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**ENERGY AND VOLUME OF VECTOR FIELDS  
ON SPHERICAL DOMAINS**

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# ENERGY AND VOLUME OF VECTOR FIELDS ON SPHERICAL DOMAINS

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We present a “boundary version” for theorems about minimality of volume and energy functionals on a spherical domain of an odd-dimensional Euclidean sphere.

## 1. Introduction

Let  $(M, g)$  be a closed,  $n$ -dimensional Riemannian manifold and  $T^1 M$  the unit tangent bundle of  $M$  considered as a closed Riemannian manifold with the Sasaki metric. Let  $X : M \rightarrow T^1 M$  be a unit vector field defined on  $M$ , regarded as a smooth section of the unit tangent bundle  $T^1 M$ . The volume of  $X$  was defined in [Gluck and Ziller 1986] by  $\text{vol } X := \text{vol } X(M)$ , where  $\text{vol } X(M)$  is the volume of the submanifold  $X(M) \subset T^1 M$ . Using an orthonormal local frame  $\{e_1, e_2, \dots, e_{n-1}, e_n = X\}$ , the volume of the unit vector field  $X$  is given by

$$\begin{aligned} \text{vol } X = \int_M & \left( 1 + \sum_{a=1}^n \|\nabla_{e_a} X\|^2 + \sum_{a < b} \|\nabla_{e_a} X \wedge \nabla_{e_b} X\|^2 + \dots \right. \\ & \left. + \sum_{a_1 < \dots < a_{n-1}} \|\nabla_{e_{a_1}} X \wedge \dots \wedge \nabla_{e_{a_{n-1}}} X\|^2 \right)^{1/2} v_M(g) \end{aligned}$$

and the energy of the vector field  $X$  is given by

$$\mathcal{E}(X) = \frac{n}{2} \text{vol } M + \frac{1}{2} \int_M \sum_{a=1}^n \|\nabla_{e_a} X\|^2 v_M(g).$$

The Hopf vector fields on  $\mathbb{S}^{2k+1}$  are unit vector fields tangent to the classical Hopf fibration  $\mathbb{S}^1 \hookrightarrow \mathbb{S}^{2k+1}$ . The following theorems gives a characterization of Hopf flows as absolute minima of volume and energy functionals:

**Theorem 1** [Gluck and Ziller 1986]. *The unit vector fields of minimum volume on the sphere  $\mathbb{S}^3$  are precisely the Hopf vector fields and no others.*

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Keywords: energy of vector fields, volume of vector fields, Hopf flow.

**Theorem 2 [Brito 2000].** *The unit vector fields of minimum energy on the sphere  $\mathbb{S}^3$  are precisely the Hopf vector fields and no others.*

We prove in this paper the following boundary version for these theorems:

**Theorem 3.** *Let  $U$  be an open set of the  $(2k+1)$ -dimensional unit sphere  $\mathbb{S}^{2k+1}$  and let  $K \subset U$  be a connected  $(2k+1)$ -submanifold with boundary of the sphere  $\mathbb{S}^{2k+1}$ . Let  $\vec{v}$  be an unit vector field on  $U$  which coincides with a Hopf flow  $H$  along the boundary of  $K$ . Then*

$$\mathcal{E}(\vec{v}) \geq \left( \frac{2k+1}{2} + \frac{k}{2k-1} \right) \text{vol } K \quad \text{and} \quad \text{vol } \vec{v} \geq \frac{4^k}{\binom{2k}{k}} \text{vol } K.$$

(Other results for higher dimensions may be found in [Brito et al. 2004; Borrelli and Gil-Medrano 2006; Chacón et al. 2001].)

## 2. Preliminaries

Let  $U \subset \mathbb{S}^{2k+1}$  be an open set of the unit sphere and let  $K \subset U$  be a connected  $(2k+1)$ -submanifold with boundary of  $\mathbb{S}^{2k+1}$ . Let  $H$  be a Hopf vector field on  $\mathbb{S}^{2k+1}$  and let  $\vec{v}$  be an unit vector field defined on  $U$ . We also consider the map  $\varphi_t^{\vec{v}} : U \rightarrow \mathbb{S}^{2k+1}(\sqrt{1+t^2})$  given by  $\varphi_t^{\vec{v}}(x) = x + t\vec{v}(x)$ . This map was introduced in [Asimov 1978; Brito et al. 1981; Milnor 1978].

**Lemma 4.** *For  $t > 0$  sufficiently small, the map  $\varphi_t^{\vec{v}}$  is a diffeomorphism.*

*Proof.* A simple application of the identity perturbation method.  $\square$

From now on, we assume that  $t > 0$  is small enough so that the map  $\varphi_t^{\vec{v}}$  is a diffeomorphism. In order to find the Jacobian matrix of  $\varphi_t^{\vec{v}}$ , we define the unit vector field  $\vec{u}$  on  $\varphi_t^{\vec{v}}(U) \subset \mathbb{S}^{2k+1}(\sqrt{1+t^2})$  by

$$\vec{u}(x) := \frac{1}{\sqrt{1+t^2}} \vec{v}(x) - \frac{t}{\sqrt{1+t^2}} x.$$

Using an adapted orthonormal frame  $\{e_1, \dots, e_{2k}, \vec{v}\}$  on a neighborhood  $V$  of  $U$ , we obtain an adapted orthonormal frame on  $\varphi_t^{\vec{v}}(V)$  given by  $\{\bar{e}_1, \dots, \bar{e}_{2k}, \vec{u}\}$ , where  $\bar{e}_i = e_i$  for all  $i \in \{1, \dots, 2k\}$ .

In this manner, we can write

$$d\varphi_t^{\vec{v}}(e_1) = \langle d\varphi_t^{\vec{v}}(e_1), e_1 \rangle e_1 + \dots + \langle d\varphi_t^{\vec{v}}(e_1), e_{2k} \rangle e_{2k} + \langle d\varphi_t^{\vec{v}}(e_1), \vec{u} \rangle \vec{u},$$

$$d\varphi_t^{\vec{v}}(e_2) = \langle d\varphi_t^{\vec{v}}(e_2), e_1 \rangle e_1 + \dots + \langle d\varphi_t^{\vec{v}}(e_2), e_{2k} \rangle e_{2k} + \langle d\varphi_t^{\vec{v}}(e_2), \vec{u} \rangle \vec{u},$$

$$\vdots$$

$$d\varphi_t^{\vec{v}}(e_{2k}) = \langle d\varphi_t^{\vec{v}}(e_{2k}), e_1 \rangle e_1 + \dots + \langle d\varphi_t^{\vec{v}}(e_{2k}), e_{2k} \rangle e_{2k} + \langle d\varphi_t^{\vec{v}}(e_{2k}), \vec{u} \rangle \vec{u},$$

$$d\varphi_t^{\vec{v}}(\vec{v}) = \langle d\varphi_t^{\vec{v}}(\vec{v}), e_1 \rangle e_1 + \dots + \langle d\varphi_t^{\vec{v}}(\vec{v}), e_{2k} \rangle e_{2k} + \langle d\varphi_t^{\vec{v}}(\vec{v}), \vec{u} \rangle \vec{u}.$$

Now, by Gauss's equation for the trivial immersion  $\mathbb{S}^{2k+1} \hookrightarrow \mathbb{R}^{2k+2}$ , we have

$$\tilde{\nabla}_Y \vec{v} = d\vec{v}(Y) = \nabla_Y \vec{v} - \langle \vec{v}, Y \rangle x$$

for every vector field  $Y$  on  $\mathbb{S}^{2k+1}$ , and then

$$\langle d\varphi_t^{\vec{v}}(e_1), e_1 \rangle = \langle e_1 + td\vec{v}(e_1), e_1 \rangle = 1 + t\langle \nabla_{e_1} \vec{v}, e_1 \rangle$$

Analogously, we can conclude that

$$\begin{aligned} \langle d\varphi_t^{\vec{v}}(e_i), e_i \rangle &= 1 + t\langle \nabla_{e_i} \vec{v}, e_i \rangle \quad \text{for } i \in \{1, \dots, 2k\}, \\ \langle d\varphi_t^{\vec{v}}(e_i), e_j \rangle &= t\langle \nabla_{e_i} \vec{v}, e_j \rangle \quad \text{for } i, j \in \{1, \dots, 2k\}, i \neq j, \\ \langle d\varphi_t^{\vec{v}}(e_i), \vec{u} \rangle &= 0 \quad \text{for } i \in \{1, \dots, 2k\}, \\ \langle d\varphi_t^{\vec{v}}(\vec{v}), \vec{u} \rangle &= \sqrt{1+t^2}. \end{aligned}$$

By employing the notation  $h_{ij}(\vec{v}) := \langle \nabla_{e_i} \vec{v}, e_j \rangle$  (where  $i, j \in \{1, \dots, 2k\}$ ), we can express the determinant of the Jacobian matrix of  $\varphi_t^{\vec{v}}$  in the form

$$\det(d\varphi_t^{\vec{v}}) = \sqrt{1+t^2} \left( 1 + \sum_{i=1}^{2k} \sigma_i(\vec{v}) t^2 \right),$$

where, by definition, the functions  $\sigma_i$  are the  $i$ -symmetric functions of the  $h_{ij}$ . For instance, if  $k = 1$ , we have

$$\begin{aligned} \sigma_1(\vec{v}) &:= h_{11}(\vec{v}) + h_{22}(\vec{v}), \\ \sigma_2(\vec{v}) &:= h_{11}(\vec{v})h_{22}(\vec{v}) - h_{12}(\vec{v})h_{21}(\vec{v}). \end{aligned}$$

### 3. Proof of the Theorem

The energy of the vector field  $\vec{v}$  (on  $K$ ) is given by

$$\mathcal{E}(\vec{v}) := \frac{1}{2} \int_K \|d\vec{v}\|^2 = \frac{2k+1}{2} \operatorname{vol} K + \frac{1}{2} \int_K \|\nabla \vec{v}\|^2$$

Using the notation above, we have

$$\mathcal{E}(\vec{v}) = \frac{2k+1}{2} \operatorname{vol} K + \frac{1}{2} \int_K \left( \sum_{i,j=1}^{2k} (h_{ij}(\vec{v}))^2 + \sum_{i=1}^{2k} \langle \nabla_{\vec{v}} \vec{v}, e_i \rangle^2 \right)$$

and then

$$(1) \quad \mathcal{E}(\vec{v}) \geq \frac{2k+1}{2} \operatorname{vol} K + \frac{1}{2} \int_K \sum_{i,j=1}^{2k} (h_{ij}(\vec{v}))^2.$$

Now observe that

$$\sum_{i < j} (h_{ii} - h_{jj})^2 = (2k - 1) \sum_i h_{ii}^2 - 2 \sum_{i < j} h_{ii} h_{jj}$$

and

$$\sum_{i < j} (h_{ij} + h_{ji})^2 = \sum_{i \neq j} h_{ij}^2 + 2 \sum_{i < j} h_{ij} h_{ji}.$$

If we sum these last two equations, we get

$$(2k - 1) \sum_i h_{ii}^2 + \sum_{i \neq j} h_{ij}^2 \geq 2\sigma_2$$

and then

$$(2) \quad \sum_i h_{ii}^2 + \frac{1}{2k - 1} \sum_{i \neq j} h_{ij}^2 \geq \frac{2}{2k - 1} \sigma_2.$$

Also, we can write

$$\sum_{i,j=1}^{2k} h_{ij}^2 = \sum_{i \neq j} h_{ij}^2 + \sum_i h_{ii}^2 \geq \sum_i h_{ii}^2 + \frac{1}{2k - 1} \sum_{i \neq j} h_{ij}^2.$$

From this and (2), we obtain

$$\sum_{i,j=1}^{2k} (h_{ij}(\vec{v}))^2 \geq \frac{2}{2k - 1} \sigma_2(\vec{v}).$$

But then, using inequality (1), we find that

$$(3) \quad \mathcal{E}(\vec{v}) \geq \frac{2k + 1}{2} \text{vol } K + \frac{1}{2k - 1} \int_K \sigma_2(\vec{v}).$$

On the other hand, by the change of variables theorem, we obtain

$$\text{vol } \varphi_t^H(K) = \int_K \sqrt{1+t^2} (1 + \sum_{i=1}^{2k} \sigma_i(H) t^i)$$

By a straightforward computation shown in [Chacón 2000] and [Brito et al. 2004], we have  $\sigma_i(H) = \eta_i$  for all  $i \in \{1, \dots, 2k\}$ , where

$$\eta_i = \begin{cases} \binom{k}{i/2} & \text{if } i \text{ is even,} \\ 0 & \text{if } i \text{ is odd.} \end{cases}$$

We know that the vector fields  $\vec{v}$  and  $H$  are the same on  $\partial K$ . Thus,  $\varphi_t^{\vec{v}}(K)$  and  $\varphi_t^H(K)$  are  $(2k + 1)$ -submanifolds of  $\mathbb{S}^{2k+1}(\sqrt{1+t^2})$  with the same boundary. We

claim that  $\varphi_t^{\vec{v}}(K) = \varphi_t^H(K)$  for all  $t$  sufficiently small. In fact, if  $p$  is an interior point of  $K$ ,

$$\lim_{t \rightarrow 0} \varphi_t^{\vec{v}}(p) = \lim_{t \rightarrow 0} \varphi_t^H(p) = p$$

and then we have necessarily

$$\varphi_t^{\vec{v}}(K) = \varphi_t^H(K)$$

for all  $t$  sufficiently small; equivalently,

$$\int_K \sqrt{1+t^2} \left( 1 + \sum_{i=1}^{2k} \sigma_i(\vec{v}) t^i \right) = \int_K \sqrt{1+t^2} \left( 1 + \sum_{i=1}^{2k} \eta_i t^i \right)$$

for all  $t > 0$  sufficiently small. Consequently, after canceling the factor  $\sqrt{1+t^2}$  and rearranging the terms, we obtain

$$\left( \int_K [\sigma_1(\vec{v}) - \eta_1] \right) t + \left( \int_K [\sigma_2(\vec{v}) - \eta_2] \right) t^2 + \dots + \left( \int_K [\sigma_{2k}(\vec{v}) - \eta_{2k}] \right) t^{2k} = 0$$

for all sufficiently small  $t$ . By identity of polynomials, we conclude

$$\int_K \sigma_i(\vec{v}) = \int_K \eta_i = \eta_i \operatorname{vol} K \quad \text{for } i \in \{1, \dots, 2k\}.$$

Using this (for  $i = 2$ ) together with (3), we get

$$\mathcal{E}(\vec{v}) \geq \frac{2k+1}{2} \operatorname{vol} K + \frac{\eta_2}{2k-1} \operatorname{vol} K = \left( \frac{2k+1}{2} + \frac{k}{2k-1} \right) \operatorname{vol} K.$$

We can obtain an analogue of this result for volumes using the following inequality (see [Brito et al. 2004] or [Chacón 2000, page 59]):

$$\operatorname{vol} \vec{v} \geq \int_K \left( 1 + \sum_{i=1}^k \frac{\binom{k}{i}}{\binom{2k}{2i}} \sigma_{2i}(\vec{v}) \right).$$

But  $\int_K \sigma_{2i} = \int_K \eta_{2i} = \eta_{2i} \operatorname{vol} K$  for all  $i \in \{1, \dots, k\}$ . Then, we have

$$\operatorname{vol} \vec{v} \geq \left( 1 + \sum_{i=1}^k \frac{\binom{k}{i}^2}{\binom{2k}{2i}} \right) \operatorname{vol} K \geq \frac{4^k}{\binom{2k}{k}} \operatorname{vol} K$$

#### 4. Final remarks

- (1) If  $K$  is a spherical cap (the closure of a connected open set with round boundary of the three unit sphere), the theorem provides a “boundary version” for

the minimization theorem of energy and volume functionals on [Brito 2000] and [Gluck and Ziller 1986].

- (2) The “Hopf boundary” hypothesis is essential. In fact, if there is no constraint for the unit vector field  $\vec{v}$  on  $\partial K$ , it is possible to construct vector fields on “small caps” such that  $\|\nabla \vec{v}\|$  is small on  $K$  (exponential maps may be used on that construction). A consequence of this is that  $\mathcal{E}(\vec{v})$  and  $\text{vol } \vec{v}$  are less than volume and energy of Hopf vector fields respectively.

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Energy and volume of vector fields on spherical domains FABIANO G. B. BRITO, ANDRÉ O. GOMES and GIOVANNI S. NUNES	1
Maps on 3-manifolds given by surgery BOLDIZSÁR KALMÁR and ANDRÁS I. STIPSICZ	9
Strong solutions to the compressible liquid crystal system YU-MING CHU, XIAN-GAO LIU and XIAO LIU	37
Presentations for the higher-dimensional Thompson groups $nV$ JOHANNA HENNIG and FRANCESCO MATUCCI	53
Resonant solutions and turning points in an elliptic problem with oscillatory boundary conditions ALFONSO CASTRO and ROSA PARDO	75
Relative measure homology and continuous bounded cohomology of topological pairs ROBERTO FRIGERIO and CRISTINA PAGLIANTINI	91
Normal enveloping algebras ALEXANDRE N. GRISHKOV, MARINA RASSKAZOVA and SALVATORE SICILIANO	131
Bounded and unbounded capillary surfaces in a cusp domain YASUNORI AOKI and DAVID SIEGEL	143
On orthogonal polynomials with respect to certain discrete Sobolev inner product FRANCISCO MARCELLÁN, RAMADAN ZEJNULLAHU, BUJAR FEJZULLAHU and EDMUNDO HUERTAS	167
Green versus Lempert functions: A minimal example PASCAL THOMAS	189
Differential Harnack inequalities for nonlinear heat equations with potentials under the Ricci flow JIA-YONG WU	199
On overtwisted, right-veering open books PAOLO LISCA	219
Weakly Krull domains and the composite numerical semigroup ring $D + E[\Gamma^*]$ JUNG WOOK LIM	227
Arithmeticity of complex hyperbolic triangle groups MATTHEW STOVER	243