# Pacific Journal of Mathematics 

IRREDUCIBLE DECOMPOSITION<br>FOR LOCAL REPRESENTATIONS OF QUANTUM TEICHMÜLLER SPACE

JÉRÉMY TOULISSE

# IRREDUCIBLE DECOMPOSITION FOR LOCAL REPRESENTATIONS OF QUANTUM TEICHMÜLLER SPACE 

JÉRÉMY TOULISSE


#### Abstract

We give an irreducible decomposition of the so-called local representations (Bai, Bonahon and Liu, 2007) of the quantum Teichmüller space $\mathcal{T}_{q}(\Sigma)$, where $\Sigma$ is a punctured surface of genus $g>0$ and $q$ is an $N$-th root of unity with $N$ odd. As an application, we construct a family of representations of the Kauffman bracket skein algebra of the closed surface $\bar{\Sigma}$.


1. Introduction 233
2. Chekhov-Fock algebra and representations of $\mathcal{T}_{q}(\Sigma) 235$
3. Decomposition of local representations 242
4. Representations of the skein algebra 248

Acknowledgment 254
References 255

## 1. Introduction

Let $\Sigma$ be the surface obtained by removing $s>0$ points $v_{1}, \ldots, v_{s}$ from the closed oriented surface $\bar{\Sigma}$ of genus $g>0$. Denote by $\mathcal{T}(\Sigma)$ the Teichmüller space of $\Sigma$, that is roughly speaking, the space of complete hyperbolic metrics on $\Sigma$. Given $\lambda$ an ideal triangulation of $\Sigma$ (that is a triangulation of the closed surface $\bar{\Sigma}$ whose vertices are exactly the $v_{i}$ ), Thurston [1986] constructed a parametrization of $\mathcal{T}(\Sigma)$ by associating a strictly positive real number to each edge $\lambda_{i}$ of the ideal triangulation, $i \in\{1, \ldots, n\}$, where $n=6 g-6+3 s$ is the number of edges of $\lambda$. These coordinates are called the shear coordinates associated to $\lambda$. In this coordinate system, the coefficients of the Weil-Petersson form on $\mathcal{T}(\Sigma)$ depend only on the combinatorics of $\lambda$ and are easy to compute.

For a parameter $q \in \mathbb{C}^{*}$, Chekhov and Fock [1999] defined the quantum Teichmüller space $\mathcal{T}_{q}(\Sigma)$ of $\Sigma$, which is a deformation quantization of the Poisson algebra of a certain class of functions over $\mathcal{T}(\Sigma)$; see also [Kashaev 1998] for a slightly different version and [Guo and Liu 2009] for a relation between the two.

[^0]This algebraic object is obtained by gluing together a collection of noncommutative algebras $\mathcal{T}_{q}(\lambda)$, called Chekhov-Fock algebras, canonically associated to each ideal triangulation of $\Sigma$. A representation of $\mathcal{T}_{q}(\Sigma)$ is then a family of representations $\left\{\rho_{\lambda}: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}(V)\right\}_{\lambda \in \Lambda(\Sigma)}$, where $\Lambda(\Sigma)$ is the space of all ideal triangulations of $\Sigma$, and $\rho_{\lambda}$ and $\rho_{\lambda^{\prime}}$ satisfy compatibility conditions whenever $\lambda \neq \lambda^{\prime}$. For $\lambda \in \Lambda(\Sigma)$, the representation $\rho_{\lambda}$ is an avatar of the representation of $\mathcal{T}_{q}(\Sigma)$ and carries almost all the information.

When $q$ is a primitive $N$-th root of unity, $\mathcal{T}_{q}(\lambda)$ admits finite-dimensional representations. In this paper, we will consider the case that $N$ is odd. The irreducible representations of $\mathcal{T}_{q}(\lambda)$ were studied in [Bonahon and Liu 2007], where it was shown that an irreducible representation of $\mathcal{T}_{q}(\lambda)$ is classified (up to isomorphism) by a weight $x_{i} \in \mathbb{C}^{*}$ assigned to each edge $\lambda_{i}$, a choice of $N$-th root $p_{j}=\left(x_{1}^{k_{j_{1}}} \ldots x_{n}^{k_{j_{n}}}\right)^{1 / N}$ associated to each puncture $v_{j}$ (where $k_{j_{i}}$ is the number of times a small simple loop around $v_{j}$ intersects $\lambda_{i}$ ) and an $N$-th root $c=\left(x_{1} \ldots x_{n}\right)^{1 / N}$. Such a representation has dimension $N^{3 g-3+s}$.

Bai, Bonahon, and Liu [Bai et al. 2007] introduced another type of representations of $\mathcal{T}_{q}(\lambda)$, called local representations, which are well behaved under cut and paste. A local representation of $\mathcal{T}_{q}(\lambda)$ is defined by an embedding into the tensorial product of triangle algebras (see definitions below). Isomorphism classes of local representations of $\mathcal{T}_{q}(\lambda)$ are classified by a weight $x_{i} \in \mathbb{C}^{*}$ associated to each edge $\lambda_{i}$ and a choice of an $N$-th root $c=\left(x_{1} \ldots x_{n}\right)^{1 / N}$. Such a representation has dimension $N^{4 g-4+2 s}$.

It follows that a local representation of $\mathcal{T}_{q}(\lambda)$ is not irreducible. In this paper, we address the question of the decomposition of a local representation into its irreducible components. We prove:

Main Theorem. Let $\lambda$ be an ideal triangulation of $\Sigma$ and $\rho$ be a local representation of $\mathcal{T}_{q}(\lambda)$ classified by weight $x_{j} \in \mathbb{C}^{*}$ associated to each edge $\lambda_{j}$ and a choice of $N$-th root $c=\left(x_{1} \ldots x_{n}\right)^{1 / N}$. Then we have the decomposition

$$
\rho=\bigoplus_{i \in \mathcal{I}} \rho^{(i)}
$$

Here, $\rho^{(i)}$ is an irreducible representation classified by the same $x_{j}$, an $N$-th root $p_{j}^{(i)}=\left(x_{1}^{k_{j_{1}}} \cdots x_{n}^{k_{j_{n}}}\right)^{1 / N}$ associated to each puncture, and the same c. Moreover, $\mathcal{I}$ is a finite set such that, given a choice of an $N$-th root $p_{j}=\left(x_{1}^{k_{j_{1}}} \cdots x_{n}^{k_{j_{n}}}\right)^{1 / N}$ for each puncture, there exists exactly $N^{g}$ elements $i \in \mathcal{I}$ with $p_{j}^{(i)}=p_{j}$ for all $j \in\{1, \ldots, s\}$.

It is proved by Bai [2007] that if $\lambda$ and $\lambda^{\prime}$ are two different triangulations of the square related by a diagonal switch, then the intertwining operators associated to two isomorphic representations $\rho: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}(V)$ and $\rho^{\prime}: \mathcal{T}_{q}\left(\lambda^{\prime}\right) \rightarrow \operatorname{End}\left(V^{\prime}\right)$ correspond to the $6 j$-symbols defined by Kashaev [1995]. These $6 j$-symbols relate
hyperbolic geometry and quantum invariants and gave birth to the famous volume conjecture; see [Murakami 2011] for an overview.

In particular, Baseilhac and Benedetti [2005] used these $6 j$-symbols to construct a $(2+1)$-dimensional topological quantum field theory (TQFT) on manifolds with $\operatorname{PSL}(2, \mathbb{C})$-character. Our result thus provides a decomposition of the vector spaces arising in the TQFT.

As an application, we adapt the construction of Bonahon and Wong [2015] to local representations of the balanced Chekhov-Fock algebra and obtain a family of representations of the Kauffman bracket skein algebra $\mathcal{S}_{A}(\bar{\Sigma})$ of the closed surface $\bar{\Sigma}$. The vector space associated to this family of representations is canonically associated to an ideal triangulation $\lambda$. In particular, it makes the computations very explicit. It also behaves well under cut and paste.

In Section 2, we recall the definition of the Chekhov-Fock algebra, the quantum Teichmüller space, the triangle algebra and the local representations. In Section 3, we prove the Main Theorem. Finally, in Section 4, we explain the connections between quantum Teichmüller theory, skein theory and construct a new family of representations of $\mathcal{S}_{A}(\bar{\Sigma})$.

## 2. Chekhov-Fock algebra and representations of $\mathcal{T}_{q}(\boldsymbol{\Sigma})$

In this section, we define the Chekhov-Fock algebra $\mathcal{T}_{q}(\lambda)$ associated to an ideal triangulation $\lambda$, describe its representations and recall the definition of the quantum Teichmüller space. Most results come from [Bonahon and Liu 2007; Bai et al. 2007].

In all this paper, for an integer $n \in \mathbb{N}$, set $\mathbb{Z}_{n}:=\mathbb{Z} / n \mathbb{Z}$ and denote by $\mathcal{U}(N)$ the group of $N$-th roots of unity.
2.1. The Chekhov-Fock algebra. In this subsection, $q$ is a formal parameter and $\Sigma$ is allowed to have boundary components with punctures on the boundary (and every boundary component has at least one puncture).

Let $\lambda$ be an ideal triangulation of $\Sigma$. We denote by $\lambda_{1}, \ldots, \lambda_{n}$ the edges of $\lambda$. The Fock's matrix associated to $\lambda$ is the skew-symmetric $n \times n$ matrix with integer coefficients $\eta_{\lambda}=\left(\sigma_{i j}\right)_{i, j=1, \ldots n}$ defined by

$$
\sigma_{i j}=a_{i j}-a_{j i}
$$

where $a_{i j}$ is the number of angular sector delimited by $\lambda_{i}$ and $\lambda_{j}$ in the faces of $\lambda$ with $\lambda_{i}$ coming before $\lambda_{j}$ counterclockwise.
Definition 2.1. The Chekhov-Fock algebra of $\lambda$ is the algebra $\mathcal{T}_{q}(\lambda)$ freely generated by the elements $X_{i}^{ \pm 1}, i \in\{1, \ldots, n\}$, subject to the relations

$$
X_{i} X_{j}=q^{2 \sigma_{i j}} X_{j} X_{i}
$$



Figure 1. The triangle $T$.

The following example is of first importance.
Example 2.2. Let $T$ be a disk with three punctures $v_{1}, v_{2}, v_{3}$ on the boundary. The boundary arcs between the punctures provides a natural triangulation $\lambda$ of $T$ (see Figure 1).

The triangle algebra is $\mathcal{T}:=\mathcal{T}_{q}(\lambda)$. It is generated by $X_{i}^{ \pm 1}, i=1,2,3$, subject to the relations

$$
X_{i} X_{i+1}=q^{2} X_{i+1} X_{i}, \quad i \in \mathbb{Z}_{3} .
$$

The algebraic structure of the Chekhov-Fock algebra is fairly simple. In particular, it is a quantum torus [Goodearl and Warfield 2004].

Given a monomial $X=X_{1}^{k_{1}} \ldots X_{n}^{k_{n}} \in \mathcal{T}_{q}(\lambda)$ for a multi-index $\mathbf{k}=\left(k_{1}, \ldots, k_{n}\right) \in$ $\mathbb{Z}^{n}$, we define the Weyl ordering of $X$ to be the monomial

$$
[X]:=q^{-\sum_{i<j} \sigma_{i j}} X_{1}^{k_{1}} \ldots X_{n}^{k_{n}} .
$$

The advantage of the Weyl ordering is its independence with respect to the order of the terms. In particular, for any permutation $\sigma:\{1, \ldots, n\} \rightarrow\{1, \ldots, n\}$, we have

$$
\left[X_{1}^{k_{1}} \ldots X_{n}^{k_{n}}\right]=\left[X_{\sigma(1)}^{k_{\sigma(1)}} \ldots X_{\sigma(n)}^{k_{\sigma(n)}}\right]
$$

For a multi-index $\mathbf{k}=\left(k_{1}, \ldots, k_{n}\right) \in \mathbb{Z}^{n}$, we define $X_{\mathbf{k}}:=\left[X_{1}^{k_{1}} \ldots X_{n}^{k_{n}}\right] \in \mathcal{T}_{q}(\lambda)$.
2.2. Finite-dimensional representations of $\mathcal{T}_{q}(\lambda)$. When the parameter $q$ is a root of unity, the structure of the Chekhov-Fock algebra is drastically different. In particular, $\mathcal{T}_{q}(\lambda)$ admits finite dimensional representations that we describe here.

In this subsection, $q \in \mathbb{C}^{*}$ is a primitive $N$-th root of unity with $N$ odd, $\Sigma$ has no boundary component and $\lambda$ is an ideal triangulation of $\Sigma$ with edges labeled $\lambda_{1}, \ldots, \lambda_{n}$.

Definition 2.3. For each puncture $v_{j}$, the puncture invariant $P_{j}=\left[X_{1}^{k_{j_{1}}} \ldots X_{n}^{k_{j_{n}}}\right] \in$ $\mathcal{T}_{q}(\lambda)$ is the monomial associated to the multi-index $\mathbf{k}_{j}=\left(k_{j_{1}}, \ldots, k_{j_{n}}\right) \in \mathbb{N}^{n}$, where $k_{j_{i}}$ is the minimum number of intersections between the edge $\lambda_{i}$ and a closed simple loop around $v_{j}$.

The puncture invariants are of main importance in the representation theory of the Chekhov-Fock algebra. In particular:

Proposition 2.4 [Bonahon and Liu 2007, Proposition 15]. The center of $\mathcal{T}_{q}(\lambda)$ is generated by:

- $X_{i}^{N}$ for each $i \in\{1, \ldots, n\}$.
- The puncture invariant $P_{j}$ associated to each puncture $v_{j} \in\left\{v_{1}, \ldots, v_{s}\right\}$.
- The element $H:=\left[X_{1} \ldots X_{n}\right]$.

Note that $\left[P_{1} \ldots P_{s}\right]=\left[H^{2}\right]$.
A representation of $\mathcal{T}_{q}(\lambda)$ is a morphism $\rho: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}(V)$ where $V$ is a vector space. Such a representation is finite-dimensional when $V$ is finite-dimensional and $\rho$ is called irreducible when there is no proper linear subspace $W \subset V$ preserved by $\rho\left(\mathcal{T}_{q}(\lambda)\right)$. Two representations $\rho: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}(V)$ and $\rho^{\prime}: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}\left(V^{\prime}\right)$ are isomorphic if there exists a linear isomorphism $L: V \rightarrow V^{\prime}$ such that

$$
\rho^{\prime}(X)=L \circ \rho(X) \circ L^{-1} \quad \text { for } X \in \mathcal{T}_{q}(\lambda)
$$

Theorem 2.5 [Bonahon and Liu 2007, Theorems 20 and 21]. Up to isomorphism, any irreducible representation

$$
\rho: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}(V)
$$

is determined by its restriction to the center of $\mathcal{T}_{q}(\lambda)$ and is classified by a nonzero complex number $x_{i}$ associated to each edges $\lambda_{i}$, a choice of an $N$-th root $p_{j}=$ $\left(x_{1}^{k_{j_{1}}} \ldots x_{n}^{k_{j_{n}}}\right)^{1 / N}$ for each puncture $v_{j}$ (where the $k_{j_{k}} \in\{1,2\}$ are as in Definition 2.3) and a choice of a square root $c=\left(p_{0} \ldots p_{s}\right)^{1 / 2}$.

Such a representation is characterized by

- $\rho\left(X_{i}^{N}\right)=x_{i} \operatorname{Id}_{V}$,
- $\rho\left(P_{j}\right)=p_{j} \operatorname{Id}_{V}$,
- $\rho(H)=c \operatorname{Id}_{V}$.

Moreover, such a representation has dimension $N^{3 g-3+s}$.
Let us come back to Example 2.2. Recall that the triangle algebra $\mathcal{T}$ is the algebra generated by $X_{i}^{ \pm 1}, i \in \mathbb{Z}_{3}$, with relations $X_{i} X_{i+1}=q^{2} X_{i+1} X_{i}$.

The center of $\mathcal{T}$ is generated by $X_{i}^{N}$ and $H=q^{-1} X_{1} X_{2} X_{3}$. One easily checks that irreducible representations of $\mathcal{T}$ have dimension $N$ and are classified (up to isomorphism) by a choice of weight $x_{i} \in \mathbb{C}^{*}$ associated to each edge $\lambda_{i}$ and a central charge, that is a choice of an $N$-th root $c=\left(x_{1} x_{2} x_{3}\right)^{1 / N}$; see [Bai et al. 2007, Lemma 2] for more details.

More precisely, let $V$ be the complex vector space generated by $\left\{e_{1}, \ldots, e_{N}\right\}$ and let $\rho$ be an irreducible representation of $\mathcal{T}$ classified by $x_{1}, x_{2}, x_{3} \in \mathbb{C}^{*}$ and $c=\left(x_{1} x_{2} x_{3}\right)^{1 / N}$. Then, up to isomorphism, the action of $\mathcal{T}$ on $V$ is defined by

$$
\rho\left(X_{1}\right) e_{i}=x_{1}^{\prime} q^{2 i} e_{i}, \quad \rho\left(X_{2}\right) e_{i}=x_{2}^{\prime} e_{i+1}, \quad \rho\left(X_{3}\right) e_{i}=x_{3}^{\prime} q^{1-2 i} e_{i-1},
$$

where $x_{i}^{\prime}$ is an $N$-th root of $x_{i}$ such that $x_{1}^{\prime} x_{2}^{\prime} x_{3}^{\prime}=c$. Note that, up to isomorphism, $\rho$ is independent of the choice of the $N$-th root $x_{i}^{\prime}$ with $x_{1}^{\prime} x_{2}^{\prime} x_{3}^{\prime}=c$.

The following lemma will be useful in the next section. Recall that $\mathcal{U}(N)$ is the group of $N$-th roots of unity.
Lemma 2.6. Let $\rho: \mathcal{T} \rightarrow \operatorname{End}(V)$ be the representation of the triangle algebra classified by $x_{1}=x_{2}=x_{3}=1$ and $c \in \mathcal{U}(N)$. For each $i \in \mathbb{Z}_{3}$ and $N$-th root $\zeta \in \mathcal{U}(N)$, the eigenspace of $\rho\left(X_{i}\right)$ of eigenvalue $\zeta$ is one-dimensional.
Proof. We use the explicit form of the representation $\rho$ in $V=\operatorname{span}\left\{e_{1}, \ldots, e_{N}\right\}$ described above. Set $x_{1}^{\prime}=x_{2}^{\prime}=1, x_{3}^{\prime}=c$ and $\zeta=q^{2 k}$ for some $k \in\{0, \ldots, N-1\}$.

For $i=1$, one checks that the eigenspace of $\rho\left(X_{1}\right)$ associated to $\zeta$ is generated by $e_{k}$.

For $i=2$, the vector $\alpha_{k}:=\sum_{i \in \mathbb{Z}_{N}} q^{-2 k i} e_{i}$ satisfies $\rho\left(X_{2}\right) \alpha_{k}=q^{2 k} \alpha_{k}$ and $\left\{\alpha_{1}, \ldots, \alpha_{k}\right\}$ form a basis of $V$. So the eigenspace of $\rho\left(X_{2}\right)$ associated to the eigenvalue $\zeta$ is generated by $\alpha_{k}$.

For $i=3$, we use the fact that $\rho\left(q^{-1} X_{1} X_{2} X_{3}\right)=c \operatorname{Id}_{V}$, where $c \in \mathcal{U}(N)$.
An ideal triangulation of $\Sigma$ is composed by $m$ faces $T_{1}, \ldots, T_{m}$. Each face $T_{j}$ determines a triangle algebra $\mathcal{T}_{j}$ whose generators are associated to the three edges of $T_{j}$. It provides a canonical embedding $\iota$ of $\mathcal{T}_{q}(\lambda)$ into $\mathcal{T}_{1} \otimes \cdots \otimes \mathcal{T}_{m}$ defined on the generators as follows:

- $\iota\left(X_{i}\right)=X_{j_{i}} \otimes X_{k_{i}}$ if $\lambda_{i}$ belongs to two distinct triangles $T_{j}$ and $T_{k}$ and $X_{j_{i}} \in$ $\mathcal{T}_{j}, X_{k_{i}} \in \mathcal{T}_{k}$ are the generators associated to the edge $\lambda_{i} \in T_{j}$ and $\lambda_{i} \in T_{k}$ respectively.
- $\iota\left(X_{i}\right)=\left[X_{j_{i_{1}}} X_{j_{i_{2}}}\right]$ if $\lambda_{i}$ corresponds to two sides of the same face $T_{j}$ and $X_{j_{i_{1}}}, X_{i_{j_{2}}} \in \mathcal{T}_{j}$ are the associated generators.
Definition 2.7. A local representation of $\mathcal{T}_{q}(\lambda)$ is a representation which factorizes as $\left(\rho_{1} \otimes \cdots \otimes \rho_{m}\right) \circ \iota$ where $\rho_{i}: \mathcal{T}_{i} \rightarrow V_{i}$ is an irreducible representation of the triangle algebra $\mathcal{T}_{i}$ and $\iota: \mathcal{T}_{q}(\lambda) \rightarrow \mathcal{T}_{1} \otimes \cdots \otimes \mathcal{T}_{m}$ is defined as above.

In particular, a local representation has dimension $N^{m}$ where $m=4 g-4+2 s$ is the number of faces of the triangulation.

Local representations were first introduced by Bai et al. [2007].
Theorem 2.8 [Bai et al. 2007, Proposition 6]. Up to isomorphism, a local representation

$$
\rho: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}(V)
$$



Figure 2. Flip of triangulation.
is classified by a nonzero complex number $x_{i}$ associated to the edge $\lambda_{i}$ and a choice of an $N$-th root $c=\left(x_{1} \ldots x_{n}\right)^{1 / N}$. Such a representation satisfies

- $\rho\left(X_{i}^{N}\right)=x_{i} \operatorname{Id}_{V}$,
- $\rho(H)=c \mathrm{Id}_{V}$.

Local representations have certain advantages over irreducible representations. First of all, these representations behave very well under cut and paste, so one can use them to construct invariant of 3-manifolds; see [Baseilhac and Benedetti 2005]. Also, the vector space associated to a local representation decomposes as a tensor product of vector spaces and each generator $X_{i} \in \mathcal{T}_{q}(\lambda)$ associated to an edge $\lambda_{i}$ only acts on the vector spaces associated to triangle adjacent to the edge $\lambda_{i}$ (that is why these representations are called local).
2.3. Quantum Teichmïller space and its representations. The quantum Teichmüller space is obtained by gluing together a family of (division algebras of) Chekhov-Fock algebras indexed by the set of ideal triangulations of $\Sigma$.

The simplex of ideal triangulations. Let $\Lambda(\Sigma)$ be the set of ideal triangulations of $\Sigma$. We say that two triangulations $\lambda, \lambda^{\prime} \in \Lambda(\Sigma)$ differ by a flip if $\lambda$ and $\lambda^{\prime}$ coincide everywhere except in a square made of two adjacent triangles where they differ as in Figure 2.

The graph of ideal triangulations is the graph whose set of vertices is $\Lambda(\Sigma)$ and two vertices $\lambda, \lambda^{\prime} \in \Lambda(\Sigma)$ are connected by an edge if and only if $\lambda$ and $\lambda^{\prime}$ differ by a flip.

The simplex of ideal triangulations is obtained from the graph of ideal triangulations by gluing a 2 -simplex on each cycle corresponding to the pentagon relation (see Figure 3).

Proposition 2.9 [Penner 1993]. The simplex of ideal triangulations is connected and simply connected. Namely, any two different ideal triangulations are connected by a sequence of flips and two sequences between two ideal triangulations differ by a sequence of pentagon relations.

Coordinate change. The Chekhov-Fock algebra $\mathcal{T}_{q}(\lambda)$ associated to an ideal triangulation $\lambda \in \Lambda(\Sigma)$ satisfies the Ore condition; see [Goodearl and Warfield 2004].


Figure 3. Pentagon relation.

In particular, $\mathcal{T}_{q}(\lambda)$ has a well-defined division algebra $\hat{\mathcal{T}}_{q}(\lambda)$ consisting of rational fractions satisfying some noncommutativity relations.

Let $\lambda, \lambda^{\prime} \in \Lambda(\Sigma)$ be two ideal triangulations related by a flip. Chekhov and Fock [1999] constructed coordinates change isomorphisms

$$
\Psi_{\lambda \lambda^{\prime}}^{q}: \hat{\mathcal{T}}_{q}\left(\lambda^{\prime}\right) \rightarrow \hat{\mathcal{T}}_{q}(\lambda)
$$

These coordinates change satisfy the pentagon relation. In particular, using the result of Penner, they extend uniquely to coordinates change $\Psi_{\lambda \lambda^{\prime}}^{q}: \hat{\mathcal{T}}_{q}\left(\lambda^{\prime}\right) \rightarrow \hat{\mathcal{T}}_{q}(\lambda)$ for any two different ideal triangulations $\lambda, \lambda^{\prime} \in \Lambda(\Sigma)$.

It was proved in [Liu 2009] that these coordinates change are the unique ones satisfying some natural relations, as for instance $\Psi_{\lambda \lambda^{\prime \prime}}^{q}=\Psi_{\lambda \lambda^{\prime}}^{q} \circ \Psi_{\lambda^{\prime} \lambda^{\prime \prime}}^{q}$ for any $\lambda, \lambda^{\prime}, \lambda^{\prime \prime} \in \Lambda(\Sigma)$. Moreover, when $q=1$, these maps reduce to the classical coordinates change in Teichmüller theory; see [loc. cit.] for more details.

Definition 2.10. The quantum Teichmüller space of $\Sigma$ is defined by

$$
\mathcal{T}_{q}(\Sigma):=\bigsqcup_{\lambda \in \Lambda(\Sigma)} \hat{\mathcal{T}}_{q}(\lambda) / \sim
$$

where $x_{\lambda} \in \hat{\mathcal{T}}_{q}(\lambda) \sim x_{\lambda^{\prime}} \in \hat{\mathcal{T}}_{q}\left(\lambda^{\prime}\right)$ if and only if $x_{\lambda}=\Psi_{\lambda \lambda^{\prime}}^{q}\left(x_{\lambda^{\prime}}\right)$.
Note that, as Since each coordinate change $\Psi_{\lambda \lambda^{\prime}}^{q}$ is an algebra isomorphism, $\mathcal{T}_{q}(\Sigma)$ inherits an algebra structure, and the $\hat{\mathcal{T}}_{q}(\lambda)$ can be thought as "global coordinates" on $\mathcal{T}_{q}(\Sigma)$.

Representations. A natural definition for a finite dimensional representation of $\mathcal{T}_{q}(\Sigma)$ would be a family of finite dimensional representations

$$
\left\{\rho_{\lambda}: \hat{\mathcal{T}}_{q}(\lambda) \rightarrow \operatorname{End}\left(V_{\lambda}\right)\right\}_{\lambda \in \Lambda(\Sigma)}
$$

such that for each pair $\lambda, \lambda^{\prime} \in \Lambda(\Sigma)$, the representations $\rho_{\lambda^{\prime}}$ and $\rho_{\lambda} \circ \Psi_{\lambda, \lambda^{\prime}}^{q}$ of $\hat{\mathcal{T}}_{q}\left(\lambda^{\prime}\right)$ are isomorphic (as representations).

However, as pointed out in [Bai et al. 2007, Section 4.2], when $V_{\lambda}$ is finitedimensional, there is no morphism $\hat{\mathcal{T}}_{q}(\lambda) \rightarrow \operatorname{End}\left(V_{\lambda}\right)$. In fact, $\hat{\mathcal{T}}_{q}(\lambda)$ is infinitedimensional as a vector space while $\operatorname{End}\left(V_{\lambda}\right)$ is finite-dimensional and so, such a homomorphism $\rho_{\lambda}$ would have nonzero kernel. In particular, there would exists elements $x \in \hat{\mathcal{T}}_{q}(\lambda)$ such that $\rho_{\lambda}(x)=0$ and so, $\rho_{\lambda}\left(x^{-1}\right)$ would make no sense.

This motivates the following definition:
Definition 2.11. A local representation (respectively an irreducible representation) of $\mathcal{T}_{q}(\Sigma)$ is a family of representations

$$
\left\{\rho_{\lambda}: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}\left(V_{\lambda}\right)\right\}_{\lambda \in \Lambda(\Sigma)}
$$

such that for each $\lambda, \lambda^{\prime} \in \Lambda(\Sigma), \rho_{\lambda}$ is a local representation (respectively an irreducible representation) of $\mathcal{T}_{q}(\lambda)$, and $\rho_{\lambda^{\prime}}$ is isomorphic to $\rho_{\lambda} \circ \Psi_{\lambda \lambda^{\prime}}^{q}$ whenever $\rho_{\lambda} \circ \Psi_{\lambda \lambda^{\prime}}^{q}$ makes sense.

We say that $\rho_{\lambda} \circ \Psi_{\lambda \lambda^{\prime}}^{q}$ makes sense, if for each Laurent polynomial $X^{\prime} \in \mathcal{T}_{q}\left(\lambda^{\prime}\right)$,

$$
\Psi_{\lambda \lambda^{\prime}}\left(X^{\prime}\right)=P Q^{-1}=Q^{\prime-1} P^{\prime} \in \hat{\mathcal{T}}_{q}(\lambda), \quad \text { for some } P, Q, P^{\prime}, Q^{\prime} \in \mathcal{T}_{q}(\lambda)
$$

In that case, we define

$$
\rho_{\lambda} \circ \Psi_{\lambda \lambda^{\prime}}\left(X^{\prime}\right):=\rho_{\lambda}(P) \rho_{\lambda}(Q)^{-1}=\rho_{\lambda}\left(Q^{\prime}\right)^{-1} \rho_{\lambda}\left(P^{\prime}\right)
$$

Proposition 2.12 [Bai et al. 2007, Proposition 10]. Let $\lambda, \lambda^{\prime} \in \Lambda(\Sigma)$ be two ideal triangulations of $\Sigma$. Then there exists a rational map

$$
\varphi_{\lambda \lambda^{\prime}}: \mathbb{C}^{n} \rightarrow \mathbb{C}^{n}
$$

such that a local representation $\rho^{\prime}: \mathcal{T}_{q}\left(\lambda^{\prime}\right) \rightarrow \operatorname{End}\left(V_{\lambda^{\prime}}\right)$ classified by $\left(x_{1}^{\prime}, \ldots, x_{n}^{\prime}\right)$ and $c^{\prime}=\left(x_{1}^{\prime} \ldots x_{n}^{\prime}\right)^{1 / N}$ is isomorphic to $\rho_{\lambda} \circ \Psi_{\lambda \lambda^{\prime}}$ (whenever it makes sense) where $\rho_{\lambda}: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}\left(V_{\lambda}\right)$ is a local representation classified by $\left(x_{1}, \ldots, x_{n}\right)$ and $c=\left(x_{1} \ldots x_{n}\right)^{1 / N}$ if and only if $c=c^{\prime}$ and

$$
\left(x_{1}^{\prime}, \ldots, x_{n}^{\prime}\right)=\varphi_{\lambda \lambda^{\prime}}\left(x_{1}, \ldots, x_{n}\right)
$$

Remark 2.13. The analogue is also proved in [Bonahon and Liu 2007] for irreducible representations. In particular, the rational maps $\varphi_{\lambda \lambda^{\prime}}$ are the same.

It turns out that the rational maps $\varphi_{\lambda \lambda^{\prime}}$ correspond to the coordinates change of the shear-bend coordinates on the character variety $\chi(\Sigma, \operatorname{SL}(2, \mathbb{C}))$.

As a result, isomorphism classes of local (respectively irreducible) representations of $\mathcal{T}_{q}(\Sigma)$ are classified, up to finitely many choices, by the character variety $\chi(\Sigma, \operatorname{SL}(2, \mathbb{C}))$; see [loc. cit.] for more details.

## 3. Decomposition of local representations

In this section, we prove the Main Theorem. Let $\rho: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}(V)$ be the local representation classified by the nonzero complex number $x_{i}$ associated to each edge and the central charge $c$. Given a puncture invariant $P_{j}=\left[X_{1}^{k_{j_{1}}} \ldots X_{n}^{k_{j_{n}}}\right]$ (see Proposition 2.4) associated to the puncture $v_{j}$, the representation $\rho$ satisfies

$$
\rho\left(P_{j}^{N}\right)=x_{1}^{k_{j_{1}}} \ldots x_{n}^{k_{j_{n}}} \operatorname{Id}_{V}
$$

It follows that if $p_{j}$ is an eigenvalue of $P_{j}$, then $p_{j}^{N}=x_{1}^{k_{j_{1}}} \ldots x_{n}^{k_{j_{n}}}$.

## Notation.

- Given $p_{j} \in \mathbb{C}^{*}$ so that $p_{j}^{N}=x_{1}^{k_{j_{1}}} \ldots x_{n}^{k_{j_{n}}}$, we denote by

$$
V_{p_{j}}\left(P_{j}\right):=\left\{x \in V: \rho\left(P_{j}\right) x=p_{j} x\right\}
$$

the associated eigenspace.

- Given $\boldsymbol{p}=\left(p_{1}, \ldots, p_{s}\right)$ so that for each $j, p_{j}^{N}=x_{1}^{k_{j_{1}}} \ldots x_{n}^{k_{j_{n}}}$, set

$$
V_{\boldsymbol{p}}:=\left\{x \in V: \rho\left(P_{j}\right) x=p_{j} x, j=1, \ldots, s\right\}=\bigcap_{j=1}^{s} V_{p_{j}}\left(P_{j}\right) .
$$

The proof of the Main Theorem will follow the next proposition:
Proposition 3.1. There exists an ideal triangulation $\lambda_{0} \in \Lambda(\Sigma)$ such that for each $\boldsymbol{p}$ as above, $V_{\boldsymbol{p}}$ has dimension $N^{4 g-3+s}$.
Proof. The dimension of $V_{\boldsymbol{p}}$ does not depend on the numbers $x_{i} \in \mathbb{C}^{*}$ characterizing $\rho$. In this proof, we will consider all the $x_{i}$ equal to 1 , which implies that the eigenvalues of $\rho\left(P_{i}\right)$ are root of unity.

Consider the one punctured surface $\Sigma^{\prime}:=\Sigma \cup\left\{v_{1}, \ldots, v_{s-1}\right\}$. As $g>0$, there exists an ideal triangulation $\lambda^{\prime}$ of $\Sigma^{\prime}$. Let $T$ be a triangle of the triangulation $\lambda^{\prime}$ and consider the triangulation of $T \backslash\left\{v_{1}, \ldots, v_{s-1}\right\}$ as in Figure 4.

The union of the triangulation $\lambda^{\prime}$ and the one of $T$ gives an ideal triangulation $\lambda_{0}$ of $\Sigma$.

Consider a local representation $\rho: \mathcal{T}_{q}\left(\lambda_{0}\right) \rightarrow \operatorname{End}(V)$. The decomposition of the ideal triangulation $\lambda_{0}$ gives the nice decomposition

$$
V=W \otimes W^{\prime}
$$



Figure 4. Triangulation of $T \cup\left\{v_{1}, \ldots, v_{s}\right\}$.
where $W^{\prime}$ is the vector space corresponding to the triangles of the triangulation $\lambda^{\prime}$ (except the triangle $T$ ), and $W$ corresponds to the triangles of $T$.

In particular, as the triangulation $\lambda^{\prime}$ contains $4 g-2$ triangles, $\operatorname{dim}\left(W^{\prime}\right)=N^{4 g-3}$ (remember that we do not consider the vector space associated to $T$ ), and $\operatorname{dim} W=$ $N^{2 s-1}$.

The interest of the triangulation $\lambda_{0}$ is clear: the puncture invariant $P_{i}$ associated to the puncture $v_{i} \neq v_{s}$ acts as the identity on $W^{\prime}$, so the eigenspaces of $\rho\left(P_{i}\right)$ has the form $E \otimes W^{\prime}$ where $E \subset W$ is an eigenspace of the restriction of $\rho\left(P_{i}\right)$ on $W$. It motivates the following notation:

## Notation.

- The vector space $W$ decomposes as

$$
W=W^{0} \otimes \cdots \otimes W^{s-1}
$$

where $W^{0}$ is associated to $T_{0}$ and $W^{j}$ to $T_{j}$ and $T_{j}^{\prime}$ for $j=1, \ldots, s-1$.

- Given a root of unity $p_{k} \in \mathcal{U}(N)$, set

$$
W_{p_{k}}^{j}\left(P_{k}\right):=\left\{x \in W^{j}: \rho\left(P_{k}\right) x=p_{k} x\right\} .
$$

- For $\boldsymbol{p}=\left(p_{1}, \ldots, p_{s-1}\right) \in \mathcal{U}(N)^{s-1}$, set

$$
W_{p}^{j}=\left\{x \in W^{j}: \rho\left(P_{k}\right) x=p_{k} x, k=1, \ldots, s-1\right\}=\bigcap_{k=1}^{s-1} W_{p_{k}}^{j}\left(P_{k}\right)
$$

- Finally, set

$$
W_{p}=\left\{x \in W: \rho\left(P_{k}\right) x=p_{k} x, k=1, \ldots, s-1\right\} .
$$



Figure 5. The generators of $\mathcal{T}_{j}$ and $\mathcal{T}_{j}^{\prime}$.
Lemma 3.2. Using the above notation, and given $\boldsymbol{p} \in \mathcal{U}(N)^{s-1}$, we have the following:
(1) $\operatorname{dim} W_{\boldsymbol{p}}^{0}= \begin{cases}1 & \text { if } \boldsymbol{p}=\left(p_{1}, 1, \ldots, 1\right), \\ 0 & \text { otherwise } .\end{cases}$
(2) For $j \in\{1, \ldots, s-2\}$,

$$
\operatorname{dim} W_{\boldsymbol{p}}^{j}= \begin{cases}1 & \text { if } \boldsymbol{p}=\left(1, \ldots, 1, p_{j}, p_{j+1}, 1, \ldots, 1\right), \\ 0 & \text { otherwise } .\end{cases}
$$

(3) $\operatorname{dim} W_{p}^{s-1}= \begin{cases}N & \text { if } \boldsymbol{p}=\left(1, \ldots, 1, p_{s-1}\right), \\ 0 & \text { otherwise. }\end{cases}$

Proof. (1) If $k \neq 1, v_{k}$ is not a vertex of $T_{0}$. It follows that $P_{k}$ acts on $W^{0}$ by the identity; so if $p_{k} \neq 1, W_{\boldsymbol{p}}^{0}=\{0\}$.

Now, if $p_{k}=1$ for all $k \neq 1$, then $W_{p}^{0}$ is the eigenspace of the action on $W^{0}$ of the edge opposite to $v_{1}$. By Lemma 2.6, it is one-dimensional.
(2) Fix $j \in\{1, \ldots, s-2\}$. For $k \notin\{j, j+1\}$, $v_{k}$ is neither a vertex of $T_{j}$ nor of $T_{j}^{\prime}$. Hence $P_{j}$ acts on $W^{j}$ as the identity, and if $p_{k} \neq 1$, then $W_{p}^{j}=\{0\}$.

Take $p_{k}=1$ for all $k \notin\{j, j+1\}$ and denote by $X^{ \pm 1}, Y^{ \pm 1}, Z^{ \pm 1}$ (respectively $X^{\prime \pm 1}, Y^{\prime \pm 1}, Z^{\prime \pm 1}$ ) the generators of the triangle algebras $\mathcal{T}_{j}$ (respectively $\mathcal{T}_{j}^{\prime}$ ) associated to the triangles $T_{j}$ (respectively $T^{\prime}{ }_{j}$ ) as in Figure 5. Set also $W^{j}=V^{j} \otimes V^{\prime j}$ where $V^{j}$ (respectively $V^{\prime j}$ ) is the vector space associated to the representation of the triangle algebra $\mathcal{T}_{j}$ (respectively $\mathcal{T}_{j}^{\prime}$ ).

Denote by $c_{j}, c_{j}^{\prime} \in \mathcal{U}(N)$ the central charges of the restriction of the representation to $\mathcal{T}_{j}$ and $\mathcal{T}_{j}^{\prime}$ respectively. Then $\rho\left(P_{j}\right)$ acts on $V^{j}:=\operatorname{span}\left\{e_{0}, \ldots, e_{N-1}\right\}$ like $c_{j} Z^{-1}$ and on $V^{\prime j}=\operatorname{span}\left\{e_{0}^{\prime}, \ldots, e_{N-1}^{\prime}\right\}$ like $c_{j}^{\prime} Z^{\prime-1}$. In the same way, $\rho\left(P_{j+1}\right)$ acts on $V_{j}$ like $c_{j} Y^{-1}$ and on $V_{j}^{\prime}$ like $c_{j}^{\prime} Y^{\prime-1}$.

Using the same action as in Example 2.2 and writing

$$
c_{j}=q^{p}, \quad c_{j}^{\prime}=q^{p^{\prime}},
$$

we get

$$
\rho\left(P_{j}\right) e_{k}=q^{2 k-1+p} e_{k+1}, \quad \rho\left(P_{j}\right) e_{l}^{\prime}=q^{1-2 l+p^{\prime}} e_{l+1}
$$

It follows that the action of $P_{j}$ on $W^{j}$ is given by

$$
P_{j} \epsilon_{k, l}=q^{2(k-l)+p+p^{\prime}} \epsilon_{k+1, l+1} \text { where } \epsilon_{k, l}:=e_{k} \otimes e_{l}^{\prime} \in W^{j}
$$

In the same way, one sees that the action of $P_{j+1}$ on $W^{j}$ is given by

$$
P_{j+1} \epsilon_{k, l}=q^{p+p^{\prime}} \epsilon_{k-1, l-1} .
$$

Now, for $m, n \in \mathbb{Z}_{N}$, set $\alpha_{m, n}:=\sum_{k=0}^{N-1} q^{2 k m} \epsilon_{k, k+n}$, an easy calculation shows that

$$
P_{j} \alpha_{m, n}=q^{-2(m+n)+p+p^{\prime}} \alpha_{m, n}, \quad P_{j+1} \alpha_{m, n}=q^{2 m+p+p^{\prime}} \alpha_{m, n}
$$

It follows that $\left\{\alpha_{n, m}: n, m \in \mathbb{Z}_{N}\right\}$ is a base of $W^{j}$ and, for all $p_{j}, p_{j+1} \in \mathcal{U}(N)$, there exists a unique couple $(m, n) \in \mathbb{Z}_{N}^{2}$ with $p_{j}=q^{-2(m+n)+p+p^{\prime}}$ and $p_{j+1}=q^{2 m+p+p^{\prime}}$.

Therefore, $\operatorname{dim} W_{\boldsymbol{p}}^{j}=1$ if and only if $p_{k}=1$ for all $k \neq j, j+1$.
(3) If $k \neq s-1$, then $v_{k}$ is neither a vertex of $T_{s-1}$ nor a vertex of $T_{s-1}^{\prime}$, so if $p_{k} \neq 1$, then $W_{\mathbf{h}}^{s}=\{0\}$.

Suppose that $p_{k}=1$ for all $k \in\{1, \ldots, s-2\}$, then

$$
W_{p}^{s-1} \supset \bigoplus_{p_{a} p_{b}=p_{s-1}} V_{p_{a}}^{s-1}\left(P_{s-1}\right) \otimes V_{p_{b}}^{s-1}\left(P_{s-1}\right)
$$

where $V_{p_{a}}^{s}\left(P_{s-1}\right)$ is the eigenspace associated to the eigenvalue $p_{a}$ of the action of $\rho\left(P_{s-1}\right)$ on the vector space associated to the triangle $T_{s-1}$, and $V_{p_{b}}^{\prime-1}\left(P_{s-1}\right)$ is defined in an analogous way.

The direct sum contains $N$ terms of dimension one, hence $\operatorname{dim}\left(W_{p}^{s-1}\right) \geq N$. On the other hand, we have

$$
\operatorname{dim}\left(W^{s-1}\right)=N^{2}
$$

and also

$$
\operatorname{dim}\left(W^{s-1}\right)=\sum_{\boldsymbol{p} \in \mathcal{U}(N)^{s-1}} \operatorname{dim}\left(W_{p}^{s-1}\right) \geq N \times N
$$

This implies that $W_{\boldsymbol{p}}^{s-1}$ has exactly dimension $N$ for $\boldsymbol{p}=\left(1, \ldots, 1, p_{s-1}\right)$.
The proof of Proposition 3.1 follows from the following elementary remark:

Remark 3.3. For all $j \in\{0, \ldots, s-1\}$, given $\boldsymbol{p}_{j} \in \mathcal{U}(N)^{s-1}$ and a nonzero vector $x_{j} \in W_{\boldsymbol{p}_{j}}^{j}$, the vector $x_{0} \otimes \cdots \otimes x_{s-1}$ is in $W_{\boldsymbol{p}}$ where $\boldsymbol{p}=\boldsymbol{p}_{0} \boldsymbol{p}_{1} \ldots \boldsymbol{p}_{s-1}$ is obtained by taking the product on each component.

We thus have the inclusion

$$
\begin{equation*}
W_{\boldsymbol{p}} \supset \bigoplus_{\boldsymbol{p}=\boldsymbol{p}_{0} \ldots \boldsymbol{p}_{s-1}} W_{\boldsymbol{p}_{0}}^{0} \otimes \cdots \otimes W_{\boldsymbol{p}_{s-1}}^{s-1} \tag{1}
\end{equation*}
$$

Writing $\boldsymbol{p}_{j}=\left(p_{0}^{(j)}, \ldots, p_{s-1}^{(j)}\right)$ and $\boldsymbol{p}=\left(p_{1}, \ldots, p_{s}\right)$, one notes that from Lemma 3.2, the only nonzero terms in the direct sum of (1) are the $\boldsymbol{p}_{j}$ satisfying

$$
\begin{align*}
& p_{1}^{(0)} p_{1}^{(1)}=p_{1} \\
& p_{2}^{(1)} p_{2}^{(2)}=p_{2} \tag{2}
\end{align*}
$$

$$
p_{s-1}^{(s-2)} p_{s-1}^{(s-1)}=p_{s-1}
$$

There exists exactly $N^{s-1}$ different choices for the $\boldsymbol{p}_{j}$ satisfying relations (2).
Moreover, for each choice of $\boldsymbol{p}_{j}$ satisfying (2), the vector space $W_{\boldsymbol{p}_{0}}^{0} \otimes \cdots \otimes W_{\boldsymbol{p}_{s-1}}^{s-1}$ has dimension $N$. It follows that for each $\boldsymbol{p} \in \mathcal{U}(N)^{s-1}$,

$$
\operatorname{dim} W_{p} \geq N^{s}
$$

On the other hand,

$$
\operatorname{dim} W=N^{2 s-1}=\sum_{\boldsymbol{p} \in \mathcal{U}(N)^{s-1}} \operatorname{dim} W_{\boldsymbol{p}}
$$

hence each $W_{\boldsymbol{p}}$ has exactly dimension $N^{s}$.
Finally, as the puncture invariants act as the identity on the vector space $W^{\prime}$, the intersection of the eigenspaces of the $\rho\left(P_{j}\right)$ for all $j=1, \ldots, s-1$ has the form $W_{\boldsymbol{p}} \otimes W^{\prime}$ for some $\boldsymbol{p} \in \mathcal{U}(N)^{s-1}$ and has dimension $N^{4 g-3+s}$. As the representation $\rho$ has fixed central charge $c$,

$$
\rho\left(\left[P_{1} P_{2} \ldots P_{s}\right]\right)=\rho\left(\left[H^{2}\right]\right)=c^{2} \operatorname{Id}_{V}
$$

It follows that the action of $\rho\left(P_{s}\right)$ on $V$ can be easily expressed as a function of the action of the $\rho\left(P_{j}\right)$ for $j=1, \ldots, s-1$, and we get the result.

Proposition 3.1 implies the decomposition of the Main Theorem for the triangulation $\lambda_{0}$. Since the dimension of the eigenspaces depends continuously on the $x_{i}$, we get the decomposition for all value of $x_{i} \in \mathbb{C}^{*}$.

Indeed, consider the local representation

$$
\rho: \mathcal{T}_{q}\left(\lambda_{0}\right) \rightarrow \operatorname{End}(V)
$$

classified by a nonzero complex number $x_{i}$ associated to each edge and central charge $c$. Let $\rho^{(i)}: \mathcal{T}_{q}\left(\lambda_{0}\right) \rightarrow \operatorname{End}\left(V^{(i)}\right)$ be an irreducible factor.

In particular,

$$
\begin{aligned}
\rho^{(i)}\left(X_{i}^{N}\right) & =\rho\left(X_{i}^{N}\right)_{\mid V^{(i)}}=x_{i} \operatorname{Id}_{V^{(i)}}, \\
\rho^{(i)}(H) & =\rho(H)_{\mid V^{(i)}}
\end{aligned}=c \operatorname{Id}_{V^{(i)}}, ~ \$
$$

so a necessary condition for $\rho^{(i)}$ to be an irreducible factor is that it is classified by the same $x_{i}$ and same central charge $c$.

For each puncture $v_{j}$, denote by $p_{j}^{(i)}$ the $N$-th root of $x_{1}^{k_{j_{1}}} \ldots x_{n}^{k_{j_{n}}}$ so that

$$
\rho^{(i)}\left(P_{j}\right)=p_{j}^{(i)} \operatorname{Id}_{V^{(i)}} .
$$

It follows that $p_{s}^{(i)}=c^{2}\left(p_{1}^{(i)} \ldots p_{s-1}^{(i)}\right)^{-1}$ and $V^{(i)} \subset V_{\boldsymbol{p}}$ where $\boldsymbol{p}=\left(p_{1}^{(i)}, \ldots, p_{s-1}^{(i)}\right)$ with

$$
V_{p}=\left\{x \in V: \rho\left(P_{j}\right) x=p_{j}^{(i)} x, j=1, \ldots, s-1\right\}
$$

Finally, as $\operatorname{dim} V_{\boldsymbol{p}}=N^{4 g-3+s}$ and the dimension of an irreducible representation of $\mathcal{T}_{q}\left(\lambda_{0}\right)$ has dimension $N^{3 g-3+s}, V_{\boldsymbol{p}}$ contains exactly $N^{g}$ irreducible factors classified by the same $x_{i}$, same central charge $c$ and $N$-the root $p_{j}^{(i)}$ associated to the puncture $v_{j}$.

Proof in the general case. Recall that, given another ideal triangulation $\lambda \in \Lambda(\Sigma)$, the "transition maps" $\varphi_{\lambda_{0} \lambda}: \mathbb{C}^{n} \rightarrow \mathbb{C}^{n}$ defined in Section 2.3 are rational, hence defined on a Zariski open set of $\mathbb{C}^{n}$.

Now, consider a local representation

$$
\rho: \mathcal{T}_{q}(\lambda) \rightarrow \operatorname{End}\left(V_{\lambda}\right)
$$

classified by a number $x_{i} \in \mathbb{C}^{*}$ associated to each edge and central charge $c$.
If there exists $\left(y_{1}, \ldots, y_{n}\right) \in \mathbb{C}^{n}$ so that $\varphi_{\lambda_{0} \lambda}\left(y_{1}, \ldots, y_{n}\right)=\left(x_{1}, \ldots, x_{n}\right)$ (which is a generic condition), then it follows from Section 2.3 that $\rho_{\lambda}$ is isomorphic to $\rho_{\lambda_{0}}: \mathcal{T}_{q}\left(\lambda_{0}\right) \rightarrow \operatorname{End}\left(V_{\lambda_{0}}\right)$. It means that there exists an isomorphism

$$
L_{\lambda_{0} \lambda}: V_{\lambda} \rightarrow V_{\lambda_{0}}
$$

so that for each $X \in \mathcal{T}_{q}(\lambda)$ we have

$$
\rho_{\lambda_{0}}\left(\Psi_{\lambda_{0} \lambda}^{q}(X)\right)=L_{\lambda_{0} \lambda} \circ \rho_{\lambda}(X) \circ L_{\lambda_{0} \lambda}^{-1} .
$$

Here $\Psi_{\lambda_{0} \lambda}^{q}: \hat{\mathcal{T}}_{q}(\lambda) \rightarrow \hat{\mathcal{T}}_{q}\left(\lambda_{0}\right)$ are the coordinates change defined in Section 2.3.
As $\rho_{\lambda_{0}}$ is a local representation of $\mathcal{T}_{q}\left(\lambda_{0}\right)$, there exists an irreducible decomposition of $\rho_{\lambda_{0}}$ given by the decomposition

$$
V_{\lambda_{0}}=\bigoplus_{i \in \mathcal{I}} V_{\lambda_{0}}^{i}
$$

In particular, each $i \in \mathcal{I}$, $V_{\lambda_{0}}^{i}$ is stable by $\rho_{\lambda_{0}}$ and has dimension $N^{3 g-3+s}$.
For each $i \in \mathcal{I}$, set $V_{\lambda}^{i}:=L_{\lambda_{0} \lambda}^{(-1)}\left(V_{\lambda_{0}}^{i}\right)$. One easily gets that each $V_{\lambda}^{i}$ is stable by $\rho_{\lambda}\left(\mathcal{T}_{q}(\lambda)\right)$, and has dimension $N^{3 g-3+s}$. In this way we get a decomposition of $\rho_{\lambda}$ into irreducible factors given by the decomposition

$$
V_{\lambda}=\bigoplus_{i \in \mathcal{I}} V_{\lambda}^{i}
$$

Finally, if $\rho_{\lambda}$ is classified by the parameters $\left(x_{1}, \ldots, x_{n}\right)$ which are not in the image of $\varphi_{\lambda_{0} \lambda}$, one can deform continuously $\left(x_{1}, \ldots, x_{n}\right)$ to get the previous decomposition and, as the decomposition does not depend of the parameters, get the result for $\rho_{\lambda}$.

## 4. Representations of the skein algebra

In this section, we use the Main Theorem to construct a nice family of representation of the Kauffman bracket skein algebra $\mathcal{S}_{A}(\bar{\Sigma})$ of the closed surface $\bar{\Sigma}=\Sigma \cup$ $\left\{v_{1}, \ldots, v_{s}\right\}$. This is done by adapting the construction of Bonahon and Wong [2015] to the case of local representations.

In Section 4.1, we describe the balanced Chekhov-Fock algebra $\mathcal{Z}_{\omega}(\lambda)$ associated to an ideal triagulation $\lambda$ of $\Sigma$ and characterize its irreducible representations. Then, in Section 4.2, we introduce the local representations of $\mathcal{Z}_{\omega}(\lambda)$ and extend the Main Theorem to decompose these local representations into irreducible factors. Finally, in Section 4.3, we use the previous decomposition to construct a family of representations of $\mathcal{S}_{A}(\bar{\Sigma})$.
4.1. Balanced Chekhov-Fock algebra. Let $q$ be a primitive $N$-th root of unity with $N$ odd and let $\omega$ be the unique fourth root of $q$ which is also a primitive $N$-th root of unity. Namely, if $q=e^{2 i \pi \frac{k}{N}}$ with $N$ and $k$ coprime, then there is a unique $l \in \mathbb{Z}_{4}$ so that $k+l N \in 4 \mathbb{Z}$, and $\omega=e^{2 i \pi \frac{k}{4 N}} e^{i \frac{l \pi}{2}}$.

Let $\lambda$ be an ideal triangulation of $\Sigma$ whose edges are $\left(\lambda_{1}, \ldots, \lambda_{n}\right)$. In order to avoid confusion, we will denote by $X_{i}$ the generators of $\mathcal{T}_{q}(\lambda)$ and by $Z_{i}$ the generators of $\mathcal{T}_{\omega}(\lambda)$.

A multi-index $\mathbf{k} \in \mathbb{Z}^{n}$ is called $\lambda$-balanced (or balanced) if for each triangle of the triangulation whose edges are $j_{1}, j_{2}, j_{3}$ we have

$$
k_{j_{1}}+k_{j_{2}}+k_{j_{3}} \in 2 \mathbb{Z}
$$

A monomial $Z \in \mathcal{T}_{\omega}(\lambda)$ is balanced if $Z$ is a scalar multiple of $Z_{\mathbf{k}}$ where $\mathbf{k} \in \mathbb{Z}^{n}$ is balanced. (Here $Z_{\mathbf{k}}$ is defined as in Section 2.1).
Definition 4.1. The balanced Chekhov-Fock algebra $\mathcal{Z}_{\omega}(\lambda)$ is the subalgebra of $\mathcal{T}_{\omega}(\lambda)$ generated (as a vector space) by balanced monomials.


Figure 6. Train track.
In particular, the image of the map

$$
\begin{aligned}
i: \mathcal{T}_{q}(\lambda) & \rightarrow \mathcal{T}_{\omega}(\lambda) \\
X_{i} & \mapsto Z_{i}^{2}
\end{aligned}
$$

lies in $\mathcal{Z}_{\omega}(\lambda)$ so we will consider $\mathcal{T}_{q}(\lambda)$ as a subalgebra of $\mathcal{Z}_{\omega}(\lambda)$.
The ideal triangulation $\lambda$ defines canonically a train track $\tau_{\lambda}$ on $\Sigma$ which looks like in Figure 6 on each triangle of the triangulation. Note that $\tau_{\lambda}$ has a switch on each edge of $\lambda$.

We denote by $\mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$ the abelian group of integer weight systems on $\tau_{\lambda}$. Namely, an element $\alpha \in \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$ is a map that associates to each edge $e$ of $\tau_{\lambda}$ an integer $\alpha(e)$ in such a way that at each switch, the sum of the weights of the incoming edges equals the sum of the weights of the outgoing edges.

A weight system $\alpha \in \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$ on $\tau_{\lambda}$ is a map that associates an integer to any edge of the train track in such a way that the sum of weights of the incoming edges equals the sum of weights of the outgoing edges. Given $\alpha \in \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$ and an edge $\lambda_{i} \in \lambda$, the sum of the weights of the edges incoming to $\lambda_{i}$ is an integer. It thus define a map

$$
\varphi: \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{Z}^{n}
$$

whose image is exactly the set of balanced multi-index. Thus, given an integer weight system $\alpha \in \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$, we define $Z_{\alpha} \in \mathcal{Z}_{\omega}(\lambda)$ to be $Z_{\varphi(\alpha)}=\left[Z_{1}^{\alpha_{1}} \ldots Z_{n}^{\alpha_{n}}\right]$ where $\varphi(\alpha)=\left(\alpha_{1}, \ldots, \alpha_{n}\right) \in \mathbb{Z}^{n}$. In particular, one gets the noncommutativity relations

$$
Z_{\alpha} Z_{\beta}=\omega^{4 \Omega(\alpha, \beta)} Z_{\beta} Z_{\alpha}
$$

where $\Omega: \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \times \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{Z}$ is the Thurston intersection form; see [Bonahon and Wong 2012, Section 2] for more details.
Definition 4.2. A twisted homomorphism is a map $\zeta: \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{C}^{*}$ such that for every $\alpha, \beta \in \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$,

$$
\zeta(\alpha+\beta)=(-1)^{\Omega(\alpha, \beta)} \zeta(\alpha) \zeta(\beta)
$$

Finally, note that each puncture $v_{j}$ defines a integer weight system $\eta_{j} \in \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$ as follow. The connected component $D_{j}$ of $\Sigma \backslash \tau_{\lambda}$ containing $v_{j}$ is bounded by a finite number of edges of $\tau_{\lambda}$. For an edge $e$ of $\tau_{\lambda}$, define $\eta_{j}(e) \in\{0,1,2\}$ to be the number of times $e$ lies in the boundary of $D_{j}$. In particular,

$$
Z_{\eta_{j}}^{2}=i\left(P_{j}\right)
$$

where $P_{j} \in \mathcal{T}_{q}(\lambda)$ is the puncture invariant associated to $v_{j}$ and $i: \mathcal{T}_{q}(\lambda) \rightarrow \mathcal{Z}_{\omega}(\lambda)$ is defined above.

Irreducible representations of $\mathcal{Z}_{\omega}(\lambda)$ were classified in [Bonahon and Wong 2015]. They proved:

Proposition 4.3 [Bonahon and Wong 2012, Proposition 14]. Up to isomorphism, an irreducible representation $\rho: \mathcal{Z}_{\omega}(\lambda) \rightarrow \operatorname{End}(V)$ has dimension $N^{3 g-3+s}$ and is classified by a twisted homomorphism $\zeta: \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{C}^{*}$ and a choice of an $N$-th root $h_{j}=\zeta\left(\eta_{j}\right)^{1 / N}$ for each puncture $v_{j}$. Such a representation satisfies:

- $\rho\left(Z_{\alpha}^{N}\right)=\zeta(\alpha) \operatorname{Id}_{V}$ for all $\alpha \in \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$.
- $\rho\left(\eta_{j}\right)=\zeta\left(\eta_{j}\right) \operatorname{Id}_{V}$ for all $j \in\{1, \ldots, s\}$.
4.2. Local representations of $\mathcal{Z}_{\omega}(\lambda)$. Here, we introduce the notion of local representation of the balanced Chekhov-Fock algebra $\mathcal{Z}_{\omega}(\lambda)$ associated to an ideal triangulation $\lambda$. We then extend the Main Theorem to give a decomposition of local representations of $\mathcal{Z}_{\omega}(\lambda)$ into its irreducible components.

Since by our choice $\omega$ is also a primitive $N$-th root of unity, there is a map

$$
\mathcal{T}_{\omega}(\lambda) \rightarrow \bigotimes_{T_{i} \in F(\lambda)} \mathcal{T}_{\omega}\left(T_{i}\right)
$$

as introduced in Section 2.2, where $F(\lambda)$ is the set of faces of $\lambda$ and $\mathcal{T}_{\omega}\left(T_{i}\right)$ is the triangle algebra associated to the face $T_{i}$. It is clear that this map restricts to a morphism

$$
\iota: \mathcal{Z}_{\omega}(\lambda) \rightarrow \bigotimes_{T_{i} \in F(\lambda)} \mathcal{Z}_{\omega}\left(T_{i}\right)
$$

Here, $\mathcal{Z}_{\omega}\left(T_{i}\right)$ is the balanced triangle algebra associated to the face $T_{i}$.
Definition 4.4. A local representation of the balanced Chekhov-Fock algebra $\mathcal{Z}_{\omega}(\lambda)$ is a representation $\rho: \mathcal{Z}_{\omega}(\lambda) \rightarrow \operatorname{End}(V)$ that can be written as $\left(\rho_{1} \otimes \cdots \otimes \rho_{m}\right) \circ \iota$ where each $\rho_{i}$ is an irreducible representation of $\mathcal{Z}_{\omega}\left(T_{i}\right)$.

In order to classify local representations of $\mathcal{Z}_{\omega}(\lambda)$, we first have to understand the irreducible representations of the balanced triangle algebra $\mathcal{Z}_{\omega}(T)$. Let $\tau$ be the train track in $T$ with edges $e_{1}, e_{2}, e_{3}$ as in Figure 6 and denote by $\mathcal{W}(\tau, \mathbb{Z})$ The group of integer weight systems on $\tau$.

Lemma 4.5. Up to isomorphism, an irreducible representation of the balanced triangle algebra $\mathcal{Z}_{\omega}(T)$ has dimension $N$ and is classified by a twisted homomorphism $\zeta: \mathcal{W}(\tau, \mathbb{Z}) \rightarrow \mathbb{C}^{*}$ together with a choice of an $N$-th root $C=(\zeta(\mu))^{1 / N}$ where $\mu \in \mathcal{W}(\tau, \mathbb{Z})$ is such that $\mu\left(e_{i}\right)=1$ for all $i \in \mathbb{Z}_{3}$. Such a representation satisfies

- $\rho\left(Z_{\alpha}^{N}\right)=\zeta(\alpha) \operatorname{Id}_{V}$.
- $\rho\left(Z_{\mu}\right)=C \operatorname{Id}_{V}$.

Proof. The group $\mathcal{W}(\tau, \mathbb{Z})$ is generated by $\left\{\alpha_{1}, \alpha_{2}, \alpha_{3}\right\}$ where

$$
\alpha_{i}\left(e_{j}\right)=\delta_{i j}, \quad i, j \in \mathbb{Z}_{3}
$$

It follows that the balanced triangle algebra $\mathcal{Z}_{\omega}(T)$ is generated by $Z_{\alpha_{1}}^{ \pm 1}, Z_{\alpha_{2}}^{ \pm 1}$ and $Z_{\alpha_{3}}^{ \pm 1}$ and the relations are

$$
Z_{\alpha_{1}} Z_{\alpha_{i+1}}=\omega^{-2} Z_{\alpha_{i+1}} Z_{\alpha_{i}}, \quad i \in \mathbb{Z}_{3}
$$

If we denote by $Z_{i}$ the generator of $\mathcal{T}_{\omega}(T)$ associated to the edge $\lambda_{i}$ (so for instance $Z_{\alpha_{1}}=\left[Z_{2} Z_{3}\right]$ ), the map

$$
\Psi: \mathcal{Z}_{\omega}(T) \rightarrow \mathcal{T}_{\omega}(T), \quad Z_{\alpha_{i}} \mapsto Z_{i}^{-1}
$$

is an isomorphism of algebras such that $\Psi\left(Z_{\mu}\right)=H^{-1}$ where $H=\left[Z_{1} Z_{2} Z_{3}\right]$.
In particular, an irreducible representation $\rho$ of $\mathcal{Z}_{\omega}(\lambda)$ has the form $\rho=\bar{\rho} \circ \Psi$ where $\bar{\rho}$ is an irreducible representation of $\mathcal{T}_{\omega}(T)$.

Using the result of Section 2.2 and the fact that a twisted homomorphism of $\mathcal{W}(\tau, \mathbb{Z})$ is fully determined by its value on the $\alpha_{i}$, we get the result.

Let $\tau_{i}$ be the restriction of the train track $\tau_{\lambda}$ to the triangle $T_{i}$ of the triangulation $\lambda$. A twisted homomorphism $\zeta: \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{C}^{*}$ induces a twisted homomorphism $\zeta_{i}$ : $\mathcal{W}\left(\tau_{i}, \mathbb{Z}\right) \rightarrow \mathbb{C}^{*}$ for each triangle $T_{i}$ of the triangulation $\lambda$. In particular, the following proposition is a straightforward extension of [Bai et al. 2007, Proposition 6]:
Proposition 4.6. A local representation $\rho: \mathcal{Z}_{\omega}(\lambda) \rightarrow \operatorname{End}(V)$ has dimension $N^{4 g-4+2 s}$ and is classified (up to isomorphism) by a twisted homomorphism $\zeta$ : $\mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{C}^{*}$ and a choice of an $N$-th root $C=\zeta(\mu)^{1 / N}$ where $\mu(e)=1$ for all edge e of $\tau_{\lambda}$. Such a representation satisfies:

- $\rho\left(Z_{\alpha}^{N}\right)=\zeta(\alpha) \operatorname{Id}_{V}$.
- $\rho\left(Z_{\mu}\right)=C \mathrm{Id}_{V}$.

Finally, the Main Theorem implies the following:
Theorem 4.7. Let $\rho: \mathcal{Z}_{\omega}(\lambda) \rightarrow \operatorname{End}(V)$ be the (isomorphism class of) representation classified by the twisted homomorphism $\zeta: \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{C}^{*}$ and the choice of an $N$-th root $C=\zeta(\mu)^{1 / N}$ (where $\mu$ is defined as above). Then $\rho=\bigoplus_{i \in \mathcal{I}} \rho^{(i)}$ where each $\rho^{(i)}$ is irreducible, classified by the same twisted homomorphism $\zeta$ and
$N$-th root $h_{k}^{(i)}=\left(\zeta\left(\eta_{k}\right)\right)^{1 / N}$ with $h_{1}^{(i)} \ldots h_{s}^{(i)}=C$ (here, the $\eta_{k}$ are defined as in Section 4.1).

Moreover, for each choice of an $N$-th root $h_{k}=\left(\zeta\left(\eta_{k}\right)\right)^{1 / N}$ for each $k \in$ $\{1, \ldots, s\}$, there are exactly $N^{g}$ indices $i \in \mathcal{I}$ such that $h_{k}^{(i)}=h_{k}$ for all $k$.
Proof. Let $\rho$ be a local representation of $\mathcal{Z}_{\omega}(\lambda)$ classified by $\zeta$ and $C$. In particular, $\rho$ induces a local representation $\bar{\rho}:=\rho \circ i$ of $\mathcal{T}_{q}(\lambda)$, where $i: \mathcal{T}_{q}(\lambda) \hookrightarrow \mathcal{Z}_{\omega}(\lambda)$. The local representation $\bar{\rho}$ is classified by the weight $\zeta\left(\beta_{i}\right)$ for all edges $\lambda_{i}$, where $\beta_{i} \in \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right)$ is defined by $Z_{\beta_{i}}=Z_{i}^{2}=i\left(X_{i}\right)$.

Let $P_{j} \in \mathcal{T}_{q}(\lambda)$ be the puncture invariant associated to the puncture $v_{j}$. The image of $P_{j}$ in $\mathcal{Z}_{\omega}(\lambda)$ is $Z_{\eta_{j}}^{2}$. We claim that the eigenspaces of $\bar{\rho}\left(P_{j}\right)$ correspond to the eigenspaces of $\rho\left(Z_{\eta_{j}}\right)$. In fact, if $V_{h_{j}}\left(Z_{\eta_{j}}\right)$ is the eigenspace of $\rho\left(Z_{\eta_{j}}\right)$ corresponding to the eigenvalue $h_{j}=\left(\zeta\left(\eta_{j}\right)\right)^{1 / N}$, then one has the inclusion

$$
V_{h_{j}}\left(Z_{\eta_{j}}\right) \subset V_{p_{j}}\left(P_{j}\right)
$$

where $V_{p_{j}}\left(P_{j}\right)$ is the eigenspace of $\bar{\rho}\left(P_{j}\right)=\rho\left(Z_{\eta_{j}}^{2}\right)$ corresponding to the eigenvalue $p_{j}=h_{j}^{2}$. Because there are only $N$ different possible eigenvalues of $\rho\left(Z_{\eta_{j}}\right)$, a dimension counting argument shows the equality.

Now, we apply the Main Theorem and get that, for each choice of $\left(h_{1}, \ldots, h_{s}\right)$ where $h_{j}=\left(\zeta\left(Z_{\eta_{j}}\right)\right)^{1 / N}$, the intersection $V_{h_{1}}\left(Z_{\eta_{1}}\right) \cap \cdots \cap V_{h_{s}}\left(Z_{\eta_{s}}\right)$ has dimension $N^{4 g-3+s}$ and hence is made of $N^{g}$ copies of the irreducible representation of $\mathcal{Z}_{\omega}(\lambda)$ classified by $\zeta$ and $h_{1}, \ldots, h_{s}$.

Bonahon and Wong [2012, Section 3] associate a character $r_{\zeta} \in \chi(\Sigma, \operatorname{SL}(2, \mathbb{C}))$ to a twisted homomorphism $\zeta: \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{C}^{*}($ here $\chi(\Sigma, \operatorname{SL}(2, \mathbb{C}))$ is the algebraic quotient of $\operatorname{Hom}\left(\pi_{1}(\Sigma), \mathrm{SL}_{2}(\mathbb{C})\right)$ by the action of $\mathrm{SL}_{2}(\mathbb{C})$ by conjugation). In particular, the (irreducible or local) representations of the balanced Chekhov-Fock algebra associated to an ideal triangulation $\lambda$ of $\Sigma$ are classified, up to finitely many choice, by a Zariski open set in $\chi(\Sigma, \operatorname{SL}(2, \mathbb{C}))$.

Note that, if $r_{\zeta}$ is the character associated to the twisted homomorphism $\zeta$, the holonomy of $r_{\zeta}$ around a puncture $v_{j}$ is parabolic exactly when $\zeta\left(\eta_{j}\right)=1$.
4.3. Representations of $\mathcal{S}_{A}(\overline{\boldsymbol{\Sigma}})$. We explain here how Theorem 4.7 gives rise to a new family of representations of the Kauffman bracket skein algebra of the closed surface $\bar{\Sigma}=\Sigma \cup\left\{v_{1}, \ldots, v_{s}\right\}$.

Skein algebra. Given an oriented 3-manifold $M$, and a nonzero complex number $A \in \mathbb{C}^{*}$, consider the complex vector space $V(M)$ freely generated by isotopy classes of framed links in $M$. The skein module $\mathcal{S}_{A}(M)$ of $M$ is the quotient of $V(M)$ by the Kauffman bracket skein relations as defined in Figure 7.

Namely, we identify three different links when differ by the previous relation in an open ball and agree everywhere else.


Figure 7. Kauffman bracket skein relations.
Given a framed link $K \subset M$, we denote by [ $K$ ] its image in the skein module $\mathcal{S}_{A}(M)$.

When $M=\Sigma \times[0,1]$ for a surface $\Sigma$, the skein module $\mathcal{S}_{A}(M)=\mathcal{S}_{A}(\Sigma)$ inherits an algebra structure given by superposition of links. Namely, given two framed links $K_{1}$ and $K_{2}$ in $\Sigma \times[0,1]$, the product [ $\left.K_{1}\right] \cdot\left[K_{2}\right]$ is defined to be the image of $K_{1} \cup K_{2}$ in $\mathcal{S}_{A}(\Sigma)$, where $K_{1} \cup K_{2}$ is given by the superposition of $K_{1}$ on top of $K_{2}$ where we rescaled so that $K_{1} \subset \Sigma \times\left[0, \frac{1}{2}\right]$ and $K_{2} \subset \Sigma \times\left[\frac{1}{2}, 1\right]$. We call $\mathcal{S}_{A}(\Sigma)$ with the product - the Kauffman bracket skein algebra of $S$.

Finite-dimensional representations of the skein algebra $\mathcal{S}_{A}(\Sigma)$ are of main importance as they appear naturally in topological quantum field theory (TQFT). For example, the Witten-Reshetekin-Turaev TQFT [Blanchet et al. 1995; Turaev 1994].

Classical shadow and quantum trace. Let $\mu: \mathcal{S}_{A}(\Sigma) \rightarrow \operatorname{End}(V)$ be an irreducible representation of the Kauffman bracket skein algebra of $\Sigma$.

Bonahon and Wong [2016] (see also [Lê 2015a] for a simpler proof) proved that if $A$ is a primitive $N$-th root of -1 , the $N$-th Chebyshev polynomial $T_{N}$ of the first kind of any skein $[K] \in \mathcal{S}_{A}(\Sigma)$ is a central element in $\mathcal{S}_{A}(\Sigma)$. In particular, the precomposition of $\rho$ by $T_{N}$ maps each skein [ $K$ ] to a multiple of the identity in $\operatorname{End}(V)$. This multiple of the identity can be interpreted as an element $r_{\mu} \in \chi\left(\Sigma, \mathrm{SL}_{2}(\mathbb{C})\right)$ in the $\operatorname{SL}(2, \mathbb{C})$ character variety of $\Sigma$. This character is called the classical shadow of the representation $\mu$.

When $A=\omega^{-2}$ (so $A$ is a primitive $N$-th root of -1 ) and $\lambda$ is an ideal triangulation of $\Sigma$, Bonahon and Wong [2011] (see also [Lê 2015b] for a more conceptual proof) constructed a quantum trace map

$$
\operatorname{tr}_{\omega}^{\lambda}: \mathcal{S}_{A}(\Sigma) \rightarrow \mathcal{Z}_{\omega}(\lambda)
$$

which turns out to be an injective algebra homomorphism.
In particular, by precomposing irreducible representations of $\mathcal{Z}_{\omega}(\lambda)$ by the quantum trace, Bonahon and Wong [2012] obtained a family of irreducible representations of the Kauffman bracket skein algebra of $S$ indexed by a Zariski open subset
of the character variety $\chi\left(\Sigma, \mathrm{SL}_{2}(\mathbb{C})\right)$. Moreover, taking the classical shadow of such an irreducible representation recovers the character.

Representations of $\mathcal{S}_{A}(\bar{\Sigma})$. The inclusion $\Sigma \hookrightarrow \bar{\Sigma}$ gives an algebra homomorphism

$$
\iota: \mathcal{S}_{A}(\bar{\Sigma}) \rightarrow \mathcal{S}_{A}(\Sigma)
$$

Let $r \in \chi(\Sigma, \operatorname{SL}(2, \mathbb{C}))$ be a character obtained from a character $r^{\prime} \in \chi(\bar{\Sigma}, \operatorname{SL}(2, \mathbb{C}))$ (namely, the holonomy of $r$ around each puncture is trivial). If $\zeta: \mathcal{W}\left(\tau_{\lambda}, \mathbb{Z}\right) \rightarrow \mathbb{C}^{*}$ is the twisted homomorphism associated to $r$, then $\zeta\left(\eta_{j}\right)=1$ for each puncture $v_{j}$.

Denote by

$$
\rho: \mathcal{Z}_{\omega}(\lambda) \rightarrow \operatorname{End}(V)
$$

the local representation of $\mathcal{Z}_{\omega}(\lambda)$ classified by $\zeta$ and the $N$-th root $C=\left(\left(-\omega^{4}\right)^{s}\right)^{1 / N}$. Let $E \subset V$ be the intersection of the eigenspaces of $\rho\left(Z_{\eta_{k}}\right)$ for $k \in\{1, \ldots, s\}$ corresponding to the eigenvalue $-\omega^{4}$.

By Theorem 4.7, the vector space $E$ is stable by $\rho\left(\mathcal{Z}_{\omega}(\lambda)\right)$, so we get an induced representation $\rho^{\prime}: \mathcal{Z}_{\omega}(\lambda) \rightarrow \operatorname{End}(E)$. Note that $\rho^{\prime}$ is made of $N^{g}$ copies of the irreducible representation of $\mathcal{Z}_{\omega}(\lambda)$ classified by $\zeta$ and puncture invariant $-\omega^{4}$.

Proposition 4.8. There is a proper linear subspace $F \subset E$ such that the composition

$$
\mu: \mathcal{S}_{A}(\bar{\Sigma}) \xrightarrow{\iota} \mathcal{S}_{A}(\Sigma) \xrightarrow{\bar{\rho}} \operatorname{End}(F)
$$

induces a representation of $\mathcal{S}_{A}(\bar{\Sigma})$. The classical shadow of each irreducible factor of $\mu$ is same. Finally, the dimension of $F$ is at least $N^{4 g-3}$ when $g>1$ and at least $N^{2}$ when $g=1$.

Proof. This is a direct consequence of the construction of Bonahon and Wong [2015]. In fact, using the decomposition of $\bar{\rho}$ into irreducible parts and considering the total off-diagonal kernel of each irreducible factor (see [op. cit., Section 4.2]), one gets the result.

The vector space $F$ is canonically associated to the triangulation $\lambda$, which makes the family of representations described above easier to handle for computations.

## Acknowledgment

I would like to thank Francis Bonahon for valuable discussions on the subject. This research was partially supported by the grant DMS-1406559 from the U.S. National Science Foundation. The author gratefully acknowledges support from the NSF grants DMS-1107452, 1107263, and 1107367 "RNMS: GEometric structures And Representation varieties" (the GEAR Network).

## References

[Bai 2007] H. Bai, "Quantum Teichmüller spaces and Kashaev's $6 j$-symbols", Algebr. Geom. Topol. 7 (2007), 1541-1560. MR Zbl
[Bai et al. 2007] H. Bai, F. Bonahon, and X. Liu, "Local representations of the quantum Teichmüller space", preprint, 2007. arXiv
[Baseilhac and Benedetti 2005] S. Baseilhac and R. Benedetti, "Classical and quantum dilogarithmic invariants of flat PSL(2, © )-bundles over 3-manifolds", Geom. Topol. 9 (2005), 493-569. MR Zbl
[Blanchet et al. 1995] C. Blanchet, N. Habegger, G. Masbaum, and P. Vogel, "Topological quantum field theories derived from the Kauffman bracket", Topology 34:4 (1995), 883-927. MR Zbl
[Bonahon and Liu 2007] F. Bonahon and X. Liu, "Representations of the quantum Teichmüller space and invariants of surface diffeomorphisms", Geom. Topol. 11 (2007), 889-937. MR Zbl
[Bonahon and Wong 2011] F. Bonahon and H. Wong, "Quantum traces for representations of surface groups in $\mathrm{SL}_{2}(\mathbb{C}) "$, Geom. Topol. 15:3 (2011), 1569-1615. MR Zbl
[Bonahon and Wong 2012] F. Bonahon and H. Wong, "Representations of the Kauffman bracket skein algebra, II: Punctured surfaces", preprint, 2012. arXiv
[Bonahon and Wong 2015] F. Bonahon and H. Wong, "Representations of the Kauffman bracket skein algebra, III: Closed surfaces and naturality", preprint, 2015. arXiv
[Bonahon and Wong 2016] F. Bonahon and H. Wong, "Representations of the Kauffman bracket skein algebra, I: Invariants and miraculous cancellations", Invent. Math. 204:1 (2016), 195-243. MR Zbl
[Fock and Chekhov 1999] V. V. Fock and L. O. Chekhov, "A quantum Teichmüller space", Theoret. and Math. Phys. 120:3 (1999), 1245-1259. MR Zbl
[Goodearl and Warfield 2004] K. R. Goodearl and R. B. Warfield, Jr., An introduction to noncommutative Noetherian rings, 2nd ed., London Mathematical Society Student Texts 61, Cambridge Univ. Press, 2004. MR Zbl
[Guo and Liu 2009] R. Guo and X. Liu, "Quantum Teichmüller space and Kashaev algebra", Algebr. Geom. Topol. 9:3 (2009), 1791-1824. MR Zbl
[Kashaev 1995] R. M. Kashaev, "A link invariant from quantum dilogarithm", Modern Phys. Lett. A 10:19 (1995), 1409-1418. MR Zbl
[Kashaev 1998] R. M. Kashaev, "Quantization of Teichmüller spaces and the quantum dilogarithm", Lett. Math. Phys. 43:2 (1998), 105-115. MR Zbl
[Lê 2015a] T. T. Q. Lê, "On Kauffman bracket skein modules at roots of unity", Algebr. Geom. Topol. 15:2 (2015), 1093-1117. MR Zbl
[Lê 2015b] T. T. Q. Lê, "Quantum Teichmüller spaces and quantum trace map", 2015. To appear in J. Inst. Math. Jussieu. arXiv
[Liu 2009] X. Liu, "The quantum Teichmüller space as a noncommutative algebraic object", J. Knot Theory Ramifications 18:5 (2009), 705-726. MR Zbl
[Murakami 2011] H. Murakami, "An introduction to the volume conjecture", pp. 1-40 in Interactions between hyperbolic geometry, quantum topology and number theory, edited by A. Champanerkar et al., Contemp. Math. 541, American Mathematical Society, Providence, RI, 2011. MR Zbl
[Penner 1993] R. C. Penner, "Universal constructions in Teichmüller theory", Adv. Math. 98:2 (1993), 143-215. MR Zbl
[Thurston 1986] W. P. Thurston, "Minimal stretch maps between hyperbolic surfaces", preprint, 1986. arXiv
[Turaev 1994] V. G. Turaev, Quantum invariants of knots and 3-manifolds, de Gruyter Studies in Mathematics 18, de Gruyter, Berlin, 1994. MR Zbl

Received October 12, 2016. Revised March 1, 2017.

JÉRÉMY TOULISSE
Department of Mathematics
University of Southern Califonia
Los Angeles, CA
United States
toulisse@usc.edu

# PACIFIC JOURNAL OF MATHEMATICS 

Founded in 1951 by E. F. Beckenbach (1906-1982) and F. Wolf (1904-1989)
msp.org/pjm
EDITORS
Don Blasius (Managing Editor)
Department of Mathematics
University of California
Los Angeles, CA 90095-1555
blasius@math.ucla.edu

Paul Balmer
Department of Mathematics
University of California
Los Angeles, CA 90095-1555
balmer@math.ucla.edu
Wee Teck Gan
Mathematics Department
National University of Singapore Singapore 119076 matgwt@nus.edu.sg

## Sorin Popa

Department of Mathematics
University of California
Los Angeles, CA 90095-1555
popa@math.ucla.edu

Vyjayanthi Chari
Department of Mathematics
University of California
Riverside, CA 92521-0135 chari@math.ucr.edu

Kefeng Liu
Department of Mathematics
University of California
Los Angeles, CA 90095-1555 liu@math.ucla.edu

## Jie Qing

Department of Mathematics
University of California
Santa Cruz, CA 95064
qing@cats.ucsc.edu

Daryl Cooper

Department of Mathematics
University of California
Santa Barbara, CA 93106-3080
cooper@math.ucsb.edu
Jiang-Hua Lu
Department of Mathematics
The University of Hong Kong
Pokfulam Rd., Hong Kong
jhlu@maths.hku.hk

## Paul Yang

Department of Mathematics
Princeton University
Princeton NJ 08544-1000
yang@math.princeton.edu

## PRODUCTION

Silvio Levy, Scientific Editor, production@msp.org

## SUPPORTING INSTITUTIONS

ACADEMIA SINICA, TAIPEI
CALIFORNIA INST. OF TECHNOLOGY
INST. DE MATEMÁTICA PURA E APLICADA KEIO UNIVERSITY
MATH. SCIENCES RESEARCH INSTITUTE NEW MEXICO STATE UNIV.
OREGON STATE UNIV.

STANFORD UNIVERSITY
UNIV. OF BRITISH COLUMBIA
UNIV. OF CALIFORNIA, BERKELEY
UNIV. OF CALIFORNIA, DAVIS
UNIV. OF CALIFORNIA, LOS ANGELES
UNIV. OF CALIFORNIA, RIVERSIDE
UNIV. OF CALIFORNIA, SAN DIEGO
UNIV. OF CALIF., SANTA BARBARA

UNIV. OF CALIF., SANTA CRUZ
UNIV. OF MONTANA
UNIV. OF OREGON
UNIV. OF SOUTHERN CALIFORNIA UNIV. OF UTAH
UNIV. OF WASHINGTON
WASHINGTON STATE UNIVERSITY

These supporting institutions contribute to the cost of publication of this Journal, but they are not owners or publishers and have no responsibility for its contents or policies.

See inside back cover or msp.org/pjm for submission instructions.
The subscription price for 2018 is US $\$ 475 /$ year for the electronic version, and $\$ 640 /$ year for print and electronic. Subscriptions, requests for back issues and changes of subscriber address should be sent to Pacific Journal of Mathematics, P.O. Box 4163, Berkeley, CA 94704-0163, U.S.A. The Pacific Journal of Mathematics is indexed by Mathematical Reviews, Zentralblatt MATH, PASCAL CNRS Index, Referativnyi Zhurnal, Current Mathematical Publications and Web of Knowledge (Science Citation Index).

The Pacific Journal of Mathematics (ISSN 0030-8730) at the University of California, c/o Department of Mathematics, 798 Evans Hall \#3840, Berkeley, CA 94720-3840, is published twelve times a year. Periodical rate postage paid at Berkeley, CA 94704, and additional mailing offices. POSTMASTER: send address changes to Pacific Journal of Mathematics, P.O. Box 4163, Berkeley, CA 94704-0163.

PJM peer review and production are managed by EditFLOw ${ }^{\circledR}$ from Mathematical Sciences Publishers.
E. mathematical sciences publishers
nonprofit scientific publishing
http://msp.org/
© 2018 Mathematical Sciences Publishers

## PACIFIC JOURNAL OF MATHEMATICS

Volume 294 No. 1 May 2018
Three-dimensional Sol manifolds and complex Kleinian groups ..... 1
Waldemar Barrera, Rene Garcia-Lara and Juan Navarrete
On periodic points of symplectomorphisms on surfaces ..... 19
Marta Batoréo
Mixing properties for hom-shifts and the distance between walks on associated ..... 41 graphsNishant Chandgotia and Brian Marcus
Simultaneous construction of hyperbolic isometries ..... 71
Matt Clay and Caglar Uyanik
A local weighted Axler-Zheng theorem in $\mathbb{C}^{n}$ ..... 89
ŽELJKo ČUČKović, SÖnMEZ ŞAhutoğLu and Yunus E. Zeytuncu
Monotonicity and radial symmetry results for Schrödinger systems with ..... 107 fractional diffusionJing Li
Moduli spaces of stable pairs ..... 123
Yinbang Lin
Spark deficient Gabor frames ..... 159
Romanos-Diogenes Malikiosis
Ordered groups as a tensor category ..... 181
Dale Rolfsen
Multiplication of distributions and a nonlinear model in elastodynamics ..... 195C. O. R. Sarrico
Some Ambrose- and Galloway-type theorems via Bakry-Émery and modified ..... 213
Ricci curvaturesHomare Tadano
Irreducible decomposition for local representations of quantum Teichmüller ..... 233spaceJÉRÉMY TOULISSE


[^0]:    MSC2010: 20G42, 57M50, 57R56.
    Keywords: quantum Teichmüller theory, skein algebras, quantum topology.

