

*Pacific
Journal of
Mathematics*

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IN LORENTZIAN SPACE FORMS**

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Volume 246 No. 2

June 2010

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We show that spacelike S-Willmore surfaces are the only spacelike Willmore surfaces with a duality in Lorentzian space forms. We obtain a classification of S-Willmore spheres in Lorentzian conformal space forms. Such a sphere must be congruent to either a complete spacelike stationary ($\vec{H} = 0$) surface in R_1^n ; a super-Willmore sphere in S^{2m+2} ; or a polar transform of a $(j - 1)$ -isotropic complete spacelike stationary ($\vec{H} = 0$) surface in R_1^{2j+2} . We also show that all Willmore spheres in Q_1^4 are conformal to a complete spacelike stationary surface in R_1^4 .

1. Introduction

The best-known results in global study of submanifolds in space forms are about minimal and CMC spheres, for example, the classical Hopf theorem and Chern's work [1970]. Later this study was extended to Willmore 2-spheres in Möbius geometry [Bryant 1984; Ejiri 1988b; Musso 1990; Montiel 2000] and Lorentzian conformal geometry [Alías and Palmer 1996; Ma and Wang 2008]. In this paper we generalize some of these results for spacelike S-Willmore spheres in Lorentzian space forms.

Blaschke [1929] and Thomsen [1923] began the study of Willmore surfaces, and derived the dual property. Bryant [1984] rediscovered the duality theorem and used it in classifying all Willmore spheres in S^3 . Then Ejiri [1988b] found that the duality theorem does not hold for general Willmore surfaces in S^n and showed that S-Willmore surfaces are the only Willmore surfaces with a dual Willmore surfaces in S^n . He also classified all S-Willmore spheres in S^n . In the case $n = 4$, Willmore spheres are automatically S-Willmore, a fact also studied in [Montiel 2000] and [Musso 1990]. When the case $n > 4$, Ma [2006] showed that there exists another category of transforms, the adjoint transforms, that generalize dual surfaces. By

MSC2000: primary 53A30; secondary 53B30.

Keywords: spacelike S-Willmore surfaces, spacelike S-Willmore sphere, stationary surfaces, duality theorem.

Supported by Program for Young Excellent Talents in Tongji University and the Tianyuan Foundation of China, grant 10926112.

using adjoint transforms, he obtained partial results concerning the classification of Willmore 2-sphere in S^n [Ma 2005].

In the Lorentzian case, Alías and Palmer [1996] showed that the duality theorem holds for spacelike Willmore surfaces in 3-dimensional Lorentzian space forms, and they proved that the only spacelike Willmore 2-sphere is the round sphere. In higher-dimensional Lorentzian space forms, the duality theorem does not hold in general either [Ma and Wang 2008]. While we can define the so-called polar transforms for surfaces in 4-dimensional Lorentzian space forms, such transforms keep the Willmore property. We also obtained the classification theorem:

Theorem A [Ma and Wang 2008]. *Any spacelike Willmore 2-sphere in Q_1^4 is either congruent to a complete spacelike stationary surface (that is, $H = 0$) in R_1^4 , or a polar surface of such a surface.*

Theorem B [Ejiri 1988b]. *Every S-Willmore 2-sphere $y : S^2 \rightarrow S^n$ must be conformally equivalent to a super S-Willmore surface or a complete minimal surface in R^n with its dual surface reducing to a point.*

For the Lorentzian case, we obtain the duality theorem too. Then we generalize the notion of polar surfaces and derive the classifications of S-Willmore spheres:

Theorem C. *Suppose $y : S^2 \rightarrow Q_1^n$ is a spacelike S-Willmore 2-sphere. Then up to conformal equivalence, it is described in one of the following:*

- (i) *y is congruent to a complete spacelike stationary surface (that is, $\vec{H} = 0$) in R_1^n , and its dual surface reduces to a point.*
- (ii) *y is a super S-Willmore surface in some Riemannian sphere $S^{2m} \subset Q_1^n$.*
- (iii) *y , together with its dual surface, is a polar surface of an $(m - 1)$ -isotropic complete spacelike stationary surface in R_1^{2m} , and all are in some subspace $Q_1^{2m} \subset Q_1^n$.*

Here the space Q_1^n is the conformal compactification of Lorentzian space forms. Theorem C is proved by constructing holomorphic forms on S^2 , which must vanish. Then it reduces to three cases. The first two cases are the same as in S^n . The third is that it is a full totally isotropic sphere in $Q_1^{2m} \subset Q_1^n$. There are two difficulties in the proof. The first is to show that the holomorphic forms are globally defined in S^2 . The second is to describe the totally isotropic sphere in Q_1^{2m} . Another problem is the existence of surfaces in the case (iii). When $n = 4$, we prove that it must be a round sphere. When $n > 4$, in R_1^6 we give some examples of surfaces belonging to case (i) and (iii). As to surfaces belonging to case (iii), we still know little about them.

This paper is organized as follows. In Section 2, we review the notion of Q_1^n , the general theory of spacelike surfaces, the description of Willmore and S-Willmore

surfaces, and the duality theorem for S-Willmore surfaces. In Section 3, we construct holomorphic forms on Willmore 2-spheres and reduce the problem to describing the totally isotropic sphere in Q_1^{2m} . In Section 4, we derive the notion of polar surfaces, showing that a totally isotropic sphere in Q_1^{2m} is either congruent to a totally isotropic minimal surface in \mathbb{R}_1^{2m} or to a polar surface of $(m - 1)$ -isotropic minimal surfaces in \mathbb{R}_1^{2m} . We also exhibit an S-Willmore sphere conformal to a totally isotropic minimal surface in R_1^6 . In the last section, we show that the spheres in the second case of Theorem A must be round spheres.

In the sequel $y : M \rightarrow Q_1^n$ will always denote a smooth spacelike immersion from an oriented surface M unless it is explicitly stated otherwise.

2. Spacelike surfaces in Lorentzian conformal geometry

Here we will review the conformal geometry of Q_1^n and the spacelike surface theory in Q_1^n . For further details, see [Alías and Palmer 1996; Ma and Wang 2008]. Our treatment here follows the surface theory in [Burstall et al. 2002].

Let \mathbb{R}_2^{n+2} be the space \mathbb{R}^{n+2} equipped with the quadric form

$$\langle x, x \rangle = \sum_1^n x_i^2 - x_{n+1}^2 - x_{n+2}^2,$$

with C_1^{n+1} its light cone. The projective light cone

$$Q_1^n = \{[x] \in \mathbb{R}P^{n+1} \mid x \in C_1^{n+1} \setminus \{0\}\}$$

is equipped with a Lorentzian metric induced from the projection $S^{n-1} \times S^1 \subset C_1^{n+1} \setminus \{0\} \rightarrow Q_1^n$ together with a Lorentzian metric $g(S^{n-1}) \oplus (-g(S^1))$, where $g(S^{n-1})$ and $g(S^1)$ are standard metrics on S^{n-1} and S^1 . So there is a conformal Lorentzian metric $[h]$ on Q_1^n . The conformal group of $(Q_1^n, [h])$ is exactly the orthogonal group $O(n, 2)/\{\pm 1\}$ of \mathbb{R}_2^{n+2} , acting on Q_1^n by

$$T([x]) = [xT] \quad \text{for } T \in O(n, 2).$$

The three n -dimensional Lorentzian space forms with constant sectional curvature $c = 0, +1, -1$ can be conformally embedded as a proper subset of Q_1^n . For more details about Lorentzian space forms, see [O'Neill 1983].

Now we focus on a spacelike surface $y : M \rightarrow Q_1^n$, with M a Riemann surface. Let $U \subset M$ be an open subset. A local lift of y is a map $Y : U \rightarrow C_1^{n+1} \setminus \{0\}$ such that $\pi \circ Y = y$. Two different local lifts differ by a scaling, so their induced metrics are conformal. We call y a conformal spacelike surface if the induced metric is conformal, that is, $\langle Y_z, Y_z \rangle = 0$ and $\langle Y_z, Y_{\bar{z}} \rangle > 0$ for any local lift Y and any complex coordinate z on M . Then there is a decomposition $M \times \mathbb{R}_2^{n+2} = V \oplus V^\perp$, where $V = \text{span}\{Y, dY, Y_{z\bar{z}}\}$ is a Lorentzian rank-4 subbundle independent of the choice

of Y and z . V^\perp is also a Lorentzian subbundle, called the conformal normal bundle of y in Q_1^n . Their complexifications are denoted respectively by $V_{\mathbb{C}}$ and $V_{\mathbb{C}}^\perp$.

Fix a local coordinate z . There is a local lift Y satisfying $|dY|^2 = |dz|^2$, called the canonical lift (with respect to z). Choose a frame $\{Y, Y_z, Y_{\bar{z}}, N\}$ of $V_{\mathbb{C}}$, where $N \in \Gamma(V)$ is uniquely determined by

$$(1) \quad \langle N, Y_z \rangle = \langle N, Y_{\bar{z}} \rangle = \langle N, N \rangle = 0 \quad \text{and} \quad \langle N, Y \rangle = -1.$$

Definition 2.1 [Bryant 1984; Ejiri 1988b; Alías and Palmer 1996; Ma and Wang 2008]. For a conformally immersed surface $y : M \rightarrow Q_1^n$ with canonical lift Y (with respect to a local coordinate z), we define

$$G := Y \wedge Y_u \wedge Y_v \wedge N = -2i \cdot Y \wedge Y_z \wedge Y_{\bar{z}} \wedge N, \quad \text{where } z = u + iv.$$

Here $N \equiv 2Y_{z\bar{z}} \pmod{Y}$ is the frame vector determined in (1). Note that $\langle G, G \rangle = 1$ and that G is well defined. We call $G : M \rightarrow G_{3,1}(\mathbb{R}_2^{n+2})$ the conformal Gauss map of y .

Given frames as above, and noting that Y_{zz} is orthogonal to Y, Y_z and $Y_{\bar{z}}$, there must be a complex function s and a section $\kappa \in \Gamma(V_{\mathbb{C}}^\perp)$ such that

$$Y_{zz} = -\frac{1}{2}sY + \kappa.$$

This defines two basic invariants κ and s depending on coordinates z , the conformal Hopf differential and the Schwarzian of y ; for discussion, see [Burstall et al. 2002; Ma 2005]. Let D denote the normal connection and $\psi \in \Gamma(V_{\mathbb{C}}^\perp)$ any section of the normal bundle. Then the structure equations can be given as

$$(2) \quad \begin{aligned} Y_{zz} &= -\frac{1}{2}sY + \kappa, \\ Y_{z\bar{z}} &= -\langle \kappa, \bar{\kappa} \rangle Y + \frac{1}{2}N, \\ N_z &= -2\langle \kappa, \bar{\kappa} \rangle Y_z - sY_{\bar{z}} + 2D_{\bar{z}}\kappa, \\ \psi_z &= D_z\psi + 2\langle \psi, D_{\bar{z}}\kappa \rangle Y - 2\langle \psi, \kappa \rangle Y_{\bar{z}}. \end{aligned}$$

The conformal Gauss, Codazzi and Ricci equations as integrable conditions are

$$(3) \quad \begin{aligned} s_{\bar{z}} &= 6\langle \kappa, D_z\bar{\kappa} \rangle + 2\langle D_z\kappa, \bar{\kappa} \rangle, \\ 0 &= \text{Im}(D_{\bar{z}}D_{\bar{z}}\kappa + \frac{1}{2}\bar{s}\kappa), \\ R_{\bar{z}z}^D &= D_{\bar{z}}D_z\psi - D_zD_{\bar{z}}\psi = 2\langle \psi, \kappa \rangle \bar{\kappa} - 2\langle \psi, \bar{\kappa} \rangle \kappa. \end{aligned}$$

These are parallel to the theory in [Burstall et al. 2002]. The conformal Hopf differentials play an important role in the research of Willmore surfaces. To see this, we first give the transformation formula of κ . For another complex coordinate w , $Y_1 = Y \cdot |dw/dz|$ is a canonical lift with respect to w . So the corresponding Hopf

differential κ_1 with respect to (Y_1, w) is

$$(4) \quad \kappa_1 = \kappa \cdot \left(\frac{dz}{dw}\right)^2 \cdot \left|\frac{dz}{dw}\right|^{-1}.$$

Direct computation using (2) and (4) shows that G induces a conformal-invariant metric

$$g := \frac{1}{4} \langle dG, dG \rangle = \langle \kappa, \bar{\kappa} \rangle |dz|^2$$

on M . This metric might be positive definite, negative definite, or degenerate, depending on the sign of $\langle \kappa, \bar{\kappa} \rangle$. We define the Willmore functional and Willmore surfaces by using this metric.

Definition 2.2. *The Willmore functional of y is defined as the area of M with respect to the metric above:*

$$W(y) := 2i \int_M \langle \kappa, \bar{\kappa} \rangle dz \wedge d\bar{z}.$$

An immersed surface $y : M \rightarrow Q_1^n$ is called a *Willmore surface* if it is a critical surface of the Willmore functional with respect to any variation of the map $y : M \rightarrow Q_1^n$.

By (4), it is easy to see that $W(y)$ is well defined. The following characterization of Willmore surfaces is similar to the one in the codimension-1 case [Alfías and Palmer 1996] as well to the one in Möbius geometry [Bryant 1984; Burstall et al. 2002; Ejiri 1988b; Wang 1998].

Theorem 2.3. *For a conformal spacelike surface $y : M \rightarrow Q_1^n$, the following four conditions are equivalent:*

- (i) y is Willmore.
- (ii) The conformal Gauss map G is a harmonic map into $G_{3,1}(\mathbb{R}_2^{n+2})$.
- (iii) The conformal Hopf differential κ of y satisfies the Willmore condition

$$(5) \quad D_{\bar{z}} D_{\bar{z}} \kappa + \frac{1}{2} \bar{s} \kappa = 0,$$

which is a stronger condition than the conformal Codazzi equation (3).

The proof of this theorem is completely the same as in Möbius geometry (we refer readers to [Ma 2005; Wang 1998]).

Now we introduce the so-called *S-Willmore surfaces*, the Willmore surfaces with dual surfaces (see the next section). See [Ejiri 1988b] and [Ma 2005] for their counterparts in Möbius geometry.

Definition 2.4. A conformal Willmore surface $y : M \rightarrow Q_1^n$ is called a *S-Willmore surface* if it satisfies $D_{\bar{z}}\kappa \parallel \kappa$, that is,

$$(6) \quad D_{\bar{z}}\kappa = -\frac{1}{2}\bar{\mu}\kappa$$

for some local function μ when $\kappa \neq 0$ (that is, when κ is nonumbilical). Here the conformal Hopf differential κ is defined in (2).

Note that a conformal surface satisfying $D_{\bar{z}}\kappa \parallel \kappa$ is not necessary Willmore; see [Ma and Wang 2008].

Theorem 2.5. Any spacelike S-Willmore surface $[Y] : M \rightarrow Q_1^n$ has a dual S-Willmore surface $[\hat{Y}]$, meaning they share the same central sphere and the same complex coordinates (and hence the same normal bundle and normal connection). The dual surface $[\hat{Y}]$ degenerates to a point if and only if $[Y]$ is congruent to a spacelike stationary ($H = 0$) surface in R_1^4 . Conversely, a spacelike Willmore surface in Q_1^n with a dual Willmore surface must be an S-Willmore surface.

Proof. If $[Y]$ is totally umbilical, the proof is trivial. If $[Y]$ is not totally umbilical, by using that $[Y]$ is real analytic (see [Ejiri 1988b; Alías and Palmer 1996]), we see that $M_0 = M \setminus \{\text{umbilical points}\}$ is an open dense subset of M .

For an S-Willmore surface $y : M \rightarrow Q_1^n$ with canonical lift Y (with respect to a local coordinate z), suppose that $D_{\bar{z}}\kappa = -\frac{1}{2}\bar{\mu}\kappa$ for some local function μ . Set

$$(7) \quad \hat{Y} = N + \bar{\mu}Y_z + \mu Y_{\bar{z}} + \frac{1}{2}|\mu|^2 Y.$$

We will show it is just the dual S-Willmore surface of y .

In fact, substituting $D_{\bar{z}}\kappa = -\frac{1}{2}\bar{\mu}\kappa$ into the Willmore equation (5) shows that

$$(8) \quad \mu_z - \frac{1}{2}\mu^2 - s = 0,$$

which leads to

$$(9) \quad \hat{Y}_z = \frac{1}{2}\mu\hat{Y} + \rho(Y_z + \frac{1}{2}\mu Y), \quad \text{where } \rho = \bar{\mu}_z - 2\langle \kappa, \bar{\kappa} \rangle.$$

So

$$\begin{aligned} \hat{Y}_{z\bar{z}} &= (\frac{1}{2}\mu\hat{Y} + \rho(Y_z + \frac{1}{2}\mu Y))_{\bar{z}} = (\dots)\hat{Y} + (\dots)\hat{Y}_z + (\dots)\hat{Y}_{\bar{z}} + (\dots)Y, \\ \text{span}_{\mathbb{C}}\{Y, Y_z, Y_{\bar{z}}, Y_{z\bar{z}}\} &= \text{span}_{\mathbb{C}}\{\hat{Y}, \hat{Y}_z, \hat{Y}_{\bar{z}}, \hat{Y}_{z\bar{z}}\}. \end{aligned}$$

Since y and $[\hat{Y}]$ have common central sphere, complex coordinates, and harmonic conformal Gauss map, $[\hat{Y}]$ is a Willmore surface by Theorem 2.3.

From the definition of ρ , using (3) and (8), we see that

$$(10) \quad \rho_{\bar{z}} = \bar{\mu}_{z\bar{z}} - 2\langle D_{\bar{z}}\kappa, \bar{\kappa} \rangle - 2\langle \kappa, D_{\bar{z}}\bar{\kappa} \rangle = \bar{\mu}\rho.$$

So either ρ must have isolated zero points or $\rho \equiv 0$; see [Chern 1970, Section 4]. If $\rho \equiv 0$, then \hat{Y} reduces to a point. If $\rho \neq 0$, then $\tilde{Y} = \hat{Y}/|\rho|$ is a canonical lift with respect to z . Direct computation leads to

$$\tilde{\kappa} = \tilde{Y}_{z\bar{z}} \text{ mod } \tilde{Y} = (\rho/|\rho|)\kappa.$$

With the same normal bundle as y it is easy to see that $\tilde{D}_z = D_z$, where \tilde{D}_z is the normal connection of \tilde{Y} . So $\tilde{D}_z\tilde{\kappa} \parallel \tilde{\kappa}$, showing that $[\hat{Y}]$ is S-Willmore.

Then Proposition 2.6 finishes the proof on M_0 .

Now consider the remaining points. The function μ is analytic satisfying (8) except at the umbilical points. So for any $p \in M \setminus \{M_0\}$, either μ extends to a finite number on p or goes to infinity on p . For the first case, we still define \hat{Y} by (7); for the second case, we define \hat{Y} by

$$[\hat{Y}_p] = \lim_{\mu \rightarrow \infty} \left[\frac{1}{|\mu|^2} (N + \bar{\mu}Y_z + \mu Y_{\bar{z}} + \frac{1}{2}|\mu|^2 Y) \right] = [Y_p].$$

So we extend the definition of \hat{Y} continuously to the umbilical points, finishing the proof of our theorem. □

Proposition 2.6. *If \hat{Y} reduces to a point, then $y = [Y]$ is congruent to a complete spacelike stationary surface (that is, $\vec{H} = 0$) in R^n , and its dual surface reduces to a point.*

Proof. If \hat{Y} reduces to a point, by applying a transformation $T \in O(n, 2)$ if necessary, we can set

$$\hat{Y} = (1, 0, \dots, 0, 1) \quad \text{and} \quad Y = (\frac{1}{2}(-1 + \langle u, u \rangle), u, \frac{1}{2}(1 + \langle u, u \rangle)),$$

where $u : U \rightarrow R^n$. Let z be an arbitrary complex coordinate. Then we have

$$Y_{z\bar{z}} = aY + \langle Y_z, Y_{\bar{z}} \rangle N \quad \text{and} \quad \hat{Y} = N + \bar{\mu}Y_z + \mu Y_{\bar{z}} + \langle Y_z, Y_{\bar{z}} \rangle |\mu|^2 Y,$$

where a, μ are two functions. It is easy to see that $\hat{Y}_z = -\mu \langle Y_z, Y_{\bar{z}} \rangle \hat{Y} + \dots$. So $\mu \equiv 0$ and $Y_{z\bar{z}} = aY + \langle Y_z, Y_{\bar{z}} \rangle \hat{Y}$. Replacing by u leads to

$$(\langle u_{z\bar{z}}, u \rangle, u_{z\bar{z}}, \langle u_{z\bar{z}}, u \rangle) = (\frac{1}{2}(-a + a \langle u, u \rangle), au, \frac{1}{2}(a + a \langle u, u \rangle)).$$

This implies $a \equiv 0$ and $u_{z\bar{z}} \equiv 0$. So u is a stationary surface in \mathbb{R}_1^n , and $y = [Y]$ belongs to class (i) in Theorem C. □

3. S-Willmore spheres in Q_1^n

Let $y : S^2 \rightarrow Q_1^n$ be an S-Willmore sphere with (Y, z) as above and with its dual surface \hat{Y} as in (7). First, we will introduce the notion of m -isotropic surfaces and give some properties of such surfaces. Then we define holomorphic forms for m -isotropic S-Willmore spheres by a lemma. With these vanishing forms, we

prove the first two cases of Theorem C and leave the discussion of the last case to Section 4.

Definition 3.1. An immersed spacelike surface $y : M \rightarrow Q_1^n$ is called an m -isotropic surface if it satisfies

$$\langle Y_z^{(i)}, Y_z^{(j)} \rangle = 0 \quad \text{for } i + j \leq 2m + 1.$$

Here $Y_z^{(i)}$ denotes the i -th derivative of Y with respect to z . If the equations above hold for any i and j , we call y totally isotropic.

It is direct to verify this notion is well defined, that is, independent of the choice of Y and z . Also this is a conformal-invariant property of surfaces. One may compare with [Ejiri 1988b; 1988a; Kokubu 2002; Ma 2005].

We can derive that, for an m -isotropic surface y with $m \geq 2$,

$$(11) \quad \langle (D_z)^i \kappa, (D_z)^j \kappa \rangle = 0 \quad \text{for } i + j \leq 2m - 3.$$

If $y : M \rightarrow Q_1^n$ is not contained in any $2m$ -dimensional spacelike Riemannian sphere S^{2m} of Q_1^n , we say that y is *full* in Q_1^n .

If y is m -isotropic and not $(m + 1)$ -isotropic, we say y is *strictly m -isotropic*.

Similarly to [Ejiri 1988b; Ma 2005; Ma and Wang 2008], we can define holomorphic forms for S-Willmore spheres:

Lemma 3.2. *Let $y : S^2 \rightarrow Q_1^n$ be a strictly m -isotropic S-Willmore sphere. On the subset $M_0 \subset S^2$ where y has no umbilic points, let Y be the canonical lift of y , and let \hat{Y} be a local lift of its dual Willmore surface satisfying $\langle Y, \hat{Y} \rangle = -1$. Write $\hat{\kappa} = \hat{Y}_{zz}^\perp$, that is, $\hat{\kappa} \perp \{\hat{Y}, \hat{Y}_z, \hat{Y}_{\bar{z}}, \hat{Y}_{z\bar{z}}\}$. Then the differential form*

$$(12) \quad \begin{aligned} \Omega_m dz^{4m+4} &:= \langle Y_z^{(m+1)}, Y_z^{(m+1)} \rangle \langle \hat{Y}_z^{(m+1)}, \hat{Y}_z^{(m+1)} \rangle dz^{4m+4} \\ &= \langle D_z^{(m-1)} \kappa, D_z^{(m-1)} \kappa \rangle \langle \hat{D}_z^{(m-1)} \hat{\kappa}, \hat{D}_z^{(m-1)} \hat{\kappa} \rangle dz^{4m+4} \end{aligned}$$

is a globally defined holomorphic $(4m + 4)$ -form on S^2 , and hence vanishes. Here D_z and \hat{D}_z denote the normal connections of Y and \hat{Y} .

Proof. On M_0 , the formula for $\Omega_m dz^{4m+4}$ is meaningful. We note that the normal connection \hat{D}_z is equal to D_z . With respect to another complex coordinate w , a canonical lift is $Y_1 = Y \cdot |dw/dz|$. So $\hat{Y}_1 = \hat{Y} \cdot |dz/dw|$ is the lift of dual surface and $\hat{\kappa}_1 = \hat{Y}_{1z\bar{z}}^\perp = \hat{\kappa} \cdot |dz/dw|^{-1} \cdot (dz/dw)^2$. Together with (4) and (12), $\Omega_m dz^{4m+4}$ is invariant when (Y, z) changes.

By using (9) and the fact that the normal connection \hat{D}_z of \hat{Y} is equal to D_z , we derive

$$\begin{aligned} \hat{\kappa} &= \rho \kappa, \\ D_z \hat{\kappa} &= \rho D_z \kappa + (\dots) \kappa, \\ D_z^{(i)} \hat{\kappa} &= \rho D_z^{(i)} \kappa + (\dots) D_z^{(i-1)} \kappa + \dots + (\dots) \kappa. \end{aligned}$$

By (11), we see that $\Omega_m d\bar{z}^{4m+4} = \rho^2 \langle D_z^{(m-1)} \kappa, D_z^{(m-1)} \kappa \rangle^2 d\bar{z}^{4m+4}$. To show Ω_m is holomorphic, with the Ricci equation (3), (11) and the S-Willmore condition, we have

$$\begin{aligned}
 \langle D_z^{(m-1)} \kappa, D_z^{(m-1)} \kappa \rangle_{\bar{z}} &= 2 \langle D_{\bar{z}} D_z^{(m-1)} \kappa, D_z^{(m-1)} \kappa \rangle \\
 &= 2 \langle D_z D_{\bar{z}} (D_z^{(m-2)} \kappa) - (\dots) \kappa, D_z^{(m-1)} \kappa \rangle = \dots \\
 (13) \quad &= 2 \langle D_z^{(m-1)} D_{\bar{z}}(\kappa) + (\dots) D_z^{(m-2)} \kappa + \dots + (\dots) \kappa, D_z^{(m-1)} \kappa \rangle \\
 &= -\bar{\mu} \langle D_z^{(m-1)} \kappa, D_z^{(m-1)} \kappa \rangle.
 \end{aligned}$$

Together with (10), we see that $\Omega_{m\bar{z}} \equiv 0$.

To show $\Omega_m (d\bar{z})^{4m+4}$ extends to S^2 as a holomorphic form, note that by (3) and (6),

$$D_{\bar{z}} D_z^{(j)} \kappa = \dots = D_z^{(j)} D_{\bar{z}} \kappa + \dots + (\dots) \kappa \in \text{span}\{\kappa, \dots, D_z^{(m-1)} \kappa\}.$$

By Chern’s lemma on pseudoholomorphic functions [1970], the zero points of $D_z^{(j)} \kappa$ for $1 \leq j \leq m - 1$ are isolated. So $\{\kappa, \dots, D_z^{(m-1)} \kappa\}$ forms a holomorphic vector bundle F^m on $M'_0 = M_0 \setminus \{\text{isolated zero points of } D_z^{(j)} \kappa \text{ for } 1 \leq j \leq m - 1\}$.

Let $\{\zeta, \dots, \zeta^{m-1}\}$ be a local frame consisting of F^m around an umbilical point $p \in S^2 \setminus M_0$ such that

$$\text{span}\{\kappa, \dots, D_z^{(i)} \kappa\} = \text{span}\{\zeta, \dots, \zeta^i\} \quad \text{for } i = 1, \dots, m - 1.$$

Set $g = \langle \zeta^{m-1}, \zeta^{m-1} \rangle$ and $D_z^{(m-1)} \kappa = f \zeta_i^{m-1} \text{ mod } \{\zeta, \dots, \zeta^{m-2}\}$. Then comparing with (13), we have $-\bar{\mu} = 2f_{\bar{z}}/f$. Together with (9) we have

$$\rho \langle D_z^{(m-1)} \kappa, D_z^{(m-1)} \kappa \rangle = -2 \langle \kappa, \bar{\kappa} \rangle f^2 g + \bar{\mu}_z f^2 g,$$

and

$$\bar{\mu}_z f^2 g = (\bar{\mu} f^2 g)_z - \bar{\mu} (2ff_z g + fg_z) = (-2f_{\bar{z}} f g)_z + 2f_{\bar{z}} (2f_z g + fg_z),$$

showing that Ω_m is a smooth function (depending on z). Then $\Omega_m (d\bar{z})^{4m+4}$ extends smoothly to $S^2 \setminus M_0$ as desired. It is holomorphic both on M'_0 and in the interior of $S^2 \setminus M_0$ (it vanishes in the latter case). For the same reason $\Omega_m (d\bar{z})^{4m+4}$ extends smoothly to $M'_0 \setminus M_0$. So it is holomorphic on all of S^2 . □

The method used here is suggested by Ma and it is indeed an improvement of [Ma 2005] an [Ma and Wang 2008].

From the proof of Lemma 3.2, we know that either y satisfies $\rho \equiv 0$, forcing the dual surface of y to reduce to a point, or that y is a totally isotropic surface in Q_1^n . In the case that y is totally isotropic, if y is in some Riemannian sphere $S^{2m} \subset Q_1^n$, we have this:

Proposition 3.3 [Ejiri 1988b]. *If y is totally isotropic and contained in some Riemannian sphere $S^{2m} \subset Q_1^n$, it must be a super Willmore surface in S^{2m} .*

If y is not contained in any Riemannian sphere, it must be full in some Q_1^{2m} . For this case, Proposition 4.3 gives the description of such surface; see Section 4 for details. Summing up, we finish the proof of Theorem C.

4. Polar surfaces of m -isotropic S-Willmore surfaces in Q_1^{2m+2}

Next we deal with m -isotropic S-Willmore surfaces that are full in Q_1^{2m+2} . Let $y : M \rightarrow Q_1^{2m+2}$ be such a surface. Then by (11), we can suppose that

$$\begin{aligned}
 \kappa &= E_1, \\
 D_z E_1 &= E_2 + \alpha_1 E_1 \quad \text{with } E_2 \perp (E_1, \bar{E}_1), \\
 &\vdots \\
 D_z E_{m-1} &= E_m + \alpha_{m-1} E_{m-1} + \dots + (\dots) E_1, \\
 &\quad \text{with } E_m \perp (E_1, \dots, E_{m-1}, \bar{E}_1, \dots, \bar{E}_{m-1}).
 \end{aligned}
 \tag{14}$$

So $\langle E_j, E_j \rangle = 0$ for $1 \leq j \leq m - 1$. We note that on the points that $E_i = 0$, by a similar treatment as before, E_i can be extended to the whole surface since it is defined on a open dense subset.

To give the definition of *polar surfaces*, we still need more preparation. By the S-Willmore condition,

$$D_{\bar{z}} E_1 = -\frac{1}{2} \mu E_1.$$

Together with the Ricci equation, we have

$$\begin{aligned}
 D_{\bar{z}} E_2 &= D_{\bar{z}} D_z E_1 - D_{\bar{z}}((\dots) E_1) = D_z D_{\bar{z}} E_1 \quad \text{mod } (E_1) \\
 &= -\frac{1}{2} \bar{\mu} E_2 \quad \text{mod } (E_1).
 \end{aligned}$$

Using Ricci equations repeatedly, we get

$$D_{\bar{z}} E_i = -\frac{1}{2} \bar{\mu} E_i \quad \text{mod } (E_1, \dots, E_{i-1}) \quad \text{for } 1 \leq i \leq m.
 \tag{15}$$

Then we have that, for all $1 \leq j \leq m - 1$,

$$\langle E_j, \bar{E}_j \rangle_z = \langle D_z E_j, \bar{E}_j \rangle + \langle E_j, D_z \bar{E}_j \rangle = (\alpha_j - \frac{1}{2}(\mu)) \langle E_j, \bar{E}_j \rangle,$$

Applying Chern’s lemma about almost holomorphic functions [1970], we see that either $\langle E_i, \bar{E}_i \rangle \equiv 0$ on a open dense subset M_0 of M , or $\langle E_i, \bar{E}_i \rangle$ has only isolated zero points. We will show that if $\langle E_j, \bar{E}_j \rangle \equiv 0$ on M_0 , then y must be contained in some Q_1^{2j+2} .

To see this, we need two lemmas. Let \mathbb{C}_1^n be the complexification of R_1^n with its inner product. The first one is about isotropic vector in \mathbb{C}_1^n .

Lemma 4.1. *Let E be a nonzero vector in \mathbb{C}_1^n with $\langle E, E \rangle = 0$, E must either equal to $\lambda(e_1 + ie_2)$ with $e_1, e_2 \in R_1^n$, $\langle e_1, e_2 \rangle = 0$, $\langle e_1, e_1 \rangle = \langle e_2, e_2 \rangle = 1$, and $\langle E, \bar{E} \rangle \neq 0$; or equal to $\lambda(e_1 - e_2)$ with $e_1, e_2 \in R_1^n$ and $\langle e_1, e_2 \rangle = 0$, $\langle e_1, e_1 \rangle = -\langle e_2, e_2 \rangle = 1$, and $\langle E, \bar{E} \rangle = 0$.*

The other concerns complex functions:

Lemma 4.2. *Let γ be a complex function such that $\gamma_{\bar{z}} = \bar{\gamma}_z$. There must be a real function $f = f(z, \bar{z})$ satisfying $\gamma = f_z$.*

Proof. Set $\gamma = \gamma_1 + i\gamma_2$ and $z = u + iv$, where γ_1, γ_2, u, v are real. Putting these in the formula $\gamma_{\bar{z}} = \bar{\gamma}_z$ and expanding it shows $\partial\gamma_1/\partial v = -\partial\gamma_2/\partial u$, which is just the integrability condition of the PDE

$$\frac{\partial f}{\partial u} = \gamma_1 \quad \text{and} \quad \frac{\partial f}{\partial v} = -\gamma_2. \quad \square$$

Proposition 4.3. *Let $y : M \rightarrow Q_1^{2m+2}$ be an $(m + 1)$ -isotropic S-Willmore surface that is full in Q_1^{2m+2} . Then y induces two m -isotropic S-Willmore surfaces $[\hat{L}]$ and $[\hat{R}]$ in Q_1^{2m+2} , called polar surfaces of y . One of these polar surfaces reduces to a point and the other is conformal to an m -isotropic complete spacelike stationary ($H = 0$) surface in R_1^{2m+2} . Both y and its dual surface are polar surfaces of the nondegenerate polar surface.*

Proof. Since y is $(m + 1)$ -isotropic and full in Q_1^{2m+2} , it is totally isotropic and its derivatives $Y_z^{(k)}$ give $(m + 1)$ independent isotropic vectors in \mathbb{C}_2^{m+2} . Using Lemma 4.1, we can suppose that $\langle E_j, \bar{E}_j \rangle = 0$ for some j . Then $E_j = \lambda_j(e_{j1} - e_{j2})$ for two orthogonal vector fields $e_{j1}, e_{j2} \in V$ such that $\langle e_{j1}, e_{j2} \rangle = 0$ and $\langle e_{j1}, e_{j1} \rangle = -\langle e_{j2}, e_{j2} \rangle = 1$. Write

$$L_j = \frac{1}{\sqrt{2}}(e_{j1} - e_{j2}), \quad R_j = \frac{1}{\sqrt{2}}(e_{j1} + e_{j2}) \quad (\text{which implies } E_j = \sqrt{2}\lambda_j L_j).$$

If $j = 1$, we have

$$(16) \quad \begin{cases} (L_1)_z = \alpha_{11}L_1 + \psi_{11}, \\ (R_1)_z = -\alpha_{11}R_1 + \psi_{12} - \bar{\mu}\langle R_1, E_1 \rangle Y - 2\langle R_1, E_1 \rangle Y_{\bar{z}}, \\ D_{\bar{z}}D_z\psi - D_zD_{\bar{z}}\psi = 0 \quad \text{for all } \psi \in V^\perp \text{ and } \psi \perp \{L_1, R_1\}. \end{cases}$$

Here $\psi_{11}, \psi_{12} \in V_{\mathbb{C}}^\perp$ and $\psi_{11}, \psi_{12} \perp \{L_1, R_1\}$. Since $D_{\bar{z}}\kappa \parallel \kappa$, we have $\psi_{11} = 0$ and $(L_1)_z = \alpha_{11}L_1$. So $[L_1]$ is a point in Q_1^{2m+2} . Together with the third formula in (16), we can choose new L_1 and $e_{13}, \dots, e_{1,2m}$ so that $\alpha_{11} = 0$ and $e_{13}, \dots, e_{1,2m}$ is a parallel orthonormal frame of $V^\perp \text{ mod } L_1, R_1$. So

$$\begin{aligned} (R_1)_z &= \sum_k \lambda_{1k}e_{1k} - \bar{\mu}\langle R_1, E_1 \rangle Y - 2\langle R_1, E_1 \rangle Y_{\bar{z}}, \\ (e_{1k})_z &= -\frac{1}{2}\lambda_{1k}L_1 \quad \text{for } k = 3, \dots, 2m, \text{ with } \lambda_{1k} \text{ functions.} \end{aligned}$$

The Ricci equations lead to $\lambda_{1k\bar{z}} = \bar{\lambda}_{1kz}$ for $k = 3, \dots, 2m$. Using Lemma 4.2, we see that there exist real functions f_{1k} such that $\lambda_{1k} = f_{1kz}$. Set

$$(17) \quad \hat{R}_1 = R_1 - \sum_k f_{1k} e_{1k} - \frac{1}{2}(\sum_k f_{1k})L_1.$$

Differentiating (17), we obtain

$$(\hat{R}_1)_z = -\bar{\mu}\langle R_1, E_1 \rangle Y - 2\langle R_1, E_1 \rangle Y_{\bar{z}}.$$

Thus Y is located in a 6-dimensional subspace $\mathbb{R}_2^6 \subset \mathbb{R}_2^{2m+4}$. By the condition that y is full, we have $j = m = 1$.

From [Ma and Wang 2008], we see that \hat{R}_1 is conformal to a stationary surface in \mathbb{R}_1^4 , and that $[Y]$ and its dual surface $[\hat{Y}]$ are just polar transforms of \hat{R}_1 . In fact, we see directly that L_1 is in the central sphere of \hat{R}_1 and dual surface of \hat{R}_1 . So \hat{R}_1 is a Willmore surface with dual surface reducing to a point and thus is S-Willmore.

If $j > 1$, by (15) and (14), we have $(L_j)_z = \alpha_{j1}L_j$ for some function α_{j1} . This means that $[L_j]$ is a point in Q_1^{2m+2} and we can choose new L_j and R_j such that $(L_j)_z = 0$. Differentiating $\langle R_j, E_i \rangle = 0$, we have

$$0 = \langle (R_j)_z, E_i \rangle + \langle R_j, (E_i)_z \rangle = \langle (R_j)_z, E_i \rangle - \sqrt{2}\lambda_{j,j-1}\delta_{i,j-1},$$

where $\lambda_{j,j-1}$ is some function. So we can assume that

$$(R_j)_z = \lambda_{j,j-1}E_{j-1} + \psi_{j+1,2},$$

where $\psi_{j+1,2} \in V_{\mathbb{C}}^\perp$ with $\psi_{j+1,2} \perp \{E_i, \bar{E}_i, L_j, R_j, 1 \leq i < j\}$. In summary, we have

$$(E_i)_z = \alpha_i E_i + E_{i+1} \quad \text{for } 1 \leq i < j,$$

$$(\bar{E}_1)_z = -\frac{1}{2}\mu\bar{E}_1 - 2\bar{\mu}\langle E_1, \bar{E}_1 \rangle Y - 2\langle E_1, \bar{E}_1 \rangle Y_{\bar{z}},$$

$$(\bar{E}_i)_z = -\frac{1}{2}\mu\bar{E}_i - \langle E_i, \bar{E}_i \rangle / \langle E_{i-1}, \bar{E}_{i-1} \rangle \bar{E}_{i-1} \quad \text{for } 1 < i < j,$$

$$(L_j)_z = 0,$$

$$(R_j)_z = \lambda_{j,j-1}E_{j-1} + \psi_{j+1,2},$$

$$\psi_z = D_z \psi \quad \text{for all } \psi \in V^\perp \text{ with } \psi \perp \{\text{Re}(E_i), \text{Im}(E_i), L_j, R_j, 1 \leq i < j\}.$$

That $D_{\bar{z}}D_z \psi = D_z D_{\bar{z}} \psi$ shows that we can choose $\{e_{j,2j+1}, \dots, e_{j,2m}\}$ as a parallel orthonormal frame of $V^\perp \text{ mod } \{\text{Re}(E_i), \text{Im}(E_i), L_j, R_j, 1 \leq i < j\}$. So

$$(R_j)_z = \lambda_{j,j-1}\bar{E}_{j-1} + \sum_k \lambda_{jk} e_{j,2j+k},$$

where $\lambda_{jk\bar{z}} = \bar{\lambda}_{jkz}$. Using Lemma 4.2, we see that there exist real functions f_{jk} such that $f_{jkz} = \lambda_{jk}$. Setting

$$\hat{R}_j = R_j - \sum_k f_{jk} e_{j,2j+k} - (\sum_k f_{jk})L_j.$$

shows that $(\hat{R}_j)_z = \lambda_{j,j-1} \bar{E}_{j-1}$.

So Y is located in a $(2j+4)$ -dimensional subspace $\mathbb{R}_2^{2j+4} \subset \mathbb{R}_2^{2m+4}$. By the condition that y is full, $j = m$.

From $(\hat{R}_m)_{z\bar{z}} = (\dots)L_m \text{ mod } \{\hat{R}, \hat{R}_z, \hat{R}_{\bar{z}}\}$, we see that L_m is in the central sphere of \hat{R}_m and dual surface of \hat{R}_1 . The methods of [Ma and Wang 2008] show that \hat{R}_m is conformal to a stationary surface in \mathbb{R}_1^{2m+2} . So \hat{R}_m is a S-Willmore surface with dual surface reducing to a point.

Following [Ma and Wang 2008], we call $[\hat{R}_m], [L_m] : S^2 \rightarrow Q_1^{2m+2}$ polar surfaces of $y = [Y] : S^2 \rightarrow Q_1^{2m+2}$ for all $m \geq 1$.

To see that \hat{R}_m is m -isotropic and y together with its dual surface $[\hat{Y}]$ are just its polar surfaces, let us calculate $\hat{R}_{zz}, \dots, \hat{R}_{mz}^{(m+2)}$:

$$\begin{aligned} (\hat{R}_m)_{zz} &= (\dots)\bar{E}_{m-2} \text{ mod } \{\hat{R}, \hat{R}_z\}, \\ &\vdots \\ (\hat{R}_m^{(m-1)})_z &= (\dots)\bar{E}_1 \text{ mod } \{\hat{R}, \hat{R}_z, \bar{E}_{m-1}, \dots, \bar{E}_2\}, \\ (\hat{R}_m^{(m)})_z &= (\dots)(Y_{\bar{z}} + \frac{1}{2}\bar{\mu}Y) \text{ mod } \{\hat{R}, \hat{R}_z, \bar{E}_{m-1}, \dots, \bar{E}_1\}, \\ (\hat{R}_m^{(m+1)})_z &= (\dots)(\rho Y + \hat{Y}) \text{ mod } \{\hat{R}, \hat{R}_z, \bar{E}_{m-1}, \dots, \bar{E}_1, Y_{\bar{z}} + \frac{1}{2}\bar{\mu}Y\}. \end{aligned}$$

So \hat{R}_m is m -isotropic. $[Y]$ and $[\hat{Y}] = [\hat{Y} = N + \bar{\mu}Y_z + \mu Y_{\bar{z}} + \frac{1}{2}|\mu|^2 Y]$ are polar surfaces of \hat{R}_m . □

From the proof of Proposition 4.3, we see that to define the polar surfaces we just need y to be m -isotropic and full in Q_1^{2m+2} . Let $y : M \rightarrow Q_1^{2m+2}$ be such an S-Willmore surface. On an open dense subset we have $\langle E_j, \bar{E}_j \rangle \neq 0$ for $1 \leq j \leq m-1$. So the complement of $\text{span}\{\text{Re } E_j, \text{Im } E_j\}$ in V^\perp is a Lorentzian 2-bundle. Let $\{L, R\}$ be a null basis of it, that is, $\langle L, L \rangle = \langle R, R \rangle = 0$ and $\langle L, R \rangle = -1$, with orientation restrictions $[L], [R] : \tilde{M} \rightarrow Q_1^{2m+2}$ well defined.

Proposition 4.4. *Let $y : M \rightarrow Q_1^{2m+2}$ be an m -isotropic S-Willmore surface that is full in Q_1^{2m+2} . Then y induces two m -isotropic S-Willmore surfaces $[L]$ and $[R]$ in Q_1^{2m+2} , called the polar surfaces of y (they may degenerate). Also y and its dual surface are polar surfaces of the nondegenerate polar surface.*

Remark 4.5. In the case $m = 1$, the normal bundle V^\perp is just a Lorentzian 2-bundle with two canonical lightlike directions, which enables us to define polar surfaces naturally. See [Ma and Wang 2008] for details.

Corollary 4.6. *Let y be a surface belonging to both case (i) and case (iii) in Theorem C. Then y is conformal to a totally isotropic complete spacelike stationary surface in R_1^{2m} .*

Remark 4.7. It is natural to ask that whether there are nontrivial S-Willmore spheres belonging to case (iii). For the case $m = 1$, the answer is negative; see the next section. For the case in R_1^6 we can give examples of spheres both of case (i) and case (iii). See the example below.

Example 4.8. Define

$$f : M = \mathbb{C} \setminus \{0, 1\} \rightarrow \mathbb{C}_1^6, \quad z \mapsto zv_0 + \frac{v_1}{z} + \frac{v_2}{z-1}$$

with $v_1 = (1, i, 0, 0, 0, 0)$, $v_2 = (0, 0, 1, i, 0, 0)$ and $v_0 = v_1 + v_2 + (0, 0, 0, 0, 1, 1)$. Then $y = 2 \operatorname{Re} f$ is just a totally isotropic complete spacelike stationary surface, and its compactification is a Willmore sphere with $W(y) = 8\pi$. Its polar surface l corresponds to a Willmore sphere with $W(y) = 8\pi$ and can be written as

$$\begin{aligned} l = 2 \operatorname{Re}(\tilde{l}), \quad \tilde{l} &= \frac{1}{3} \left(\left(\frac{1}{z} + \frac{1}{z-1} \right) \tilde{v}_0 + \left(z + 1 + \frac{1}{z-1} \right) \tilde{v}_1 - \left(z - 2 + \frac{1}{z} \right) \tilde{v}_2 \right) \\ &= \frac{1}{3} \left(z(\tilde{v}_1 - \tilde{v}_2) + \frac{1}{z}(\tilde{v}_0 - \tilde{v}_2) + \frac{1}{z-1}(\tilde{v}_0 + \tilde{v}_1) + (\tilde{v}_1 + 2\tilde{v}_2) \right), \end{aligned}$$

with $\tilde{v}_0 = (-1, 0, 0, 0, 0, 1)$, $\tilde{v}_1 = (0, 1, -i, 0, 0, 0)$ and $\tilde{v}_2 = (0, 0, 0, 1, -i, 0)$. So they are both spacelike stationary surfaces of the type

$$x = \operatorname{Re} \left(\frac{v_1}{z - \lambda_1} + \frac{v_1}{z - \lambda_2} + \frac{v_1}{z - \lambda_3} \right),$$

where the v_j are orthogonal isotropic vectors and the λ_j are distinct complex numbers such that the surfaces are complete.

Remark 4.9. It is desirable to have examples that fall under case (iii) and but not case (i). One way to find these follows Bryant’s methods [1984] and is used in Example 4.8. Alternatively, see [Ejiri 1993] and [Kokubu 2002] for two iterative constructions of m -isotropic minimal surfaces. Both methods require detailed technical computations, and so far we have not succeeded in using them.

5. Null-umbilic S-Willmore spheres in Q_1^4

Let $y : S^2 \rightarrow Q_1^4$ be a null-umbilic S-Willmore spheres in Q_1^4 . Recall that y is called null-umbilic if κ is a complexified lightlike section, that is, $\langle \kappa, \kappa \rangle = 0$. Then the structure equations simplify to

$$\begin{aligned} Y_{zz} &= -\frac{1}{2}sY + \lambda_1L, & N_z &= -sY_{\bar{z}} + 2\lambda_{1\bar{z}}L, & L_z &= 0, \\ Y_{z\bar{z}} &= \frac{1}{2}N, & R_z &= \lambda_{1\bar{z}}Y + 2\lambda_1Y_{\bar{z}} \end{aligned}$$

for some local function λ_1 and normal frame L, R . The Gauss equation shows that $s_{\bar{z}} = 0$. Let the complex coordinate z transform projectively. Then $s(dz)^2$ is a globally defined holomorphic 2-form on S^2 , and $s \equiv 0$. The only nontrivial invariant left is λ_1 , and the only remaining integrability condition to be satisfied is

$(\lambda_1)_{\bar{z}\bar{z}} \equiv 0$ on the whole $S^2 = \mathbb{C}\mathbb{P}^1$. It is easy to see that on \mathbb{C} the general solution to this equation takes the form $\lambda_1(z, \bar{z}) = A(z) \cdot \bar{z} + B(z)$, where $A(z)$ and $B(z)$ are holomorphic functions. With respect to the new coordinate $w = 1/z$, we find that

$$\lambda_1 = \left(A\left(\frac{1}{w}\right) \frac{1}{\bar{w}} + B\left(\frac{1}{w}\right) \right) \cdot \frac{\bar{w}}{w^3}.$$

The poles will show up unless $A(z) = B(z) \equiv 0$ and $\lambda_1 \equiv 0$ identically. Hence we have proved this:

Proposition 5.1. *Let $y : M \rightarrow Q_1^4$ be an immersed spacelike Willmore 2-sphere that is also null-umbilic. Then $y(M)$ is congruent to a round 2-sphere.*

Remark 5.2. Leitner [2000] showed that compact null-umbilic surfaces in R_1^4 , S_1^4 or H_1^4 must be topological 2-spheres and be certain deformations of round 2-spheres.

Our conclusion does not mean that there are no nontrivial compact *twistorial surfaces* in Q_1^4 . It should be understood that there exist many them, but all have branch points. Note that there are many stationary surfaces in R_1^4 with planar ends, and polar transforms of surfaces provide null-umbilic Willmore surfaces (which may have branch points).

Acknowledgments

The author is thankful to Professor Xiang Ma for valuable discussions and suggestions, and also thanks the referees for valuable suggestions.

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Received May 15, 2009. Revised September 5, 2009.

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The Pacific Journal of Mathematics (ISSN 0030-8730) at the University of California, c/o Department of Mathematics, 969 Evans Hall, Berkeley, CA 94720-3840, is published monthly except July and August. Periodical rate postage paid at Berkeley, CA 94704, and additional mailing offices. POSTMASTER: send address changes to Pacific Journal of Mathematics, P.O. Box 4163, Berkeley, CA 94704-0163.

PJM peer-review and production is managed by EditFlow™ from Mathematical Sciences Publishers.

PUBLISHED BY PACIFIC JOURNAL OF MATHEMATICS

at the University of California, Berkeley 94720-3840

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Typeset in L^AT_EX

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PACIFIC JOURNAL OF MATHEMATICS

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0030-8730(201006)246:2;1-B