_1 < i_2 < \dots < i_d\}$ can be obtained within this scheme also by setting

$$i_k = \min\{1 \leq j \leq m : \mathfrak{g}_j \cap \mathfrak{b}_{k-1}(l) \not\subset \mathfrak{r}(l, \mathfrak{b}_{k-1}(l))\},$$

but in this context that is not necessary.) Thus

$$\mathfrak{b}_1(l) = \mathfrak{g}_{i_1}^l, \mathfrak{b}_2(l) = (\mathfrak{g}_{i_2} \cap \mathfrak{g}_{i_1}^l)^l,$$

and so on. One has $\mathfrak{g} = \mathfrak{b}_0(l) \supset \mathfrak{b}_1(l) \supset \mathfrak{b}_2(l) \supset \dots \supset \mathfrak{b}_d(l) = \mathfrak{b}(l)$. In our case here it is easily seen from what we have done that for each k ,

$$\mathfrak{b}_k(l) = (\text{span}\{Z_{i_1}, Z_{i_2}, \dots, Z_{i_k}\})^l.$$

This gives the pairing of j_k with $i_k, 1 \leq k \leq d$, and hence a pairing of the basis elements Z_{i_k} and $Z_{j_k}, 1 \leq k \leq d$. For simplicity of notation, set $V_k = Z_{i_k}, W_k = Z_{j_k}, 1 \leq k \leq d$. Note that by Corollary 1.3, the V_k lie in \mathfrak{n} , but the W_k may or may not lie in \mathfrak{n} . Of course if $W_k \notin \mathfrak{n}$, then W_k is one of the basis elements of \mathfrak{h} . Let $\mathcal{R} = \{1 \leq k \leq d : W_k \in \mathfrak{n}\}$.

Lemma 1.4. *For each $k, 1 \leq k \leq d$, if $k \notin \mathcal{R}$, then $W_k = A_k$.*

Proof. Suppose that $W_k \notin \mathfrak{n}$, that is, $j_k > n$. Since $[A_r, \mathfrak{g}_{i_k}] = (0)$ for $r > k$, then $A_r \in \mathfrak{b}_k(l), k \leq r \leq d$. On the other hand, $[A_k, \mathfrak{g}_{i_{k-1}}] = (0)$

and $A_k \notin \mathfrak{g}_{i_k}^l$. By definition of the algebras $\mathfrak{b}_r(l), 1 \leq r \leq d$, this implies $A_k \in \mathfrak{b}_{k-1}(l) \sim \mathfrak{b}_k(l)$, hence $W_k = A_k$. \square

Lemma 1.5. *For each $k, 1 \leq k \leq d, [W_k, V_k] \neq 0, [\mathfrak{g}_{j_k-1}, V_k] = 0, [W_k, \mathfrak{g}_{i_k-1}] = 0$, and $[W_k, V_r] = 0$ for $r \neq k$. Also if $k \in \mathcal{R}$, then $\lambda_{j_k}(A_k) = -1$ and $\lambda_{j_k}(A_r) = 0$ for $r > k$.*

Proof. We proceed by induction on $k, 1 \leq k \leq d$. Suppose $k = 1$. Then by definition of i_1 and $j_1, [W_1, V_1] \neq 0$, and for every $l \in \Omega, j < j_1, l([Z_j, V_1]) = 0$. Since Ω is dense in \mathfrak{n}^* , this means $[Z_j, V_1] = 0, 1 \leq j < j_1$. Suppose that $1 \in \mathcal{R}$. For $j_1 \leq j \leq n$, we have $[Z_j, V_1] \in R(A_1, \lambda_j(A_1) + \lambda_{i_1}(A_1)) \cap \mathfrak{g}_{i_1-1}$, so that if $[Z_j, V_1] \neq 0$, then $\lambda_j(A_1) = -1$. In particular $\lambda_{j_1}(A_1) = -1$. For $r > 1, 0 \neq [W_1, V_1] \in R(A_r, \lambda_{j_1}(A_r)) \cap \mathfrak{g}_{i_1-1}$ yields $\lambda_{j_1}(A_r) = 0$. But now

$$[W_1, V_r] \in R(A_r, \lambda_{j_1}(A_r) + \lambda_{i_r}(A_r)) \cap \mathfrak{g}_{i_r-1} = R(A_r, 1) \cap \mathfrak{g}_{i_r-1} = (0).$$

Similarly $[W_1, \mathfrak{g}_{i_1-1}] \subset r(A_1, -1) \cap \mathfrak{g}_{i_1-1} = (0)$. On the other hand, if $1 \notin \mathcal{R}$, then $W_1 = A_1$ and $[W_1, V_r] = 0$ if and only if $r > 1$, and $[W_1, \mathfrak{g}_{i_1-1}] = (0)$, by definition of A_1 . This proves the lemma for the case $k = 1$.

Suppose that $k > 1$, and that the lemma is true for each $h, 1 \leq h < k$. Now by definition of j_k , for each $l \in \Omega$ and for each $j, 1 \leq j \leq j_k$, we have real numbers $c_{h,j}(l), 1 \leq h < k$, such that

$$Z_j(l) = Z_j + \sum_{1 \leq h < k} c_{h,j}(l) W_h \in \mathfrak{g}_{j_k} \cap \mathfrak{b}_{k-1}(l),$$

and again by definition of $j_k, l([Z_j(l), V_k]) \neq 0$. But by induction, $[W_h, V_k] = 0, 1 \leq h < k$, hence we have $[W_k, V_k] \neq 0$, and for $j < j_k, l([Z_j, V_k]) = 0$ for all $l \in \Omega$ so that $[Z_j, V_k] = 0$. If $k \in \mathcal{R}$, then arguing as in the case $k = 1$ we find that $\lambda_{j_k}(A_k) = -1, \lambda_{j_k}(A_r) = 0, r > k$, and $[W_k, \mathfrak{g}_{i_k-1}] \in R(A_k, -1) \cap \mathfrak{g}_{i_k-1} = (0)$. In particular, $[W_k, V_r] = 0$ for $r < k$, since $i_r < i_k$, and for $r > k, [W_k, V_r] \in R(A_r, 1) \cap \mathfrak{g}_{i_r-1} = (0)$. If $k \notin \mathcal{R}$, then $W_k = A_k$, and in this case, $[W_k, \mathfrak{g}_{i_k-1}] = 0$ and $[W_k, V_r] = 0$ if and only if $r \neq k$ just because of the definition of A_k . This proves the lemma. \square

Lemma 1.6. *If both k and r belong to \mathcal{R} , then $[W_k, W_r] = 0$.*

Proof. Assume $k > r$, and set $Z = [W_k, W_r]$. It is easily seen from Lemma 1.5 that $Z \in \mathfrak{b}(l)$ (for every $l \in \Omega$), and that $[A_k, Z] = -Z$. Let $i_0 = \min \{1 \leq i \leq n : Z \in \mathfrak{n}_i\}$, and set

$$a_k(l) = l([A_k, V_k]) / l([W_k, V_k]) = l(V_k) / l([W_k, V_k]), l \in \Omega.$$

The function $a_k(l)$ depends only on $l|_{\mathfrak{n}_{i_k}}$, and by Corollary 1.3 and Lemma 1.5, $A_k - a_k(l)W_k$ belong to $\mathfrak{b}(l)$. Set

$$P(l) = l([A_k - a_k(l)W_k, Z_k]) = -l(Z) - a_k(l)l([W_k, Z]).$$

Then P is identically zero on Ω , hence on all of \mathfrak{n}^* . Now if $i_0 < i_k$, then $P(l) = -l(Z)$ so $Z = 0$, while if $i_0 > i_k$, then we have $P(l) = -l(Z) + P_0(l)$ where $P_0(l)$ depends only on $l|_{\mathfrak{n}_{i_0-1}}$, so again $Z = 0$. But $i_0 = i_k$ is impossible, since if $i_0 = i_k$, say $Z = cV_k + Z_0$ with $c \neq 0$ and $Z_0 \in \mathfrak{n}_{i_0-1}$, then we find that

$$P(l) = -2cl(V_k) - l(Z_0) - l([W_k, Z_0]) / l([W_k, V_k]) = -2cl(V_k) + P_0(l)$$

where $P_0(l)$ depends only on $l|_{\mathfrak{n}_{i_0-1}}$. If this were the case then $P(l)$ could not be identically zero. Hence $Z = 0$. □

Proposition 1.7. *If $k \in \mathcal{R}$, then $[W_k, V_k]$ belongs to the center of \mathfrak{n} , $[W_k, Z_j] = 0$ for all $j \neq i_k, j \leq n$, and $\lambda_j(A_k) = 0$ for all $j \neq i_k$ or j_k . Moreover, for each $k, 1 \leq k \leq d$, we have $[Z_j, V_k] = 0$ for all $j \neq j_k$, and for any $i \notin e, j \notin e, 1 \leq i, j \leq n$, we have $[Z_j, Z_i] = 0$. In particular, \mathfrak{n} is two-step.*

Proof. Let $k \in \mathcal{R}$ and set $Z = [W_k, V_k]$. For any $j \leq n$, $[Z_j, W_k] \in \mathfrak{g}_{j_k-1}, [Z_j, V_k] \in \mathfrak{g}_{i_k-1}$, and so by Lemma 1.5, $[Z_j, Z] = [[Z_j, W_k], V_k] + [W_k, [Z_j, V_k]] = 0$.

Secondly we show that if $k \in \mathcal{R}$, and $j \neq i_k$ or j_k , then $\lambda_j(A_k) = 0$. Note that if $j \in i^0$, or if $j < i_k$, then we have $\lambda_j(A_k) = 0$ just by definition of A_k . Suppose that $j \in j^0$, say $j = j_r, Z_j = W_r$. We may assume that $r \in \mathcal{R}$, and by Lemma 1.4, we may assume that $r > k$. Now set $Z = [W_r, V_r]$. We have $[A_k, Z] = \lambda_j(A_k)Z$. For $l \in \Omega$, we have $A_k - a_k(l)W_k \in \mathfrak{b}(l)$ as in Lemma 1.6, and $Z \in \text{cent}(\mathfrak{n}) \subset \mathfrak{b}(l)$, so

$$P(l) = l([A_k - a_k(l)W_k, Z])$$

is therefore identically zero. But $P(l) = \lambda_j(A_k)l(Z)$, and since $Z \neq 0$ we must have $\lambda_j(A_k) = 0$. Next suppose that $j \notin e, j > j_k$, and for $1 \leq r \leq d$ set

$$c_r(l) = l([Z_j, V_r]) / l([W_r, V_r]).$$

Then $Z_j - \sum c_r(l)W_r$ belongd to $\mathfrak{b}(l)$ Now by Lemma 1.5, $c_r(l) = 0$ unless $j_r < j$, hence $l \rightarrow c_r(l)$ depends only on $l|_{\mathfrak{n}_{j-1}}$. Thus

$$P(l) = l\left([A_k - a_k(l)W_k, Z_j - \sum c_r(l)W_r]\right)$$

is identically zero, but $P(l) = \lambda_j(A_k)l(Z_j) + P_0(l)$ with $P_0(l)$ depending only on $l|_{\mathfrak{n}_{j-1}}$, which gives $\lambda_j(A_k) = 0$.

Now fix $k, 1 \leq k \leq d$, and $j \neq j_k$. To show that $Z = [Z_j, V_k] = 0$, from above results we may assume $j \notin e$ and $n \geq j > j_k$. Thus $k \in \mathcal{R}$ and so by the above $\lambda_j(A_k) = 0$. But then $Z \in R(A_k, 1) \cap \mathfrak{n}_{i_k-1} = (0)$.

Next we show that $[Z_j, W_k] = 0$ for $k \in \mathcal{R}$ and $j \neq i_k, j \leq n$, and here we can assume that $j \notin e$ and $j > j_k$. Note that from the above we have $[Z_j, V_k] = 0, 1 \leq k \leq d$, hence $Z_j \in \mathfrak{b}(l)$, for every $l \in \Omega$. We claim that $Z = [Z_j, W_k] \in \mathfrak{b}(l)$. Let $1 \leq r \leq d$; if $r \neq k$, then $[V_r, W_k] = [V_r, Z_j] = 0$ by above results and so $[V_r, Z] = 0$. On the other hand $[V_k, Z] = [Z_j, [V_k, W_k]]$, and since $[V_k, W_k] \in \text{cent}(\mathfrak{n})$ we have $l([V_k, Z]) = 0$. This proves the claim. Now $P(l) = l([A_k - a_k(l)W_k, Z])$ is identically zero. But $P(l) = -l(Z) - a_k(l)l([W_k, Z])$, and now an argument exactly like that of Lemma 1.6 gives $Z = 0$.

Finally, for $i \notin e, j \notin e, 1 \leq i, j \leq n$, the above shows that both Z_i and Z_j belong to $\mathfrak{b}(l)$ for every $l \in \Omega$. So $l([Z_i, Z_j]) = 0$ for every $l \in \Omega$ and this implies $[Z_i, Z_j] = 0$. This completes the proof. \square

Write $\mathcal{R} = \{r_1 < r_2 < \dots < r_u\}$ and

$$\{1, 2, \dots, d\} \sim \mathcal{R} = \{s_1 < s_2 < \dots < s_\nu\}$$

(here \sim denotes “set minus”). For the remainder of the paper we change notation for the basis elements of \mathfrak{h} : set $A_h = A_{r_h}, 1 \leq h \leq u$, and write $B_k = A_{s_k}, 1 \leq k \leq \nu$. We will use the coordinates

$$(t, s) = \exp(t_1 A_1) \dots \exp(t_u A_u) \exp(s_1 B_1) \dots \exp(s_\nu B_\nu)$$

for H freely, e.g., $q_{B,G}(t, s), \Lambda_j(t, s)$, etc.

Let us summarize what we know about the structure of G . Set $X_h = W_{r_h}, Y_h = V_{r_h}, 1 \leq h \leq u$.

Theorem 1.8. *Let G be the semi-direct product NH , with N nilpotent and normal in G , and with H abelian, $\text{Ad } H$ consisting of semi-simple transformations. If $H \backslash G$ is finite multiplicity, then N is two-step. Moreover, there are elements $X_1, X_2, \dots, X_u, Y_1, Y_2, \dots, Y_u$ in \mathfrak{n} such that*

- (i) $[X_k, Y_r] = 0$ if and only if $r \neq k$, and $[X_k, Y_k]$ is central in $\mathfrak{n}, 1 \leq k \leq u$,
- (ii) for every $r, k, [X_k, X_r] = [Y_k, Y_r] = 0$,
- (iii) $\mathfrak{n} = \text{cent}(\mathfrak{n}) + \text{span}\{X_1, X_2, \dots, X_u, Y_1, Y_2, \dots, Y_u\}$, and
- (iv) each X_k and Y_k is an eigenvector for $\text{Ad}(h), h \in H$.

It is clear that for each $l \in \Omega, \mathfrak{n}/\ker(l) \cap \text{cent}(\mathfrak{n})$ is Heisenberg, and

$$\begin{aligned} \mathfrak{p} &= \mathfrak{b}(l) \cap \mathfrak{n} = \text{span}\{Z_i : 1 \leq i \leq n, i \notin j^0\} = \\ &= \text{cent}(\mathfrak{n}) + \text{span}\{Y_1, Y_2, \dots, Y|_u\} \end{aligned}$$

is an abelian ideal in \mathfrak{g} (and of course a polarization at $p(l) = l|_{\mathfrak{n}}$).

It follows that

$$\mathfrak{b}(l) = \mathfrak{p} + \text{span}\{A_1 + a_1(l)X_1, A_2 + a_2(l)X_2, \dots, A_u + a_u(l)X_u\},$$

where the $a_h(l)$ are defined as in Lemma 1.6. Each A_k commutes with $\text{cent}(\mathfrak{n})$ and satisfies $[A_k, Y_k] = Y_k$, $[A_k, X_k] = -X_k$, and $[A_k, X_r] = [A_k, Y_k] = 0$, $r \neq k$. In particular $B(l)$ is a semi-direct product of $P = \exp(\mathfrak{p})$ with a vector group $W(l)$ of dimension u .

It is immediate from the above that there is a *single* subset Γ which is a cross-section for all of the coset spaces $B(l) \backslash G$.

Corollary 1.9. *For $(x_1, x_2, \dots, x_u) \in \mathbb{R}^u$ and $(s_1, s_2, \dots, s_\nu) \in \mathbb{R}^\nu$, set $\gamma(x, s) = \exp(x_1 X_1) \dots \exp(x_u X_u) \exp(s_1 B_1) \dots \exp(s_\nu B_\nu)$. For any $l \in \Omega$, the set $\Gamma = \{\gamma(x, s) : (x, s) \in \mathbb{R}^u \times \mathbb{R}^\nu\}$ is a cross-section for $B(l) \backslash G$.*

Remark 1.10. As mentioned in the introduction, our class of homogeneous spaces $H \backslash G$ has also been studied in [13]. There the irreducibles are constructed by means of the Mackey machine, and the spectrum of the quasi-regular representation τ is described by “Mackey parameters”. Let G be as in the hypothesis of Theorem 1.8, but without the assumption that $H \backslash G$ is finite multiplicity. Let $\sigma \in \hat{N}$, and let H_σ be the stabilizer in H of σ . One of the main ideas of [13] is to choose $f \in \mathfrak{n}^*$, belonging to the coadjoint orbit corresponding to σ , such that H_σ coincides with the stabilizer $H(f)$ of f in H [13, Theorem 3.2]; such a linear functional is said to be *aligned*. Then the natural map $\alpha : H_\sigma \rightarrow \text{Sp}(\mathfrak{n}/\mathfrak{n}(f))$ is considered. A result of this (though not explicitly stated there) is that $H \backslash G$ has uniform multiplicity 2^u if

$$(1.1) \quad \dim(\alpha(H_\sigma)) = \dim(\mathfrak{n}/\mathfrak{n}(f)) / 2$$

holds for generic $\sigma \in \hat{N}$, where $u = \dim(\alpha(H_\sigma))$.

On the other hand, suppose that $H \backslash G$ is finite multiplicity, and set $\mathfrak{h}_0 = \text{span}\{A_k : k \in \mathcal{R}\}$, $H_0 = \exp(\mathfrak{h}_0)$. Let $l \in \Omega_0 \subset \mathfrak{g}^*$, and let $\sigma \in \hat{N}$ be the irreducible representation (equivalence class) corresponding to the N -orbit of $f = p(l)$. It is easily seen that $\mathfrak{n} + \mathfrak{b}(l) = \mathfrak{n} + \mathfrak{n}^l = \mathfrak{n} + \mathfrak{b}_0$, and hence $H_0 = H_\sigma$. As we said above, a natural choice for $l \in \Omega_0$ when using the Mackey machine is one for which f is aligned, and in this setting that means $H_0 = H(f)$, hence $f(Y_k) = f(X_k) = 0$, $1 \leq k \leq u$. Thus the aligned linear functionals which are used in [13] are *not* in Ω . In the present work we shall construct the irreducibles as monomial representations by means of polarizations, and we shall use generic H -orbit parameters for the concrete Plancherel formula. For linear functionals in generic H -orbits we have $f(Y_k) \neq 0$. Thus for H -orbit parameters we use linear functionals in Ω , while for Mackey parameters one uses linear functionals that are not in Ω .

There is a strong parallel between the present work and the theory of Gelfand pairs (H, N) where N is nilpotent and H acts on N by automorphisms, but now H is a compact Lie group [1, 2]. To begin with, in [1]

it is shown that if (H, N) is a Gelfand pair, then N is two-step. Secondly, we can relate our situation to a result of Carcano [4] concerning Gelfand pairs. Let $\sigma \in \hat{N}$; realizing σ in a Hilbert space \mathcal{H}_σ , the Weil representation associated with σ is a representation ω_σ of H_σ acting in \mathcal{H}_σ that “extends” σ . It is well-known that in the setting of the present paper, ω_σ is quasi-equivalent to the regular representation of H_σ and has uniform multiplicity, and that (1.1) holds if and only if ω_σ has finite multiplicity, in which case that multiplicity is 2^u [13, Prop. 3.4]. Thus from the work of [13] and the present work, we can say the following, which parallels the above-mentioned result of Carcano. If ω_σ has finite multiplicity for almost every $\sigma \in \hat{N}$ (with respect to Plancherel measure), then $H \backslash G$ is finite multiplicity (and in this case both multiplicities are 2^u). Conversely, if $H \backslash G$ is finite multiplicity, then from our structural results on N it is easily seen that for every $\sigma \in \hat{N}$, the multiplicity of ω_σ is $2^{u'}$, where $u' \leq u$.

We conclude this section with the observation that Theorem 1.8 provides a coordinate-free description of all nilpotent groups that can arise in the class of homogeneous spaces we are considering. For $X \in \mathfrak{n}$ let $\mathfrak{c}(X)$ be the centralizer of X in \mathfrak{n} . If \mathfrak{n} is as in the Theorem 1.8, then $\dim(\mathfrak{n}/\mathfrak{c}(X)) \leq 1$, for every $X \in \mathfrak{n}$. On the other hand if \mathfrak{n} is a nilpotent Lie algebra such that $\dim(\mathfrak{n}/\mathfrak{c}(X)) \leq 1$, for every $X \in \mathfrak{n}$, then \mathfrak{n} is two-step (or abelian) and there are elements $X_1, X_2, \dots, X_u, Y_1, Y_2, \dots, Y_u$ in \mathfrak{n} that satisfy conditions (i), (ii), (iii) of the theorem. If \mathfrak{n} does have this form, it is clear that there is H as in the theorem such that $H \backslash NH$ is finite multiplicity. Hence we have the following.

Corollary 1.10. *Let N be a connected, simply connected nilpotent Lie group with Lie algebra \mathfrak{n} . Then the following are equivalent.*

- (i) *There is a vector subgroup H of $\text{Aut}(N)$ whose derived group in $\text{Aut}(\mathfrak{n})$ consists of semi-simple transformations and so that if G is the resulting semi-direct product, then $H \backslash G$ is finite multiplicity.*
- (ii) *For every $X \in \mathfrak{n}$, $\dim(\mathfrak{n}/\mathfrak{c}(X)) \leq 1$.*

We remark that the above class of two-step nilpotent groups is very different from the class of nilpotent groups known as Heisenberg-type (or H-type) groups [12] (that occur naturally in the study of Gelfands pairs.) A two-step nilpotent Lie algebra is H-type if $\dim(\mathfrak{n}/\mathfrak{c}(X)) = \dim(\text{cent}(\mathfrak{n}))$ holds for every $X \notin \text{cent}(\mathfrak{n})$. Hence if \mathfrak{n} is H-type and satisfies (ii) above then \mathfrak{n} is a Heisenberg Lie algebra. There seems to be no simple description of the class of two-step N that can arise in a Gelfand pair (H, N) .

2. Smooth vectors and generalized vectors.

Given a subalgebra \mathfrak{k} of \mathfrak{g} , let dk be a right Haar measure on $K = \exp(\mathfrak{k})$. Let Δ_K be the modular function of K (the derivative of right Haar measure with respect to a left Haar measure). In particular, one can take $\Delta_G(g) = \prod_{1 \leq i \leq n} \Lambda_i(g)$. For exponential solvable groups G , it is well known that there is a positive character q on G such that $q(k) = \Delta_K(k) / \Delta_G(k)$, $k \in K$, and that the space $K \backslash G$ carries a relatively invariant measure $d\gamma$ with modulus q^{-1} , that is, a measure $d\gamma$ which satisfies

$$\int_{K \backslash G} f(\gamma g) d\gamma = \int_{K \backslash G} f(\gamma) q(g)^{-1} d\gamma$$

for compactly-supported f on $K \backslash G$. We want to make natural choices of $dk, d\gamma$ for $K = B(l)$, but before addressing that issue, we make some more general comments. Let χ be a unitary character of K . Let $C^\infty(G, K, \chi)$ denote the space of smooth functions f on G which satisfy $f(kg) = \chi(k) f(g)$, and let $C_c^\infty(G, K, \chi)$ be the subspace of $C^\infty(G, K, \chi)$ consisting of those f which are compactly supported mod K . The Hilbert space $L^2(G, K, \chi)$ is the completion of $C_c^\infty(G, K, \chi)$ under the norm $\|f\|_2 = \left[\int_{K \backslash G} |f(\gamma)|^2 d\gamma \right]^{1/2}$. Let π_χ be the irreducible representation induced from the character χ of K , so that π_χ acts in the space $\mathcal{H}_\chi = L^2(G, K, \chi)$ by the formula

$$\pi_\chi(s) f(g) = f(gs) q(s)^{1/2}.$$

Let \mathcal{H}_χ^∞ be the Fréchet space of smooth vectors for π_χ in \mathcal{H}_χ , and let $\mathcal{H}_\chi^{-\infty}$ denote its antidual. It is well-known that $\mathcal{H}_\chi^\infty \subset C^\infty(G, K, \chi)$ [20].

Fix $l \in \Omega$, and let $B = B(l) = \exp(\mathfrak{b}(l))$. We have seen that $B = PW$, where $P = \exp(\mathfrak{p})$ is the polarization in \mathfrak{n} at $p(l)$, and $W = \exp(\mathfrak{w}(l))$ is an abelian group of dimension u . A basis for \mathfrak{p} is $\{Z_j : 1 \leq j \leq n, j \notin j^0\}$ and for $\mathfrak{w}(l)$ is $\{A_k - a_k(l) X_k : 1 \leq k \leq u\}$. Letting dp and dw be the Lebesgue measures on P and W resp. obtained from these coordinates, a natural choice for right Haar measure on B is just $db = dpdw$. Recall that we have the index set $\mathcal{R} = \{1 \leq k \leq d : j_k \leq n\}$. Define a positive character $q_{B,G}$ on G by

$$q_{B,G} = \prod_{k \in \mathcal{R}} \Lambda_{j_k}^{-1}(g).$$

Then for $Y \in \mathfrak{b}$,

$$q_{B,G}(\exp Y) = e^{-\text{tr ad}_{\mathfrak{g}/\mathfrak{b}} Y} = \Delta_B(\exp Y) / \Delta_G(\exp Y).$$

Note that this is not the only choice for $q_{B,G}$ that we could have taken (one can extend Δ_B / Δ_G in many ways), and the choice of $q_{B,G}$ affects the

relatively invariant measure $d\gamma$ as well as the growth properties of functions in the resulting space \mathcal{H}_χ . The above choice is natural and more importantly, will result in manageable growth properties.

Recall that we have the coordinates $\gamma(x, s)$ on $B \setminus G$ given in the previous section.

Lemma 2.1. *Let $d\gamma$ be the measure on $B \setminus G$ defined by*

$$\int_{B \setminus G} f(\gamma) d\gamma = \int_{\mathbb{R}^d} f(\gamma(x, s)) q_{B,G}(\gamma(0, s)) dx ds$$

where $dx ds$ denotes Lebesgue measure on $\mathbb{R}^u \times \mathbb{R}^\nu$. Then $d\gamma$ is relatively invariant with modulus $q_{B,G}^{-1}$.

Proof. Let $g \in G$, and define the diffeomorphism $T_g : \mathbb{R}^u \times \mathbb{R}^\nu \rightarrow \mathbb{R}^u \times \mathbb{R}^\nu$ by $B\gamma(x, s)g = B\gamma(T_g(x, s))$. If $g \in P$, then normality of P gives that $T_g = \text{Id}$. Let t be any real number; we compute T_g in the cases (a) $g = \exp(tB_h)$ for some $h, 1 \leq h \leq \nu$, (b) $g = \exp(tX_k)$, for some $k, 1 \leq k \leq u$, and (c) $g = \exp(t(A_k - a_k(l)X_k))$, for $1 \leq k \leq u$.

(a) Here we have

$$\begin{aligned} \gamma(x, s)g &= \exp(x_1 X_1) \dots \exp(x_u X_u) \times \\ &\quad \times \exp(s_1 B_1) \dots \exp([t + s_h] B_h) \dots \exp(s_\nu B_\nu) \end{aligned}$$

hence $T_g(x, s) = (x, s_1, \dots, t + s_h, \dots, s_\nu)$.

(b) In this case

$$\begin{aligned} \gamma(x, s)g &= \exp(x_1 X_1) \dots \exp([e^a t + x_k] X_k) \dots \exp(x_u X_u) \times \\ &\quad \times \exp(s_1 B_1) \dots \exp(s_\nu B_\nu) \end{aligned}$$

where $a = \sum_r \lambda_{j_k}(B_r)$, so $T_g(x, s) = (x_1, x_2, \dots, e^a t + x_k, \dots, x_u, s)$.

(c) Here $g = x(t) \exp(tA_k)$ and $\exp(tA_k) = gy(t)$ where

$$x(t) = \exp(a_k(l)(e^{-t} - 1)X_k)$$

and

$$y(t) = \exp(a_k(l)(e^t - 1)X_k).$$

We have

$$\begin{aligned} \gamma(x, s) \exp(tA_k) &= \exp(tA_k) \times \\ &\quad \times \gamma((x_1, \dots, e^t x_k, \dots, x_u), s) = \\ &= gy(t) \gamma((x_1, \dots, e^t x_k, \dots, x_u), s) = \\ &\quad = g\gamma((x_1, \dots, e^t x_k + a_k(l)(e^t - 1), \dots, x_u), s) \end{aligned}$$

hence

$$\begin{aligned} \gamma(x, s) g &= \gamma(x, s) \exp(t(A_k - a_k(l)X_k)) = \gamma(x, s) x(t) \exp(tA_k) = \\ &= \gamma\left(\left(x_1, \dots, x_k + a_k(l)e^{\alpha(s)}(e^{-t} - 1), \dots, x_u\right), s\right) \exp(tA_k) = \\ &= \exp(tA_k) \gamma\left(\left(x_1, \dots, e^t x_k + a_k(l)e^{\alpha(s)}(1 - e^t), \dots, x_u\right), s\right) = \\ &= g\gamma\left(\left(x_1, \dots, e^t x_k + a_k(l)(1 - e^t)(e^{\alpha(s)} - 1), \dots, x_u\right), s\right). \end{aligned}$$

Thus

$$T_g(x, s) = \left(\left(x_1, \dots, e^t x_k + a_k(l)(1 - e^t)(e^{\alpha(s)} - 1), \dots, x_u\right), s\right).$$

To finish the proof one need only check that in each case,

$$q_{B,G}(\gamma(T_g(x, s))) J_g(x, s) = q_{B,G}(g),$$

where $J_g(x, s)$ is the Jacobian determinant of $T_g(x, s)$. We leave this to the reader. \square

Let $\chi = \chi_l$ be the character of B defined by $\chi_l(\exp Y) = e^{i(Y)}$, $Y \in B$. Set $\pi_l = \pi_\chi$, $\mathcal{H}_l = \mathcal{H}_\chi$, etc. Since BH is an open subset of G , H may be regarded as an open subset of $B \setminus G$. Using the coordinates (t, s) for H and the coordinates (x, s) for $B \setminus G$, we compute that the map $\varphi : H \rightarrow \mathbb{R}^u \times \mathbb{R}^v$ defined by

$$\varphi(t, s) = (a_1(l)(e^{t_1} - 1), \dots, a_u(l)(e^{t_u} - 1), s)$$

satisfies $B\gamma(\varphi(t, s)) = B(t, s)$.

We want to construct an appropriately covariant generalized vector for π_l , that is, an element of

$$(\mathcal{H}_l^{-\infty})^{q_{H,G}^{-1/2}} = \left\{ \beta \in \mathcal{H}_l^{-\infty} : \pi_l(h)\beta = q_{H,G}(h)^{-1/2}\beta, \text{ for every } h \in H \right\}.$$

Following Fujiwara and Yamagami [11], and Lipsman [16], we define *formally*

$$(2.1) \quad \beta_l(f) = \int_H \bar{f}(h) q_{B,G}^{1/2} q_{H,G}^{-1/2} \chi_f(h) dh, f \in \mathcal{H}_l^\infty.$$

It is not at all obvious that (2.1) is convergent for all $f \in \mathcal{H}_l^\infty$. Note for example that if $f \in C_c^\infty(B, G, \chi)$, then $f|_H$ may not be compactly supported (if $\mathcal{R} \neq \emptyset$, then the image of H in $B \setminus G$ is not closed). Hence it is not

immediate that (2.1) is finite even for $f \in C_c^\infty(G, B, \chi)$. However we shall prove the following.

Theorem 2.2. *The integral (2.1) is absolutely convergent for every $f \in \mathcal{H}_l^\infty$, and β_l is continuous on \mathcal{H}_l^∞ .*

This result is a generalization of the proof of convergence in [18], in which it is assumed the N is abelian. We have seen that N is abelian if and only if $\mathcal{R} = \emptyset$, so in what follows we assume that $\mathcal{R} \neq \emptyset$. To prove the result we need information about the growth properties of f on H . For simplicity of notation we shall write $f(t, s)$ for $f|_{H(t,s)}$. We make a couple of observations. First, $f \in \mathcal{H}_l$ implies that $f|_H$ is square integrable on H with respect to the measure $q_{B,G}(t, s) dt ds$, or in other words, $(f|_H)(q_{B,G})^{1/2} \in L^2(H)$, and $\left\| (f|_H)(q_{B,G})^{1/2} \right\|_{L^2(H)} \leq \|f\|_{\mathcal{H}_l}$. Second, for $f \in \mathcal{H}_l^\infty$, the differential operators $\pi(Z)$, $Z \in \mathbf{U}(\mathfrak{g}_c)$ (=the enveloping algebra of the complexification \mathfrak{g}_c of \mathfrak{g}), act on $f|_H$, and $\left\| (q_{B,G})^{1/2} \pi(Z)(f|_H) \right\|_{L^2(H)} \leq \|\pi(Z)f\|_{\mathcal{H}_l}$ for every $Z \in \mathbf{U}(\mathfrak{g}_c)$. We can compute $\pi(Z)$ as an operator on $\varphi(H)$ (in the G -coordinates (x, s)) or as an operator on H (in the H -coordinates (t, s)). In the latter coordinates the algebra $\pi(\mathbf{U}(\mathfrak{g}_c))$ is more easily described and provides us with more useful information about \mathcal{H}_l^∞ .

First we set some notation that will be convenient. For $J \subset \{1, 2, \dots, n\}$, set $\Lambda_J(t, s) = \prod_{j \in J} \Lambda_j(t, s)$. Recall we have written $\mathcal{R} = \{r_1 < r_2 < \dots < r_u\}$; set $q_k(t, s) = \Lambda_{j_{r_k}}(t, s)$, $1 \leq k \leq u$, and for $K \subset \{1, 2, \dots, u\}$, set $q_K(t, s) = \prod_{k \in K} q_k(t, s)$ (so that for $K = \{1, 2, \dots, u\}$, $q_K = q_{B,G}^{-1}$). Denote by $C_0(H)$ the space of continuous functions on H that vanish at infinity. We recall (a weak form of) a standard regularity result that if f and its partial derivatives of all orders belong to $L^2(H)$, then $f \in C_0(H)$. In fact if we choose a fixed constant-coefficient partial differential operator D on H such that the reciprocal of its ‘‘symbol’’ $P = \hat{D}$ belongs to $L^2(H)$, then we have $\|f\|_\infty \leq \|1/P\|_{L^2(H)} \|Df\|_{L^2(H)}$.

Lemma 2.3. *Let $f \in \mathcal{H}_l^\infty$, let $J \subset \{1, 2, \dots, n\}$, $J' \subset \{1, 2, \dots, n\} \sim j^0$, and let $K \subset \{1, 2, \dots, u\}$. Set $D_K = \prod_{k \in K} \partial_{t_k}$. Then*

(a) *the function*

$$\phi(t, s) = q_{B,G}^{1/2}(t, s) \Lambda_J(0, s) \Lambda_{J'}(0, s) f(t, s) \in C_0(H),$$

and there is $V \in \mathbf{U}(\mathfrak{g}_c)$ and a constant M , depending only on J and J' , such that $\|\phi\|_\infty \leq M \|\pi(V)f\|_{\mathcal{H}_l}$ and

(b) *the function*

$$\phi(t, s) = q_{B,G}^{1/2}(t, s) \Lambda_J(0, s) \Lambda_{J'}(t, s) q_K(t, s) D_K f(t, s) \in C_0(H),$$

and there is $W \in \mathbf{U}(\mathfrak{g}_c)$ and a constant M , depending only on J, J' , and K , such that $\|\phi\|_\infty \leq M \|\pi(W)f\|_{\mathcal{H}_j}$.

Proof. We begin by computing $\pi(Z)$ as an operator on H , for certain $Z \in \mathfrak{g}$. First, consider a basis element Z_j which belongs to the center of \mathfrak{n} . We have $Z_j \in \mathfrak{p}$ and by Proposition 1.7, $\lambda_j(A_k) = 0, 1 \leq k \leq u$, so one finds that $\pi(Z_j) = i\Lambda_j(t, s) = i\Lambda_j(0, s)$. Next, let $r = r_k \in \mathcal{R}$. If $j = i_r$, then $\pi(Z_j) = i\Lambda_j(t, s) = e^{t_k}$ and $\Lambda_j(0, s) = \Lambda_j(0, 0) = 1$. Suppose that $j = j_r$. Then $[A_h, [X_k, Y_k]] = 0, 1 \leq h \leq u$, and since $[B_h, Y_k] = 0, [B_k, [X_k, Y_k]] = \lambda_{j_r}(B_h)[X_k, Y_k], 1 \leq h \leq \nu$. Hence $\pi([X_k, Y_k]) = i\Lambda_{j_r}(0, s)$. Thus by taking the appropriate element U of $\mathbf{U}(\mathfrak{g}_c)$, we have $\pi(U) = \Lambda_J(0, s)\Lambda_{J'}(t, s)$ as an operator on H and

$$q_{B,G}^{1/2}(t, s)\Lambda_J(0, s)\Lambda_{J'}(t, s)f(t, s) \in L^2(H).$$

Now $\pi(A_k) = \partial_{t_k}, 1 \leq k \leq u$ and $\pi(B_k) = \partial_{s_k}, 1 \leq k \leq \nu$, so if D_α is any mixed partial of order $|\alpha|$, then $D_\alpha \in \pi(\mathbf{U}(\mathfrak{h}))$ and so

$$q_{B,G}^{1/2}(t, s)\Lambda_J(0, s)\Lambda_{J'}(t, s)D_\alpha f(t, s)$$

belongs to $L^2(H)$ also. But since the function $q_{B,G}^{1/2}(t, s)\Lambda_{J'}(t, s)\Lambda_J(0, s)$ involves only exponentials in t and s , then $D_\alpha(q_{B,G}^{1/2}(t, s)\Lambda_J(0, s)f(t, s))$ can be written as a sum of terms of the form

$$c_\beta q_{B,G}^{1/2}(t, s)\Lambda_J(0, s)\Lambda_{J'}(t, s)D_\beta f(t, s), |\beta| < |\alpha|, c_\beta \in \mathbb{R}.$$

Hence all partials of $\phi(t, s) = q_{B,G}^{1/2}(t, s)\Lambda_J(0, s)\Lambda_{J'}(t, s)f(t, s)$ also belong to $L^2(H)$, and so $\phi \in C_0(H)$. Now choose $Z \in \mathbf{U}(\mathfrak{h}_c)$ for which the reciprocal of the Fourier transform $\pi(\hat{Z})$ belongs to $L^2(H)$ and we have $\pi(Z)\phi(t, s) = q_{B,G}^{1/2}(t, s)\sum_{|\beta| < |\alpha|} \pi(Z_\beta)f(t, s)$ where in this case α is the order of $\pi(Z)$. Hence

$$\|\phi\|_\infty \leq M \|\pi(Z)\phi\|_{L^2(H)} = M \left\| \sum_{|\beta| < |\alpha|} \pi(Z_\beta)f(t, s) \right\|_{\mathcal{H}_i}.$$

As for the function (b), we compute that for $1 \leq k \leq u, \pi(X_k) = q_k(t, s)\partial_{t_k}$. Now by Lemma 1.5 and Proposition 1.7, $q_k(t, s) = e^{-t_k}q_k(0, s), 1 \leq k \leq u$, so if $X_K = \prod_{k \in K} X_k$, then $\pi(X_K) = q_K(t, s)D_K$, and we have

$$q_{B,G}^{1/2}(t, s)\Lambda_J(0, s)\Lambda_{J'}(t, s)q_K(t, s)D_K f(t, s) \in L^2(H).$$

Now in a similar manner as before we find that

$$q_{B,G}^{1/2}(t, s)\Lambda_J(0, s)\Lambda_{J'}(t, s)q_K(t, s)Df(t, s)$$

belongs to $C_0(H)$ and the indicated estimate is obtained. □

Note that by taking $K = \{k\}$, Lemma 2.3 tells us something about the growth of $\partial_{t_k} f$ as $T_k \rightarrow -\infty, 1 \leq k \leq u$, (in particular, if all other variables are held constant, then $\partial_{t_k} f(t, s) \rightarrow 0$ rapidly as $t \rightarrow -\infty$). In the next lemma we derive information about the growth of f itself as $t_k \rightarrow -\infty$.

Lemma 2.4. *Let $f \in \mathcal{H}_l^\infty$, let $P \subset \{1, 2, \dots, u\}$, and let*

$$U(P) = \{(t, s) \in H : \log(q_k(t, s)) > 0, \text{ for every } k \in P\}.$$

Let $\tilde{P} = \{1, 2, \dots, u\} \sim P$ and let K be any subset of $P, \tilde{K} = P \sim K$. Write $K = \{k_1, k_2, \dots, k_a\}, \tilde{K} = \{h_1, h_2, \dots, h_b\}$, and write $t \in \mathbb{R}^u$ as $t = (t_K, t_{\tilde{K}}, t_{\tilde{P}})$ (with the obvious meaning). Finally write $Q_{\tilde{K}}(s) = (\log(q_{h_1}(0, s)), \log(q_{h_2}(0, s)), \dots, \log(q_{h_b}(0, s)))$.

Then for any $J \subset \{1, 2, \dots, n\}, J' \subset \{i_{r_k} : k \in \tilde{P}\}$, the function

$$\phi(t, s) = \Lambda_J(0, s) \Lambda_{J'}(t, s) q_{\tilde{P}}(t, s)^{-1/2} f((t_K, Q_{\tilde{K}}(s), t_{\tilde{P}}), s)$$

is bounded on $U(P)$. Moreover, there is a finite set of positive constants $\{M_1, M_2, \dots\}$, and elements $\{W_1, W_2, \dots\}$ in $\mathbf{U}(\mathfrak{g}_c)$, depending only on J, J' and K , such that

$$\sup_{U(P)} |\phi(t, s)| \leq \sum_{\beta} M_{\beta} \|\pi(W_{\beta}) f\|_{\mathcal{H}_l}.$$

Proof. Note that if $P = \emptyset$, then $t = t_{\tilde{P}}$ and $q_{\tilde{P}}^{-1/2} = q_{B,G}^{1/2}$, so in this case we are done by Lemma 2.3. Assume that $P \neq \emptyset$. We proceed by induction on $a = \#(K)$. If $a = 0$, then

$$\phi(t, s) = \Lambda_J(0, s) \Lambda_{J'}(t, s) q_{B,G}(t, s)^{1/2} f(t, s) |_{t_k = \log(q_k(0, s)), k \in P}$$

so again by Lemma 2.3 we are done. Suppose that $a > 0$, and that the lemma holds for all K' with $\#(K') < a$. Now for each $k = 1, 2, \dots, u$, $q_k(t, s) = e^{-t_k} q_k(0, s)$ so $\log(q_k(0, s)) > 0$ means $t_k < \log(q_k(0, s))$. For each $(t, s) \in U(P)$ let $E = E(t, s)$ be the subset of \mathbb{R}^a defined by

$$E(t, s) = \{\tau \in \mathbb{R}^a : t_{k_{\alpha}} < \tau_{\alpha} < \log(q_{k_{\alpha}}(0, s)), 1 \leq \alpha \leq a\}$$

and set $D_K = \prod_{k \in K} \partial_{t_k}$.

Replacing $t_K = (t_{k_1}, t_{k_2}, \dots, t_{k_a})$ by $\tau = (\tau_1, \tau_2, \dots, \tau_a)$ in $f((t_K, Q_{\tilde{K}}(s), t_{\tilde{P}}), s)$ and integrating $D_K f$ over E , repeated application of the fundamental theorem of calculus gives

$$\int_E D_K f((\tau, Q_{\tilde{K}}(s), t_{\tilde{P}}), s) d\tau = \sum (-1)^{\#(K \sim K')} f((t_{K'}, Q_{\tilde{K}'}(s), t_{\tilde{P}}), s)$$

where the sum is taken over all subsets K' of K . Now multiply both sides of the above by $\Lambda_J(0, s) \Lambda_{J'}(t, s) q_{\tilde{P}}(t, s)^{-1/2}$, and we get

$$\begin{aligned} |\phi(t, s)| &= \\ &= \left| \int_E \Lambda_J(0, s) \Lambda_{J'}(t, s) q_{\tilde{P}}(t, s)^{-1/2} D_K f((\tau, Q_{\tilde{K}}(s), t_{\tilde{P}}), s) d\tau \right| + \\ &\quad + \sum_{K' \neq K} \left| \Lambda_J(0, s) \Lambda_{J'}(t, s) q_{\tilde{P}}(t, s)^{-1/2} f((t_{K'}, Q_{\tilde{K}'}(s), t_{\tilde{P}}), s) \right|. \end{aligned}$$

Let $g_{K'}$ be a term in the right hand sum with $K' \neq K$. Then by induction, there are finitely many constants $M_{\beta, K'}$ and elements $W_{\beta, K'} \in \mathbf{U}(\mathfrak{g}_c)$ depending only on J, J' , and K' , such that

$$\sup_{U(P)} |g_{K'}(t, s)| \leq \sum_{\beta} M_{\beta, K'} \|\pi(W_{\beta, K'}) f\|_{\mathcal{H}_i}.$$

Therefore it remains to show that the function

$$I(t, s) = \int_E \Lambda_J(0, s) \Lambda_{J'}(t, s) q_{\tilde{P}}(t, s)^{-1/2} D_K f((\tau, Q_{\tilde{K}}(s), t_{\tilde{P}}), s) d\tau$$

is bounded on $U(P)$ in a similar way. To see this, note that

$$\begin{aligned} q_{B, G}((\tau, Q_{\tilde{K}}(s), t_{\tilde{P}}), s) &= q_K(\tau, s)^{-1} q_{\tilde{P}}(t, s)^{-1} \text{ and} \\ q_{B, G}((\tau, Q_{\tilde{K}}(s), Q_{\tilde{P}}(s)), s) &= q_K(\tau, s)^{-1} \end{aligned}$$

where $Q_{\tilde{P}}(s)$ means we have replaced t_k by $\log(q_k(0, s))$ for all $k \in \tilde{P}$. Hence

$$\begin{aligned} \Lambda_J(0, s) \Lambda_{J'}(t, s) q_{\tilde{P}}^{-1/2} D_K f((\tau, Q_{\tilde{K}}(s), t_{\tilde{P}}), s) &= \\ \left[q_{B, G}^{1/2}((\tau, Q_{\tilde{K}}(s), t_{\tilde{P}}), s) \Lambda_J(0, s) \Lambda_{J'}(t, s) q_K(\tau, s) D_K \times \right. \\ \left. \times f((\tau, Q_{\tilde{K}}(s), t_{\tilde{P}})) \right] q_{B, G}^{1/2}((\tau, Q_{\tilde{K}}(s), Q_{\tilde{P}}(s)), s). \end{aligned}$$

But we can apply Lemma 2.4 to the function inside the brackets above, and in so doing obtain $M > 0$ and $W \in \mathbf{U}(\mathfrak{g}_c)$ such that

$$\begin{aligned} I(t, s) &\leq M \|\pi(W) f\|_{\mathcal{H}_i} \int_E q_{B, G}^{1/2}((\tau, Q_{\tilde{K}}(s), Q_{\tilde{P}}(s)), s) d\tau = \\ &= M \|\pi(W) f\|_{\mathcal{H}_i} \int_E q_K(\tau, s)^{-1} d\tau. \end{aligned}$$

But it is easily seen that $\int_E q_K(\tau, s)^{-1} d\tau \leq 1$ on $U(P)$, and this proves the lemma. \square

Proof of Theorem 2.2. Let $S = \{1 \leq j \leq n : \Lambda_j(t, s) \neq 1 \text{ for some } (t, s) \in H\}$. For each $j \in S$, let ε_j be a choice of sign, $\varepsilon_j = \pm 1$. Set

$$U_\varepsilon = \{(t, s) \in H : \log \Lambda_j(t, s) \varepsilon_j > 0, \text{ for every } j \in S\}.$$

(Of course some of the U_ε may be empty). H is the disjoint union of the sets U_ε , and for each ε , set $P = \{1 \leq k \leq u : \varepsilon_{j_{r_k}} = +1\}$ so that $U_\varepsilon \subset U(P)$.

Fix $U_\varepsilon \neq \emptyset$; we need to show that $f q_{B,G}^{1/2} q_{H,G}^{-1/2}$ is integrable on U_ε . We begin by noting that from Lemma 1.5 and Proposition 1.7, we have $q_{H,G}^{-1/2}(t, s) = q_{H,G}^{-1/2}(0, s)$ for every $(t, s) \in H$. Recall that for each $k = 1, 2, \dots, u$, $\Lambda_{i_{r_k}}(t, s) = e^{t_k}$ and $\Lambda_{j_{r_k}}(t, s) = q_k(t, s) = e^{-t_k} q_k(0, s)$, and recall also that for all other j , $\Lambda_j(t, s) = \Lambda_j(0, s)$, $(t, s) \in H$. Next we observe that for each k the sign of $\log(q_k(0, s))$ is constant on U_ε , for since $[X_k, Y_k] \in \text{cent}(\mathfrak{n})$ and $\text{Ad}(t, s)([X_k, Y_k]) = q_k(0, s)[X_k, Y_k]$, there is some j (with $Z_j \in \text{cent}(\mathfrak{n})$) such that $\Lambda_j(t, s) = \Lambda_{j_{r_k}}(0, s) = q_k(0, s)$ for all $(t, s) \in H$. Hence the sign of $\log(\Lambda_j(0, s))$ is constant on U_ε for each $j \in S, j \neq i_{r_k}$. Let $I = S \sim \{i_{r_k} : 1 \leq k \leq u\}$ and let $I^+ = \{j \in I : \log \Lambda_j(0, s) > 0 \text{ on } U_\varepsilon\}$, $I^- = I \sim I^+$. Then

$$\begin{aligned} q_{H,G}^{-1/2}(t, s) &= q_{H,G}^{-1/2}(0, s) = \prod_{j \in I} \Lambda_j(0, s)^{1/2} = \\ &= \prod_{j \in I^+} \Lambda_j(0, s) \prod_{j \in I^+} \Lambda_j(0, s)^{-1/2} \prod_{j \in I^-} \Lambda_j(0, s)^{1/2}. \end{aligned}$$

On the other hand

$$q_{B,G}^{1/2}(t, s) = q_P(t, s)^{-1/2} q_{\tilde{P}}(t, s)^{-1/2}.$$

We partition \tilde{P} : let

$$Q = \{k \in \tilde{P} : \log(\Lambda_{i_{r_k}}) > 0 \text{ on } U_\varepsilon\} = \{k \in \tilde{P} : t_k > 0 \text{ on } U_\varepsilon\}$$

and let $\tilde{Q} = \tilde{P} \sim Q$. Applying Lemma 2.4 with $K = P, J = I^+$, and $J' = \{i_{r_k} : k \in Q\}$, we have that

$$\phi(t, s) = \left(\prod_{j \in I^+} \Lambda_j(0, s) \right) \left(\prod_{k \in Q} e^{t_k} \right) q_{\tilde{P}}(t, s)^{-1/2} f(t, s)$$

is bounded on U_ε . We claim that the function

$$\Psi(t, s) = q_P^{-1/2}(t, s) \prod_{k \in Q} e^{-t_k} \prod_{j \in I^+} \Lambda_j(0, s)^{-1/2} \prod_{j \in I^-} \Lambda_j(0, s)^{1/2}$$

is integrable on U_ϵ . If this is so then

$$f(t, s) q_{B,G}^{1/2}(t, s) q_{H,G}^{-1/2}(t, s) = \phi(t, s) \Psi(t, s)$$

is integrable on U_ϵ and we are done. To prove the claim, we partition Q : set $Q^+ = \{k \in Q : \log(q_k(0, s)) > 0\}$, $Q^- = Q \sim q^+$. Set

$$\begin{aligned} \Psi_{Q^+}(t, s) &= \prod_{k \in Q^+} e^{-t_k} q_k(0, s)^{-1/2} \\ \Psi_{Q^-}(t, s) &= \prod_{k \in Q^-} e^{-t_k} q_k(0, s)^{1/2} \\ \Psi_{\tilde{Q}}(s) &= \prod_{k \in \tilde{Q}} q_k(0, s)^{1/2} \\ \Psi_+(s) &= \prod_{j \in I^+ \sim j^0} \Lambda_j(0, s)^{-1/2} \\ \Psi_-(s) &= \prod_{j \in I^- \sim j^0} \Lambda_j(0, s)^{1/2}. \end{aligned}$$

Then $\Psi(t, s) = q_P(t, s)^{-1/2} \Psi_{Q^+}(t, s) \Psi_{Q^-}(t, s) \Psi_{\tilde{Q}}(s) \Psi_+(t, s) \Psi_-(s)$. Note that $I \sim j^0$ consists of indices j for which $Z_j \in \text{cent}(\mathfrak{n})$. Next we describe the set U_ϵ : define

$$V_\epsilon = \{s : \log(\Lambda_j(0, s)) \epsilon_j > 0 \text{ for } j \in I \sim j^0\}$$

and for $s \in V_\epsilon$, set

$$W_\epsilon(s) = \left\{ t : t_k \epsilon_{i_{r_k}} > 0 \text{ and } (-t_k + \log(q_k(0, s))) \epsilon_{j_{r_k}} > 0, 1 \leq k \leq u \right\},$$

so that $U_\epsilon = \{(t, s) : s \in V_\epsilon, t \in W_\epsilon(s)\}$ and

$$\int_{U_\epsilon} \Psi(t, s) dt ds = \int_{V_\epsilon} \left[\int_{W_\epsilon(s)} \Psi(t, s) dt \right] ds.$$

It is enough to show that the function $s \rightarrow \int_{W_\epsilon(s)} \Psi(t, s) dt$ is exponentially decreasing on V_ϵ . To see this, note first that the function $\Psi_+(s) \Psi_-(s)$ is exponentially decreasing on V_ϵ . Now for each $s \in V_\epsilon$, $W_\epsilon(s)$ is simply a u -dimensional cube in \mathbb{R}^u , and we consider each interval which makes up $W_\epsilon(s)$. If $k \in \tilde{Q}$, then $\log(q_k(0, s)) < t_k < 0$, and

$$\int_{(\log(q_k(0, s)), 0)} q_k(0, s) dt_k = \log(q_k(0, s)) q_k(0, s).$$

If $k \in Q^+$ then $0 < \log(q_k(0, s)) < t_k$, and

$$\int_{(\log(q_k(0, s)), +\infty)} e^{-t_k} q_k(0, s)^{-1/2} dt_k < q_k(0, s)^{-1/2},$$

while if $k \in Q^-$, we have $0 < t_k < +\infty$ and $\int_{(0,+\infty)} e^{-t_k} q_k(0, s)^{1/2} dt_k = q_k(0, s)^{1/2}$. Finally if $k \in P$, then $t_k < \log(q_k(0, s))$ and

$$\int_{(-\infty, \log(q_k(0, s)))} e^{-t_k/2} q_k(0, s)^{1/2} dt_k = 2.$$

Thus for each $s \in V_\epsilon$,

$$\begin{aligned} \int_{W_\epsilon(s)} \Psi(t, s) dt &= \\ &= \int_{W_\epsilon(s)} q_P(t, s)^{-1/2} \Psi_{Q^+}(t, s) \Psi_{Q^-}(t, s) \Psi_{\tilde{Q}}(s) \Psi_+(s) \Psi_-(s) dt \leq \\ &\leq 2^{\#(P)} \left[\prod_{k \in Q^+} q_k(0, s)^{-1/2} \right] \left[\prod_{k \in Q^-} q_k(0, s)^{1/2} \right] \times \\ &\quad \times \left[\prod_{k \in \tilde{Q}} \log(q_k(0, s)) q_k(0, s)^{1/2} \right] \Psi_+(s) \Psi_-(s). \end{aligned}$$

From the definition of the index sets Q^+ , Q^- and \tilde{Q} we now see that the function $s \rightarrow \int_{W_\epsilon(s)} \Psi(t, s) dt$ is exponentially decreasing on V_ϵ , and the integral (2.1) is convergent.

Now to show that β_l is continuous on \mathcal{H}_l^∞ , we apply the estimate of Lemma 2.4, and the above analysis, to each set U_ϵ . We have ϕ_ϵ and Ψ_ϵ as above, and Lemma 2.4 gives $\sup_{U_\epsilon} |\phi_\epsilon(t, s)| \leq \sum_\beta M_{\beta, \epsilon} \|\pi(W_{\beta, \epsilon}) f\|_{\mathcal{H}_l}$, for some constants $M_{\beta, \epsilon}$ and elements $W_{\beta, \epsilon} \in \mathbf{U}(\mathfrak{g}_c)$ independent of f . Hence

$$|\beta_l(f)| \leq \sum_\epsilon \|\Psi_\epsilon\|_1 \sup_{U_\epsilon} |\phi_\epsilon(t, s)| \leq \sum_\epsilon \sum_\beta \|\Psi_\epsilon\|_1 M_{\beta, \epsilon} \|\pi(W_{\beta, \epsilon}) f\|_{\mathcal{H}_l}.$$

By definition of the topology on \mathcal{H}_l^∞ , this finishes the proof. □

3. The Plancherel formula.

Now that we have generalized vectors $\beta_l \in (\mathcal{H}_l^{-\infty})^{q_{H, G}^{-1/2}}$, $l \in \Omega$, the results of [16] show that we have a Plancherel formula. Here we shall derive it by simple Fourier inversion. To do so we must choose some l in each H -orbit, and we want to do this in a smooth natural way. Hence the first task is to compute a nice cross-section for H -orbits in $\Omega \cap (f + \mathfrak{h}^\perp)$. In our scenario here, since $H(f + l) = f + Hl$ and H acts only by “dilations”, a cross-section is easy to find. Specifically, given $l \in \mathfrak{h}^\perp$, set $l_j = l(Z_j)$, and set $\epsilon_r(l) = \text{sign}(l_{i_r})$, $1 \leq i \leq d$. We have analytic functions $Q_j(w, l)$, $1 \leq j \leq n$, such that

$$Hl = \left\{ \sum_j Q_j(w, l) Z_j^* : w = (w_1, w_2, \dots, w_d) \in ((0, +\infty))^d \right\}.$$

For each j , if $j = i_r \in \Psi$, then $Q_j(w, l) = \varepsilon_j(l) w_r$, and if $j \notin \Psi$, say $i_r < j < i_{r+1}$, then $Q_j(w, l) = \rho_j(w_1, w_2, \dots, w_r) l_j$, where ρ_j is an analytic function of the form $\rho_j(w_1, w_2, \dots, w_r) = w_1^{\alpha_1} \dots w_r^{\alpha_r}$, $\alpha_h \in \mathbb{R}$. For each w , the function $Q_j(w, \cdot)$ is H -invariant on $\Omega \cap \mathfrak{h}^\perp$. A nice cross-section Σ for H -orbits in $\Omega \cap \mathfrak{h}^\perp$ is given by simply putting $w_r = 1, 1 \leq r \leq d$, that is $\Sigma = \{l \in \Omega \cap \mathfrak{h}^\perp : |l_{i_r}| = 1, 1 \leq r \leq d\}$. The cross-section in $f + \mathfrak{h}^\perp$ is $f + \Sigma$. Fixing a choice of signs $\varepsilon = (\varepsilon_1, \varepsilon_2, \dots, \varepsilon_d), \varepsilon_r = \pm 1$, we have $\Sigma = \bigcup \Sigma_\varepsilon$ where Σ_ε is the flat variety

$$\Sigma_\varepsilon = \{l \in \Omega \cap \mathfrak{h}^\perp : l_{i_r} = \varepsilon_r\}.$$

We choose coordinates for each Σ_ε . The center of \mathfrak{g} has as a basis

$$\{Z_j : j \notin e, j \leq n\} = \{C_1, C_2, \dots, C_a\},$$

and we set $l(C_h) = \zeta_h, 1 \leq h \leq a$, and $l(X_k) = \mu_k, 1 \leq k \leq u$. Then there is a dense, open subset D_ε of $\mathbb{R}^a \times \mathbb{R}^u$ such that

$$\Sigma_\varepsilon = \left\{ \sum \zeta_h C_h^* + \sum \varepsilon_h V_r^* + \sum \mu_k X_k^* : (\zeta, \mu) \in D_\varepsilon \right\}.$$

Given a function Θ on $f + \mathfrak{h}^\perp$, we shall write

$$\begin{aligned} \int_{f+\Sigma} \Theta(l) dl &= \\ &= \sum_{\varepsilon \in \{1, -1\}^d} \int_{\mathbb{R}^a \times \mathbb{R}^u} \Theta \left(f + \sum \zeta_h C_h^* + \sum \varepsilon_r V_r^* + \sum \mu_k X_k^* \right) d\zeta d\mu. \end{aligned}$$

Before deriving the Plancherel formula we may as well compute multiplicities, which amounts to just counting the number of H -orbits in each G -orbit intersection with $\Omega \cap (f + \mathfrak{h}^\perp)$. Fix $l_0 = \sum \zeta_h C_h^* + \sum \varepsilon_r V_r^* + \sum \mu_k X_k^* \in \Sigma_\varepsilon$, and let $l = f + l_0$. Set $f(A_k) = \alpha_k, 1 \leq k \leq u$. An ordered ‘‘co-exponential’’ basis for $\mathfrak{g} \bmod \mathfrak{g}(l)$ is $\{Z_{e_1}, Z_{e_2}, \dots, Z_{e_{2d}}\}$ (where we have written $e = \{e_1 < e_2 < \dots < e_{2d}\}$) and consists of the V_r 's, the X_k 's, and the B_h 's. Using the methods of [7, 8], we find that the G -orbit of l is the set of all l' of the form

$$\begin{aligned} l' = \sum \zeta_h C_h^* + \sum \varepsilon_{s_h} w_{s_h} V_{s_h}^* + \sum y_k Y_k^* + \sum x_k X_k^* + \\ + \sum P_k(w, x_k, y_k, l') A_k^* + \sum z_h B_h^* \end{aligned}$$

where W_{s_h} runs through $(0, +\infty), 1 \leq h \leq \nu$, x_k and y_k run through $\mathbb{R}, 1 \leq k \leq u$, and z_h runs through $\mathbb{R}, 1 \leq h \leq \nu$. Recall that A_k commutes with every basis element except X_k and Y_k , that X_k commutes with every element except A_k, Y_k and possibly some of the $B_h, 1 \leq h \leq \nu$, and Y_k commutes with

every element except X_k and A_k . This is why P_k depends only on x_k, y_k , and the $w_{s_h}, 1 \leq h \leq \nu$. In fact the function $P_k(w, x_k, y_k, l')$ can be computed as

$$P_k(w, x_k, y_k, l') = \alpha_k + (\sigma_k(w) x_k y_k - \varepsilon_{r_k} \mu_k) / l' ([X_k, Y_k]),$$

where $\sigma_k(w)$ is a positive analytic function in the positive variables $w_{s_1}, w_{s_2}, \dots, w_{s_\nu}$. (The function $\rho_{j_{r_k}}$ above involves only the variables $w_{s_1}, w_{s_2}, \dots, w_{s_\nu}$ and w_{r_k} , and $\sigma_k(w_{s_1}, w_{s_2}, \dots, w_{s_\nu}) = \rho_{j_{r_k}}(w_{s_1}, w_{s_2}, \dots, w_{s_\nu}, w_{r_k})^{-1} |_{w_{r_k}=1}$). Thus

$$\begin{aligned} Gl \cap (f + \mathfrak{h}^\perp) &= f + \{l' \in \mathfrak{p}(Gl) : P_k(w, x_k, y_k, l') = \alpha_k, 1 \leq k \leq u\} = \\ &= f + \{l' \in \mathfrak{p}(Gl) : \sigma_k(w) x_k y_k = \varepsilon_{r_k} \mu_k, 1 \leq k \leq u\}. \end{aligned}$$

Now $\#((Gl \cap (f + \mathfrak{h}^\perp)) / H) = \#((Gl \cap (\mathfrak{h}^\perp) \cap (f + \Sigma)))$. If l' belongs to this intersection, we must put each $w_{s_h} = 1, 1 \leq h \leq \nu$, and so every coordinate of l' is fixed except the X_k and Y_k coordinates, where we are allowed $l'(Y_k) = \varepsilon_{r_k}(l') = \pm 1, 1 \leq k \leq u$ while $l'(X_k) = \mu'_k$ is determined by

$$\varepsilon_{r_k}(l') \mu'_k = \varepsilon_{r_k} \mu_k, 1 \leq k \leq u.$$

Hence the intersection $Gl \cap (f + \mathfrak{h}^\perp) \cap (f + \Sigma)$ consists of 2^u elements corresponding to the possible choices of signs for $\varepsilon_{r_1}(l'), \varepsilon_{r_2}(l'), \dots, \varepsilon_{r_u}(l')$, and we have proved the following.

Proposition 3.1. *For any $f \in \mathfrak{h}^*$, the representation τ_f is uniform multiplicity 2^u , where $u = \dim(\text{cent}(N) \sim N) / 2$.*

We turn now to the Plancherel formula and the intertwining operator. Let $f \in \mathfrak{h}^*, \tau = \tau_f$ the representation induces by χ_f , and let α_τ be the canonical cyclic generalized vector for τ , that is, $\alpha_\tau(\phi) = \phi(e), \phi \in \mathcal{H}_\tau^\infty$. Then for any test function $\omega \in \mathcal{D}(G), \tau(\omega) \alpha_\tau$ belongs to \mathcal{H}_τ^∞ , in fact to $C_c^\infty(G, H, \chi_f)$, and is given by the formula (cf. [16])

$$\begin{aligned} \tau(\omega) \alpha_\tau(g) &= \omega_{H,f}(g) = \\ &= \Delta_G^{-1}(g) q_{H,G}^{-1/2}(g) \int_H \omega(g^{-1}h^{-1}) \Delta_G^{-1}(h) q_{H,G}^{-1/2}(h) \chi_f(h)^{-1} dh. \end{aligned}$$

One also has

$$\langle \tau(\omega) \alpha_\tau, \alpha_\tau \rangle = \omega_{H,f}(e) = \int_H \omega(h^{-1}) \Delta_G^{-1}(h) q_{H,G}^{-1/2}(h) \chi_f(h)^{-1} dh.$$

Let $l \in f + \mathfrak{h}^\perp$, and let π_l and β_l be as in Section 2. Then $\pi_l(\omega) \beta_l \in \mathcal{H}_l^\infty$ is given by

$$\pi_l(\omega) \beta_l(g) = \int_B \omega_{H,f}(bg) \chi_l(b)^{-1} q_{B,G}^{-1/2}(bg) q_{H,G}^{1/2}(bg) \Delta_B(b) db$$

and

$$\begin{aligned} \langle \pi_l(\omega) \beta_l, \beta_l \rangle &= \int_H \int_B \omega_{H,f}(bh) \chi_l(b)^{-1} q_{B,G}^{-1/2}(b) \times \\ &\quad \times q_{H,G}^{1/2}(b) \Delta_B(b) \chi_f(h)^{-1} dbdh \\ &= \int_H \int_B \omega_{H,f}(h^{-1}bh) \chi_l(b)^{-1} q_{B,G}^{-1/2}(b) \times \\ &\quad \times q_{H,G}^{1/2}(b) \Delta_B(b) dbdh \end{aligned}$$

(cf. [16] or [17] for the computations). For any $\phi \in C_c(G, H, \chi)$ (where χ is any unitary character of H), we set

$$I_l(\phi) = \int_H \int_B \phi(h^{-1}bh) \chi_l(b)^{-1} q_{B,G}^{-1/2}(b) q_{H,G}^{1/2}(b) \Delta_B(b) dbdh$$

so that $I_l(\omega_{H,f}) = \langle \pi_l(\omega) \beta_l, \beta_l \rangle$ when $l \in f + \mathfrak{h}^\perp$.

Theorem 3.2. *Let χ be any unitary character of H , and let $\phi \in C_c(G, H, \chi)$. Then the integral $\int_{f+\Sigma} I_l(\phi) |R(l)| dl$ is independent of the choice of $f \in \mathfrak{h}^*$ and we have*

$$\phi(e) = \int_{f+\Sigma} I_l(\phi) |R(l)| dl$$

where $R(l) = ((2\pi)^n l([X_1, Y_1]) l([X_2, Y_2]) \dots l([X_u, Y_u]))^{-1}$. In particular

$$\langle \tau(\omega) \alpha_\tau, \alpha_\tau \rangle = \int_{f+\Sigma} \langle \pi_l(\omega) \beta_l, \beta_l \rangle |R(l)| dl.$$

Proof. Let $f \in \mathfrak{h}^*$ and $l \in f + \Sigma$, writing $l = f + \sum \zeta_h C_h^* + \sum \varepsilon_\tau V_\tau^* + \sum \mu_k X_k^* \in f + \Sigma_\varepsilon$ as above. We use the following coordinates on $B(l)$: an element

$$\begin{aligned} b &= \prod \exp(c_h C_h) \prod \exp(z_h V_{s_h}) \prod \exp(y_k Y_k) \times \\ &\quad \times \prod \exp(w_k (A_k - a_k(l)) X_k) \end{aligned}$$

is identified with $(c, z, y, w) \in \mathbb{R}^a \times \mathbb{R}^\nu \times \mathbb{R}^u \times \mathbb{R}^u$. Recalling the formula for $a_k(l)$ we have $a_k(l) = \varepsilon_{\tau_k} / l([X_k, Y_k])$ in this case, and so $|R(l)| = (2\pi)^{-n} \prod |a_k(l)|$. We denote $\prod e^{t_k}$ by e^t , $\prod e^{s_k}$ by e^s , $(e^{s_1} z_1, e^{s_2} z_2, \dots, e^{s_\nu} z_\nu)$ by $e^s z$, etc. We compute explicitly that

$$q_{B,G}^{-1/2}(b) q_{H,G}^{1/2}(b) \Delta_B(b) = e^{w_1/2} e^{w_2/2} \dots e^{w_u/2} = e^{w/2}.$$

We have

$$\begin{aligned} I_l(\phi) &= \iint_H \iiint_B \phi\left(c, e^{-s}z, e^{-t}y, \text{Ad}^{-1}(t, s)w\right) e^{i[\sum w_k(a_k(l)\mu_k - \alpha_k)]} \\ &\quad \times e^{-i[\sum y_k \varepsilon_{r_k} + \sum z_h \varepsilon_{s_h} + \sum c_k \zeta_k]} e^{w/2} dcdzdydwdsdt \\ &= \iint_H \iiint_B \phi\left(c, z, y, \text{Ad}^{-1}(t, s)w\right) e^{i[\sum w_k(a_k(l)\mu_k - \alpha_k)]} \\ &\quad \times e^{-i[\sum y_k e^{t k} \varepsilon_{r_k} + \sum z_h e^{s h} \varepsilon_{s_h} + \sum c_k \zeta_k]} e^t e^s e^{w/2} dcdzdydwdsdt. \end{aligned}$$

It is easy to check that for $t, s, c, z,$ and y fixed, the function

$$w \rightarrow \phi\left(c, z, y, \text{Ad}^{-1}(t, s)w\right) e^{w/2}$$

is rapidly decreasing. By the change of variables $\mu_k \rightarrow a_k(l)(\mu_k + \alpha_k), 1 \leq k \leq u,$ and Fourier inversion in the variables $w, \mu,$ we get

$$\begin{aligned} (2\pi)^{n-u} \int_{f+\Sigma} I_l(\phi) |R(l)| dl &= \sum_{\varepsilon} \int_{\mathbb{R}^a} \int_{\mathbb{R}^u} \int_{\mathbb{R}^\nu} \int_{\mathbb{R}^u} \int_{\mathbb{R}^\nu} \int_{\mathbb{R}^a} \phi(c, z, y, 0) \\ &\quad \times e^{-i[\sum y_k e^{t k} \varepsilon_{r_k} + \sum z_h e^{s h} \varepsilon_{s_h} + \sum c_k \zeta_k]} e^t e^s dcdzdydsdtd\zeta. \end{aligned}$$

Note that the above is independent of the choice of $f \in \mathfrak{h}^*.$ Set $\nu_k = e^{t k}, 1 \leq k \leq u, \rho_h = e^{s h}, 1 \leq h \leq \nu,$ and set $I = (0, +\infty).$ Note that $a + \nu + u = n - u.$ A simple computation gives

$$\begin{aligned} \int_{f+\Sigma} I_l(\phi) |R(l)| dl &= (2\pi)^{u-n} \sum_{\varepsilon} \int_{\mathbb{R}^a} \int_{\mathbb{R}^u} \int_{\mathbb{R}^\nu} \int_{\mathbb{R}^a} \int_{I^\nu} \int_{I^u} \phi(c, z, y, 0) \\ &\quad \times e^{-i[\sum y_k \nu_k \varepsilon_{r_k} + \sum z_h \rho_h \varepsilon_{s_h} + \sum c_k \zeta_k]} d\nu d\rho dcdzdyd\zeta \\ &= (2\pi)^{u-n} \int_{\mathbb{R}^a} \int_{\mathbb{R}^u} \int_{\mathbb{R}^\nu} \int_{\mathbb{R}^a} \int_{\mathbb{R}^\nu} \int_{\mathbb{R}^u} \phi(c, z, y, 0) \\ &\quad \times e^{-i[\sum y_k \gamma_k + \sum z_h \xi_h + \sum c_k \zeta_k]} d\gamma d\xi dcdzdyd\zeta \\ &= \phi(c). \end{aligned}$$

□

Define $T : C_c(G, H, \chi_f) \rightarrow \int_{f+\Sigma} \mathcal{H}_l |R(l)| dl$ by

$$T(\omega_{H,f}) = \{\pi_l(\omega) \beta_l(g)\}_{l \in f+\Sigma}.$$

By [16, Prop 3.2], T extends to an intertwining operator

$$L^2(G, H, \chi_f) \rightarrow \int_{f+\Sigma} \mathcal{H}_l |R(l)| dl.$$

Explicitly,

$$\begin{aligned}
 T(\omega_{H,f})_l(g) &= \pi_l(\omega) \beta_l(g) \\
 &= \iiint\limits_B \phi((c, z, y, w,) g) e^{i[\sum w_k(a_k(l)\mu_k - \alpha_k)]} \\
 &\quad \times e^{-i[\sum y_k \varepsilon_{r_k} + \sum z_k \varepsilon_{s_k} + \sum c_k \zeta_k]} e^{w/2} dc dz dy dw \\
 &\quad \times q_{B,G}^{-1/2}(g) q_{H,G}^{1/2}(g).
 \end{aligned}$$

Identifying \mathcal{H}_l with $L^2(\mathbb{R}^u \times \mathbb{R}^v)$ via the mapping $\gamma(x, s), l \in f + \Sigma$, it is clear that for any $\Phi \in L^2(\mathbb{R}^u \times \mathbb{R}^v)$, the function

$$l \rightarrow \langle \pi_l(\omega) \beta_l, \Phi \rangle$$

is C^∞ , so $\{\pi_l(\omega) \beta_l(g)\}_{l \in f + \Sigma}$ is a smooth section of $\{\mathcal{H}_l\}_{l \in f + \Sigma}$.

4. Examples.

We provide two examples. The first will be the basic split oscillator group, wherein most of the essential difficulties of the subject are already exhibited. In the second example we let \mathbb{R} act semi-simply on the split oscillator so as to be non-trivial on the center, in order to show the differences created by a non-trivial action of H on the center of the nilradical. By the results of Section 1, in some sense the general case just amounts to taking higher dimensional analogues of these examples (or of the $ax + b$ group), where the commutator $[n, n]$ is allowed to be arbitrarily large, and where the portion of H acting non-trivially on $\text{cent}(\mathfrak{n})$ can act on \mathfrak{n} in a fairly arbitrary manner. In an attempt to make the techniques of Section 2 more transparent, we have related results of that section to computations in these examples.

1. $\mathfrak{g} = \text{span}\{A, X, Y, Z\}$ with non-vanishing brackets $[A, X] = -X, [A, Y] = Y, [X, Y] = Z$. Here $\mathfrak{h} = \mathbb{R}A, \mathfrak{n} = \text{span}\{X, Y, Z\}$. G is the semi-direct product of the 3-dimensional Heisenberg group with \mathbb{R} , and is diffeomorphic with \mathbb{R}^4 by identifying (z, y, x, t) with $\exp(zZ) \exp(yY) \exp(xX) \exp(tA)$. The multiplication is

$$(z, y, x, t)(z', y', x', t') = (z + z' + e^t xy', y + e^t y', x + e^{-t} x', t + t').$$

The Jordan-Holder sequence is given by $\mathfrak{g}_1 = \mathbb{R}Z, \mathfrak{g}_2 = \text{span}\{Z, Y\}, \mathfrak{g}_3 = \text{span}\{Z, Y, X\}, \mathfrak{g}_4 = \mathfrak{g}$. For $l \in \mathfrak{g}^*$, write $l = (\lambda, \gamma, \mu, \alpha)$, where $\lambda = l(Z), \gamma = l(Y), \mu = l(X), \alpha = l(A)$. $\Omega_0 = \{l \in \mathfrak{g}^* : \lambda \neq 0\}$ and $\Omega_1 = \{l \in \mathfrak{g}^* : \gamma \neq 0\}$. Fix $l \in \Omega = \Omega_0 \cap \Omega_1$; then $\mathfrak{b}(l) = \text{span}\{A_l, Y, Z\}$ where

$A_l = A - a(l)X, a(l) = \gamma/\lambda$. We have

$$B(l) = \{ \exp(zZ) \exp(yY) \exp(wA_l) \\ = (z, y, a(l)(e^{-w} - 1)w) : w \in \mathbb{R}, y \in \mathbb{R}, z \in \mathbb{R} \}.$$

Haar measure on G is $dg = dzdydxdt$, on B is $db = dzdydt$, and on $H, dh = dt$. G is unimodular but $\Delta_B(z, y, a(l)(e^{-w} - 1), w) = e^{-w}$ and $q_{B,G}(z, y, x, t) = e^t$. Let $f \in \mathcal{H}_l^\infty$, and write $f(t)$ for $f|_H$. Then the generalized vector β_l is defined formally by

$$(4.1) \quad \beta_l(f) = \int_H \bar{f}(t) e^{t/2} dt.$$

How do we see that this integral converges absolutely and defines a generalized vector? Using H -coordinates, we compute that $\pi(A)f(t) = f'(t)$, $\pi(Y)f(t) = ie^t f(t)$, and $\pi(X)f(t) = e^{-t} f'(t)$. Set

$$\phi_1(t) = q_{B,G}(t) \pi(Y)f(t).$$

Since all derivatives of ϕ_1 are square integrable on H , then $\phi_1(t) \in C_0(H)$ and we have $\|\phi_1\|_\infty \leq M_1 \|\pi(V)f\|_{\mathcal{H}_l}$ (as in Lemma 2.3, part (a)). In particular $f(t)e^{t/2} = \phi_1(t)e^{-t}$ is absolutely integrable over $0 \leq t \leq +\infty$. Let $\phi_2(t) = e^{-t/2} f'(t) = q_{B,G}(t) \pi(X)f'(t)$; then all derivatives of ϕ_2 are square integrable on H so $\phi_2 \in C_0(H)$ and $\|\phi_2\|_\infty \leq M_2 \|\pi(W)f\|_{\mathcal{H}_l}$ (this is Lemma 2.3, part (b)). Hence as $t \rightarrow -\infty, f' \rightarrow 0$ faster than $e^{t/2}$, and we apply the fundamental theorem of calculus to see that f is bounded on $-\infty \leq t \leq 0$, and that $\sup\{|f(t)| : -\infty \leq t \leq 0\} \leq 2M_2 \|\pi(W)f\|_{\mathcal{H}_l} + |f(0)|$. (This is a special case of Lemma 2.4.) Now as in the proof of Theorem 2.2, it follows that (4.1) converges absolutely and

$$|\beta_l(f)| \leq \int_{(0,+\infty)} |f(t)| e^{t/2} dt + \int_{(-\infty,0)} |f(t)| e^{t/2} dt \\ \leq \sup\{|\phi_1(t)| : 0 \leq t \leq +\infty\} \int_{(0,+\infty)} e^{-t} dt + \\ + \sup\{|f(t)| : -\infty \leq t \leq 0\} \int_{(-\infty,0)} e^{t/2} dt \\ \leq M_1 \|\pi(V)f\|_{\mathcal{H}_l} + 4M_2 \|\pi(W)f\|_{\mathcal{H}_l} + 2|f(0)|.$$

The cross-section for H -orbits in $\mathfrak{h}^\perp \cap \Omega$ is

$$\Sigma = \{(\lambda, \varepsilon, \mu, 0) : \lambda \in \mathbb{R} \sim \{0\}, \mu \in \mathbb{R}, \varepsilon = \pm 1\} = \bigcup_{\varepsilon=\pm 1} \Sigma_\varepsilon$$

and Theorem 3.2 says that the Plancherel measure is given on Σ_ε by $(2\pi)^{-3} \lambda^{-1} d\lambda d\mu$. The matrix elements for $\pi_l, l \in \Sigma$, are

$$\begin{aligned} \langle \pi_l(\omega) \beta_l, \beta_l \rangle &= \int_H \iiint_B \phi \left(\exp(zZ) \exp(e^{-t}yY) \exp(\text{Ad}^{-1}(t)w) \right) \\ &\quad \times e^{i\omega a(l)\mu} e^{-i[y\varepsilon+z\lambda]} e^{w/2} dzdydwdt = \\ &= \int_H \iiint_B \phi \left(z, y, \text{Ad}^{-1}(t)w \right) \\ &\quad \times e^{i\omega\mu} e^{-i[y\varepsilon^{-t}\varepsilon+z\lambda]} e^t e^{w/2} |a(l)|^{-1} dzdydwdt \end{aligned}$$

where $\omega \in \mathcal{D}(G)$ and $\phi = \tau(\omega) \alpha_\tau$. Now $R(l) = (2\pi)^{-3} a(l)$ here. Thus

$$\begin{aligned} &\int_\Sigma \langle \pi_l(\omega) \beta_l, \beta_l \rangle |R(l)| dl \\ &= \sum_{\varepsilon=\pm 1} \iint_{\Sigma_\varepsilon} \int_H \iiint_B \phi \left(z, y, \text{Ad}^{-1}(t)w \right) e^{i\omega\mu} e^{-i[y\varepsilon^{-t}\varepsilon+z\lambda]} \\ &\quad \times e^t e^{w/2} dzdydwdt (2\pi)^{-3} d\lambda d\mu \\ &= (2\pi)^{-2} \sum_{\varepsilon=\pm 1} \iiint \phi(z, y, 0) e^{-i[y\varepsilon^t\varepsilon+z\lambda]} dy e^t dt dz d\lambda \\ &= (2\pi)^{-2} \iint \left[\int_{(-\infty, 0)} + \int_{(0, +\infty)} \right] \int \phi(z, y, 0) \\ &\quad \times e^{-i[y\gamma+z\lambda]} dy d\gamma dz d\lambda = \\ &= \phi(e) = \langle \tau(\omega) \alpha_\tau, \alpha_\tau \rangle. \end{aligned}$$

2. $\mathfrak{g} = \text{span}\{B, A, X, Y, Z\}$ where $\mathfrak{g}_4 = \text{span}\{A, X, Y, Z\}$ is the split oscillator, and $[B, X] = X, [B, Z] = Z, [B, A] = [B, Y] = 0$. G_4 is realized as above, and $G = G_4 \exp(\mathbb{R}B)$ so that the multiplication is

$$\begin{aligned} (z, y, x, t, s) (z', y', x', t', s') \\ = (z + e^s z' + e^t x y', y + e^t y', x + e^{s-t} x', t + t', s + s'). \end{aligned}$$

The Jordan-Holder sequence is the obvious extension of that which was chosen above for \mathfrak{g}_4 , the set Ω of generic linear functionals is the same as before, and for $l \in \Omega$, the polarization $\mathfrak{b}(l)$ is the same as before, but now $q_{B,G}(s, t) = e^{t-s}$. G is no longer unimodular: $\Delta_G(z, y, x, t, s) = q_{H,G}(z, y, x, t, s)^{-1} = e^{2s}$. Thus the formula for β_l now is

$$(4.2) \quad \beta_l(f) = \int_H \bar{f}(s, t) e^s e^{(t-s)/2} ds dt$$

where we have written $f(s, t)$ for $f|_H$. One computes that, for smooth vectors restricted to H , $\pi(A) = \partial_t, \pi(B) = \partial_s, \pi(X) = e^{s-t}\partial_t, \pi(Y) = ie^t$. By Lemma 2.3, for each $k, j \geq 0$, there are constants $M_{k,j}$ and $N_{k,j}$ and elements $V_{k,j}, W_{k,j}$ in $\mathbf{U}(\mathfrak{g}_c)$ such that

$$(i) \left\| e^{ks} e^{jt} e^{(t-s)/2} f(s, t) \right\|_\infty \leq M_{k,j} \|\pi(V_{k,j}) f\|_{\mathcal{H}_l}, \text{ and}$$

$$(ii) \left\| e^{ks} e^{jt} e^{(t-s)/2} e^{s-t} \partial_t f(s, t) \in C_0(H) \right\|_\infty \leq N_{k,j} \|\pi(W_{k,j}) f\|_{\mathcal{H}_l}.$$

Let $U(++) = \{(s, t) : s > 0, t > s\}$, $U(+-) = \{(s, t) : s > 0, t < s\}$, $U(-+) = \{(s, t) : s < 0, t > s\}$, $U(--) = \{(s, t) : s < 0, t < s\}$. Using (i) and (ii) and the fundamental theorem of calculus we get (Lemma 2.4)

$$(iii) \sup \{|e^{ks} f(s, t)| : (s, t) \in U(+-) \cup U(--)\} \leq M_{k,0} \|\pi(V_{k,0}) f\|_{\mathcal{H}_l} + 2N_{k,0} \|\pi(W_{k,0}) f\|, k \geq 0.$$

For each of the above four subsets U of H , we write $\bar{f}(s, t)e^s e^{(t-s)/2}$ as a product of a function ϕ for which one of the above estimates (i), (ii) or (iii) holds, and a function Ψ which is absolutely integrable over U .

$$U(++) : \phi(s, t) = f(s, t) e^{2s} e^t e^{(t-s)/2}, \Psi(s, t) = e^{-s} e^{-t},$$

$$U(+-) : \phi(s, t) = f(s, t) e^{2s}, \Psi(s, t) = e^{-s} e^{(t-s)/2},$$

$$U(-+) : \phi(s, t) = f(s, t) e^t e^{(t-s)/2}, \Psi(s, t) = e^{s-t},$$

$$U(--) : \phi(s, t) = f(s, t), \Psi(s, t) = e^s e^{(t-s)/2}.$$

Now in a manner similar to example (1) (cf. also the proof of Theorem 2.2) we see that (4.2) is absolutely convergent and defines a generalized vector for π .

The cross-section for H -orbits in $\mathfrak{h}^\perp \cap \Omega$ is

$$\Sigma = \{(\varepsilon_1, \varepsilon_2, \mu, 0, 0) : \mu \in \mathbb{R}, \varepsilon_1 = \pm 1, \varepsilon_2 = \pm 1, \} = \bigcup_\varepsilon \Sigma_\varepsilon$$

where $\varepsilon = (\varepsilon_1, \varepsilon_2)$ runs through $\{-1, 1\}^2$. Here the Plancherel measure is given on each Σ_ε by $(2\pi)^{-3} d\mu$. For $\omega \in \mathcal{D}(G)$ and $\phi = \tau(\omega) \alpha_\tau$, we have

$$\begin{aligned} \langle \pi_l(\omega) \beta_l, \beta_l \rangle &= \langle \pi_{\varepsilon_1, \varepsilon_2, \mu}(\omega) \beta_{\varepsilon_1, \varepsilon_2, \mu}, \beta_{\varepsilon_1, \varepsilon_2, \mu} \rangle \\ &= \iiint_H \iiint_B \phi(\exp(e^s z Z) \exp(e^{-t} y Y) \exp(Ad^{-1}(s, t) w)) \\ &\quad \times e^{iw(\varepsilon_2/\varepsilon_1)\mu} e^{-i[y\varepsilon_2 + z\varepsilon_1]} e^{w/2} dz dy dw dt ds, \end{aligned}$$

and a computation like that of Example (1) shows that

$$(2\pi)^{-3} \sum_\varepsilon \int \langle \pi_{\varepsilon_1, \varepsilon_2, \mu}(\omega) \beta_{\varepsilon_1, \varepsilon_2, \mu}, \beta_{\varepsilon_1, \varepsilon_2, \mu} \rangle d\mu = \phi(e) = \langle \tau(\omega) \alpha_\tau, \alpha_\tau \rangle.$$

5. Concluding Remark.

Suppose that G is any connected, simply connected nilpotent Lie group, H a closed connected subgroup, and τ an associated finite multiplicity monomial representation. Given choices of appropriate basis for \mathfrak{g} and \mathfrak{h} , there is a unique construction of a flat cross-section Σ for generic H -orbits in \mathfrak{h}^\perp . Attaching the Vergne polarizations to each $l \in \Sigma$ one has a natural, explicit algorithm for deriving Fujiwara's Plancherel formula, and hence for constructing an explicit, smooth decomposition of τ over Σ , as described in the introduction. Can one give an explicit description of the Plancherel measure on Σ ? One could even hope that this decomposition diagonalizes the differential operators on $H \backslash G$ that commute with τ , in the manner of [5]. Of course one could also entertain such questions for G completely solvable, once the technical difficulties surrounding the construction of the generalized vectors β_l are overcome.

References

- [1] C. Benson, J. Jenkins and G. Ratcliff, *On Gelfand pairs associated with solvable Lie groups*, Trans. Amer. Math. Soc., **321** (1990), 85-116
- [2] C. Benson, J. Jenkins and G. Ratcliff, *Bounded K -spherical functions on Heisenberg groups*, J. Funct. Anal., **105** (1992), 409-443
- [3] P. Bernat et al., *Représentations des groupes de Lie résolubles*, Dunod, Paris, 1972
- [4] G. Carcano, *A commutativity condition for algebras of invariant functions*, Bull. U.M.I., **7** (1987), 1091-1105
- [5] L. Corwin and F. Greenleaf, *Commutativity of invariant differential operators on nilpotent homogeneous spaces $K \backslash G$ with finite multiplicity*, Comm. Pure Appl. Math., Vol. XLV, No. 6(1992), 681-748
- [6] L. Corwin, F. Greenleaf and G. Grelaud, *Direct integral decompositions and multiplicities for induced representations of nilpotent Lie groups*, Trans. Amer. Math. Soc., **304** (1987), 549-583
- [7] B. Currey and R. Penney, *The structure of the space of coadjoint orbits of a completely solvable Lie group*, Michigan Math. Jour., **36** (1989), 309-320
- [8] B. Currey, *The structure of the space of coadjoint orbits of an exponential solvable Lie group*, Trans. Amer. Math. Soc., **332** (1992), 241-269
- [9] H. Fujiwara, *Représentations monomiales d'un groupe de Lie nilpotent*, Pacific J. Math., **127** (1987), 329-352
- [10] H. Fujiwara, *Représentations monomiales des groupes de Lie résolubles exponentiels*, Progress in Math., **82** (1990), 61-84
- [11] H. Fujiwara and S. Yamagami, *Certaines représentations monomiales d'un groupe de Lie résoluble exponentiel*, Adv. Stud. Pure Math., **14** (1988), 153-190
- [12] A. Kaplan and F. Ricci, *Harmonic analysis on groups of Heisenberg type*, Lecture Notes in Math., **992** (1983), 416-435

- [13] R. Lipsman, *Harmonic analysis on exponential solvable homogeneous spaces: the algebraic or symmetric cases*, Pacific J. Math., **140**, No. 1 (1989), 117-147
- [14] R. Lipsman, *Induced representations of completely solvable Lie groups*, Ann Scuola Norm. Sup. Pisa, **17** (1990), 127-165
- [15] R. Lipsman, *The multiplicity function on exponential and completely solvable homogeneous spaces*, Geom. Dedicata, **39** (1991), 155-161
- [16] R. Lipsman, *The Penney-Fujiwara Plancherel formula for abelian symmetric spaces and completely solvable homogeneous spaces*, Pacific J. Math., **151**, No. 2 (1991), 265-295
- [17] R. Lipsman, *The Penney-Fujiwara Plancherel formula for homogeneous spaces*, Representation Theory of Lie Groups and Lie Algebras (Fuji-Kawaguchiko, 1990), 120-139, World Sci. Pub., River Edge, NJ, 1992
- [18] R. Lipsman, *The Penney-Fujiwara Plancherel formula for symmetric spaces*, The Orbit Method in Representation Theory (Copenhagen, 1988), 135-145, Birkhäuser, Boston, 1990
- [19] R. Penney, *Abstract Plancherel Theorems and a Frobenius reciprocity theorem*, J. Funct. Anal., **18** (1975), 177-190
- [20] N. Poulsen, *On C^∞ -vectors and intertwining bilinear forms for representations of Lie groups*, J. Funct. Anal., **9** (1972), 87-120

Received April 19, 1993 and revised April 4, 1994.

SAINT LOUIS UNIVERSITY
ST. LOUIS, MO 63130

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0030-8730(1995)170:2;1-H