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Using basic homotopy constructions, we show that isomorphism classes of string structures on spin bundles are naturally given by certain degree 3 cohomology classes, which we call string classes, on the total space of the bundle. Using a Hodge isomorphism, we then show that the harmonic representative of a string class gives rise to a canonical 3-form on the base space, refining the associated differential character. We explicitly calculate this 3-form for homogeneous metrics on 3-spheres, and we discuss how the cohomology theory tmf could potentially encode obstructions to positive Ricci curvature metrics.

1. Introduction

Degree four characteristic classes arise as obstructions in several ways in math and theoretical physics. This is analogous to the way the Stiefel–Whitney classes w_1 and w_2 encode obstructions to orientations and spin structures on a manifold M . One usually encounters the degree four classes when considering structures analogous to the spin structure, but on mapping spaces $\text{Map}(\Sigma, M)$, where Σ is a 1- or 2-dimensional manifold. It is common to say that $\frac{1}{2}p_1(M) = 0 \in H^4(M; \mathbb{Z})$ is the obstruction to forming a *string structure* on a manifold M .

In this paper, we only deal with a homotopy-theoretic version of string structures. While geometric notions, such as [Coquereaux and Pilch 1989; Stolz and Teichner 2004; Waldorf 2009], are necessary for applications, we show we can recover some of this geometric information from the topological data for free. When dealing with these degree 4 classes, one usually must also deal with the associated differential characters. We naturally obtain globally defined forms representing these characters. We also speculate on the possibility that the string orientation of tmf may encode obstructions to positive Ricci curvature metrics. This would be analogous to the obstructions for positive scalar curvature metrics encoded in the spin orientation of KO .

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Throughout, all manifolds will be compact, connected, oriented, smooth, and without boundary. We now set up notation. Let G be a compact, simply connected, simple Lie group and $\lambda \in H^4(BG; \mathbb{Z})$ a universal characteristic class. Let $P \xrightarrow{\pi} M$ be a principal G -bundle with connection Θ , and let

$$\lambda(P) \in H^4(M; \mathbb{Z}), \quad \lambda(\Theta) \in \Omega^4(M), \quad \check{\lambda}(\Theta) \in \check{H}^4(M),$$

respectively be the naturally induced characteristic class, Chern–Weil form, and Cheeger–Simons differential character. The differential character is closely related to the Chern–Simons form $CS_\lambda(\Theta) \in \Omega^3(P)$. Finally, g will be a Riemannian metric on M . (In Section 2, we also consider more general G and λ .)

To the characteristic class

$$\lambda \in H^4(BG; \mathbb{Z}) \cong H^3(G; \mathbb{Z}) \cong \pi_3(G) \cong \mathbb{Z},$$

one can associate a topological group \widetilde{G}_λ and homomorphism $\widetilde{G}_\lambda \rightarrow G$ killing off the corresponding element in $\pi_3(G)$ and inducing isomorphisms in all higher homotopy groups [Stolz 1996; Stolz and Teichner 2004; Baez et al. 2007; Henriques 2008; Schommer-Pries 2009]. When $G = \text{Spin}(k)$ and $\lambda = \frac{1}{2}p_1$, the resulting group is commonly known as $\text{String}(k)$. While the actual groups \widetilde{G}_λ are not easy to describe, their homotopy type is clearly fixed. Therefore, we base our constructions only on the homotopy type. While more concrete models of \widetilde{G}_λ lead to more geometric definitions of \widetilde{G}_λ -structures, we only consider the problem of lifting the classifying map from BG to \widetilde{BG}_λ . We call a specific choice of lift a trivialization of the cohomology class λ .

In Section 2, we show that, up to homotopy, such trivializations of λ are naturally equivalent to cohomology classes $\mathcal{S} \in H^3(P; \mathbb{Z})$ that restrict to $\Omega\lambda \in H^3(G; \mathbb{Z})$ on the fibers. Here, $\Omega\lambda$ is the class that universally transgresses to λ . These classes \mathcal{S} are referred to as λ -trivialization classes. In the case where $G = \text{Spin}(k)$ and $\lambda = \frac{1}{2}p_1$, we see that the homotopy class of a string structure is equivalent to its string class \mathcal{S} . An important consequence is that one can describe an element in string bordism by a spin manifold M and string class $\mathcal{S} \in H^3(\text{Spin}(TM); \mathbb{Z})$.

In fact, these statements hold in greater generality, and Section 2 considers the more general case where G is a topological group and $\lambda \in H^n(BG; H)$. Homotopy classes of lifts to \widetilde{BG}_λ still induce canonical classes in $H^{n-1}(P; H)$, and there is an equivalence when $\check{H}^i(BG; H) = 0$ for $i < n$.

In Section 3, we analyze the harmonic representative of a λ -trivialization class \mathcal{S} on $P \xrightarrow{\pi} M$. The metric on P is naturally induced by a connection Θ on P and a Riemannian metric g on M . The harmonic 3-forms on P , in an adiabatic limit, were previously analyzed in [Redden 2008]. In Theorem 3.7, we see that the induced Hodge isomorphism $H^3(P; \mathbb{R}) \xrightarrow{\cong} \Omega^3(P)$ sends the string class \mathcal{S} to the 3-form $CS_\lambda(\Theta) - \pi^* H_{\mathcal{S}, g, \Theta}$. We call the form $H_{\mathcal{S}, g, \Theta} \in \Omega^3(M)$ the canonical 3-form

associated to the λ -trivialization class, metric, and connection. Proposition 3.12 states

$$d^* H_{\mathcal{S},g,\Theta} = 0 \in \Omega^2(M) \quad \text{and} \quad \check{H}_{\mathcal{S},g,\Theta} = \check{\lambda}(\Theta) \in \check{H}^4(M),$$

where $\check{H}_{\mathcal{S},g,\Theta}$ is the induced differential character. Thus, the form $H_{\mathcal{S},g,\Theta}$ lifts $\check{\lambda}(\Theta)$ to take values in \mathbb{R} instead of \mathbb{R}/\mathbb{Z} . This lift is independent of the metric on M ; the metric picks out the forms with smallest norm lifting $\check{\lambda}(\Theta)$.

We note that any time one encounters $\check{\lambda}(\Theta)$ and $\lambda(P) = 0$, the form $H_{\mathcal{S},g,\Theta}$ is relevant because it gives a purely local version of $\check{\lambda}(\Theta)$. This situation arises in theoretical physics under the guise of anomaly cancellation. It also arises when constructing the loop group extension bundle $\widehat{LP} \rightarrow LP$ restricting to $\widehat{LG} \rightarrow LG$.

In Sections 4–6, we deal exclusively with string structures on the frame bundle $\text{Spin}(TM) \rightarrow M$ of a manifold with spin structure. Given a string class \mathcal{S} and Riemannian metric g , we use the Levi-Civita connection to produce the canonical 3-form $H_{\mathcal{S},g}$.

Section 4 is largely motivational and provides background information on how string structures arise and why they are important. In particular, we discuss the string orientation

$$M\text{String} \xrightarrow{\sigma} \text{tmf}$$

of the cohomology theory of topological modular forms [Hopkins 2002] and its tentative relationship to index theory on loop spaces. This is analogous to the well-understood relationship between KO-theory and index theory. A theorem of Hitchin shows that the spin orientation of KO encodes obstructions to positive scalar curvature metrics. In the hope of an analogous theorem, Question 4.3 asks, If (M, \mathcal{S}, g) is a closed Riemannian n -manifold with string class \mathcal{S} satisfying both $\text{Ric}(g) > 0$ and $H_{\mathcal{S},g} = 0$, does this imply that $\sigma[M, \mathcal{S}] = 0 \in \text{tmf}^{-n}(pt)$?

In Section 5, we give an equivalent reformulation of Question 4.3. One can use the canonical 3-form $H_{\mathcal{S},g}$ to modify the Levi-Civita connection, inducing a metric connection $\nabla^{\mathcal{S},g}$ with torsion. Since the metric is used to “raise an index” of $H_{\mathcal{S},g}$, the global rescaling of M determines a canonical 1-parameter family of connections associated to $H_{\mathcal{S},g}$. This converges to the Levi-Civita connection in the large volume limit. Proposition 5.4 states that the simultaneous condition $(\text{Ric}(g) > 0, H_{\mathcal{S},g} = 0)$ is equivalent to the modified connection having positive Ricci curvature in a small volume limit. An interesting side note is the alternate description of the Levi-Civita connection. For a fixed metric g , the Levi-Civita connection is the unique metric connection maximizing the Ricci curvature.

In Section 6, we examine Question 4.3 in the case where $M = S^3$ with a homogeneous metric (under the left or right action of $S^3 \cong \text{SU}(2)$). In this case, the answer to Question 4.3 is yes, but not if either of the conditions $(\text{Ric}(g) > 0, H_{\mathcal{S},g} = 0)$ are weakened in an obvious way. The only (\mathcal{S}, g) satisfying both conditions is the

string class and round metric induced from D^4 . In this case, the string bordism class is obviously 0. However, there is a 1-parameter family of left-invariant metrics g satisfying $\text{Ric}(g) \geq 0$ and $H_{\mathcal{R},g} = 0$, where \mathcal{R} is induced by the right-invariant framing. Since $\sigma[S^3, \mathcal{R}] = 1/24 \in \text{tmf}^{-3}(pt)$, we see that our question would have a negative answer if one were to weaken the curvature condition. Also, one can find Ricci positive metrics g such that $H_{\mathcal{R},g}$ is arbitrarily small, so one cannot easily weaken the condition $H_{\mathcal{G},g} = 0$ either.

We close by noting that Section 3 is part of a more general story. The results of [Redden 2008] and Section 2 imply that the adiabatic-harmonic representative of a spin^c class gives a canonical 2-form refining the flat differential character $\check{W}_3(\Theta)$. Similarly, the harmonic representative of an SU-class on a $U(n)$ -bundle canonically gives a 1-form refining the character $\check{c}_1(\Theta)$. It appears there is a very general relationship between certain cohomology classes on a bundle P , their harmonic representatives, and the associated differential characters. The author is currently attempting to prove and properly understand these relations.

2. Trivializations of characteristic classes

In this section we make some observations on the general theory of trivializing a characteristic class, and we apply it to the Pontrjagin class $\frac{1}{2}p_1 \in H^4(B\text{Spin}; \mathbb{Z})$ to obtain results in subsequent sections. In the case of spin structures, or trivializations of w_2 , the results in this section are quite standard. In fact, this section is essentially a rewriting of [Lawson and Michelsohn 1989, Chapter 2.1] so that it applies in greater generality.

Since we will frequently use the notions of homotopy fibers and Eilenberg–Mac Lane spaces, we recall a couple of key facts. If H is an abelian group,¹ then an Eilenberg–Mac Lane space of type $K(H, n)$ is a space with the only nontrivial homotopy group being $\pi_n K(H, n) \cong H$. The space $K(H, n)$ is unique up to homotopy and is the classifying space for ordinary cohomology; that is, for a CW-complex X ,

$$H^n(X; H) \cong [X, K(H, n)],$$

where the right side is homotopy classes of based maps $X \rightarrow K(H, n)$. Furthermore, the loop-space functor Ω induces a homotopy equivalence

$$\Omega K(H, n) \simeq K(H, n-1) \quad \text{for } n > 1.$$

The homotopy fiber of a map is defined as the pullback of the pathspace fibration. Given a space Y with basepoint y_1 , we obtain the pathspace

$$PY = \{\gamma : [0, 1] \rightarrow Y \mid \gamma(1) = y_1\}.$$

¹ H is unrelated to the canonical forms in subsequent sections.

The natural map $PY \rightarrow Y$ given by $\gamma(0)$ is a fibration whose fiber is homotopic to ΩY . In fact, ΩY acts on the total space of this fibration. The homotopy fiber \tilde{X}_f of a map $X \xrightarrow{f} Y$ is then the actual pullback of PY . If the homotopy fiber construction is repeated, one obtains a sequence of fibrations homotopic to

$$\dots \Omega \tilde{X}_f \rightarrow \Omega X \xrightarrow{\Omega f} \Omega Y \rightarrow \tilde{X}_f \rightarrow X \xrightarrow{f} Y.$$

Now, let G be a connected topological group (of CW type so that standard classifying space constructions apply). Then, BG is the classifying space for G -bundles, and $H^*(BG; H)$ is the cohomology of BG with coefficients in H . The examples we will be concerned with are when G is a classical Lie group such as $SO(n)$ or $Spin(n)$, and $H = \mathbb{Z}$ or $\mathbb{Z}/2$. Consider a universal characteristic class $\lambda \in H^n(BG; H)$, equivalent to a homotopy class of maps

$$BG \xrightarrow{\lambda} K(H, n).$$

We fix a specific map λ and will not distinguish notationally between the map and the cohomology class.

Let \widetilde{BG}_λ be the homotopy fiber of $BG \xrightarrow{\lambda} K(H, n)$. This gives rise to the sequence

$$\dots \rightarrow G \xrightarrow{\Omega \lambda} K(H, n-1) \rightarrow \widetilde{BG}_\lambda \rightarrow BG \xrightarrow{\lambda} K(H, n)$$

of fibrations up to homotopy. Let $P \xrightarrow{\pi} M$ be a principal G -bundle over the space M ; that is, P has a free continuous (right) G -action with quotient map $\pi : P \rightarrow P/G \cong M$. Any such bundle P can be obtained as the pullback of the universal bundle

$$\begin{array}{ccc} P & \xrightarrow{f^*} & EG \\ \downarrow \pi & & \downarrow \pi \\ M & \xrightarrow{f} & BG. \end{array}$$

Consequently, any G -bundle has a natural characteristic class

$$\lambda(P) := f^* \lambda \in H^n(M; H).$$

Definition 2.1. A trivialization of the characteristic class λ on P is a lift of the classifying map to \widetilde{BG}_λ , that is, a lift \tilde{f}

$$\begin{array}{ccc} & & \widetilde{BG}_\lambda \\ & \nearrow \tilde{f} & \downarrow \\ M & \xrightarrow{f} & BG. \end{array}$$

We say two trivializations \tilde{f}_0 and \tilde{f}_1 are homotopic if they are homotopic through

the space of lifts, that is, if there exists a homotopy $\tilde{F} : [0, 1] \times M \rightarrow \widetilde{BG}_\lambda$ such that $\tilde{F}|_0 = \tilde{f}_0$, $\tilde{F}|_1 = \tilde{f}_1$, and $\tilde{F}|_t$ is a lift of f for all $t \in [0, 1]$.

Proposition 2.2. *Let $P \xrightarrow{\pi} M$ be a G -bundle classified by the map $f : M \rightarrow BG$.*

- (1) *There exists a trivialization of λ on P if and only if $\lambda(P) = 0 \in H^n(M; H)$.*
- (2) *If $\lambda(P) = 0$, the set of trivializations of λ up to homotopy has a free and transitive action of $H^{n-1}(M; H)$; that is, it is an $H^{n-1}(M; H)$ -torsor.*

Proof. Part (1) follows from the definition of the homotopy fiber. A lift \tilde{f} is precisely the choice of a nullhomotopy of $\lambda \circ f : M \rightarrow K(H, n)$, and $\lambda \circ f$ is nullhomotopic precisely when the cohomology class $\lambda(P) = 0$.

For part (2), assume an initial trivialization \tilde{f}_0 . This is equivalent to a global section $\tilde{f}_0 : M \rightarrow f^*\widetilde{BG}_\lambda$, and $\widetilde{BG}_\lambda \rightarrow BG$ is a fibration with fibers of type $\Omega K(H, n) \simeq K(H, n - 1)$. In fact the H-space $\Omega K(H, n)$ acts fiberwise on \widetilde{BG}_λ , so a global section \tilde{f}_0 induces a fiber homotopy equivalence

$$\begin{array}{ccc}
 M \times \Omega K(H, n) & \xrightarrow{\cong} & f^*\widetilde{BG}_\lambda \\
 & \searrow & \swarrow \\
 & & M.
 \end{array}$$

Therefore, the homotopy class of any other section $\tilde{f}_1 : M \rightarrow f^*\widetilde{BG}_\lambda$ is equivalent to the homotopy class of a function $M \rightarrow \Omega K(H, n) \simeq K(H, n - 1)$. □

Note that the connectedness of G implies that BG is simply connected, so we don't have to use local coefficients when dealing with the cohomology of fibers. The cohomology of any fiber is canonically isomorphic to $H^*(G; H)$, and we have a well-defined "restriction to fibers" map in cohomology, given by

$$i^* : H^*(P; H) \rightarrow H^*(G; H).$$

Proposition 2.3. (1) *A trivialization \tilde{f} of $\lambda(P)$ gives a canonical cohomology class in $H^{n-1}(P; H)$ that restricts on fibers to the class $\Omega\lambda \in H^{n-1}(G; H)$.*

- (2) *The cohomology class in (1) only depends on the homotopy class of \tilde{f} .*
- (3) *Furthermore, $H^{n-1}(M; H)$ acts equivariantly on the homotopy classes of λ -trivializations and $H^{n-1}(P; H)$ via π^* .*

Proof. For part (1), consider the universal pullback bundle

$$\begin{array}{ccc}
 EG & \xleftarrow{\pi^*} & \Pi^* EG \\
 \downarrow & & \downarrow \\
 BG & \xleftarrow{\pi} & \widetilde{BG}_\lambda
 \end{array}$$

Then, a lift $\tilde{f} : M \rightarrow \widetilde{BG}_\lambda$ such that $\Pi \circ \tilde{f} = f$ is equivalent to a G -equivariant map $\tilde{f}^* : P \rightarrow \Pi^*EG$ such that $\Pi^* \circ \tilde{f}^* = f^*$.

Since EG is contractible, Π^*EG is a $K(H, n - 1)$ space, as evidenced by the natural homotopy equivalence of fibrations given by

$$\begin{array}{ccc}
 G & \longrightarrow & \Pi^*EG \\
 \downarrow \simeq & & \downarrow \simeq \\
 \Omega BG & \xrightarrow{\Omega\lambda} & \Omega K(H, n)
 \end{array}
 \begin{array}{c}
 \nearrow \\
 \searrow
 \end{array}
 \widetilde{BG}_\lambda.$$

Therefore, any lift \tilde{f} is equivalent to $\tilde{f}^* : P \rightarrow \Pi^*EG \simeq K(H, n - 1)$. When restricted to a fiber, $\tilde{f}^* : G \rightarrow \Pi^*EG$ is equivalent to $\Omega\lambda : G \rightarrow K(H, n - 1)$. This is shown in the following commutative diagram:

$$\begin{array}{ccccc}
 & & & & K(H, n - 1) \\
 & & & \tilde{f}^* \text{ (dashed)} & \downarrow \\
 & & & \nearrow & \Pi^* \\
 P & \xrightarrow{\tilde{f}^*} & EG & \xrightarrow{\tilde{f}} & \widetilde{BG}_\lambda \\
 \downarrow & \nearrow & \downarrow & \searrow & \downarrow \\
 M & \xrightarrow{f} & BG & \xrightarrow{\Pi} & \widetilde{BG}_\lambda
 \end{array}$$

For part (2), a homotopy \tilde{F} between any two trivialisations \tilde{f}_0 and \tilde{f}_1 naturally lifts to an equivariant homotopy \tilde{F}^* between the bundle maps \tilde{f}_0^* and \tilde{f}_1^* . Therefore, the cohomology class $\tilde{f}^* \in H^{n-1}(P; H)$ of a trivialisaton only depends on the homotopy class of \tilde{f} .

For part (3), the fiberwise action of $\Omega K(H, n)$ on $\widetilde{BG}_\lambda \xrightarrow{\Pi} BG$ naturally pulls back via π^* to an action on Π^*EG . If $\tilde{f}_1 = \phi \cdot \tilde{f}_0$, where $\phi : M \rightarrow \Omega K(H, n)$, then

$$\tilde{f}_1^* = \pi^* \phi \cdot \tilde{f}_0^*.$$

Therefore, if two homotopy classes trivialisations $[\tilde{f}_0]$ and $[\tilde{f}_1]$ differ by $[\phi] \in H^{n-1}(M, H)$, their natural cohomology classes $[\tilde{f}_0^*], [\tilde{f}_1^*] \in H^{n-1}(P, H)$ differ by $\pi^*[\phi]$. □

The previous proposition gives a map

$$(2.4) \quad \{\lambda\text{-trivialisations}\} / \sim \longrightarrow \{\mathcal{S} \in H^{n-1}(P; H) \mid i^*\mathcal{S} = \Omega\lambda \in H^{n-1}(G; H)\}$$

that is equivariant under the natural $H^{n-1}(M; H)$ action. Here, \sim denotes equivalence up to homotopy. In general, this map is neither injective nor surjective. We will refer to such a cohomology class \mathcal{S} as λ -trivialisaton class.

Proposition 2.5. *Suppose $\tilde{H}^i(G; H) = 0$ for $i < n - 1$. Then (2.4) is a bijection.*

Proof. The connectedness of G implies the E_2 term in the Leray–Serre cohomology spectral sequence for $EG \rightarrow BG$ is

$$E_2^{r,s} \cong H^r(BG; H^s(G; H)),$$

and the contractibility of EG implies that $E_\infty^{r,s} = 0$ for $(r, s) \neq (0, 0)$. This, combined with the vanishing of $H^i(G; H)$ for $i < n - 1$, implies that the transgression

$$d_n : E_n^{0,n-1} \cong H^{n-1}(G; H) \rightarrow E_n^{n,0} \cong H^n(BG; H)$$

is an isomorphism. In fact, Lemma 2.6 says that $d_n(\Omega\lambda) = \lambda$.

The Leray–Serre cohomology spectral sequence for $P \xrightarrow{\pi} M$ is pulled back from the sequence for the universal bundle. This results in the exact sequence

$$0 \rightarrow H^{n-1}(M; H) \xrightarrow{\pi^*} H^{n-1}(P; H) \xrightarrow{i^*} H^{n-1}(G; H) \xrightarrow{d_n} H^n(M; H) \\ \Omega\lambda \mapsto \lambda(P).$$

If $\lambda(P) = 0$, then the action of $H^{n-1}(M; H)$ is free and transitive on classes in $H^{n-1}(P; H)$ restricting to $\Omega\lambda$. Since (2.4) is an equivariant map, and both sides are torsors for $H^{n-1}(M; H)$, it must be a bijection. \square

Lemma 2.6. *Suppose that $\tilde{H}^i(X) = 0$ for $i < n$. Then, the cohomology transgression for the pathspace fibration $\Omega X \hookrightarrow PX \rightarrow X$ is the inverse of the loop functor; that is, $d_n^{-1} = \Omega$ in*

$$H^{n-1}(\Omega X; H) \begin{array}{c} \xrightarrow{d_n} \\ \cong \\ \xleftarrow{\Omega} \end{array} H^n(X; H).$$

Proof. For the fibration $\Omega X \hookrightarrow PX \rightarrow X$, the transgression and loop functor are related by

$$\begin{array}{ccc} H^n(X; H) & \xrightarrow{\Omega} & H^{n-1}(\Omega X; H) \\ \downarrow & & \uparrow \\ E_n^{n,0} & \xleftarrow{d_n} & E_n^{0,n-1}. \end{array}$$

This follows from the general relationship between the transgression and cohomology loop suspension [Serre 1951]. If $\tilde{H}^i(X; H) = 0$ for $i < n$, then there is no room for any nontrivial differentials in the Serre spectral sequence until d_n . Therefore, $E_n^{n,0} \cong H^n(X; H)$, $E_n^{0,n-1} \cong H^{n-1}(\Omega X; H)$ and d_n is an isomorphism with inverse Ω . \square

Finally, we wish to make a general note about *stable* cohomology classes. The usual examples are Chern classes, Pontryagin classes, and Stiefel–Whitney classes, and they correspond to the stable cohomology of classifying spaces for the groups

$U(k)$ and $O(k)$. In general, assume one has a sequence of groups $\{G(k)\}$ and natural inclusions $G(k) \hookrightarrow G(k+1)$ inducing maps

$$\dots \rightarrow BG(k) \rightarrow BG(k+1) \rightarrow BG(k+2) \rightarrow \dots$$

such that the cohomology stabilizes. We then refer to the cohomology of $BG = \lim_{k \rightarrow \infty} BG(k)$. Any cohomology class $\lambda \in H^n(BG; H)$ is stable and defines a sequence of cohomology classes $\lambda_k \in H^n(BG(k); H)$ for all k :

$$H^n(BG; H) \rightarrow H^n(BG(k); H), \quad \lambda \mapsto \lambda_k,$$

though the k -subscript is usually unnecessary and dropped. Given a $G(k)$ -bundle $P(k)$ classified by $f : M \rightarrow BG(k)$, one can stably extend to a $G(k+l)$ -bundle $P(k+l)$ by $M \xrightarrow{\tilde{f}} BG(k) \rightarrow BG(k+l)$. It is obvious that the characteristic class is stable in that $\lambda_{k+l}(P(k+l)) = \lambda_k(P(k)) \in H^n(M; H)$.

Proposition 2.7. *Consider $\lambda \in H^n(BG; H)$. A trivialization of λ_k on any $G(k)$ -bundle naturally induces a trivialization of λ on any stable extension of P .*

Proof. This follows from the naturality of homotopy fibers. If we consider the inclusion map $\iota : BG(k) \rightarrow BG(k+1)$, then

$$\widetilde{BG(k)}_\lambda = \lambda_k^* PK(H, n) = (\lambda_{k+1} \circ \iota)^* PK(H, n) = \iota^* \widetilde{BG(k+1)}_\lambda.$$

Drawing this bundle map, we have

$$\begin{array}{ccc} & \widetilde{BG(k)}_\lambda & \longrightarrow & \widetilde{BG(k+1)}_\lambda \\ & \nearrow \tilde{f} & & \downarrow \\ M & \xrightarrow{f} & BG(k) & \longrightarrow & BG(k+1). \end{array}$$

Any trivialization of $\widetilde{BG(k)}_\lambda$ naturally extends to a trivialization of $\widetilde{BG(k+1)}_\lambda$ by composition, and this process can be continued indefinitely. □

To the $G(k_1)$ -bundle $P_1 \xrightarrow{\pi_1} M$ and $G(k_2)$ -bundle $P_2 \xrightarrow{\pi_2} M$ we can associate the $G(k_1) \times G(k_2)$ -bundle $P_1 \times_M P_2 \rightarrow M$. If there are inclusions

$$BG(k_1) \times BG(k_2) \xrightarrow{\iota_1 \times \iota_2} BG(k_1 + k_2),$$

the bundle $P_1 \times_M P_2$ is also naturally a $G(k_1 + k_2)$ -bundle.

Suppose that $H^i(BG; H) = 0$ for $i < n$. The Kunneth formula then implies the additivity of $\lambda \in H^n(BG; H)$:

$$(2.8) \quad \lambda(P_1 \times_M P_2) = \lambda(P_1) + \lambda(P_2) \in H^n(M; H).$$

The bottom square of the diagram below then commutes up to homotopy, implying the existence of the dotted arrow map.

$$\begin{array}{ccc}
 \widetilde{BG}(k_1)_\lambda \times \widetilde{BG}(k_2)_\lambda & \dashrightarrow & \widetilde{BG}(k_1+k_2)_\lambda \\
 \downarrow & & \downarrow \\
 BG(k_1) \times BG(k_2) & \longrightarrow & BG(k_1+k_2) \\
 \downarrow \lambda \times \lambda & & \downarrow \lambda \\
 K(H, n) \times K(H, n) & \longrightarrow & K(H, n)
 \end{array}$$

Therefore, a trivialization of λ on the bundles P_1 and P_2 induces a trivialization of λ on $P_1 \times_M P_2$ when viewed as a $G(k_1+k_2)$ -bundle (at least up to homotopy). This can also be seen explicitly in terms of cohomology classes.

Proposition 2.9. *For $l = 1, 2$, let $P_l \xrightarrow{\pi_l} M$ be a $G(k_l)$ -bundle. Let $P_1 \times_M P_2 \rightarrow M$ be the $G(k_1) \times G(k_2)$ -bundle and $P \rightarrow M$ the induced $G(k_1+k_2)$ -bundle. Assume $H^i(BG(k); H) = 0$ for $i < n$ (here $k = k_1, k_2, k_1+k_2$) and $\lambda \in H^n(BG; H)$. Then up to homotopy, a λ -trivialization on any two of $\{P, P_1, P_2\}$ induces a λ -trivialization on the third.*

Proof. Equation (2.8) implies the existence of a λ -trivialization on the third bundle if the other two admit λ -trivializations. Proposition 2.5 states the choice of a trivialization, up to homotopy, is equivalent to a λ -trivialization class $\mathcal{G}_i \in H^{n-1}(P_i; H)$ restricting to $\Omega\lambda$ on the fibers. We now show that the choice of λ -trivialization class on any two bundles determines one on the third bundle.

Note that there are natural bundle maps

$$\begin{array}{ccc}
 P_1 \times_M P_2 & \xrightarrow{\iota_1 \times \iota_2} & P \\
 \swarrow \pi_1 & & \searrow \pi_2 \\
 P_1 & & P_2
 \end{array}$$

We seek solutions to the equation

$$(\iota_1 \times \iota_2)^* \mathcal{G} = \pi^* \mathcal{G}_1 + \pi^* \mathcal{G}_2.$$

Just as in Proposition 2.5, the following commutative diagram is obtained from the Serre spectral sequences for the bundles P and $P_1 \times_M P_2$:

$$\begin{array}{ccccccc}
 0 \rightarrow H^{n-1}(M) & \longrightarrow & H^{n-1}(P) & \longrightarrow & H^{n-1}(G(k_1+k_2)) & \longrightarrow & H^n(M) \\
 & & \downarrow (\iota_1 \times \iota_2)^* & & \downarrow (\iota_1 \times \iota_2)^* & & \parallel \\
 0 \rightarrow H^{n-1}(M) & \rightarrow & H^{n-1}(P_1 \times_M P_2) & \rightarrow & H^{n-1}(G(k_1)) \oplus H^{n-1}(G(k_2)) & \rightarrow & H^n(M)
 \end{array}$$

The cohomology coefficients are all H but suppressed for spacing purposes.

We know that $(\iota_1 \times \iota_2)^* \lambda_{k_1+k_2} = \lambda_{k_1} \oplus \lambda_{k_2}$. For any three classes $\mathcal{S}, \mathcal{S}_1, \mathcal{S}_2$ in the respective bundles, the exact sequence implies

$$(\iota_1 \times \iota_2)^* \mathcal{S} - \pi_1^* \mathcal{S}_1 - \pi_2^* \mathcal{S}_2 = \pi^* \phi$$

for a unique $\phi \in H^{n-1}(M; H)$. If we fix two of the classes $\mathcal{S}, \mathcal{S}_1, \mathcal{S}_2$, modifying the third by ϕ gives us a solution to our desired equation. \square

The previous proposition is useful when dealing with cobordism theories. In the Pontryagin–Thom construction, the relevant extra structure takes place on the stable normal bundle. Suppose the m -manifold M already has a G -structure on the stable normal bundle $\nu(M)$. A lift of the classifying map to \widetilde{BG}_λ induces maps on the Thom spaces, which in turn give an element in the \widetilde{G}_λ -bordism group $M\widetilde{G}_\lambda^{-m}(pt)$.

However, it is often easier or more desirable to describe structures on the tangent bundle. For any manifold M , $TM \oplus \nu(M)$ is canonically isomorphic to the trivial bundle, so Proposition 2.9 often allows us to construct cobordism classes while only dealing with TM , or $G(TM)$.

Corollary 2.10. *Let $\lambda \in H^n(BG; H)$ be a stable class and suppose $H^i(BG; H)$ vanishes for $i < n$. Then, an m -manifold M with G -structure and λ -trivialization class $\mathcal{S} \in H^{n-1}(G(TM); H)$ canonically determines a \widetilde{G}_λ -bordism class $[M, \mathcal{S}]$ in $M\widetilde{G}_\lambda^{-m}(pt)$.*

We now apply Propositions 2.2, 2.3, and 2.5 and Corollary 2.10 to recover standard information on spin and spin^c structures as well as a convenient description of string structures.

2a. Spin structures. For $k > 2$, $\pi_1(\text{SO}(k)) \cong \mathbb{Z}/2$, and the nontrivial double cover is known as $\text{Spin}(k)$. The Hurewicz image of the generator of $\pi_1(\text{SO}(k))$ is the generator of $H^1(\text{SO}(k); \mathbb{Z}/2)$, which transgresses to $w_2 \in H^2(\text{BSO}(k); \mathbb{Z}/2)$. It is then clear that

$$\widetilde{\text{BSO}(k)}_{w_2} \simeq \text{BSpin}(k).$$

Moreover, there is a spin orientation of KO-theory $\alpha : M\text{Spin} \rightarrow \text{KO}$. Propositions 2.2, 2.3, 2.5, and Corollary 2.10 imply the following.

Proposition 2.11. *Let $P \xrightarrow{\pi} M$ be a principal $\text{SO}(k)$ -bundle.*

- *P admits a spin structure if and only if $w_2(P) = 0 \in H^2(M; \mathbb{Z}/2)$.*
- *The set of spin structures up to isomorphism is naturally equivalent to the set of spin classes $\mathcal{S} \in H^1(P; \mathbb{Z}/2)$ that restrict to the nontrivial class in $H^1(\text{SO}(k); \mathbb{Z}/2)$.*
- *The set of spin structures up to isomorphism is a torsor for $H^1(M; \mathbb{Z}/2)$.*

- An oriented m -manifold M with spin class $\mathcal{S} \in H^1(\mathrm{SO}(TM); \mathbb{Z}/2)$ gives rise to the bordism class $[M, \mathcal{S}] \in \mathrm{MSpin}^{-m}(pt)$ and the KO-theory class $\alpha[M, \mathcal{S}] \in \mathrm{KO}^{-m}(pt)$.

Geometrically, the statements can be understood by interpreting $H^1(\cdot; \mathbb{Z}/2)$ in terms of double covers; see [Lawson and Michelsohn 1989, Chapter 2.1]. Then, a spin structure on P is an equivariant double cover of P restricting fiberwise to the nontrivial double cover of $\mathrm{SO}(k)$.

2b. Spin^c structures. For $k > 2$,

$$H^1(\mathrm{SO}(k); \mathbb{Z}) = 0 \quad \text{and} \quad H^2(\mathrm{SO}(k); \mathbb{Z}) \cong H_1(\mathrm{SO}(k); \mathbb{Z}) \cong \pi_1(\mathrm{SO}(k)) \cong \mathbb{Z}/2.$$

The group $\mathrm{Spin}^c(k) = \mathrm{Spin}(k) \times_{\mathbb{Z}/2} S^1$ is a nontrivial S^1 -bundle over $\mathrm{SO}(k)$ and hence classified by the generator of $H^2(\mathrm{SO}(k); \mathbb{Z})$; this generator transgresses to $W_3 \in H^3(\mathrm{BSO}(k); \mathbb{Z}) \cong \mathbb{Z}/2$. Therefore,

$$\widetilde{\mathrm{BSO}(k)}_{W_3} \simeq \mathrm{BSpin}^c(k).$$

Furthermore, there is a spin^c orientation $\mathrm{MSpin}^c \rightarrow K$ of K -theory. Propositions 2.2, 2.3, 2.5, and Corollary 2.10 imply the following.

Proposition 2.12. *Let $P \xrightarrow{\pi} M$ be a principal $\mathrm{SO}(k)$ -bundle.*

- P admits a spin^c structure if and only if $W_3(P) = 0 \in H^3(M; \mathbb{Z})$.
- The set of spin^c structures up to homotopy is naturally equivalent to the set of classes $\mathcal{S} \in H^2(P; \mathbb{Z})$ that restrict to the nontrivial class in $H^2(\mathrm{SO}(k); \mathbb{Z})$.
- The set of spin^c structures up to homotopy is a torsor for $H^2(M; \mathbb{Z})$.
- An oriented m -manifold M with spin^c class $\mathcal{S} \in H^2(\mathrm{SO}(TM); \mathbb{Z})$ gives rise to the bordism class $[M, \mathcal{S}] \in \mathrm{MSpin}^c{}^{-m}(pt)$ and K -theory class $\in K^{-m}(pt)$.

Again, the statements above all have direct geometric interpretations based on $K(\mathbb{Z}, 2) \simeq BS^1$. A spin^c structure is an equivariant S^1 -extension of P restricting to the nontrivial extension on fibers. One can always tensor an S^1 -bundle over P with the pullback of an S^1 -bundle on M .

2c. String structures. Let G be any compact simple simply connected Lie group. Then, $\pi_2(G) = 0$ and $\pi_3(G) \cong H^3(G; \mathbb{Z}) \cong \mathbb{Z}$. What happens when you kill $\pi_3(G)$? The 3-connected cover $G\langle 4 \rangle \rightarrow G$ cannot be a finite-dimensional Lie group, since any connected nonabelian Lie group has nontrivial π_3 . However, there do exist topological groups $\widetilde{G} \rightarrow G$ that are 3-connected coverings. Various constructions can be found in [Stolz 1996; Stolz and Teichner 2004; Baez et al. 2007; Henriques 2008; Schommer-Pries 2009]. The results of all these imply the following (and usually one only needs G to be semisimple):

Choose a “level” $\lambda \in H^4(BG; \mathbb{Z}) \cong H^3(G; \mathbb{Z})$. Then, there exists a topological group and continuous homomorphism $G\langle\lambda\rangle \rightarrow G$ such that $G\langle\lambda\rangle$ has the homotopy type of the fiber of $G \xrightarrow{\Omega\lambda} K(\mathbb{Z}, 3)$. Applying the classifying space functor gives

$$BG\langle\lambda\rangle \rightarrow BG \xrightarrow{\lambda} K(\mathbb{Z}, 4).$$

When this construction is applied to $G = \text{Spin}(k)$ with $\lambda = \frac{1}{2}p_1 \in H^4(B\text{Spin}(k); \mathbb{Z})$, the resulting topological group is known as $\text{String}(k)$. Trivializations of $\frac{1}{2}p_1$ are commonly referred to as string structures. Applying the classifying space functor gives us $B\text{String}(k)$, and it is clear that

$$B\text{String}(k) \simeq \widetilde{B\text{Spin}(k)}_{\frac{1}{2}p_1}.$$

Remark 2.13. While multiple models for $G\langle\lambda\rangle$ exist, there is no “easy” model like the one Clifford algebras provide for the Spin groups. One must deal with some combination of higher categories, von Neumann algebras, or gerbes, each of which have particular subtleties. In this paper, we avoid these subtleties by only considering the homotopy type of $G\langle\lambda\rangle$. While we lose some information, we can characterize lifts of structure groups purely in terms of ordinary cohomology classes.

Remark 2.14. One should be careful when talking about spin, spin^c , and string structures up to homotopy. In addition to ignoring geometric considerations, these structures are naturally categories and have automorphisms; we only deal with isomorphism classes. The automorphisms play an important role, especially if one wishes to talk about structures locally or glue together manifolds with structures. See [Stolz and Teichner 2004; Waldorf 2009] for more concrete and categorical models of string structures.

For $k \geq 3$, $\text{Spin}(k)$ is simply connected and compact; thus $\widetilde{H}^i(\text{Spin}(k); \mathbb{Z}) = 0$ for $i < 3$. Also, there is a generalized cohomology theory tmf that has a string orientation $M\text{String} \xrightarrow{\sigma} \text{tmf}$, as discussed more in Section 4. Propositions 2.2, 2.3 and 2.5 and Corollary 2.10 then imply the following statements, which can obviously be rewritten for arbitrary $\lambda \in H^4(BG; \mathbb{Z})$ (except for the string orientation).

Definition 2.15. Let $P \xrightarrow{\pi} M$ be a principal $\text{Spin}(k)$ -bundle for $k \geq 3$.

- A string structure on a principal $\text{Spin}(k)$ -bundle $P \rightarrow M$ is a lift of the classifying map to $B\text{String}(k)$, that is, a lift \tilde{f}

$$\begin{array}{ccc}
 & & B\text{String}(k) \\
 & \tilde{f} \nearrow & \downarrow \\
 M & \xrightarrow{f} & B\text{Spin}(k).
 \end{array}$$

- A string class $\mathcal{S} \in H^3(P; \mathbb{Z})$ is a cohomology class that restricts fiberwise to the stable generator of $H^3(\text{Spin}(k); \mathbb{Z})$.

Proposition 2.16. *Let $P \xrightarrow{\pi} M$ be a principal $\text{Spin}(k)$ -bundle for $k \geq 3$.*

- *P admits a string structure if and only if $\frac{1}{2}p_1(P) = 0 \in H^4(M; \mathbb{Z})$.*
- *Up to homotopy, the choice of a string structure is equivalent to the choice of a string class $\mathcal{S} \in H^3(P; \mathbb{Z})$.*
- *If \mathcal{S} is a string class, then so is $\mathcal{S} + \pi^*\phi$ for $\phi \in H^3(M; \mathbb{Z})$. This natural action of $H^3(M; \mathbb{Z})$ on string classes is free and transitive; that is, the set of string classes is a torsor for $H^3(M; \mathbb{Z})$.*
- *A spin m -manifold M with string class $\mathcal{S} \in H^3(\text{Spin}(TM); \mathbb{Z})$ determines canonical classes $[M, \mathcal{S}] \in M\text{String}^{-m}(pt)$ and $\sigma[M, \mathcal{S}] \in \text{tmf}^{-m}(pt)$.*

Remark 2.17. The cohomology $H^3(\text{Spin}(k); \mathbb{Z})$ does not stabilize until $k = 5$, so we briefly describe the stable generator in dimensions 3 and 4. Under the low-dimensional isomorphisms with the symplectic groups, the groups $\text{Spin}(k)$ for $k = 3, 4, 5$ are related through the diagram

$$\begin{array}{ccccc}
 \text{Spin}(3) & \hookrightarrow & \text{Spin}(4) & \hookrightarrow & \text{Spin}(5) \\
 \downarrow \cong & & \downarrow \cong & & \downarrow \cong \\
 \text{Sp}(1) & \xrightarrow{\text{Id} \times \text{Id}} & \text{Sp}(1) \times \text{Sp}(1) & \hookrightarrow & \text{Sp}(2) \\
 \downarrow & & \downarrow & & \downarrow \\
 \mathbb{H} & \xrightarrow{\text{Id} \times \text{Id}} & \mathbb{H} \oplus \mathbb{H} & \hookrightarrow & \text{GL}(\mathbb{H}, 2).
 \end{array}$$

The $\text{Spin}(4)$ decomposition is induced by left and right multiplications of the unit quaternions. The second inclusion $\text{Spin}(4) \hookrightarrow \text{Spin}(5)$ is isomorphic to the matrix inclusion $\text{Sp}(1) \times \text{Sp}(1) \hookrightarrow \text{Sp}(2)$ along the diagonal. Since $H^3(\text{Sp}(k); \mathbb{Z})$ stabilizes at $k = 1$, we denote by 1 a generator of $H^3(\text{Sp}(1); \mathbb{Z}) \cong H^3(\text{SU}(2); \mathbb{Z})$. Then

$$\begin{array}{ccccc}
 H^3(\text{Spin}(5); \mathbb{Z}) & \longrightarrow & H^3(\text{Spin}(4); \mathbb{Z}) & \longrightarrow & H^3(\text{Spin}(3); \mathbb{Z}) \\
 1 & \longmapsto & (1, 1) & \longmapsto & 2.
 \end{array}$$

We originally defined $p_1 = -c_2$. Therefore, we see that $\Omega \frac{1}{2}p_1 \in H^3(\text{Spin}(3); \mathbb{Z})$ is twice a generator, and in fact $\Omega \frac{1}{2}p_1 = -2\Omega c_2 \in H^3(S^3; \mathbb{Z})$, where Ωc_2 is the usual generator.

3. Harmonic representative of a string class

In Section 2 we showed that, up to homotopy, a string structure on a principal $\text{Spin}(k)$ -bundle $P \rightarrow M$ is equivalent to a string class $\mathcal{S} \in H^3(P; \mathbb{Z})$. In this section we consider the harmonic representative of a string class. This will depend on the

choice of a metric on a closed manifold M , a connection on P , and it involves taking an adiabatic limit. We also must pass from \mathbb{Z} coefficients to \mathbb{R} coefficients and lose torsion information. The harmonic representative of \mathcal{S} is the Chern–Simons 3-form associated to the connection, minus a 3-form on M representing the differential cohomology class $\frac{1}{2}\check{p}_1(\Theta)$. The corresponding result holds for arbitrary $\lambda \in H^4(BG; \mathbb{Z})$, where G is a compact, simple, simply connected Lie group.

3a. Background: Differential characters. The canonical 3-form of Theorem 3.7 is best understood in the language of differential characters, originally developed in [Cheeger and Simons 1985]. See also [Freed 2002]. Let $C_i(M)$ and $Z_i(M)$ denote the group of smooth i -chains and cycles on M , respectively. Let $\Omega_{\mathbb{Z}}^i(M)$ denote the closed differential i -forms with integral periods; that is, their image in $H^i(M; \mathbb{R})$ lies in the image of $H^i(M; \mathbb{Z}) \rightarrow H^i(M; \mathbb{R})$. The group of differential characters $\check{H}^i(M)$ is defined as certain homomorphisms satisfying a transgression property:

$$\check{H}^i(M) := \{ \chi : Z_{i-1}(M) \rightarrow \mathbb{R}/\mathbb{Z} \mid \text{there exists } \omega \in \Omega^i(M) \text{ satisfying} \\ \int_{\Sigma} c^* \omega = \chi(\partial c) \pmod{\mathbb{Z}} \text{ for all } c : \Sigma \rightarrow M \in C_i(M) \}$$

The form ω associated to a character χ must be unique, and in fact $\omega \in \Omega_{\mathbb{Z}}^i(M)$. The character σ also determines a cohomology class in $H^i(M; \mathbb{Z})$ whose image in $H^i(M; \mathbb{R})$ is the same as $[\omega]$. These two maps induce the short exact sequences

$$(3.1) \quad 0 \rightarrow \frac{\Omega^{i-1}(M)}{\Omega_{\mathbb{Z}}^{i-1}(M)} \rightarrow \check{H}^i(M) \rightarrow H^i(M; \mathbb{Z}) \rightarrow 0,$$

$$(3.2) \quad 0 \rightarrow H^{i-1}(M; \mathbb{R}/\mathbb{Z}) \rightarrow \check{H}^i(M) \rightarrow \Omega_{\mathbb{Z}}^i(M) \rightarrow 0,$$

$$(3.3) \quad 0 \rightarrow \frac{H^{i-1}(M; \mathbb{R})}{H^{i-1}(M; \mathbb{Z})} \rightarrow \check{H}^i(M) \rightarrow H^i(M; \mathbb{Z}) \times_{H^i(M; \mathbb{R})} \Omega_{\mathbb{Z}}^i(M) \rightarrow 0.$$

In fact, these exact sequences uniquely characterize the groups $\check{H}^i(M)$ [Simons and Sullivan 2008]; one can refer to $\check{H}^*(M)$ as the differential cohomology of M without specifying the exact model being used, just as one refers to ordinary cohomology without specifying the model.

The importance of differential cohomology is due to the natural factoring of the Chern–Weil homomorphism through $\check{H}^*(M)$. Any compact Lie group G and universal class $\lambda \in H^{2i}(BG; \mathbb{Z})$ determine the following for any G -bundle $P \rightarrow M$ with connection Θ :

Characteristic class	$\lambda(P) \in H^{2i}(M; \mathbb{Z})$
Chern–Weil form	$\lambda(\Theta) \in \Omega^{2i}(M)$
Chern–Simons form	$\text{CS}_{\lambda}(\Theta) \in \Omega^{2i-1}(P)$
Differential character	$\check{\lambda}(\Theta) \in \check{H}^{2i}(M)$

The integral class and form associated to $\check{\lambda}(\Theta)$ are $\lambda(P)$ and $\lambda(\Theta)$, respectively.

Suppose G is compact, semisimple and simply connected. Let $\lambda \in H^4(BG; \mathbb{Z})$. Then, as discussed in [Freed 1995], the associated Chern–Weil form is

$$\lambda(\Theta) = \langle \Omega \wedge \Omega \rangle \in \Omega^4(M),$$

where Ω is the curvature of Θ , and $\langle \cdot, \cdot \rangle$ is a suitably normalized Ad-invariant inner product on \mathfrak{g} . In this case, the Chern–Simons form is

$$CS_\lambda(\Theta) = \langle \Theta \wedge \Omega \rangle - \frac{1}{6} \langle \Theta \wedge [\Theta \wedge \Theta] \rangle \in \Omega^3(P),$$

and

$$(3.4) \quad [i^* CS_\lambda(\Theta)] = \Omega\lambda \in H^3(G; \mathbb{R}).$$

Suppose that $c : X \rightarrow M$ is a 3-cycle. The assumptions on G imply that $c^*P \rightarrow X$ admits a global section p . Then

$$\check{\lambda}(\Theta)(c) = \int_X p^*(c^* CS_\lambda(\Theta)) \pmod{\mathbb{Z}}.$$

Hence, the information contained in $\check{\lambda}(\Theta) \in \check{H}^4(M)$ is simply the \mathbb{R}/\mathbb{Z} -periods of the Chern–Simons 3-form. One is forced to only consider the \mathbb{R}/\mathbb{Z} -periods because different global sections will give different \mathbb{R} -periods. Note that when M is a connected oriented 3-manifold, (3.3) implies

$$\check{H}^4(M) \cong H^3(M; \mathbb{R})/H^3(M; \mathbb{Z}) \cong \mathbb{R}/\mathbb{Z},$$

and the isomorphism is given by evaluating on the fundamental cycle $[M]$. On a 3-manifold, the element $\check{\lambda}(\Theta) \in \check{H}^4(M) \cong \mathbb{R}/\mathbb{Z}$ is often called the Chern–Simons [1974] invariant or number of the connection Θ . This invariant motivated the theory of differential characters.

3b. Hodge isomorphism on P . A Riemannian metric g on an n -manifold M induces the Hodge star $*$: $\Lambda^i TM \rightarrow \Lambda^{n-i} TM$, creating the codifferential

$$d^* := (-1)^{n(i+1)+1} * d * : \Omega^i(M) \rightarrow \Omega^{i-1}(M).$$

The Hodge Laplacian is the operator

$$\Delta_g = dd^* + d^*d = (d + d^*)^2 : \Omega^i(M) \rightarrow \Omega^i(M).$$

When M is closed (compact with no boundary), classical Hodge theory states that there is a canonical isomorphism

$$H^i(M; \mathbb{R}) \xrightarrow[\simeq]{\Pi_{\text{Ker } \Delta_g}} \text{Ker } \Delta_g \subset \Omega^i(M).$$

We will later denote $\text{Ker } \Delta_g$ by $\mathcal{H}^i(M)$, though the forms in $\mathcal{H}^3(P)$ will only be harmonic in a limit.

Let (M, g) be a closed Riemannian manifold, and let $P \xrightarrow{\pi} M$ be a principal G -bundle with connection Θ (G a compact, simple, simply connected Lie group). This naturally gives rise to a one-parameter family of right-invariant Riemannian metrics on P :

$$g_\delta := \delta^{-2} \pi^* g \oplus g_G \quad \text{for } \delta > 0,$$

where g_G is any biinvariant metric on G . (The metric g_G exists since G is compact, and it is unique up to a scaling constant because G is simple.) Conceptually, g_δ is given by using the connection to decompose TP into horizontal and vertical spaces; the metrics on M and G determine metrics on the horizontal and vertical components, respectively.

For any $\delta > 0$, we have the harmonic forms $\text{Ker } \Delta_{g_\delta}^3 \subset \Omega^3(P)$. In general this finite-dimensional subspace varies with δ , and we will not be concerned with $\text{Ker } \Delta_{g_\delta}$ for any particular δ . Instead, we analyze the *adiabatic limit*, the limit as $\delta \rightarrow 0$. Note that we had to choose the metric g_G . For this reason, it seems natural to introduce the scaling factor δ and take a limit, thus removing the dependence on the initial choice of g_G . Indeed, this is supported by concrete calculations, where the adiabatic limit appears to be of most interest.

Theorem 3.5 [Mazzeo and Melrose 1990; Dai 1991; Forman 1995]. *The 1-parameter space $\text{Ker } \Delta_{g_\delta}^i \subset \Omega^i(P)$ smoothly extends to $\delta = 0$. Furthermore, there is a spectral sequence computing $\lim_{\delta \rightarrow 0} \text{Ker } \Delta_{g_\delta}$ that is isomorphic to the Serre spectral sequence.*

This theorem holds in greater generality, and the context of each cited paper applies to the principal G -bundles with metric that we are considering. The spectral sequence mentioned is a Hodge-theoretic sequence, the details of which are given in [Forman 1995] and also summarized in [Redden 2008]. The fact that $\text{Ker } \Delta_{g_\delta}$ extends continuously to $\delta = 0$ (as a path in Grassmannian space) implies that there is still a Hodge isomorphism

$$H^i(P; \mathbb{R}) \xrightarrow[\simeq]{\Pi_{\text{Ker } \Delta_0}} \lim_{\delta \rightarrow 0} \text{Ker } \Delta_{g_\delta} \subset \Omega^i(P).$$

We now introduce the notation

$$\begin{aligned} \mathcal{H}^i(M) &:= \text{Ker } \Delta_g \subset \Omega^i(M), \\ \mathcal{H}^i(P) &:= \lim_{\delta \rightarrow 0} \text{Ker } \Delta_{g_\delta} \subset \Omega^i(P), \\ \mathcal{H}^i(G) &:= \text{Ker } \Delta_{g_G} \subset \Omega^i(G). \end{aligned}$$

In [Redden 2008], the spectral sequence interpretation of $\mathcal{H}^3(P)$ was used to give the following description of harmonic 3-forms on P in the adiabatic limit.

Theorem 3.6 [Redden 2008, Proposition 4.5 and Theorem 4.6]. *Consider the set $(P \xrightarrow{\pi} M, g, \Theta)$, where G is a compact simple Lie group. If $\lambda(P) = 0 \in H^4(M; \mathbb{R})$, then*

$$\mathcal{H}^3(P) = \mathbb{R}[\text{CS}_\lambda(\Theta) - \pi^*h] \oplus \pi^*\mathcal{H}^3(M),$$

where $h \in \Omega^3(M)$ is the unique coexact form satisfying $dh = \lambda(\Theta)$.

When G is also simply connected, the Serre spectral sequence gives the following exact sequence, as seen in Proposition 2.5:

$$\begin{array}{ccccccc} 0 & \longrightarrow & H^3(M; \mathbb{Z}) & \xrightarrow{\pi^*} & H^3(P; \mathbb{Z}) & \xrightarrow{i^*} & H^3(G; \mathbb{Z}) \xrightarrow{d_4} H^4(M; \mathbb{Z}) \\ & & & & \mathcal{G} \mapsto & \xrightarrow{?} & \Omega\lambda \longmapsto \lambda(P) \end{array}$$

Theorem 3.7. *Consider $(P \xrightarrow{\pi} M, g, \Theta)$ where G is a simply connected compact simple Lie group. Suppose that $\lambda(P) = 0 \in H^4(M; \mathbb{Z})$ and that $\mathcal{G} \in H^3(P; \mathbb{Z})$ is a λ -trivialization class, that is, $i^*\mathcal{G} = \Omega\lambda \in H^3(G; \mathbb{Z})$. Then, the image of \mathcal{G} under the Hodge isomorphism is of the form*

$$\begin{array}{ccc} H^3(P; \mathbb{Z}) \rightarrow H^3(P; \mathbb{R}) & \xrightarrow[\simeq]{\Pi_{\text{Ker } \Delta_0}} & \mathcal{H}^3(P) \subset \Omega^3(P) \\ \mathcal{G} \mapsto & & \text{CS}_\lambda(\Theta) - \pi^*H_{\mathcal{G},g,\Theta}, \end{array}$$

where $H_{\mathcal{G},g,\Theta} \in \Omega^3(M)$. Alternatively, $\Pi_{\text{Ker } \Delta_0}\mathcal{G} - \text{CS}_\lambda(\Theta) \in \pi^*\Omega^3(M)$.

Proof. The orthogonal decomposition of $\mathcal{H}^3(P)$ in Theorem 3.6 corresponds to a splitting $H^3(P; \mathbb{R}) \cong H^3(G; \mathbb{R}) \oplus H^3(M; \mathbb{R})$. We know $\pi^*H^3(M; \mathbb{R})$ restricts to $0 \subset H^3(G; \mathbb{R})$. As mentioned in (3.4), both $\text{CS}_\lambda(\Theta) - \pi^*h$ and \mathcal{G} cohomologically restrict to $\Omega\lambda \in H^3(G; \mathbb{R})$, so

$$\Pi_{\text{Ker } \Delta_0}\mathcal{G} - (\text{CS}_\lambda(\Theta) - \pi^*h) \in \pi^*\mathcal{H}^3(M).$$

Therefore, the harmonic representative of \mathcal{G} must be of the form

$$\text{CS}_\lambda(\Theta) - \pi^*h - \pi^*h',$$

with $h' \in \mathcal{H}^3(M)$, and we define $H_{\mathcal{G},g,\Theta} := h + h' \in \Omega^3(M)$. □

Remark 3.8. The theorem does not hold without taking an adiabatic limit. For a general $\delta > 0$,

$$\Pi_{\text{Ker } \Delta_{g\delta}}\mathcal{G} - \text{CS}_\lambda(\Theta) \notin \pi^*\Omega^3(M),$$

but instead will contain forms with bidegree (2,1) and (1,2) in the (horizontal, vertical) decomposition of $\Omega^3(P)$ given by the connection.

Remark 3.9. When restricted to the fibers, the Chern–Simons form is the standard harmonic (biinvariant) form representing $\Omega\lambda$; that is, $i^*\text{CS}_\lambda(\Theta) \in \mathcal{H}^3(G)$. Just

as \mathcal{S} is a cohomological extension of $\Omega\lambda$ to all of P , we see that $\Pi_{\text{Ker } \Delta_0} \mathcal{S} = \text{CS}_\lambda(\Theta) - \pi^* H_{\mathcal{S},g,\Theta}$ is a harmonic extension of $\Omega\lambda$ to all of P .

3c. Properties of canonical 3-form. Theorem 3.7 gives a canonical construction

$$(3.10) \quad \{\lambda\text{-triv classes}\} \times \text{Met}(M) \times \mathcal{A}(P) \rightarrow \Omega^3(M), \quad \mathcal{S}, g, \Theta \mapsto H_{\mathcal{S},g,\Theta}.$$

We call $H_{\mathcal{S},g,\Theta}$ the canonical 3-form associated to (\mathcal{S}, g, Θ) . While Theorem 3.7 only uses information about \mathcal{S} as a class in $H^3(P; \mathbb{R})$, the integrality becomes necessary when understanding $H_{\mathcal{S},g,\Theta}$ in terms of differential characters. The exact sequence (3.1) gives rise to

$$(3.11) \quad 0 \rightarrow \Omega_{\mathbb{Z}}^3(M) \rightarrow \Omega^3(M) \rightarrow \check{H}^4(M) \rightarrow H^4(M; \mathbb{Z}) \rightarrow 0$$

$$H_{\mathcal{S},g,\Theta} \mapsto \check{H}_{\mathcal{S},g,\Theta},$$

where the character $\check{H}_{\mathcal{S},g,\Theta}$ obtained via $\Omega^3(M) \rightarrow \check{H}^4(M)$ is given by simply by integrating $H_{\mathcal{S},g,\Theta}$ on cycles and reducing mod \mathbb{Z} . Also, note that $H^3(M; \mathbb{Z})$ acts naturally on $\{\lambda\text{-triv classes}\}$, and it also acts on $\Omega^3(M) \times \text{Met}(M)$ by adding a harmonic representative.

Proposition 3.12. *The construction (3.10) is equivariant with respect to the natural action of $H^3(M; \mathbb{Z})$; that is, $H_{\mathcal{S}+\pi^*\phi,g,\Theta} = H_{\mathcal{S},g,\Theta} + \Pi_{\text{Ker } \Delta_g} \phi$.*

Furthermore, the forms $H_{\mathcal{S},g,\Theta}$ satisfy the following:

- $d^* H_{\mathcal{S},g,\Theta} = 0 \in \Omega^2(M)$,
- $d H_{\mathcal{S},g,\Theta} = \lambda(\Theta) \in \Omega^4(M)$,
- $\check{H}_{\mathcal{S},g,\Theta} = \check{\lambda}(\Theta) \in \check{H}^4(M)$.

Proof. The action of $H^3(M; \mathbb{Z})$ on λ -trivialization classes is given by addition under π^* , and the action on $\Omega^3(M)$ is given by adding the harmonic representative (with respect to a fixed metric M). Theorem 3.6 implies that for $\phi \in H^3(M; \mathbb{Z})$ with harmonic representative $\Pi_{\text{Ker } \Delta_g} \phi \in \mathcal{H}^3(M)$,

$$\pi^*(\Pi_{\text{Ker } \Delta_g} \phi) = \Pi_{\text{Ker } \Delta_0}(\pi^* \phi) \in \Omega^3(P).$$

The property $d^* H_{\mathcal{S},g,\Theta} = 0$ also follows directly from Theorem 3.6. That $\Pi_{\text{Ker } \Delta_0} \mathcal{S}$ is closed implies

$$d(\text{CS}_\lambda(\Theta) - \pi^* H_{\mathcal{S},g,\Theta}) = 0,$$

$$\pi^* \lambda(\Theta) - \pi^* d H_{\mathcal{S},g,\Theta} = 0,$$

$$d H_{\mathcal{S},g,\Theta} = \lambda(\Theta),$$

with the last equality following from π^* being injective on forms.

Finally, suppose that $X \xrightarrow{c} M$ is a smooth 3-cycle on M . Then the value of $\check{\lambda}(\Theta)$ on (X, c) is

$$\check{\lambda}(\Theta)(c) = c^*\check{\lambda}(\Theta) \in \check{H}^4(X) \cong \mathbb{R}/\mathbb{Z}.$$

Now, standard obstruction theory implies that $c^*P \rightarrow X$ admits a global section $p : X \rightarrow c^*P$, and it is easy to see that

$$p^*c^*\check{CS}_\lambda(\Theta) = p^*\check{CS}_\lambda(c^*\Theta) = c^*\check{\lambda}(\Theta) \in \check{H}^4(X).$$

Because $CS_\lambda(\Theta) - \pi^*H_{\mathcal{G},g,\Theta} \in \Omega_{\mathbb{Z}}^3(P)$, we have $\check{CS}_\lambda(\Theta) = \pi^*\check{H}_{\mathcal{G},g,\Theta} \in \check{H}^3(P)$ and hence

$$\begin{aligned} c^*\check{\lambda}(\Theta) &= p^*c^*\check{CS}_\lambda(\Theta) = p^*c^*\pi^*\check{H}_{\mathcal{G},g,\Theta} \\ &= p^*\pi^*c^*\check{H}_{\mathcal{G},g,\Theta} = c^*\check{H}_{\mathcal{G},g,\Theta} \in \check{H}^4(X). \end{aligned}$$

This implies that for all 3-cycles $X \xrightarrow{c} M$

$$\check{H}_{\mathcal{G},g,\Theta}(c) = \check{\lambda}(\Theta)(c),$$

and hence $\check{\lambda}(\Theta) = \check{H}_{\mathcal{G},g,\Theta} \in \check{H}^4(M)$. (This also implies $dH_{\mathcal{G},g,\Theta} = \lambda(\Theta)$.) □

Integrating the form $H_{\mathcal{G},g,\Theta}$ naturally gives values in \mathbb{R} , and Proposition 3.12 says that reducing mod \mathbb{Z} gives the same values as $\check{\lambda}(P)$. In other words, the choice of a λ -trivialization class naturally gives a lift

$$(3.13) \quad \begin{array}{ccc} & & \mathbb{R} \\ & \nearrow^{H_{\mathcal{G},g,\Theta}} & \downarrow \\ Z_3(M) & \xrightarrow{\check{\lambda}(\Theta)} & \mathbb{R}/\mathbb{Z} \end{array}$$

and the action of $H^3(M; \mathbb{Z})$ modifies the lift by the induced map $Z_3(M) \rightarrow \mathbb{Z}$. While the actual form $H_{\mathcal{G},g,\Theta}$ depends on the choice of a metric, this lift does not.

Proposition 3.14. *The lift in (3.13) is independent of the choice of metric g .*

Proof. If g_0 and g_1 are two different metrics, then (3.11) implies

$$H_{\mathcal{G},g_1,\Theta} - H_{\mathcal{G},g_0,\Theta} \in \Omega_{\mathbb{Z}}^3(M).$$

The space of Riemannian metrics is contractible, so

$$H_{\mathcal{G},g_1,\Theta} - H_{\mathcal{G},g_0,\Theta} \in d\Omega^2(M). \quad \square$$

The role of the metric in (3.10) is to pick out the forms $H_{\mathcal{G},g,\Theta}$ with smallest norm still satisfying $\check{H}_{\mathcal{G},g,\Theta} = \check{\lambda}(\Theta)$. We denote the lift by $H_{\mathcal{G},\Theta} \in \check{H}_{\mathbb{R}}^4(M)$. Here we use the nonstandard notation of $\check{H}_{\mathbb{R}}^4(M)$ to denote characters $Z_3(M) \rightarrow \mathbb{R}$ satisfying the usual transgression assumption.

In summary, the construction (3.10) induces lifts of the standard differential character construction, which are encoded in the following diagram:

$$\begin{array}{ccc}
 \{\lambda\text{-triv classes}\} \times \text{Met}(M) \times \mathcal{A}(P) & \xrightarrow{H_{\mathcal{G},g,\Theta}} & \Omega^3(M) \\
 \downarrow & & \downarrow \\
 \{\lambda\text{-triv classes}\} \times \mathcal{A}(P) & \xrightarrow{H_{\mathcal{G},\Theta}} & \check{H}_{\mathbb{R}}^4(M) \\
 \downarrow & & \downarrow \\
 \mathcal{A}(P) & \xrightarrow{\check{\lambda}(\Theta)} & \check{H}^4(M).
 \end{array}$$

Remark 3.15. Stolz and Teichner [2004] define a *geometric* trivialization of $\lambda(P)$ as a trivialization of the extended Chern–Simons field theory on $P \rightarrow M$. This includes defining a lift of the differential character $\check{\lambda}(\Theta)$ to take values in \mathbb{R} , and it aligns nicely with the construction above. In fact, if $H^3(M; \mathbb{Z})$ has no torsion, then the choice of a lift of $\check{\lambda}(\Theta)$ to $\check{H}_{\mathbb{R}}^4(M)$ is equivalent to the choice of a λ -trivialization class. Waldorf [2009] gives an explicit model for string structures in terms of trivializations of a Chern–Simons 2-gerbe, and shows that a string structure produces a 3-form on M . The 3-forms obtained in our construction are a proper subset of those he obtained, analogous to the relationship between forms representing a de Rham class and harmonic forms.

Note that one can also directly define the lift $H_{\mathcal{G},\Theta}$ without using the Hodge isomorphism. On a 3-cycle $c : X \rightarrow M$,

$$(3.16) \quad H_{\mathcal{G},\Theta}(c) = \int_X p^*(\text{CS}_{\lambda}(c^*\Theta) - c^*S),$$

where p is any global section, and S is any de Rham representative of \mathcal{G} . This is a simple consequence of $S = \text{CS}_{\lambda}(\Theta) - H_{\mathcal{G},g,\Theta} + d\beta$. It is also easy to verify that the integral on the right side is independent of p . In cases like Lemma 3.18, this allows us to calculate the form $H_{\mathcal{G},g,\Theta}$ without solving a differential equation.

Suppose the G -bundle $P \xrightarrow{\pi} M$ is topologically trivial. Then, the choice of a global section $p : M \rightarrow P$ is equivalent to a trivialization $P \cong M \times G$. The canonical λ -trivialization on $M \times G$ induces one on P , and the corresponding cohomology class is given by the Kunneth isomorphism

$$(3.17) \quad H^3(P; \mathbb{Z}) \cong H^3(M; \mathbb{Z}) \oplus H^3(G; \mathbb{Z}), \quad \mathcal{G} \leftrightarrow (0, \Omega\lambda).$$

Lemma 3.18. *Suppose $P \xrightarrow{\pi} M$ is a trivial bundle with λ -trivialization class \mathcal{G} induced by the trivialization $p : M \rightarrow P$. Then,*

$$H_{\mathcal{G},g,\Theta} - p^* \text{CS}_{\lambda}(\Theta) \in d\Omega^2(M).$$

In particular, $p^* \text{CS}_\lambda(\Theta) = H_{\mathcal{G}, \Theta}$ as elements of $\check{H}^4_{\mathbb{R}}(M)$. If $d^* p^* \text{CS}_\lambda(\Theta) = 0$, then $p^* \text{CS}_\lambda(\Theta) = H_{\mathcal{G}, g, \Theta} \in \Omega^3(M)$.

Proof. As seen in (3.17), $p^* \mathcal{S} = 0 \in H^3(M; \mathbb{Z})$. Therefore, (3.16) simplifies to

$$\int_X c^* H_{\mathcal{G}, g, \Theta} = \int_X c^* p^* \text{CS}_\lambda(\Theta)$$

for all 3-cycles $c : X \rightarrow M$, so $[H_{\mathcal{G}, g, \Theta} - p^* \text{CS}_\lambda(\Theta)] = 0 \in H^3(M; \mathbb{R})$. □

One usually chooses $\lambda \in H^4(BG; \mathbb{Z}) \cong \mathbb{Z}$ to be the generator. This is because \widehat{BG}_λ is the universal extension. This universality is also reflected in the associated canonical 3-forms.

Proposition 3.19. *If $\mathcal{S} \in H^3(P; \mathbb{Z})$ is a λ -trivialization class and $\ell \in \mathbb{Z}$, then $\ell \mathcal{S}$ is an $\ell \lambda$ -trivialization class, and $H_{\ell \mathcal{S}, g, \Theta} = \ell H_{\mathcal{S}, g, \Theta} \in \Omega^3(P)$.*

Proof. The first statement is obvious, and the second follows from the linearity of the Hodge isomorphism. □

We now apply the construction above to $G = \text{Spin}(k)$ for $k \geq 3$ with $\lambda = \frac{1}{2} p_1 \in H^4(B\text{Spin}(k); \mathbb{Z})$ to canonically produce 3-forms associated to string structures. Since $\text{Spin}(4) \cong \text{SU}(2) \times \text{SU}(2)$ is not simple, we define the canonical 3-form when $k = 4$ to be the one obtained by stabilizing to $\text{Spin}(5)$, a process that does not affect $\frac{1}{2} \check{p}_1(\Theta)$.²

Theorem 3.20. *Let $P \xrightarrow{\pi} M$ be a principal $\text{Spin}(k)$ -bundle ($k \geq 3$) with connection Θ over the Riemannian manifold (M, g) . Under the Hodge isomorphism (in an adiabatic limit), a string class $\mathcal{S} \in H^3(P; \mathbb{Z})$ is represented by*

$$\Pi_{\text{Ker } \Delta_0} \mathcal{S} = \text{CS}_{\frac{1}{2} p_1}(\Theta) - \pi^* H_{\mathcal{S}, g, \Theta} \in \Omega^3(P).$$

The canonical form $H_{\mathcal{S}, g, \Theta} \in \Omega^3(M)$ is such that

- $d^* H_{\mathcal{S}, g, \Theta} = 0 \in \Omega^2(M)$,
- $\check{H}_{\mathcal{S}, g, \Theta} = \frac{1}{2} \check{p}_1(\Theta) \in \check{H}^4(M)$, and
- the construction of $H_{\mathcal{S}, g, \Theta}$ is equivariant with respect to the natural action of $H^3(M; \mathbb{Z})$.

In particular, consider the case where (M, g) is a Riemannian manifold with spin structure satisfying $\frac{1}{2} p_1(M) = 0 \in H^4(M; \mathbb{Z})$. Then, we can let $P = \text{Spin}(TM)$, and we call a string structure on $\text{Spin}(TM)$ a string structure on M . Letting Θ be the Levi-Civita connection, this gives a map

$$(3.21) \quad \{\text{String classes on } M\} \times \text{Met}(M) \rightarrow \Omega^3(M), \quad \mathcal{S}, g \mapsto H_{\mathcal{S}, g}.$$

²The arguments in [Redden 2008] can be extended to semisimple groups, and Theorem 3.20 also holds for $k = 4$ without stabilizing.

4. Canonical 3-forms and the string orientation of tmf

We now review how string structures arise and give a possible new application of the canonical 3-forms $H_{\mathcal{G},g}$ from (3.21). First recall some classical results from index theory; an excellent source is [Lawson and Michelsohn 1989]. Suppose M is an oriented closed manifold. A priori, one cannot form a spinor bundle $\mathrm{SO}(M) \times_{\mathrm{SO}(n)} S^\pm \rightarrow M$, because the spinor representations $\mathrm{SO}(n) \rightarrow \mathrm{GL}(S^\pm)$ are only projective. The choice of a spin structure, discussed in Section 2a, allows one to define the spinor bundle $S_M^\pm := \mathrm{Spin}(M) \times_{\mathrm{Spin}(n)} S^\pm$ and Dirac operator $\mathcal{D}_M : \Gamma(S^\pm) \rightarrow \Gamma(S^\mp)$.

While the Fredholm operator \mathcal{D}_M depends on the spin structure, the Atiyah–Singer index theorem states that its index does not, and in fact

$$\mathrm{index}(\mathcal{D}_M) = \hat{A}(M) \in \mathbb{Z}.$$

Here, $\hat{A}(M)$ is a topological invariant determined by a manifold’s Pontryagin classes and is defined for any oriented manifold. In general $\hat{A}(M) \in \mathbb{Q}$, but $\hat{A}(M) \in \mathbb{Z}$ when $w_2(M) = 0$. There is also a refinement of $\hat{A}(M)$ given by the spin orientation $\alpha : \mathrm{MSpin} \rightarrow \mathrm{KO}$. This refinement can be thought of as the Clifford-linear index, and it does depend on the spin structure.

$$(4.1) \quad \begin{array}{ccc} & & \mathrm{KO}^{-n}(pt) \\ & \nearrow \alpha & \downarrow \\ \mathrm{MSpin}^{-n}(pt) & \xrightarrow{\hat{A}} & \mathbb{Z}. \end{array}$$

The KO -invariants usually appear in family index theorems, but they also contain interesting information for a single manifold due to the torsion in $\mathrm{KO}^{-*}(pt)$.

Index theory is now a central part of mathematics, and one of its powerful applications is to the problem of when a closed manifold admits positive scalar curvature metrics. The Lichnerowicz–Weitzenböck formula $\mathcal{D}_M^2 = \nabla^* \nabla + \frac{1}{4}s$, which relates \mathcal{D}_M to a positive operator and the scalar curvature s , implies the following: If a closed spin manifold M admits a metric of positive scalar curvature, then $\mathrm{index}(\mathcal{D}_M) = \hat{A}(M) = 0$ [Lichnerowicz 1962]. Furthermore, $\alpha[M] = 0 \in \mathrm{KO}^{-n}(pt)$ for all spin structures [Hitchin 1974]. In fact, for simply connected spin manifolds of dimension ≥ 5 , all the α -invariants vanish if and only if M admits a metric of positive scalar curvature [Stolz 1992].

There is an analogous story, though not fully developed, involving the Witten genus, index theory on loop spaces, and elliptic cohomology. Witten [1988] used intuition from theoretical physics and defined a topological invariant $\varphi_W(M)$ known as the Witten genus. He claimed it should be the S^1 -equivariant index of the Dirac

operator on the free loop space LM ; that is,

$$\text{“index}^{S^1} \widehat{\mathcal{D}}_{LM}\text{”} = \varphi_W(M).$$

We place the left side in quotes because of analytic difficulties in defining a good theory of Fredholm operators on infinite-dimensional manifolds. However, the Witten genus (and other elliptic genera) are well defined, and one can make formal sense of index theory on LM by using localization formulas or the representation theory of loop groups. For a good overview on these ideas, see [Liu 1996]. While $\varphi_W(M) \in \mathbb{Q}[[q]][[q^{-1}]]$ for any oriented manifold, for a string manifold, $\varphi_W(M)$ is the q -expansion of a modular form (MF) with integer coefficients and weight $n/2$, and we say $\varphi_W(M) \in MF_n$. The intuitive reason is that when $\frac{1}{2}p_1(M) = 0$, one can define the spinor bundle on LM [Coquereaux and Pilch 1989]. We wish to form $L\text{Spin}(M) \times_{L\text{Spin}(n)} S \rightarrow LM$, where S is a positive energy representation of $L\text{Spin}(n)$. However, these representations are all projective, so one must pass to an S^1 -extension $L\widehat{\text{Spin}}(n) \rightarrow L\text{Spin}(n)$. Topologically, our string class $\mathcal{S} \in H^3(\text{Spin}(M); \mathbb{Z})$ transgresses to a class in $H^2(L\text{Spin}(M); \mathbb{Z})$ that defines an isomorphism class of S^1 -extension $L\widehat{\text{Spin}}(M) \rightarrow L\text{Spin}(M) \rightarrow LM$. We say that a string structure on M transgresses to a spin structure on LM (though in this paper we have only discussed isomorphism classes of such structures). This led to the following conjecture.

Conjecture 4.2 (Höhn and Stolz [Stolz 1996]). *Let M be a closed oriented n -manifold admitting spin and string structures. If M admits a metric of positive Ricci curvature, then the Witten genus $\varphi_W(M)$ vanishes.*

Stolz’s heuristic argument comes from the hope that there is some Weizenböck-type formula such that positive Ricci curvature on M implies positive scalar curvature on LM , which in turn implies $\text{Ker}(\widehat{\mathcal{D}}_{LM}) = \varphi_W(M) = 0$. Though this reasoning is far from rigorous, there are no known counterexamples, and the conjecture holds true for homogeneous spaces and complete intersections. To the author’s knowledge, there are no known examples of simply connected closed manifolds admitting metrics of positive scalar curvature, but not metrics of positive Ricci curvature. If the conjecture is true, it would provide examples of such manifolds.

Just as KO-theory refines the \hat{A} -genus, there is a cohomology theory tmf , or topological modular forms, with string-orientation refining the Witten genus (see [Hopkins 2002]):

$$\begin{array}{ccc}
 & & \text{tmf}^{-n}(pt) \\
 & \nearrow \sigma & \downarrow \\
 M\text{String}^{-n}(pt) & \xrightarrow{\varphi_W} & MF_n
 \end{array}$$

The map $\mathrm{tmf}^{-*}(pt) \rightarrow MF_*$ is a rational isomorphism, but it is not integrally surjective or injective. In particular, $\mathrm{tmf}^{-*}(pt)$ contains a great deal of torsion. While defining tmf is a subtle process, informally tmf is the universal elliptic cohomology theory, or the elliptic cohomology theory associated to the universal moduli stack of elliptic curves. Despite several attempts [Baas et al. 2004; Hu and Kriz 2004; Segal 1988; Stolz and Teichner 2004], there is still no geometric description of tmf . However, it is believed that tmf should provide a natural home for family index theorems on loop spaces.

One might hope that all the refined invariants in tmf also vanish for string manifolds admitting positive Ricci curvature metrics, giving an analogy of Hitchin’s theorem. However, there exist a fair number of compact nonabelian Lie groups (thus admitting positive Ricci curvature metrics) that are sent to torsion elements in $\mathrm{tmf}^{-*}(pt)$ via their left-invariant framing [Hopkins 2002]. In Section 6, we investigate the case where $M = S^3$.

Conceptually, this is still compatible with the analogy to classical index theory. The group $\mathrm{Spin}(n)$ is a discrete cover of $\mathrm{SO}(n)$, so there are no local differences between the bundles $\mathrm{Spin}(M)$ and $\mathrm{SO}(M)$ and their connections. However, $\mathrm{Spin}^c(n) \rightarrow \mathrm{SO}(n)$ is an S^1 -extension, and one must choose a connection on the S^1 -bundle $\mathrm{Spin}^c(M) \rightarrow \mathrm{SO}(M)$. The curvature of this connection appears in the Weizenböck formula for the spin^c Dirac operator. Since $\mathrm{String}(M) \rightarrow \mathrm{Spin}(M)$ has $K(\mathbb{Z}, 2)$ -fibers, string structures are more analogous to spin^c structures. When constructing the S^1 -extension $L\widehat{\mathrm{Spin}}(M) \rightarrow L\mathrm{Spin}(M)$, one really needs an S^1 -extension with connection [Coquereaux and Pilch 1989]. The form $\mathrm{CS}_{\frac{1}{2}p_1}(g) - \pi^*H_{\mathcal{G},g} \in \Omega^3(\mathrm{Spin}(M))$ representing \mathcal{G} transgresses to the curvature (minus a canonically defined term) of this connection on $L\mathrm{Spin}(M)$. One would reasonably expect any Weizenböck-type formula for \not{D}_{LM} to also involve the form $H_{\mathcal{G},g}$. We ask the following question in an attempt to formulate a connection between tmf and obstructions for certain types of curvature.

Question 4.3. Let M be a closed n -dimensional manifold with spin structure such that $\frac{1}{2}p_1(M) = 0 \in H^4(M; \mathbb{Z})$, and let \mathcal{G} be a specified string class. Suppose there exists a metric g such that

$$\mathrm{Ric}(g) > 0 \quad \text{and} \quad H_{\mathcal{G},g} = 0 \in \Omega^3(M).$$

Does this imply that

$$\sigma[M, \mathcal{G}] = 0 \in \mathrm{tmf}^{-n}(pt)?$$

Remark 4.4. The condition $H_{\mathcal{G},g} = 0$ for some string class \mathcal{G} is equivalent to $\check{p}_1(g) = 0 \in \check{H}^4(M)$. This is a strong condition and is not usually satisfied for generic metrics. While a great deal of information about the characters $\check{p}_1(g)$

is known for certain manifolds, the author is not aware of any general results guaranteeing the existence or nonexistence of such metrics.

Remark 4.5. The condition $H_{\mathcal{F},g} = 0$ is conformally invariant; if $H_{\mathcal{F},g} = 0$, then $H_{\mathcal{F},e^f g} = 0$ for any conformally related metric $e^f g$. This follows from the conformal invariance of $\frac{1}{2}\check{p}_1(g)$ and the fact that $0 \in \mathcal{H}^3(M)$ for all metrics.

We close this discussion by noting that \mathcal{D}_M and \mathcal{D}_{LM} can both be thought of as partition functions of certain 1- and 2-dimensional supersymmetric nonlinear sigma models [Witten 1999]. These sigma models require spin and string structures, respectively. In the 2-dimensional sigma models, the form $H_{\mathcal{F},g}$ is used to trivialize the natural connection on a certain determinant line bundle [Witten 1999; Alvarez and Singer 2002]. Sometimes, terms in the action of these sigma models are combined and written as the connection $\nabla^{\mathcal{F},g}$ discussed in Section 5.

Stolz and Teichner [2004] have shown that KO^{-n} is homotopy equivalent to the space of supersymmetric 1-dimensional Euclidean field theories of degree n , and the spin orientation is (up to homotopy) given by the previously mentioned sigma model. The hope is that the analogous statement should hold for 2-dimensional field theories with the string orientation σ given by these sigma models. In this context, Question 4.3 is essentially asking, If one does not have to add in the terms $H_{\mathcal{F},g}$, does positivity of the Ricci curvature imply that the corresponding sigma model is qualitatively trivial?

5. Metric connections with torsion

Question 4.3 can be reformulated in terms of the Ricci curvature of a metric connection with torsion. Given a string class and metric (\mathcal{F}, g) , we define the torsion tensor $T^{\mathcal{F},g}$ by

$$T^{\mathcal{F},g} := g^{-1} H_{\mathcal{F},g} \in \Omega^1(M; \mathfrak{gl}(TM)),$$

where $H_{\mathcal{F},g}$ is the canonical 3-form from (3.21). This is simply a case of “raising indices” and is equivalent to saying $g(T_X^{\mathcal{F},g} Y, Z) = H_{\mathcal{F},g}(X, Y, Z)$, or in coordinates $T_{ij}^k = g^{rk} H_{ijr}$. Then

$$\nabla^{\mathcal{F},g} := \nabla^g + \frac{1}{2} T^{\mathcal{F},g}$$

is a metric connection with torsion $T^{\mathcal{F},g}$, where ∇^g is the Levi-Civita connection.

In general, torsion tensor T of a connection is called *totally skew-symmetric* if $gT \in \Omega^3(M)$, that is, $g(T(\cdot, \cdot), \cdot)$ is skew-symmetric in all three variables. By construction, $\nabla^{\mathcal{F},g}$ is a metric connection with totally skew-symmetric torsion. We also note that the connection $\nabla^{\mathcal{F},g}$ still preserves the geodesics of the Levi-Civita connection. In general for a fixed metric g , we have the following equalities of

subsets of connections on TM :

$$\{\text{Metric connections}\} = \left\{ \begin{array}{l} \text{Metric connections} \\ \text{with } \nabla^g\text{-geodesics} \end{array} \right\} = \left\{ \begin{array}{l} \text{Metric connections with} \\ \text{totally skew-symmetric} \\ \text{torsion} \end{array} \right\}$$

One can easily prove this by writing any connection ∇ as $\nabla^g + A$ and plugging into the geodesic equation $\nabla_X X = 0$ and metric equation $g(\nabla_X Y, Z) = -g(Y, \nabla_X Z)$.

For a torsion connection $\nabla^T = \nabla^g + \frac{1}{2}T$, we can still define the curvature tensor

$$R_{X,Y}^T Z := (\nabla_X^T \nabla_Y^T - \nabla_Y^T \nabla_X^T - \nabla_{[X,Y]}^T) Z,$$

and Ricci tensor

$$\text{Ric}^T(X, Y) := \sum_i g(R_{e_i, X}^T Y, e_i),$$

where $\{e_i\}$ is any orthonormal basis. We let Ric^g denote the Ricci tensor of the Levi-Civita connection.

Lemma 5.1. *Suppose that $\nabla^T = \nabla^g + \frac{1}{2}T$ is a metric connection with totally skew-symmetric torsion satisfying $gT = H \in \Omega^3(M)$. Then the Ricci tensor satisfies*

$$\text{Ric}^T(X, Y) = \text{Ric}^g(X, Y) - \frac{1}{4} \sum_i g(T_{e_i} X, T_{e_i} Y) - \frac{1}{2} d^* H(X, Y).$$

Proof. Let $\langle \cdot, \cdot \rangle$ denote $g(\cdot, \cdot)$. Simply expanding using $\nabla^T = \nabla^g + \frac{1}{2}T$, we get

$$\begin{aligned} \langle R_{e_i, X}^T Y, e_i \rangle &= \langle \nabla_{e_i}^T \nabla_X^T Y - \nabla_X^T \nabla_{e_i}^T Y - \nabla_{[e_i, X]}^T Y, e_i \rangle \\ &= \langle R_{e_i, X}^g Y, e_i \rangle - \frac{1}{4} \langle T_X T_{e_i} Y, e_i \rangle + \frac{1}{2} \langle \nabla_{e_i}^g T_X Y - T_X \nabla_{e_i}^g Y - T_{\nabla_{e_i}^g X} Y, e_i \rangle \\ &\quad + \frac{1}{2} \langle T_{\nabla_{e_i}^g X - \nabla_X^g e_i - [e_i, X]} Y, e_i \rangle \\ &= \langle R_{e_i, X}^g Y, e_i \rangle - \frac{1}{4} \langle T_{e_i} X, T_{e_i} Y \rangle + \frac{1}{2} \langle (\nabla_{e_i}^g T)(X, Y), e_i \rangle. \end{aligned}$$

The last term is easily seen to be a tensor. Using a normal orthonormal frame $\{e_i\}$ at a point (that is, $\nabla_{e_i} e_j = 0$), one easily calculates that

$$\sum_i \langle (\nabla_{e_i}^g T)(e_j, e_k), e_i \rangle = \sum_i \partial_i T_{jk}^i = \sum_i \partial_i H_{ijk} = -d^* H(e_j, e_k). \quad \square$$

The usual Ricci tensor Ric^g is symmetric, and Lemma 5.1 shows that the skew-symmetric part of Ric^T is $-\frac{1}{2}d^*H$. Because the canonical form $H_{\mathcal{F},g}$ satisfies $d^*H_{\mathcal{F},g} = 0$, it gives rise to a metric connection $\nabla^{\mathcal{F},g}$ with symmetric Ricci tensor. For an arbitrary metric connection ∇^T , we refer to the Ricci curvature $\text{Ric}^T(X) := \text{Ric}^T(X, X)$ as the symmetric component of Ric^T , which satisfies

$$\text{Ric}^g(X) - \text{Ric}^T(X) = \frac{1}{4} \sum_i \|T_{e_i} X\|^2 \geq 0,$$

with equality for all X precisely when $T = 0$. This gives an alternative description of the Levi-Civita connection.

Corollary 5.2. *For a fixed Riemannian metric g , the Levi-Civita connection is the unique metric connection that maximizes the Ricci curvature.*

One convenient property of both the Levi-Civita connection and the usual Ricci tensor is the invariance under a global scaling. A quick check shows that for $\epsilon > 0$,

$$\text{Ric}^{\epsilon g}(X) = \sum_i \epsilon g(R_{\epsilon^{-1/2}e_i, X}^{\epsilon g} X, \epsilon^{-1/2}e_i) = \sum_i g(R_{e_i, X}^g X, e_i) = \text{Ric}^g(X).$$

The form $H_{\mathcal{F}, g}$ was constructed using a Hodge isomorphism in an adiabatic limit, giving us the scale invariance $H_{\mathcal{F}, \epsilon g} = H_{\mathcal{F}, g}$. However, we use the metric to change $H_{\mathcal{F}, g}$ into a torsion tensor. Therefore,

$$T^{\mathcal{F}, \epsilon g} = (\epsilon g)^{-1} H_{\mathcal{F}, \epsilon g} = \epsilon^{-1} T^{\mathcal{F}, g}.$$

It is more natural then to consider the 1-parameter family of connections $\nabla^{\mathcal{F}, \epsilon g}$ than any fixed $\nabla^{\mathcal{F}, g}$. In the large volume limit, as $\epsilon \rightarrow \infty$, this connection converges to the Levi-Civita connection ∇^g . In the small volume limit, as $\epsilon \rightarrow 0$, the terms $T^{\mathcal{F}, \epsilon g}$ blow up and $\nabla^{\mathcal{F}, \epsilon g}$ does not converge to a connection unless $H_{\mathcal{F}, g} = 0$.

Question 5.3. Let (M, g, \mathcal{F}) be an n -dimensional Riemannian manifold with string class. Suppose that the Ricci tensor of the modified connection $\nabla^{\mathcal{F}, g}$ is strictly positive in the small volume scaling limit; that is

$$\lim_{\epsilon \rightarrow 0} \text{Ric}(\nabla^{\mathcal{F}, \epsilon g}) > 0.$$

Does this imply $\sigma[M, \mathcal{F}] = 0 \in \text{tmf}^{-n}(pt)$?

Proposition 5.4. *Question 4.3 is equivalent to Question 5.3.*

Proof. This follows directly from the description of the Ricci tensor in Lemma 5.1, which implies

$$\text{Ric}^{\epsilon g, \mathcal{F}}(X) = \text{Ric}^g(X) - \frac{1}{4}\epsilon \sum_i \|T_{e_i} X\|^2.$$

Consequently, if $H_{\mathcal{F}, g} \neq 0$, then $\text{Ric}^{\epsilon g, \mathcal{F}}(X) \xrightarrow{\epsilon \rightarrow 0} -\infty$ for some X . The simultaneous conditions $\text{Ric}(g) > 0$ and $H_{\mathcal{F}, g} = 0$ are equivalent to $\text{Ric}(\nabla^{\mathcal{F}, \epsilon g}) > 0$ for arbitrarily small ϵ . □

6. Homogeneous metrics on S^3

We now examine the canonical 3-forms obtained when $M = S^3$ with a homogeneous metric, and we compare the results with Question 4.3. We see that it has an affirmative answer in this special situation, but it would not if the conditions were weakened. In particular, there exists a 1-dimensional family of left-invariant metrics g with nonnegative Ricci curvature such that the right-invariant framing \mathcal{R} produces $H_{\mathcal{R}, g} = 0$ and $\sigma[S^3, \mathcal{R}] \neq 0 \in \text{tmf}^{-3}(pt)$. The previous sentence is also true with left and right swapped.

6a. String structures on S^3 . Using the isomorphism $S^3 \cong \text{SU}(2) \cong \text{Sp}(1)$, the left- and right-invariant framings induce two string classes, which we denote \mathcal{L} and \mathcal{R} . The disc D^4 inherits a standard framing from its inclusion $D^4 \subset \mathbb{R}^4$, and this restricts to a framing of the stable tangent bundle for $\partial D^4 = S^3$. We denote the induced string class by ∂D^4 and note that, by construction, the string bordism class $[S^3, \partial D^4] = 0 \in M\text{String}^{-3}(pt)$.

The set of string classes is a torsor for $H^3(S^3; \mathbb{Z}) \cong \mathbb{Z}$, an affine copy of \mathbb{Z} . In other words, the difference between any two string classes is naturally an integer. We now determine where the three previously defined string classes live on this affine line, and we use $\Omega c_2 \in H^3(S^3; \mathbb{Z})$ as our standard generator. The left and right framings are related by

$$S^3 \times \text{Spin}(3) \xrightarrow{L} \text{Spin}(S^3) \xleftarrow{R} S^3 \times \text{Spin}(3),$$

and the composition $R^{-1} \circ L$ is the Adjoint representation lifted to Spin :

$$S^3 \cong \text{SU}(2) \xrightarrow{\text{Ad}} \text{Spin}(\mathfrak{su}(2)) \cong \text{Spin}(3).$$

The difference $\mathcal{L} - \mathcal{R}$ is equal to $\pi^*(\text{Ad}^* \Omega \frac{1}{2} p_1)$. The Adjoint representation here is an isomorphism of Lie groups and hence an isomorphism on cohomology. As mentioned in Remark 2.17, there is a factor of 2 and minus sign at work: The class $\Omega \frac{1}{2} p_1$ is twice a generator of $H^3(S^3; \mathbb{Z})$, and stably $p_1 = -c_2$. Hence $\Omega \frac{1}{2} p_1$ is mapped to $-2\Omega c_2$, or $-2 \in \mathbb{Z} \cong H^3(S^3; \mathbb{Z})$, and we use the shorthand $\mathcal{L} + 2 = \mathcal{R}$.

Similarly, we examine the difference between the left-framing and the bounding string structure, and in doing so reference Remark 2.17. The string structure induced from D^4 is a framing of the stable tangent bundle. The normal bundle $\nu \rightarrow S^3$ is trivial, and we have the standard isomorphisms of bundles over S^3 :

$$\text{Spin}(TS^3 \oplus \mathbb{R}) \cong \text{Spin}(TS^3 \oplus \nu) \cong \text{Spin}(D^4) \cong \text{Spin}(4).$$

The difference in framing of the two stable bundles differs by the left-multiplication map $S^3 \rightarrow \text{Spin}(4)$ given by considering S^3 as the unit quaternions. Under the standard isomorphisms $S^3 \cong \text{SU}(2)$ and $\text{Spin}(4) \cong \text{SU}(2) \times \text{SU}(2)$, this left-multiplication map is the inclusion into the first factor:

$$\text{SU}(2) \xrightarrow{\text{Id} \times \{1\}} \text{SU}(2) \times \text{SU}(2) \cong \text{Spin}(4).$$

The induced map on cohomology sends $\Omega \frac{1}{2} p_1$ to $-\Omega c_2$, or $-1 \in \mathbb{Z} \cong H^3(S^3; \mathbb{Z})$. Therefore,

$$\mathcal{L} + 1 = \partial D^4 \quad \text{and} \quad \mathcal{L} + 2 = \partial D^4 + 1 = \mathcal{R}.$$

The Adams e -invariant gives an isomorphism $\pi_3^s \xrightarrow{\cong} \mathbb{Z}/24$ and sends the left and right framings to the two generators [Atiyah and Smith 1974]. Our calculations also verify this explicitly. On a framed $(4k-1)$ -dimensional manifold M , the e -invariant

can be computed as follows. Choose a spin manifold W such that $\partial W = M$ as spin manifolds; such a manifold exists because $M\text{Spin}^{4k-1}(pt) = 0$. Using the framing of TM , define the Pontryagin classes $p_i(W, M)$ as relative classes in $H^*(W, M)$. We then obtain $\hat{A}(W, M)$ by evaluating $\hat{A}(TW, TM)$ on the fundamental class of W , where $\hat{A}(TW, TM)$ is the \hat{A} -polynomial with relative Pontryagin classes. Then,

$$e[M] = \begin{cases} \hat{A}(W, M) \pmod{\mathbb{Z}} & \text{for } k \text{ even,} \\ \frac{1}{2}\hat{A}(W, M) \pmod{\mathbb{Z}} & \text{for } k \text{ odd.} \end{cases}$$

The e -invariant is well-defined as an element of \mathbb{Q}/\mathbb{Z} , since choosing a different W' will give $\hat{A}(W', M) - \hat{A}(W, M) = \hat{A}(W' \cup_M (-W))$, which is an integer (or even integer) by the Atiyah–Singer index theorem.

If we include metrics so that (W, \tilde{g}) is a Riemannian spin manifold with boundary (M, g) , then we naturally have the Pontryagin forms $p_i(\tilde{g}) \in \Omega^{4k}(W)$.

Proposition 6.1. *If (M, g, \mathcal{F}) is a Riemannian spin 3-manifold with string class, then*

$$e(M, \mathcal{F}) = -\frac{1}{48} \int_W p_1(\tilde{g}) + \frac{1}{24} \int_M H_{\mathcal{F}, g} \pmod{\mathbb{Z}}.$$

Proof.

$$e[M, \mathcal{F}] = \frac{1}{2} \int_W \hat{A}(W, M) = \frac{1}{2} \int_W (1 - \frac{1}{24} p_1(W, M) + \dots) = -\frac{1}{48} \int_W p_1(W, M).$$

We now construct a de Rham representative of $p_1(W, M)$. If $\partial W = M$, then consider the bordism $W \cup_M ([0, 1] \times M)$ obtained by gluing ∂W to $\{0\} \times M$. The string class \mathcal{F} gives a stable trivialization p of $\text{Spin}(TM)$ up to homotopy, and we let Θ_p denote the induced flat connection. Denoting the Levi-Civita connection on $\text{Spin}(TM)$ by Θ_g , we have the connection $\Theta(t)$ on $[0, 1] \times M$, where

$$\Theta(t) = t\Theta_p + (1 - t)\Theta_g.$$

Finally, define $\tilde{\Theta}$ to be the connection on $\text{Spin}(W \cup_M ([0, 1] \times M))$ induced by $\Theta_{\tilde{g}}$ and $\Theta(t)$. The form $p_1(\tilde{\Theta})$ is a de Rham representative of $p_1(W, M)$, and

$$\begin{aligned} \int_{\tilde{W}} p_1(\tilde{\Theta}) &= \int_W p_1(\tilde{g}) + \int_{M^3} \int_{[0, 1]} p_1(\Theta(t)) \\ &= \int_W p_1(\tilde{g}) + \int_M \text{CS}_{p_1}(\Theta_p, \Theta_g) = \int_W p_1(\tilde{g}) - 2 \int_M \text{CS}_{\frac{1}{2}p_1}(\Theta_g, \Theta_p), \end{aligned}$$

where $\text{CS}_\lambda(\Theta_g, \Theta_p)$ is the general Chern–Simons transgression between two connections. Lemmas 6.3 and 3.18 together imply

$$\int_M \text{CS}_{\frac{1}{2}p_1}(\Theta_g, \Theta_p) = \int_M p^* \text{CS}(\Theta_g) = \int_M H_{\mathcal{F}, g}.$$

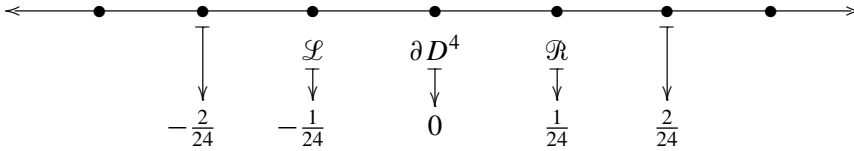
Therefore,

$$\begin{aligned}
 -\frac{1}{48} \int_{\tilde{W}} p_1(\tilde{\Theta}) &= -\frac{1}{48} \int_W p_1(\tilde{g}) + \frac{1}{24} \int_M \text{CS}_{\frac{1}{2}p_1}(\Theta_g, \Theta_p) \\
 &= -\frac{1}{48} \int_W p_1(\tilde{g}) + \frac{1}{24} \int_M H_{\mathcal{G},g}. \quad \square
 \end{aligned}$$

Corollary 6.2. *When $M = S^3$ and g is the standard round metric,*

$$e(S^3, \mathcal{G}) = \frac{1}{24} \int_{S^3} H_{\mathcal{G},g} \pmod{\mathbb{Z}}.$$

In the next subsection, we calculate $H_{\mathcal{G},g}$ for all left-invariant metrics on S^3 . Equation (6.8) and the corollary above imply that $e[S^3, \mathcal{L}] = -\frac{1}{24}$, $e[S^3, \partial D^4] = 0$, and $e[S^3, \mathcal{R}] = \frac{1}{24}$. Below is a pictorial description of the space of string classes on S^3 and their corresponding string bordism class under $e : M\text{String}^{-3} \xrightarrow{\cong} \mathbb{Z}/24$.



Lemma 6.3. *If $p : M \rightarrow P$ is a global section and Θ_p the induced flat connection, then*

$$\text{CS}_\lambda(\Theta, \Theta_p) = p^* \text{CS}_\lambda(\Theta) \in \Omega^3(M).$$

Proof. This lemma is essentially a tautology. Using the notation of [Freed 2002], in general $\text{CS}_\lambda(\Theta_1, \Theta_0) := \int_{[0,1]} \lambda(\Theta_t) \in \Omega^{2i-1}(M)$, where $\Theta_t := t\Theta_1 + (1-t)\Theta_0$ is a connection on $[0, 1] \times P \rightarrow [0, 1] \times M$. Then,

$$\text{CS}_\lambda(\Theta) := \text{CS}_\lambda(\pi^*\Theta, \Theta_{\text{taut}}) \in \Omega^{2i-1}(P),$$

where Θ_{taut} is the trivial connection induced by the canonical section of π^*P . Since one can compute these transgression forms via local frames, and by definition $p^*\Theta_p = 0$, we easily see

$$\begin{aligned}
 \text{CS}_\lambda(\Theta, \Theta_p) &= \int_{[0,1]} \lambda(p^*(t\Theta + (1-t)\Theta_p)) = \int_{[0,1]} \lambda(p^*t\Theta) \\
 &= p^* \int_{[0,1]} \lambda(t\Theta) = p^* \text{CS}_\lambda(\Theta). \quad \square
 \end{aligned}$$

6b. Calculation of canonical 3-forms. We now investigate Question 4.3 by considering left-invariant metrics on $S^3 \cong \text{SU}(2)$; that is, metrics g on $\text{SU}(2)$ such that left multiplication is an isometry. As noted in Proposition 6.11, the calculations for right-invariant metrics only differ from those for left-invariant metrics by a sign. Any such left-invariant metric is determined by its behavior on the tangent space at

the identity, so we are considering metrics on the Lie algebra $\mathfrak{su}(2)$ of left-invariant vector fields. A global rescaling of g leaves the Ricci tensor and canonical form $H_{\mathcal{L},g}$ invariant; hence it does not affect the outcome of Question 4.3. The space of left-invariant metrics, up to change of oriented basis and global rescaling, is the 2-dimensional space $\text{Sym}_{>0}^2(\mathbb{R}^3)/(\text{SO}(\mathbb{R}^3) \times \mathbb{R}_+)$, where $\text{Sym}_{>0}^2(\mathbb{R}^3)$ denotes the 6-dimensional space of positive-definite 3×3 -matrices.

We now give a more computationally explicit description of this space. Let $\{e_1, e_2, e_3\}$ be the standard basis for $\mathfrak{su}(2)$ satisfying

$$[e_1, e_2] = 2e_3, [e_2, e_3] = 2e_1, [e_3, e_1] = 2e_2.$$

When $\{e_1, e_2, e_3\}$ is an orthonormal basis, the metric is biinvariant and equal to the standard round metric on $S^3 \subset D^4$. For any $\alpha_1, \alpha_2 \in \mathbb{R}_{>0}$, define the left-invariant metric g_{α_1, α_2} by declaring $\{\alpha_1 e_1, \alpha_2 e_2, e_3\}$ to be an orthonormal basis. In the case where $\alpha_2 = 1$, we recover the 1-parameter family of Berger metrics on S^3 . Based on knowledge from [Milnor 1976], it suffices to consider the 2-parameter family of metrics $\{g_{\alpha_1, \alpha_2}\}$.

Lemma 6.4. *If g is a left-invariant metric on $SU(2)$, then there exists $\alpha_1, \alpha_2 \in \mathbb{R}_{>0}$ such that g_{α_1, α_2} is isometric to a constant multiple of g .*

Proof. Lemma 4.1 [Milnor 1976] implies that there exists an orthonormal basis $\{E_1, E_2, E_3\}$ for g such that

$$[E_1, E_2] = \lambda_3 E_3, \quad [E_2, E_3] = \lambda_1 E_1, \quad [E_3, E_1] = \lambda_2 E_2,$$

where $\lambda_i \in \mathbb{R}_{>0}$. (Milnor's e_i correspond to our E_i .) For any $(\lambda_1, \lambda_2, \lambda_3)$, it is clear that the orthonormal basis

$$\left\{ \frac{1}{2} \sqrt{\lambda_2 \lambda_3} e_1, \frac{1}{2} \sqrt{\lambda_3 \lambda_1} e_2, \frac{1}{2} \sqrt{\lambda_1 \lambda_2} e_3 \right\}$$

defines a left-invariant metric isometric to the original g . Finally, we normalize so that the coefficient of e_3 is 1. Hence, there is a surjective map

$$\mathbb{R}_{>0}^2 \rightarrow \{\text{Left-invariant metrics}\} / \{\text{Isom} \times \text{Scale}\}, \quad \alpha_1, \alpha_2 \mapsto g_{\alpha_1, \alpha_2}. \quad \square$$

We first calculate the Ricci curvature for g_{α_1, α_2} . A straightforward computation gives the covariant derivative of the Levi-Civita connection in our invariant frame. The nonzero components are

$$\begin{aligned} \langle \nabla_{\alpha_1 e_1} \alpha_2 e_2, e_3 \rangle &= \alpha_1 \alpha_2 + \alpha_1 / \alpha_2 - \alpha_2 / \alpha_1, \\ \langle \nabla_{\alpha_2 e_2} \alpha_3 e_3, \alpha_1 e_1 \rangle &= \alpha_1 \alpha_2 - \alpha_1 / \alpha_2 + \alpha_2 / \alpha_1, \\ \langle \nabla_{\alpha_3 e_3} \alpha_1 e_1, \alpha_1 e_2 \rangle &= -\alpha_1 \alpha_2 + \alpha_1 / \alpha_2 + \alpha_2 / \alpha_1. \end{aligned}$$

The Ricci tensor is then diagonalized with eigenvalues

$$\begin{aligned} \text{Ric}(\alpha_1 e_1) &= 2(\alpha_1 \alpha_2 - \alpha_1/\alpha_2 + \alpha_2/\alpha_1)(-\alpha_1 \alpha_2 + \alpha_1/\alpha_2 + \alpha_2/\alpha_1), \\ \text{Ric}(\alpha_1 e_2) &= 2(-\alpha_1 \alpha_2 + \alpha_1/\alpha_2 + \alpha_2/\alpha_1)(\alpha_1 \alpha_2 + \alpha_1/\alpha_2 - \alpha_2/\alpha_1), \\ \text{Ric}(e_3) &= 2(\alpha_1 \alpha_2 + \alpha_1/\alpha_2 - \alpha_2/\alpha_1)(\alpha_1 \alpha_2 - \alpha_1/\alpha_2 + \alpha_2/\alpha_1). \end{aligned}$$

Solving inequalities tells us that the Ricci curvature is strictly positive if and only if (α_1, α_2) is in the interior of the region bounded by the three curves

$$(6.5) \quad \alpha_2 = \sqrt{\frac{\alpha_1^2}{1 + \alpha_1^2}}, \quad \alpha_2 = \sqrt{\frac{\alpha_1^2}{-1 + \alpha_1^2}}, \quad \alpha_2 = \sqrt{\frac{-\alpha_1^2}{-1 + \alpha_1^2}}.$$

This region is shown in Figure 1, left. The Ricci curvature is nonnegative with one zero eigenvalue on the three boundary curves.

Now we calculate the canonical 3-form $H_{\mathcal{L}, g_{\alpha_1, \alpha_2}} \in \Omega^3(S^3)$. For dimensional reasons, $H_{\mathcal{L}, g_{\alpha_1, \alpha_2}}$ is harmonic and therefore

$$H_{\mathcal{L}, g_{\alpha_1, \alpha_2}} \in \mathcal{H}^3(S^3) \cong H^3(S^3; \mathbb{R}) \cong \mathbb{R}$$

with its value in \mathbb{R} determined by integrating over S^3 . Lemma 3.18 says that we can calculate $H_{\mathcal{L}, g_{\alpha_1, \alpha_2}}$ by simply calculating the Chern–Simons 3-form $\text{CS}_{\frac{1}{2}p_1}(g_{\alpha_1, \alpha_2})$ on the global frame $\{\alpha_1 e_1, \alpha_2 e_2, e_3\}$. This is a straightforward, though lengthy, calculation.

For the class $\frac{1}{2}p_1$, the Chern–Simons form is

$$\text{CS}_{\frac{1}{2}p_1}(\Theta) = -\frac{1}{16\pi^2} \text{Tr}(\Omega \wedge \Theta - \frac{1}{6}\Theta \wedge [\Theta \wedge \Theta]),$$

with Tr being the ordinary matrix trace. The normalization constant can be seen from $\frac{1}{2}p_1(\Theta) = -\frac{1}{2}c_2(\Theta) = -\frac{1}{2} \frac{1}{8\pi^2} \text{Tr}(\Omega \wedge \Omega)$. The frame $\{e_i\}$ gives rise to the dual frame $\{e^i\}$ on $\mathfrak{su}(2)^*$. In our global frame, the Chern–Simons form is a constant multiple of $e^1 \wedge e^2 \wedge e^3$, the standard volume form for $\text{SU}(2) \cong S^3 \subset D^4$. Using a direct calculation along with $\int_{S^3} e^1 \wedge e^2 \wedge e^3 = 2\pi^2$, we obtain

$$(6.6) \quad \begin{aligned} \int_{S^3} H_{\mathcal{L}, g_{\alpha_1, \alpha_2}} &= -\frac{1}{16\pi^2} \int_{S^3} \text{Tr}(\Theta \wedge \Omega - \frac{1}{6}\Theta \wedge [\Theta \wedge \Theta]) \\ &= -\frac{\alpha_1^6 \alpha_2^6 - \alpha_1^6 \alpha_2^4 - \alpha_1^4 \alpha_2^6 - \alpha_1^6 \alpha_2^2 - \alpha_1^2 \alpha_2^6 - \alpha_1^4 \alpha_2^2 - \alpha_1^2 \alpha_2^4 + 4\alpha_1^4 \alpha_2^4 + \alpha_1^6 + \alpha_2^6}{\alpha_1^4 \alpha_2^4}. \end{aligned}$$

See Figure 1, right, for a graph of this function. If we set $\alpha_2 = 1$ and only consider the usual Berger metrics, we obtain

$$(6.7) \quad \int_{S^3} H_{\mathcal{L}, g_{\alpha_1, 1}} = -2 + \frac{2\alpha_1^2 - 1}{\alpha_1^4}.$$

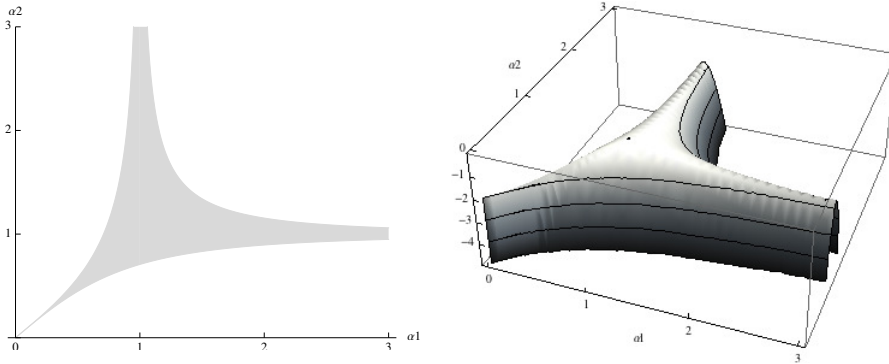


Figure 1. At left: Region with positive Ricci curvature. At right: Values of $\int_{S^3} H_{\mathcal{L},g_{\alpha_1,\alpha_2}}$.

These values are graphed in Figure 2, left. Note that when reduced mod \mathbb{Z} , (6.7) coincides with the calculation performed in the original [Chern and Simons 1974]. If we set $\alpha_1 = \alpha_2 = 1$, we obtain the standard biinvariant metric and see that

$$(6.8) \quad \int_{S^3} H_{\mathcal{L},g_{1,1}} = -1, \quad \int_{S^3} H_{\partial D^4,g_{1,1}} = 0, \quad \int_{S^3} H_{\mathcal{R},g_{1,1}} = 1.$$

We now analyze (6.6) on the region $\text{Ric} \geq 0$. The only critical point occurs at $\alpha_1 = \alpha_2 = 1$, where $\int_{S^3} H_{\mathcal{L},g_{\alpha_1,\alpha_2}} = -1$ is a maximal value. Furthermore, $\int_{S^3} H_{\mathcal{L},g_{\alpha_1,\alpha_2}} = -2$ identically on the three curves bounding the region of positive Ricci curvature. So, we have the range of values

$$\{\int_{S^3} H_{\mathcal{L},g} \mid \text{Ric}(g) > 0, \text{ } g \text{ left-invariant}\} = (-2, -1].$$

Figure 1, right, demonstrates this with the help of Mathematica; the level curves for -2 are precisely the three functions from (6.5).

Due to the equivariance of the canonical 3-form under change of string class (see Proposition 3.12), our calculation using \mathcal{L} gives us $H_{\mathcal{G},g_{\alpha_1,\alpha_2}}$ for any other string class \mathcal{G} by

$$\int_{S^3} H_{\mathcal{G}+j,g_{\alpha_1,\alpha_2}} = j + \int_{S^3} H_{\mathcal{L},g_{\alpha_1,\alpha_2}} \quad \text{for any } j \in \mathbb{Z} \cong H^3(S^3; \mathbb{Z}).$$

Therefore,

$$(6.9) \quad \{\int_{S^3} H_{\mathcal{G}+j,g} \mid \text{Ric}(g) > 0, \text{ } g \text{ left-invariant}\} = (-2 + j, -1 + j].$$

To graphically demonstrate this, Figure 2, left, shows the canonical 3-forms for various string classes on the 1-parameter family of left-invariant Berger metrics.

The entire previous discussion was based on left-invariant Riemannian metrics. What if we had decided to use right-invariant metrics? Given an inner product g_e

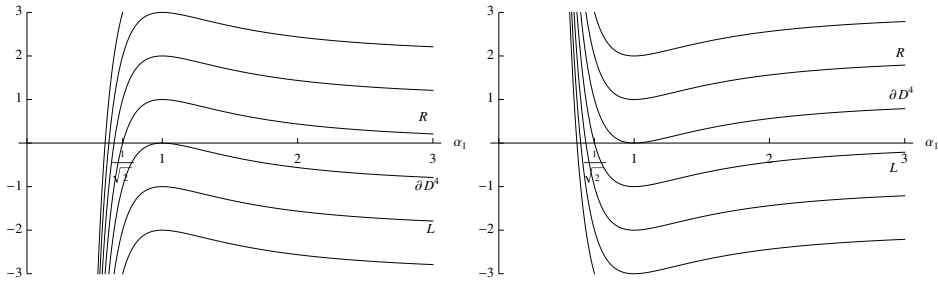


Figure 2. $\int_{S^3} H_{\mathcal{G}, g_{\alpha_1, \alpha_2}}$ on Berger metrics, for left- and right-invariant metrics, respectively.

on $T_e \text{SU}(2)$, we can form a left-invariant metric g^L and a right-invariant metric g^R by left or right multiplying g_e . The canonical 3-forms are related by the following easy lemma, whose proof is at the end of this section.

Lemma 6.10. $H_{\mathcal{G}, g^L} = -H_{\mathcal{R}, g^R}$.

This fact is graphically demonstrated in Figure 2, right. In the case of the Berger metrics, note that the Ricci curvature is positive for all $\alpha_1 > 1/\sqrt{2}$, and the Ricci curvature is nonnegative with a 0 eigenvalue at $\alpha_1 = 1/\sqrt{2}$.

Proposition 6.11. *Suppose the string class and (left or right)-invariant Riemannian metric (\mathcal{S}, g) on S^3 satisfy*

$$\text{Ric}(g) > 0 \quad \text{and} \quad H_{\mathcal{S}, g} = 0.$$

Then $\mathcal{S} = \partial D^4$ and g is the biinvariant round metric. Consequently,

$$\sigma[S^3, \mathcal{S}] = 0 \in \text{tmf}^{-3}(pt) \cong \mathbb{Z}/24.$$

Proof. If g is a left-invariant metric with positive Ricci curvature and $H_{\mathcal{S}, g} = 0$, then (6.9) implies that $\mathcal{S} = \mathcal{L} + 1 = \partial D^4$ with g the biinvariant metric $g_{1,1}$.

If g is a right-invariant metric, Lemma 6.10 and (6.9) imply that

$$(6.12) \quad \left\{ \int_{S^3} H_{\mathcal{R}+j, g} \mid \text{Ric}(g) > 0, \quad g \text{ right-invariant} \right\} = [1 + j, 2 + j].$$

If $H_{\mathcal{S}, g} = 0$, then $\mathcal{S} = \mathcal{R} - 1 = \partial D^4$ and $g = g_{1,1}$. Finally, $[S^3, \partial D^4] = 0 \in M\text{String}^{-3}$, so $\sigma[S^3, \partial D^4] = 0 \in \text{tmf}^{-3}$. □

We conclude that in this case, Question 4.3 has a very nontrivial affirmative answer. In particular, there are 1-dimensional families of left- and right-invariant metrics that are Ricci nonnegative and satisfy $H_{\mathcal{R}, g} = 0$ and $H_{\mathcal{L}, g} = 0$, respectively. Furthermore, as evidenced by Figure 2, one can find Ricci positive metrics with

$H_{\mathcal{S},g}$ arbitrarily small but nonzero. Finally, we point out that for any string class \mathcal{S} , the lift of the Chern–Simons invariant

$$\text{Met}(S^3) \xrightarrow{\int H_{\mathcal{S},g}} \mathbb{R}$$

is surjective. The 1-parameter families of left- and right-invariant Berger metrics in Figure 2 show this.

Proof of Lemma 6.10. In a left- or right-invariant frame, the connection is computed purely in terms of the Lie bracket on vector fields. On a Lie group G , one can define two Lie algebra structures $[\cdot, \cdot]_L$ and $[\cdot, \cdot]_R$ corresponding to the usual Lie bracket on left- or right-invariant vector fields. For $X, Y \in T_e G$, these are related by

$$[X, Y]_L = -[X, Y]_R.$$

If Θ_L, Θ_R denote the connections in the two frames, we have $\Theta_L = -\Theta_R$ and $\Omega_L = \Omega_R$, so

$$\text{Tr}(\Theta_L \wedge \Omega_L - \frac{1}{6} \Theta_L \wedge [\Theta_L \wedge \Theta_L]) = -\text{Tr}(\Theta_R \wedge \Omega_R - \frac{1}{6} \Theta_R \wedge [\Theta_R \wedge \Theta_R]). \quad \square$$

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