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**INEQUALITIES OF ALEXANDROV-FENCHEL TYPE  
FOR CONVEX HYPERSURFACES IN HYPERBOLIC SPACE  
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## INEQUALITIES OF ALEXANDROV–FENCHEL TYPE FOR CONVEX HYPERSURFACES IN HYPERBOLIC SPACE AND IN THE SPHERE

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**By applying the unit normal flow to well-known inequalities in hyperbolic space  $\mathbb{H}^{n+1}$  and in the sphere  $\mathbb{S}^{n+1}$ , we derive some new inequalities of Alexandrov–Fenchel type for closed convex hypersurfaces in these spaces. We also use the inverse mean curvature flow in the sphere to prove an optimal Sobolev-type inequality for closed convex hypersurfaces in the sphere.**

### 1. Introduction

Let  $N^{n+1}(c)$  be the simply connected space form of constant sectional curvature  $c$  and  $\psi : \Sigma^n \rightarrow N^{n+1}(c)$  be a closed hypersurface. Denote the  $k$ -th order mean curvature of  $\Sigma$  by  $p_k$  (see Section 2A). Inequalities about the integrals  $\int_{\Sigma} p_k d\mu$  have attracted much attention for a long time. Among them the most famous one is the classical Minkowski inequality for closed convex surfaces  $\Sigma \subset \mathbb{R}^3$ , which can be written as

$$(1-1) \quad \left( \frac{1}{\omega_2} \int_{\Sigma} p_1 d\mu \right)^2 \geq \frac{|\Sigma|}{\omega_2},$$

with equality if and only if  $\Sigma$  is a sphere. Here  $\omega_n$  is the area of  $\mathbb{S}^n(1)$  and  $|\Sigma| = \int_{\Sigma} d\mu$  is the area of  $\Sigma$  with respect to the induced metric from  $\mathbb{R}^3$ . The general inequality is the Alexandrov–Fenchel inequality [Alexandrov 1937; 1938; Fenchel 1936] which states that for convex hypersurfaces in the Euclidean space  $\mathbb{R}^{n+1}$ ,

$$(1-2) \quad \frac{1}{\omega_n} \int_{\Sigma} p_k d\mu \geq \left( \frac{1}{\omega_n} \int_{\Sigma} p_l d\mu \right)^{(n-k)/(n-l)} \quad \text{for } 0 \leq l < k \leq n,$$

with equality if and only if  $\Sigma$  is a sphere. See [Chang and Wang 2011; Guan and Li 2009; McCoy 2005; Schneider 1993] for other references on Alexandrov–Fenchel inequalities for closed hypersurfaces in Euclidean space  $\mathbb{R}^{n+1}$ .

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It is natural to generalize the Minkowski inequality and Alexandrov–Fenchel inequalities to the hypersurfaces in space forms. See, for example, [Borisenko and Miquel 1999; Gallego and Solanes 2005; Natário 2015]. Recently, the following optimal inequalities of Alexandrov–Fenchel type in  $\mathbb{H}^{n+1}$  were obtained (see [Ge et al. 2013; 2014b; Li et al. 2014; Wang and Xia 2014]): for  $1 \leq k \leq n$  and any closed horospherical convex hypersurface  $\Sigma \subset \mathbb{H}^{n+1}$ ,

$$(1-3) \quad \frac{1}{\omega_n} \int_{\Sigma} p_k d\mu \geq \left( \left( \frac{|\Sigma|}{\omega_n} \right)^{2/k} + \left( \frac{|\Sigma|}{\omega_n} \right)^{2(n-k)/kn} \right)^{k/2},$$

with equality if and only if  $\Sigma$  is a geodesic sphere in  $\mathbb{H}^{n+1}$ . In particular, when  $k = 2$ , Li, Wei and Xiong [Li et al. 2014] proved that (1-3) holds under the weaker condition that  $\Sigma$  is star-shaped and 2-convex. In the proof of (1-3), the geometric flow was used and was an important tool. However, so far there is no inequality comparing  $\int_{\Sigma} p_k d\mu$  and  $\int_{\Sigma} p_l d\mu$  in  $\mathbb{H}^{n+1}$  like (1-2) in  $\mathbb{R}^{n+1}$ . And one also wants to know whether there exist other inequalities of Alexandrov–Fenchel type in  $\mathbb{H}^{n+1}$  for closed hypersurfaces under a weaker condition than horospherical convex. Besides, in space forms, the integrals  $\int_{\Sigma} p_k d\mu$  are essentially the so-called quermassintegrals from convex geometry and integral geometry (see, e.g., [Solanes 2006] for the transformation formula) and many attempts have been devoted to establishing the relationships for quermassintegrals. See [Santaló 1976; Solanes 2003] and the references therein. So in this paper we are interested in obtaining new inequalities between the integrals  $\int_{\Sigma} p_k d\mu$ .

The Minkowski inequality and the Alexandrov–Fenchel inequalities can be viewed as the generalizations of the classical isoperimetric inequality, which compares the area of the hypersurface  $\Sigma$  and the volume of the domain enclosed by  $\Sigma$ . The Minkowski inequality (1-1) was used by Minkowski himself to prove the isoperimetric inequality for closed convex surfaces (see [Minkowski 1903; Osserman 1978]). Recently, J. Natário [2015] reversed Minkowski’s idea and derived a new Minkowski-type inequality for closed convex surfaces in the hyperbolic space  $\mathbb{H}^3$  from the isoperimetric inequality by using the unit normal flow. In this paper, first, we deal with the higher dimensional case by adapting Natário’s method [2015]. We will derive some new inequalities of Alexandrov–Fenchel type for closed convex hypersurfaces in  $\mathbb{H}^{n+1}$  and in  $\mathbb{S}^{n+1}$ , starting from the isoperimetric inequality.

Let  $\Sigma$  be a closed and convex hypersurface in  $\mathbb{H}^{n+1}$ . We say a hypersurface  $\Sigma$  is convex if all the principal curvatures of  $\Sigma$  are nonnegative everywhere. Then by the well-known result of do Carmo and Warner [1970],  $\Sigma$  is embedded and bounds a convex body in  $\mathbb{H}^{n+1}$ . Inspired by [Natário 2015], we flow the initial hypersurface  $\Sigma$  by its unit outer normal  $\nu$ . The resulting hypersurfaces are  $\Sigma_t = \psi_t(\Sigma)$ , where  $\psi_t(x) = \exp_{\psi(x)}(t\nu(x))$ ,  $x \in \Sigma$ . The  $\Sigma_t$  are also called the parallel hypersurfaces of  $\Sigma$ . From Steiner’s formula [Allendoerfer 1948], we can compute the area of  $\Sigma_t$  and

the volume of the domain  $\Omega_t$  enclosed by  $\Sigma_t$ . In Natário’s paper, the area of  $\Sigma_t$  was obtained by using the first and second variation formulas with the help of the Gauss–Bonnet formula. Steiner’s formula can also be obtained by using the precise expressions of the geodesics in space forms (see Section 2C). Since  $\mathbb{H}^{n+1}$  has constant negative curvature  $c = -1$  and  $\Sigma$  is convex, it follows that  $\Sigma_t$  can be well-defined for all  $t \geq 0$ . Define a function  $r(t)$  such that  $|\Sigma_t| = |S_{r(t)}|$ . Then the isoperimetric inequality (see [Schmidt 1940; Ros 2005]) implies that  $\text{Vol}(\Omega_t) \leq \text{Vol}(B_{r(t)})$ , where  $S_{r(t)}$  and  $B_{r(t)}$  are the geodesic sphere and geodesic ball of radius  $r(t)$  in  $\mathbb{H}^{n+1}$ , respectively. Applying the isoperimetric inequality to  $\Sigma_t$  for sufficiently large  $t$ , we obtain the following inequalities of Alexandrov–Fenchel type in  $\mathbb{H}^{n+1}$ .

**Theorem 1.1.** *Let  $\Sigma^n$  be a closed and convex hypersurface in  $\mathbb{H}^{n+1}$  with  $n \geq 3$ . Then*

$$(1-4) \quad \sum_{k=0}^n \frac{2k-n}{n\omega_n} \int_{\Sigma} C_n^k p_k d\mu \geq \left( \frac{1}{\omega_n} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \right)^{(n-2)/n}.$$

A direct calculation shows that if  $\Sigma$  is a geodesic sphere, then the equality in (1-4) holds. However, we do not obtain the rigidity (i.e., we don’t know whether the equality in (1-4) implies that  $\Sigma$  is a geodesic sphere). In Remark 3.2, we note that when the hypersurface  $\Sigma \subset \mathbb{H}^{n+1}$  is sufficiently small, the inequality (1-4) reduces to one of the Alexandrov–Fenchel inequalities in Euclidean space.

Besides the isoperimetric inequality, there are many other known inequalities in hyperbolic space. If we use the warped product model for the hyperbolic space  $\mathbb{H}^{n+1}$ , i.e.,  $\mathbb{H}^{n+1} = \mathbb{R}^+ \times \mathbb{S}^n$  with the metric  $g = dr^2 + \sinh^2 r g_{\mathbb{S}^n}$ , then there are two important functions on the hypersurface  $\Sigma$  in  $\mathbb{H}^{n+1}$ . One is the weight function  $f = \cosh r$ , and the other one is the support function  $u = \langle Df, \nu \rangle$ . Recently, the following inequality of Alexandrov–Fenchel type with weight  $f$  was proved by de Lima and Girão [2015]: for any mean convex and star-shaped closed hypersurface  $\Sigma$  in  $\mathbb{H}^{n+1}$ ,

$$(1-5) \quad \frac{1}{\omega_n} \int_{\Sigma} f p_1 d\mu \geq \left( \frac{|\Sigma|}{\omega_n} \right)^{(n+1)/n} + \left( \frac{|\Sigma|}{\omega_n} \right)^{(n-1)/n},$$

with equality if and only if  $\Sigma$  is a geodesic sphere centered at the origin in  $\mathbb{H}^{n+1}$ . For more weighted inequalities of Alexandrov–Fenchel type in different ambient spaces, readers can refer to the recent papers [Brendle et al. 2014; Ge et al. 2014a; 2015]. We remark that in [Ge et al. 2014a], the weighted Alexandrov–Fenchel-type inequalities were used to prove the Penrose-type inequality for the new Gauss–Bonnet–Chern mass in asymptotically hyperbolic graphs. Thus it is an interesting question to establish new inequalities with weight.

Applying the same method as in Theorem 1.1 to inequality (1-5), we can obtain a new inequality as follows:

**Theorem 1.2.** *Let  $\Sigma^n$  be a closed and convex hypersurface in  $\mathbb{H}^{n+1}$ . Then*

$$(1-6) \quad \frac{1}{\omega_n} \int_{\Sigma} (f + u) \sum_{k=0}^n C_n^k p_k d\mu \geq \left( \frac{1}{\omega_n} \int_{\Sigma} \sum_{k=0}^n C_n^k p_k d\mu \right)^{(n+1)/n}.$$

We remark that if  $\Sigma$  is a geodesic sphere centered at the origin, then the equality in (1-6) holds. But as before we do not obtain the rigidity.

Next we will use the same method to derive inequalities for closed convex hypersurfaces in  $\mathbb{S}^{n+1}$ . In this case, we can prove the rigidity result.

**Theorem 1.3.** *Let  $\Sigma^n$  be a closed and convex hypersurface in  $\mathbb{S}^{n+1}$  with  $n \geq 2$ . Then*

$$(1-7) \quad \omega_n \leq \sum_{s=\frac{1-(-1)^n}{2}, +2}^n \sqrt{(E(s))^2 + (F(s))^2},$$

where “+2” means that the step-length of the summation for  $s$  is 2 and

$$E(s) = \sum_{\substack{p+q=(n\pm s)/2 \\ p,q \geq 0}} \sum_{\substack{q \leq k \leq n-p \\ 2|k}} C_n^k \frac{1}{2^n} C_{n-k}^p C_k^q (-1)^{\frac{k}{2}+k-q} \int_{\Sigma} p_k d\mu,$$

$$F(s) = \sum_{\substack{p+q=(n\pm s)/2 \\ p,q \geq 0}} \sum_{\substack{q \leq k \leq n-p \\ 2 \nmid k}} C_n^k \frac{1}{2^n} C_{n-k}^p C_k^q (-1)^{[\frac{k}{2}]+k-q} \\ \times (-1)^{\chi\{2(p+q)-n \leq 0\}} \int_{\Sigma} p_k d\mu,$$

Moreover, the equality holds in (1-7) if and only if  $\Sigma^n$  is a geodesic sphere.

When  $n = 2$ , it is easy to check that

$$E(0) = 2\pi, \quad F(0) = 0,$$

$$E(2) = |\Sigma| - 2\pi, \quad F(2) = \int_{\Sigma} p_1 d\mu,$$

using the Gauss–Bonnet theorem  $|\Sigma| + \int_{\Sigma} p_2 d\mu = 4\pi$  (see Section 2B). So (1-7) implies the Minkowski-type inequality in the sphere

$$(1-8) \quad \left( \int_{\Sigma} p_1 d\mu \right)^2 \geq |\Sigma|(4\pi - |\Sigma|),$$

which is just Theorem 0.2 in [Natário 2015]. See also [Blaschke 1938; Knothe 1952; Santaló 1963]. Makowski and Scheuer [2013] proved (1-8) by using the inverse curvature flow in sphere. To get a better feeling of the inequality (1-7), we also give the precise expressions of (1-7) in the case of  $n = 3$  and  $n = 4$ ; see Remark 3.3.

Finally, in the last part of this paper, we use the inverse mean curvature flow in the sphere [Makowski and Scheuer 2013; Gerhard 2015] to prove the following optimal inequalities for strictly convex hypersurfaces in sphere  $\mathbb{S}^{n+1}$ .

**Theorem 1.4.** *Let  $\Sigma^n$  be a closed and strictly convex hypersurface in  $\mathbb{S}^{n+1}$ . Then we have the optimal inequality*

$$(1-9) \quad \int_{\Sigma} L_k d\mu \geq C_n^{2k} (2k)! \omega_n^{2k/n} |\Sigma|^{(n-2k)/n} \quad \text{for } k \leq n/2.$$

*Equality holds in (1-9) if and only if  $\Sigma$  is a geodesic sphere. Here  $L_k$  is the Gauss–Bonnet curvature of the induced metric on  $\Sigma$  (see Section 2B for details).*

The proof of Theorem 1.4 uses a similar idea as in [Brendle et al. 2014; de Lima and Girão 2015; Guan and Li 2009; Ge et al. 2013; 2014b; Li et al. 2014]. We define a curvature quantity  $Q(t)$  which is monotone nonincreasing under the inverse mean curvature flow in the sphere. Then we obtain the inequality (1-9) by comparing the initial value  $Q(0)$  with the limit  $\lim_{t \rightarrow T^*} Q(t)$ . We remark that since  $\Sigma$  is a closed and strictly convex hypersurface in  $\mathbb{S}^{n+1}$ , a well-known result due to do Carmo and Warner [1970] implies that  $\Sigma$  is embedded and is contained in an open hemisphere.

When  $k = 1$ , the inequality (1-9) reduces to

$$(1-10) \quad \int_{\Sigma} p_2 d\mu + |\Sigma| \geq \omega_n^{2/n} |\Sigma|^{(n-2)/n},$$

which was already proved by Makowski and Scheuer [2013]. One can compare (1-10) with the case  $k = 2$  of the Alexandrov–Fenchel-type inequality (1-3) in  $\mathbb{H}^{n+1}$ ; that is,

$$(1-11) \quad \int_{\Sigma} p_2 d\mu - |\Sigma| \geq \omega_n^{2/n} |\Sigma|^{(n-2)/n},$$

which was proved by Li, Wei and Xiong [Li et al. 2014] for star-shaped and 2-convex hypersurfaces in  $\mathbb{H}^{n+1}$ . For  $k > 1$ , inequalities of the same type as (1-9) were proved by Ge, Wang and Wu [Ge et al. 2013; 2014b] for horospherical convex hypersurfaces in the hyperbolic space  $\mathbb{H}^{n+1}$ .

## 2. Preliminaries

**2A.  $k$ -th order mean curvature.** Let  $\Sigma$  be a closed hypersurface in  $N^{n+1}(c)$  with unit outward normal  $\nu$ . The second fundamental form  $h$  of  $\Sigma$  is defined by

$$h(X, Y) = \langle \bar{\nabla}_X \nu, Y \rangle$$

for any two tangent vector fields  $X, Y$  on  $\Sigma$ . For an orthonormal basis  $\{e_1, \dots, e_n\}$  of  $\Sigma$ , the components of the second fundamental form are given by  $h_{ij} = h(e_i, e_j)$  and  $h_i^j = g^{jk} h_{ki}$ , where  $g$  is the induced metric on  $\Sigma$ . The principal curvatures

$\kappa = (\kappa_1, \dots, \kappa_n)$  are the eigenvalues of  $h$  with respect to  $g$ . The  $k$ -th order mean curvature of  $\Sigma$  for  $1 \leq k \leq n$  is defined as

$$(2-1) \quad p_k = \frac{1}{C_n^k} \sigma_k(\kappa) = \frac{1}{C_n^k} \sum_{i_1 < i_2 < \dots < i_k} \kappa_{i_1} \cdots \kappa_{i_k},$$

or equivalently as

$$(2-2) \quad p_k = \frac{1}{C_n^k} \sigma_k(h_i^j) = \frac{1}{C_n^k k!} \delta_{j_1 \dots j_k}^{i_1 \dots i_k} h_{i_1}^{j_1} \cdots h_{i_k}^{j_k},$$

where  $\delta_{j_1 \dots j_k}^{i_1 \dots i_k}$  is the generalized Kronecker delta given by

$$\delta_{j_1 \dots j_k}^{i_1 \dots i_k} = \det \begin{pmatrix} \delta_{j_1}^{i_1} & \delta_{j_1}^{i_2} & \cdots & \delta_{j_1}^{i_k} \\ \delta_{j_2}^{i_1} & \delta_{j_2}^{i_2} & \cdots & \delta_{j_2}^{i_k} \\ \vdots & \vdots & \ddots & \vdots \\ \delta_{j_k}^{i_1} & \delta_{j_k}^{i_2} & \cdots & \delta_{j_k}^{i_k} \end{pmatrix}.$$

We have the following Newton–MacLaurin inequalities (see, e.g., [Guan 2006]).

**Lemma 2.1.** For  $\kappa \in \bar{\Gamma}_k^+$ ,  $1 \leq k \leq n$ , where  $\bar{\Gamma}_k^+$  is the closure of the Garding cone

$$\Gamma_k^+ = \{\kappa \in \mathbb{R}^n \mid p_j(\kappa) > 0, \forall j \leq k\},$$

we have the following Newton–MacLaurin inequalities

$$p_1 p_{k-1} \geq p_k, \\ p_1 \geq p_2^{1/2} \geq \cdots \geq p_k^{1/k}.$$

Moreover, equalities above hold for some  $\kappa \in \Gamma_k^+$  if and only if  $\kappa = c(1, \dots, 1)$ , where  $c$  is a constant.

**2B. Gauss–Bonnet curvature  $L_k$ .** Given an  $n$ -dimensional Riemannian manifold  $(M, g)$ , the Gauss–Bonnet curvature  $L_k$ , where  $k \leq n/2$ , is defined by (see, e.g., [Ge et al. 2014b; 2014c])

$$(2-3) \quad L_k = \frac{1}{2^k} \delta_{j_1 j_2 \dots j_{2k-1} j_{2k}}^{i_1 i_2 \dots i_{2k-1} i_{2k}} R_{i_1 i_2}^{j_1 j_2} \cdots R_{i_{2k-1} i_{2k}}^{j_{2k-1} j_{2k}}.$$

For a closed hypersurface  $\Sigma^n \subset \mathbb{R}^{n+1}$ , recall the Gauss equation

$$R_{ij}{}^{kl} = h_i^k h_j^l - h_i^l h_j^k.$$

Then the Gauss–Bonnet curvature of the induced metric on  $\Sigma^n \subset \mathbb{R}^{n+1}$  is

$$(2-4) \quad L_k = \delta_{j_1 j_2 \dots j_{2k-1} j_{2k}}^{i_1 i_2 \dots i_{2k-1} i_{2k}} h_{i_1}^{j_1} h_{i_2}^{j_2} \cdots h_{i_{2k-1}}^{j_{2k-1}} h_{i_{2k}}^{j_{2k}} \\ = (2k)! C_n^{2k} p_{2k}.$$



For a closed hypersurface  $\Sigma^n \subset \mathbb{S}^{n+1}$ , the Gauss equations are

$$(2-5) \quad R_{ij}{}^{kl} = (h_i^k h_j^l - h_i^l h_j^k) + (\delta_i^k \delta_j^l - \delta_i^l \delta_j^k).$$

Then by a straightforward calculation, we have

$$\begin{aligned} L_k &= \delta_{j_1 j_2 \dots j_{2k-1} j_{2k}}^{i_1 i_2 \dots i_{2k-1} i_{2k}} (h_{i_1}^{j_1} h_{i_2}^{j_2} + \delta_{i_1}^{j_1} \delta_{i_2}^{j_2}) \cdots (h_{i_{2k-1}}^{j_{2k-1}} h_{i_{2k}}^{j_{2k}} + \delta_{i_{2k-1}}^{j_{2k-1}} \delta_{i_{2k}}^{j_{2k}}) \\ &= \sum_{i=0}^k C_k^i (n-2k+1)(n-2k+2) \cdots (n-2k+2i)(2k-2i)! \sigma_{2k-2i} \\ &= \sum_{i=0}^k C_k^i (n-2k+1)(n-2k+2) \cdots (n-2k+2i)(2k-2i)! C_n^{2k-2i} p_{2k-2i} \\ &= \sum_{i=0}^k C_k^i \frac{n!}{(n-2k)!} p_{2k-2i} \\ &= C_n^{2k} (2k)! \sum_{i=0}^k C_k^i p_{2k-2i}. \end{aligned}$$

Similarly, for a closed hypersurface  $\Sigma^n \subset \mathbb{H}^{n+1}$ , its Gauss–Bonnet curvature is

$$(2-6) \quad L_k = C_n^{2k} (2k)! \sum_{i=0}^k C_k^i (-1)^i p_{2k-2i}.$$

Finally, note that throughout our paper, we assume that the hypersurface  $\Sigma \subset N^{n+1}(c)$  is closed and convex. It follows that  $\Sigma$  is homeomorphic to the  $n$ -sphere (see [do Carmo and Warner 1970]). Then if the dimension of  $\Sigma$  is even, the Gauss–Bonnet–Chern theorem [Chern 1944] implies that

$$(2-7) \quad \int_{\Sigma} L_{n/2} d\mu = n! \omega_n.$$

Equation (2-7) will be used in the following sections. Also (2-7) shows that when  $2k = n$ , the inequality (1-9) is an equality.

**2C. Unit normal flow and Steiner's formula.** Let  $\psi : \Sigma \rightarrow N^{n+1}(c)$  be a closed and convex hypersurface in the simply connected space form  $N^{n+1}(c)$  of constant sectional curvature  $c$ . Denote by  $\Omega$  the domain enclosed by  $\Sigma$ . The area of  $\Sigma$  is denoted by  $|\Sigma|$  and the volume of  $\Omega$  is denoted by  $|V|$ . As we mentioned in Section 1, we flow the initial hypersurface  $\Sigma$  by its unit outer normal  $\nu$ . The resulting hypersurfaces are  $\Sigma_t = \psi_t(\Sigma)$ , where  $\psi_t(x) = \exp_{\psi(x)}(t\nu(x))$ ,  $x \in \Sigma$ . The  $\Sigma_t$  are also called the parallel hypersurfaces of  $\Sigma$ . Denote by  $\Omega_t$  the domain

bounded by  $\Sigma_t$ . The convexity assumption of  $\Sigma$  and the curvature of  $N^{n+1}(c)$  guarantee that the  $\Sigma_t$  are well-defined in the following range:

$$t \in \left[0, \frac{\pi}{2}\right) \quad \text{for } c = 1,$$

$$t \geq 0 \quad \text{for } c = 0 \text{ or } -1.$$

Further, denote the area of  $\Sigma_t$  and the volume of  $\Omega_t$  by  $|\Sigma_t|$  and  $|V_t|$ , respectively. Then Steiner's formula [Allendoerfer 1948] implies that

$$(2-8) \quad |\Sigma_t| = \begin{cases} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu t^k & \text{if } c = 0, \\ \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \cosh^{n-k} t \sinh^k t & \text{if } c = -1, \\ \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \cos^{n-k} t \sin^k t & \text{if } c = 1. \end{cases}$$

and

$$(2-9) \quad |V_t| = \begin{cases} |V| + \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \frac{1}{k+1} t^{k+1} & \text{if } c = 0, \\ |V| + \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \int_0^t \cosh^{n-k} s \sinh^k s ds & \text{if } c = -1, \\ |V| + \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \int_0^t \cos^{n-k} s \sin^k s ds & \text{if } c = 1. \end{cases}$$

We give a simple proof of (2-8) and (2-9) here. First, when  $c = 0$ , the parallel hypersurface can be expressed as  $\psi_t = \psi + tv$  (see [Montiel and Ros 1991]). So  $(\psi_t)_*(e_i) = (1 + t\kappa_i)e_i$ . Therefore the area element of  $\Sigma_t$  is

$$d\mu_t = (1 + t\kappa_1) \cdots (1 + t\kappa_n) d\mu,$$

which implies that the areas  $|\Sigma_t|$  of the parallel hypersurfaces  $\Sigma_t$  are equal to

$$|\Sigma_t| = \int_{\Sigma} (1 + t\kappa_1) \cdots (1 + t\kappa_n) d\mu = \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu t^k.$$

Note that  $\Sigma_t = \psi_t(\Sigma)$  are parallel hypersurfaces of  $\Sigma$  given by

$$\psi_t(x) = \exp_{\psi(x)}(tv(x))$$

for  $x \in \Sigma$ . By integrating and using the co-area formula, we obtain

$$|V_t| = |V| + \int_0^t |\Sigma_s| ds = |V| + \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \frac{1}{k+1} t^{k+1}.$$

Similarly, when  $c = -1$ ,  $\psi_t = \cosh t \psi + \sinh t v$  (see [Montiel and Ros 1991]) and so  $(\psi_t)_*(e_i) = (\cosh t + \sinh t \kappa_i)e_i$ . Therefore the area element of  $\Sigma_t$  is

$$d\mu_t = (\cosh t + \sinh t \kappa_1) \cdots (\cosh t + \sinh t \kappa_n) d\mu,$$

which implies

$$\begin{aligned} |\Sigma_t| &= \int_{\Sigma} (\cosh t + \sinh t \kappa_1) \cdots (\cosh t + \sinh t \kappa_n) d\mu \\ &= \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \cosh^{n-k} t \sinh^k t. \end{aligned}$$

Then by integrating, we obtain

$$|V_t| = |V| + \int_0^t |\Sigma_s| ds = |V| + \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \int_0^t \cosh^{n-k} s \sinh^k s ds.$$

Finally, the case  $c = 1$  can also be proved by noting that  $\psi_t = \cos t \psi + \sin t \nu$ , where  $t \in [0, \frac{\pi}{2})$ .

### 3. The results by the method of unit normal flow

**3A. The Euclidean case.** To demonstrate the method which will be used to prove Theorems 1.1–1.3, in this subsection we first consider the simple case that  $\Sigma$  is a closed and convex hypersurface in  $\mathbb{R}^{n+1}$ . Let  $\Sigma_t$  be the parallel hypersurfaces of  $\Sigma$  and  $\Omega_t$  be the domain enclosed by  $\Sigma_t$ . Then  $\Sigma_t$  is well-defined for all  $t \geq 0$ . For all  $t \geq 0$ , the isoperimetric inequality (see [Osserman 1978]) in Euclidean space  $\mathbb{R}^{n+1}$  implies

$$(3-1) \quad \left( \frac{|\Sigma_t|}{\omega_n} \right)^{n+1} \geq \left( (n+1) \frac{|V_t|}{\omega_n} \right)^n.$$

Substitute Steiner’s formulas (2-8), (2-9) into (3-1). If  $n$  is odd, then comparing the coefficient of  $t^{n(n+1)}$  in (3-1) yields

$$(3-2) \quad \int_{\Sigma} p_n d\mu \geq \omega_n,$$

which is a special Alexandrov–Fenchel inequality.

If  $n$  is even, (2-4) and the Gauss–Bonnet–Chern theorem (2-7) imply that  $\int_{\Sigma} p_n d\mu = \omega_n$ . Thus expanding the two sides of (3-1) and comparing the coefficients of  $t^{n(n+1)}$ ,  $t^{n(n+1)-1}$  and  $t^{n(n+1)-2}$ , we can get

$$(3-3) \quad \left( \frac{1}{\omega_n} \int_{\Sigma} p_{n-1} d\mu \right)^2 \geq \frac{1}{\omega_n} \int_{\Sigma} p_{n-2} d\mu,$$

which is also an Alexandrov–Fenchel inequality. In particular, when  $n = 2$ , (3-3) reduces to the classical Minkowski inequality (1-1).

**3B. The hyperbolic case, I.** In this subsection, we prove Theorem 1.1. Assume that  $\Sigma$  is a closed and convex hypersurface in  $\mathbb{H}^{n+1}$ . Then the parallel hypersurfaces  $\Sigma_t$  are well-defined for all  $t \geq 0$ . Recall that the area of a geodesic sphere  $S_r$  and the volume of a geodesic ball  $B_r$  with radius  $r$  in the hyperbolic space  $\mathbb{H}^{n+1}$  are

$$S(r) := |S_r| = \omega_n \sinh^n r,$$

$$V(r) := \text{Vol}(B_r) = \omega_n \int_0^r \sinh^n s \, ds.$$

Now define a function  $r(t)$  such that  $|\Sigma_t| = S(r(t))$ . That is,

$$(3-4) \quad \sum_{k=0}^n \cosh^{n-k} t \sinh^k t \int_{\Sigma} C_n^k p_k \, d\mu = \omega_n \sinh^n r(t).$$

Then the isoperimetric inequality (see [Schmidt 1940; Ros 2005]) implies

$$(3-5) \quad |V_t| \leq V(r(t)) \quad \text{for } t \geq 0.$$

From this inequality, we can get some information for  $\Sigma$ .

First we get a rough estimate for  $r(t)$ . When  $t \rightarrow +\infty$ ,  $\cosh^{n-k} t \sinh^k t = \sinh^n t (1 + o(1))$ . Thus from  $|\Sigma_t| = S(r(t))$ , we get

$$\sinh^n t (1 + o(1)) \sum_{k=0}^n \int_{\Sigma} C_n^k p_k \, d\mu = \omega_n \sinh^n r(t),$$

which implies

$$(3-6) \quad r(t) = t + \frac{1}{n} \ln \left( \frac{1}{\omega_n} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k \, d\mu \right) + o(1).$$

However, this estimate for  $r(t)$  is not enough. For our purposes, we should make better use of  $|\Sigma_t| = S(r(t))$  as follows. The case of  $n = 2$  was considered by Natário [2015], so we assume that  $n \geq 3$  in the following calculation. Since we will examine (3-5) for sufficiently large  $t$ , we only care about the terms involving  $e^{nt}$  and  $e^{(n-2)t}$ . The other terms are  $o(e^{(n-2)t})$ . It is straightforward to check that

$$\cosh^{n-k} t \sinh^k t = \frac{1}{2^n} e^{nt} + \frac{1}{2^n} (n-2k) e^{(n-2)t} + \dots.$$

Consequently (3-4) implies

$$(3-7) \quad \frac{1}{2^n} \sum_{k=0}^n (e^{nt} + (n-2k)e^{(n-2)t} + \dots) \int_{\Sigma} C_n^k p_k \, d\mu$$

$$= \omega_n \left( \frac{1}{2^n} e^{nr} - \frac{1}{2^n} n e^{(n-2)r} + \dots \right).$$

On the other hand, from Steiner's formula (2-9), we have

$$\begin{aligned}
 |V_t| &= |V| + \sum_{k=0}^n \int_0^t \cosh^{n-k} s \sinh^k s \, ds \int_{\Sigma} C_n^k p_k \, d\mu \\
 &= |V| + \frac{1}{2^n} \sum_{k=0}^n \int_0^t (e^{ns} + (n-2k)e^{(n-2)s} + \dots) \, ds \int_{\Sigma} C_n^k p_k \, d\mu \\
 &= |V| + \frac{1}{2^n} \sum_{k=0}^n \left( \frac{1}{n} e^{nt} + \frac{n-2k}{n-2} e^{(n-2)t} + \dots \right) \int_{\Sigma} C_n^k p_k \, d\mu \\
 &= \frac{1}{2^n} \frac{1}{n} e^{nt} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k \, d\mu + \frac{1}{2^n} e^{(n-2)t} \sum_{k=0}^n \frac{n-2k}{n-2} \int_{\Sigma} C_n^k p_k \, d\mu + \dots,
 \end{aligned}$$

and

$$\begin{aligned}
 V(r(t)) &= \omega_n \int_0^r \sinh^n s \, ds \\
 &= \omega_n \int_0^r \left( \frac{1}{2^n} e^{ns} - \frac{1}{2^n} n e^{(n-2)s} + \dots \right) \, ds \\
 &= \frac{\omega_n}{2^n} \frac{1}{n} e^{nr} - \frac{\omega_n}{2^n} \frac{n}{n-2} e^{(n-2)r} + \dots.
 \end{aligned}$$

Now taking (3-7) into account yields

$$\begin{aligned}
 V(r(t)) &= \frac{\omega_n}{2^n} e^{(n-2)r} + \frac{1}{2^n} \frac{1}{n} \sum_{k=0}^n (e^{nt} + (n-2k)e^{(n-2)t} + \dots) \int_{\Sigma} C_n^k p_k \, d\mu \\
 &\qquad\qquad\qquad - \frac{\omega_n}{2^n} \frac{n}{n-2} e^{(n-2)r} \\
 &= \frac{\omega_n}{2^n} \left( \frac{-2}{n-2} \right) e^{(n-2)r} + \frac{1}{2^n} \frac{1}{n} e^{nt} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k \, d\mu \\
 &\qquad\qquad\qquad + \frac{1}{2^n} \frac{1}{n} e^{(n-2)t} \sum_{k=0}^n (n-2k) \int_{\Sigma} C_n^k p_k \, d\mu + \dots.
 \end{aligned}$$

Noting (3-6), we have

$$\begin{aligned}
 V(r(t)) &= \frac{\omega_n}{2^n} \left( \frac{-2}{n-2} \right) e^{(n-2)t} \left( \frac{1}{\omega_n} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k \, d\mu \right)^{(n-2)/n} \\
 &\quad + \frac{1}{2^n} \frac{1}{n} e^{nt} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k \, d\mu + \frac{1}{2^n} \frac{1}{n} e^{(n-2)t} \sum_{k=0}^n (n-2k) \int_{\Sigma} C_n^k p_k \, d\mu + \dots.
 \end{aligned}$$

Now  $|V_t| \leq V(r(t)), t \rightarrow +\infty$  gives us

$$\frac{1}{2^n} \sum_{k=0}^n \left( \frac{n-2k}{n-2} - \frac{n-2k}{n} \right) \int_{\Sigma} C_n^k p_k d\mu \leq \frac{\omega_n}{2^n} \frac{-2}{n-2} \left( \frac{1}{\omega_n} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \right)^{(n-2)/n},$$

or equivalently

$$(3-8) \quad \sum_{k=0}^n \frac{2k-n}{n} \int_{\Sigma} C_n^k p_k d\mu \geq \omega_n \left( \frac{1}{\omega_n} \sum_{k=0}^n \int_{\Sigma} C_n^k p_k d\mu \right)^{(n-2)/n} \quad \text{for } n \geq 3.$$

Hence we complete the proof of Theorem 1.1.

**Remark 3.1.** It is easy to check that for a geodesic sphere in  $\mathbb{H}^{n+1}$ , the equality holds in (3-8). However this method can not yield the rigidity result; i.e., we cannot conclude that  $\Sigma$  is a geodesic sphere if the equality holds in (3-8).

**Remark 3.2.** We also remark that for a small hypersurface  $\Sigma \subset \mathbb{H}^{n+1}$  (i.e., with small diameter), the inequality (3-8) can reduce to the Euclidean inequalities (3-2) and (3-3). For example, we first assume  $n = 4$ . For 4-dimensional hypersurface  $\Sigma \subset \mathbb{H}^5$ , the Gauss–Bonnet–Chern formula (2-7) implies

$$(3-9) \quad \int_{\Sigma} (p_4 - 2p_2 + 1) d\mu = \frac{1}{4!} \int_{\Sigma} L_2 d\mu = \omega_4.$$

Substituting (3-9) into the inequality (3-8) gives that

$$\left( 1 + \frac{2}{\omega_4} \int_{\Sigma} (p_3 + p_2 - p_1 - 1) d\mu \right)^2 \geq 1 + \frac{4}{\omega_4} \int_{\Sigma} (p_3 + 2p_2 + p_1) d\mu.$$

Expanding the left-hand side of the above inequality, and comparing both sides by orders (note that  $\Sigma$  is a small hypersurface), we obtain that

$$(3-10) \quad \left( \frac{1}{\omega_4} \int_{\Sigma} p_3 d\mu \right)^2 \geq \frac{1}{\omega_4} \int_{\Sigma} p_2 d\mu.$$

This is just the inequality (3-3) for hypersurfaces in Euclidean space  $\mathbb{R}^5$ . For the general even-dimensional case, by using the Gauss–Bonnet–Chern formula,

$$\int_{\Sigma} \sum_{k=0}^{n/2} C_{n/2}^k (-1)^k p_{n-2k} d\mu = \frac{1}{n!} \int_{\Sigma} L_{n/2} d\mu = \omega_n.$$

We can also reduce the inequality (3-8) to the Euclidean version (3-3) for small hypersurfaces  $\Sigma \subset \mathbb{H}^{n+1}$ . For the odd-dimensional case, the argument is similar.

**3C. The hyperbolic case, II.** In this subsection, we will prove Theorem 1.2. Since the method is similar to that of the last subsection, we just sketch it.

Here we need the following model of the hyperbolic space. Let  $\mathbb{R}_1^{n+2}$  be the Minkowski space with the Lorentzian metric

$$\langle x, y \rangle = x_1 y_1 + \dots + x_{n+1} y_{n+1} - x_{n+2} y_{n+2}.$$

Then the  $(n+1)$ -dimensional hyperbolic space can be defined by

$$\mathbb{H}^{n+1} = \{x \in \mathbb{R}_1^{n+2} \mid \langle x, x \rangle = -1, x_{n+2} \geq 1\}$$

with the induced metric from  $\mathbb{R}_1^{n+2}$ .

Fix a point  $a = (0, \dots, 0, -1)$ . Then it is easy to check that the weight function and the support function can be written down as

$$\begin{aligned} f &= \langle \psi, a \rangle, \\ u &= \langle \nu, a \rangle. \end{aligned}$$

Next define a family of parallel hypersurfaces  $\Sigma_t = \psi_t(\Sigma)$ , where  $\psi_t(x) = \exp_{\psi(x)}(t\nu(x))$ ,  $x \in \Sigma$ , and  $\nu(x)$  is the outward unit normal of  $\Sigma$ . In fact,  $\psi_t = \cosh t \psi + \sinh t \nu$ . And since the initial hypersurface is convex,  $\Sigma_t$  is well-defined for all  $t \geq 0$ . Then  $(\psi_t)_*(e_i) = (\cosh t + \kappa_i \sinh t)e_i$  and

$$\kappa_i(t) = \frac{\tanh t + \kappa_i}{1 + \kappa_i \tanh t}.$$

For convenience, we define a function  $Q(t)$  by

$$Q_n(t) = (1 + t\kappa_1) \cdots (1 + t\kappa_n) = 1 + C_n^1 p_1 t + \dots + C_n^n p_n t^n.$$

Then the mean curvature of  $\Sigma_t$  is

$$p_1(t) = \frac{n \cosh t \sinh t Q_n(\tanh t) + Q'_n(\tanh t)}{n \cosh^2 t Q_n(\tanh t)}.$$

Note that  $p_1(t) \rightarrow 1$  as  $t \rightarrow +\infty$ . So for sufficiently large  $t$ ,  $\Sigma_t$  is mean convex. And  $\langle \nu_t, a \rangle = \langle \sinh t \psi + \cosh t \nu, a \rangle \geq 0$  for sufficiently large  $t$ , which implies  $\Sigma_t$  is star-shaped for these  $t$ . Thus, we can apply (1-5) to  $\Sigma_t$ :

$$\begin{aligned} &\frac{1}{\omega_n} \int_{\Sigma} \langle \cosh t \psi + \sinh t \nu, a \rangle p_1(t) \cosh^n t Q_n(\tanh t) d\mu \\ &\geq \left( \frac{1}{\omega_n} \int_{\Sigma} \cosh^n t Q_n(\tanh t) d\mu \right)^{\frac{n+1}{n}} + \left( \frac{1}{\omega_n} \int_{\Sigma} \cosh^n t Q_n(\tanh t) d\mu \right)^{\frac{n-1}{n}}. \end{aligned}$$

Let  $t \rightarrow +\infty$ . Taking into account that  $\tanh t \rightarrow 1$ ,  $p_1(t) \rightarrow 1$  and  $\sinh t = \cosh t (1 + o(1))$ , we obtain (1-6). So we have finished the proof.

**3D. The spherical case.** We now prove Theorem 1.3. Assume that  $\Sigma$  is a closed and convex hypersurface in  $\mathbb{S}^{n+1}$ . Then the parallel hypersurface  $\Sigma_t$  is well-defined for  $t \in [0, \frac{\pi}{2})$ . Recall that the area of a geodesic sphere  $S_r$  and the volume of a geodesic ball  $B_r$  with radius  $r$  in the sphere  $\mathbb{S}^{n+1}$  are

$$S(r) = \omega_n \sin^n r,$$

$$V(r) = \omega_n \int_0^r \sin^n s \, ds.$$

Now since  $|V_t|$  is increasing in  $t$ , when  $t$  satisfies

$$|V_t| = V\left(\frac{\pi}{2}\right) = \omega_n \int_0^{\pi/2} \sin^n r \, dr,$$

the isoperimetric inequality (see [Ros 2005]) implies  $|\Sigma_t| \geq S(\frac{\pi}{2}) = \omega_n$  for this  $t$ . Therefore, a weaker requirement is

$$(3-11) \quad \max_{t \in [0, \pi/2)} |\Sigma_t| \geq \omega_n.$$

Then the key point is to estimate  $\max_{t \in [0, \pi/2)} |\Sigma_t|$ . Direct computation shows that

$$\begin{aligned} \cos^{n-k} t \sin^k t &= \left(\frac{e^{it} + e^{-it}}{2}\right)^{n-k} \left(\frac{e^{it} - e^{-it}}{2i}\right)^k \\ &= \frac{1}{2^n} \sum_{p=0}^{n-k} \sum_{q=0}^k C_{n-k}^p C_k^q \cos\left((2(p+q)-n)t - \frac{k\pi}{2}\right) (-1)^{k-q}. \end{aligned}$$

Then Steiner's formula (2-8) implies

$$\begin{aligned} |\Sigma_t| &= \sum_{k=0}^n C_n^k \cos^{n-k} t \sin^k t \int_{\Sigma} p_k \, d\mu \\ &= \sum_{k=0}^n C_n^k \frac{1}{2^n} \sum_{p=0}^{n-k} \sum_{q=0}^k C_{n-k}^p C_k^q \cos\left((2(p+q)-n)t - \frac{k\pi}{2}\right) (-1)^{k-q} \int_{\Sigma} p_k \, d\mu \\ &= \sum_{\substack{0 \leq k \leq n \\ 2 \nmid k}} C_n^k \frac{1}{2^n} \sum_{p=0}^{n-k} \sum_{q=0}^k C_{n-k}^p C_k^q (-1)^{\frac{k}{2}} \cos((2(p+q)-n)t) (-1)^{k-q} \int_{\Sigma} p_k \, d\mu \\ &+ \sum_{\substack{0 \leq k \leq n \\ 2 \nmid k}} C_n^k \frac{1}{2^n} \sum_{p=0}^{n-k} \sum_{q=0}^k C_{n-k}^p C_k^q (-1)^{[\frac{k}{2}]} \sin((2(p+q)-n)t) (-1)^{k-q} \int_{\Sigma} p_k \, d\mu. \end{aligned}$$



Next let  $2(p + q) - n = \pm s$  and sum up in terms of  $s$  first. We get

$$\begin{aligned}
 (3-12) \quad |\Sigma_t| &= \sum_{s=\frac{1-(-1)^n}{2}, +2}^n \sum_{\substack{p+q=(n\pm s)/2 \\ p,q \geq 0}} \sum_{\substack{q \leq k \leq n-p \\ 2|k}} C_n^k \frac{1}{2^n} C_{n-k}^p C_k^q \\
 &\quad \times (-1)^{\frac{k}{2}+k-q} \cos(st) \int_{\Sigma} p_k d\mu \\
 &+ \sum_{s=\frac{1-(-1)^n}{2}, +2}^n \sum_{\substack{p+q=(n\pm s)/2 \\ p,q \geq 0}} \sum_{\substack{q \leq k \leq n-p \\ 2 \nmid k}} C_n^k \frac{1}{2^n} C_{n-k}^p C_k^q (-1)^{[\frac{k}{2}]+k-q} \\
 &\quad \times (-1)^{\chi\{2(p+q)-n \leq 0\}} \sin(st) \int_{\Sigma} p_k d\mu \\
 &\leq \sum_{s=\frac{1-(-1)^n}{2}, +2}^n \sqrt{(E(s))^2 + (F(s))^2},
 \end{aligned}$$

in the notation of Theorem 1.3.

Next we show that for the geodesic sphere with radius  $r \in [0, \frac{\pi}{2})$ , the equality holds. For this special hypersurface,  $\int_{\Sigma} p_k d\mu = \omega_n \sin^n r \cot^k r = \omega_n \sin^{n-k} r \cos^k r$ . Thus

$$\begin{aligned}
 |\Sigma_t| &= \omega_n \sum_{k=0}^n C_n^k (\cos t \sin r)^{n-k} (\sin t \cos r)^k = \omega_n \sin^n(r + t) \\
 &= \omega_n \frac{1}{2^n} \sum_{q=0}^n C_n^q \cos\left((2q - n)(r + t) - \frac{n\pi}{2}\right) (-1)^{n-q}.
 \end{aligned}$$

For simplicity, we assume  $n$  is even. Then

$$\begin{aligned}
 |\Sigma_t| &= \omega_n \frac{1}{2^n} \sum_{q=0}^n C_n^q \cos((2q - n)(r + t)) (-1)^{n/2+n-q} \\
 &= \omega_n \sum_{s=0,2,\dots,n} \sum_{2q-n=\pm s} \frac{1}{2^n} C_n^q \cos(s(r + t)) (-1)^{3n/2-q} \\
 &= \omega_n \sum_{s=0,2,\dots,n} \sum_{2q-n=s} 2 \frac{1}{2^n} C_n^q \cos(s(r + t)) (-1)^{3n/2-q},
 \end{aligned}$$

where we note that the coefficients of  $\cos(s(r + t))$  for the two choices of  $q$  are the same.

Now expand  $\cos(s(r + t)) = \cos sr \cos st - \sin sr \sin st$ . We find that all the inequalities in (3-12) become equalities for  $t = \frac{\pi}{2} - r$  and  $|\Sigma_t| = \omega_n \sin^n \frac{\pi}{2} = \omega_n$ . Thus for the geodesic sphere, the equality in (3-12) holds.

On the other hand, assume the equality holds. Then when some  $t$  satisfies  $|V_t| = V(\frac{\pi}{2}) = \omega_n \int_0^{\pi/2} \sin^n r dr$ , we must have  $|\Sigma_t| = S(\frac{\pi}{2}) = \omega_n$  for this  $t$ . So

the isoperimetric inequality implies that  $\Sigma_t = \mathbb{S}^n(1)$ . Then the initial hypersurface must be a geodesic sphere.

Thus Theorem 1.3 is proved.

**Remark 3.3.** In Section 1, we discussed the special case  $n = 2$  of (1-7), which is just the Minkowski-type inequality for convex surfaces in  $\mathbb{S}^3$ . Here, to get a better feeling of the inequality (1-7), we give the precise expressions for  $n = 3$  and  $n = 4$ . For  $n = 3$ , we have

$$\begin{aligned} \omega_3 \leq & \sqrt{\left(\frac{1}{4}\left(|\Sigma| - 3 \int_{\Sigma} p_2 d\mu\right)\right)^2 + \left(\frac{1}{4}\left(3 \int_{\Sigma} p_1 d\mu - \int_{\Sigma} p_3 d\mu\right)\right)^2} \\ & + \sqrt{\left(\frac{3}{4}\left(|\Sigma| + \int_{\Sigma} p_2 d\mu\right)\right)^2 + \left(\frac{3}{4}\left(\int_{\Sigma} p_1 d\mu + \int_{\Sigma} p_3 d\mu\right)\right)^2}. \end{aligned}$$

And for  $n = 4$ , we have

$$\begin{aligned} (3-13) \quad \omega_4 \leq & \sqrt{\left(\frac{1}{8}\left(|\Sigma| - 6 \int_{\Sigma} p_2 d\mu + \int_{\Sigma} p_4 d\mu\right)\right)^2 + \left(\frac{1}{2}\left(\int_{\Sigma} p_1 d\mu - \int_{\Sigma} p_3 d\mu\right)\right)^2} \\ & + \sqrt{\left(\frac{1}{2}\left(|\Sigma| - \int_{\Sigma} p_4 d\mu\right)\right)^2 + \left(\int_{\Sigma} p_1 d\mu + \int_{\Sigma} p_3 d\mu\right)^2} \\ & + \frac{3}{8}\left(|\Sigma| + 2 \int_{\Sigma} p_2 d\mu + \int_{\Sigma} p_4 d\mu\right). \end{aligned}$$

For a 4-dimensional hypersurface  $\Sigma$  in  $\mathbb{S}^5$ , we have the Gauss–Bonnet–Chern formula

$$(3-14) \quad \int_{\Sigma} (p_4 + 2p_2 + 1) d\mu = \frac{1}{4!} \int_{\Sigma} L_2 d\mu = \omega_4.$$

Therefore, the inequality (3-13) can be further simplified by using the formula (3-14).

**Remark 3.4.** As in the hyperbolic case, when the hypersurface  $\Sigma \subset \mathbb{S}^5$  is small, the inequality (3-13) reduces to the Euclidean version (3-3). This can be seen using a similar argument to that in Remark 3.2.

#### 4. The results by the method of inverse mean curvature flow

In this section we give the proof of Theorem 1.4 using a different method from the one in the previous section.

**4A. Evolution equations.** Considering  $\Sigma$  as the initial hypersurface, we flow  $\Sigma$  in  $\mathbb{S}^{n+1}$  under the flow equation  $X : \Sigma \times [0, T^*) \rightarrow \mathbb{S}^{n+1}$ ,

$$\partial_t X = F\nu,$$

where  $F$  is a curvature function and  $\nu$  is the unit normal to the flow hypersurfaces  $\Sigma_t$ . First we recall the following evolution equations.

**Lemma 4.1** [Makowski and Scheuer 2013]. *Under the curvature flow  $\partial_t X = F\nu$  in  $\mathbb{S}^{n+1}$ , we have*

$$(4-1) \quad \frac{d}{dt} |\Sigma_t| = n \int_{\Sigma_t} F p_1 d\mu_t,$$

$$(4-2) \quad \frac{d}{dt} \int_{\Sigma_t} p_m d\mu_t = (n-m) \int_{\Sigma_t} F p_{m+1} d\mu_t - m \int_{\Sigma_t} F p_{m-1} d\mu_t.$$

To simplify the notation, in the following we define

$$(4-3) \quad \tilde{L}_k = \frac{1}{C_n^{2k} (2k)!} L_k = \sum_{i=0}^k C_k^i p_{2k-2i}.$$

Using Lemma 4.1, we obtain the following.

**Lemma 4.2.** *Under the curvature flow  $\partial_t X = F\nu$  in  $\mathbb{S}^{n+1}$ , we have*

$$\frac{d}{dt} \int_{\Sigma_t} \tilde{L}_k d\mu_t = (n-2k) \sum_{i=0}^k C_k^i \int_{\Sigma_t} F p_{2k-2i+1} d\mu_t.$$

*Proof.* The proof is by a direct calculation:

$$\begin{aligned} \frac{d}{dt} \int_{\Sigma_t} \tilde{L}_k d\mu_t &= \sum_{i=0}^k C_k^i \frac{d}{dt} \int_{\Sigma_t} p_{2k-2i} d\mu_t \\ &= \sum_{i=0}^k C_k^i \int_{\Sigma_t} ((n-2k+2i) F p_{2k-2i+1} - 2(k-i) F p_{2k-2i-1}) d\mu_t \\ &= \sum_{i=0}^k C_k^i \int_{\Sigma_t} (n-2k+2i) F p_{2k-2i+1} d\mu_t \\ &\quad - \sum_{i=1}^k C_k^{i-1} \int_{\Sigma_t} 2(k-i+1) F p_{2k-2i+1} d\mu_t \\ &= (n-2k) \sum_{i=0}^k C_k^i \int_{\Sigma_t} F p_{2k-2i+1} d\mu_t. \quad \square \end{aligned}$$

**4B. Proof of Theorem 1.4.** Recently, Makowski and Scheuer [2013] and Gerhardt [2015] studied the curvature flows in the sphere. If the initial hypersurface  $\Sigma \subset \mathbb{S}^{n+1}$  is closed and strictly convex, then under the inverse mean curvature flow

$$\partial_t X = \frac{1}{p_1} \nu,$$

there exists a finite time  $T^* < \infty$  such that the flow hypersurface  $\Sigma_t$  converges to an equator in  $\mathbb{S}^{n+1}$  and the mean curvature of  $\Sigma_t$  converges to zero almost everywhere in the sense of (see Theorem 1.4 in [Makowski and Scheuer 2013])

$$(4-4) \quad \lim_{t \rightarrow T^*} \int_{\Sigma_t} p_1^\alpha d\mu_t = 0 \quad \text{for all } 1 \leq \alpha < \infty.$$

For each  $t \in [0, T^*)$ , define the quantity  $Q(t)$  by

$$(4-5) \quad Q(t) = |\Sigma_t|^{-(n-2k)/n} \int_{\Sigma_t} \tilde{L}_k d\mu_t.$$

On the one hand, by Lemmas 4.2 and 2.1 (note that strictly convex implies all principal curvatures of  $\Sigma_t$  are positive, and certainly belong to  $\Gamma_k^+$ ), we have

$$\begin{aligned} \frac{d}{dt} \int_{\Sigma_t} \tilde{L}_k d\mu_t &= (n-2k) \sum_{i=0}^k C_k^i \int_{\Sigma_t} \frac{p_{2k-2i+1}}{p_1} d\mu_t \\ &\leq (n-2k) \sum_{i=0}^k C_k^i \int_{\Sigma_t} p_{2k-2i} d\mu_t \\ &= (n-2k) \int_{\Sigma_t} \tilde{L}_k d\mu_t. \end{aligned}$$

Equality holds if and only if  $\Sigma_t$  is totally umbilical. On the other hand, the area of the flow hypersurface evolves as

$$\frac{d}{dt} |\Sigma_t| = n |\Sigma_t|.$$

Therefore we obtain that the quantity  $Q(t)$  is monotone nonincreasing in  $t$ ; i.e.,

$$(4-6) \quad \frac{d}{dt} Q(t) \leq 0.$$

Since under the inverse mean curvature flow, the flow hypersurfaces converge to an equator in  $\mathbb{S}^{n+1}$  and the mean curvature of  $\Sigma_t$  converges to zero almost everywhere in the sense of (4-4), we have

$$(4-7) \quad \lim_{t \rightarrow T^*} Q(t) = \omega_n^{2k/n}.$$

Combining (4-6) and (4-7), we have

$$Q(0) = |\Sigma|^{-(n-2k)/n} \int_{\Sigma} \tilde{L}_k d\mu \geq \lim_{t \rightarrow T^*} Q(t) = \omega_n^{2k/n}.$$

Hence noting (4-3), we obtain that

$$(4-8) \quad \int_{\Sigma} L_k d\mu \geq C_n^{2k} (2k)! \omega_n^{2k/n} |\Sigma|^{(n-2k)/n}.$$

Equality holds in (4-8) if and only if  $Q(t)$  is constant in  $t$ . Then  $\Sigma_t$  is totally umbilical for each  $t \in [0, T^*)$ , and, in particular,  $\Sigma$  is totally umbilical and hence a geodesic sphere. The inequality (4-8) says that the induced metric of convex hypersurfaces in  $\mathbb{S}^{n+1}$  satisfies the optimal Sobolev inequalities. See [Ge et al. 2014b] for further information about the Sobolev inequalities of the same type.

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
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