

Tunisian Journal of Mathematics

an international publication organized by the Tunisian Mathematical Society

**On the ultimate energy bound of solutions
to some forced second-order evolution equations
with a general nonlinear damping operator**

Alain Haraux

2019

vol. 1

no. 1



On the ultimate energy bound of solutions to some forced second-order evolution equations with a general nonlinear damping operator

Alain Haraux

Under suitable growth and coercivity conditions on the nonlinear damping operator g which ensure nonresonance, we estimate the ultimate bound of the energy of the general solution to the equation $\ddot{u}(t) + Au(t) + g(\dot{u}(t)) = h(t)$, $t \in \mathbb{R}^+$, where A is a positive selfadjoint operator on a Hilbert space H and h is a bounded forcing term with values in H . In general the bound is of the form $C(1 + \|h\|^4)$, where $\|h\|$ stands for the L^∞ norm of h with values in H and the growth of g does not seem to play any role. If g behaves like a power for large values of the velocity, the ultimate bound has quadratic growth with respect to $\|h\|$ and this result is optimal. If h is antiperiodic, we obtain a much lower growth bound and again the result is shown to be optimal even for scalar ODEs.

1. Introduction

We investigate a specific quantitative aspect of solutions to the equation

$$\ddot{u}(t) + Au(t) + g(\dot{u}(t)) = h(t),$$

where V is a real Hilbert space, $A \in L(V, V')$ is a symmetric, positive, coercive operator, $g \in C(V, V')$ is monotone and h is a forcing term. This equation has been intensively studied in the literature when g is a local damping term, covering the following topics: existence of almost periodic solutions, asymptotic behavior of the general solution, rate of decay to 0 of the difference of two solutions in the energy space in the best cases; see, e.g., [Amerio and Prouse 1969; Prouse 1965a; 1965b; 1965c; 1965d; Biroli 1973; Biroli and Haraux 1980; Haraux 1981; 1982; 1985; 1987; 1991; Haraux and Zuazua 1988; Zuazua 1988]. In the more recent paper [Aloui et al. 2013], a result generalizing the theorems of [Haraux 1987] on boundedness and compactness has been proved for possibly nonlocal damping terms. However when looking at the arguments of those two papers and trying to extract an estimate of the solutions for t large, we find immediately that

MSC2010: 34A34, 34D20, 35B40, 35L10, 35L90.

Keywords: second-order equation, nonlinear damping, energy bound, antiperiodic.

the methodology cannot be adapted to that purpose. The present article aims at improving the situation. Actually we devise a new technique which allows us to “forget” the influence of the initial data from the very beginning of the estimates, thus dropping all unnecessary terms related to transient behavior. The plan of the paper is as follows: In Section 2, we introduce the basic tools used in the statements and proofs of the main results. Section 3 is devoted to a very general case. Section 4 covers a still rather general case where the damping operator behaves like a power for large values of the velocity, this for instance allows us to encompass any polynomial map, and we give a short list of examples in the field of PDEs of the second order in t for which our result is optimal. In Section 5 we establish two partial results when the forcing is antiperiodic, a situation which is known (see, e.g., [Haraux 1989]) to prevent resonance under weaker conditions on g than the general periodic case. We obtain a better estimate which is optimal in finite dimensions, but in the infinite-dimensional setting we can only slightly improve the general estimate and we do not reach what one might expect to be the optimal result.

2. Functional framework and the initial value problem

We now recall the exact functional framework that shall be used in the formulation as well as in the proofs of our new results. We follow the presentation from [Aloui et al. 2013] at the exception of a small difference for the approximation of weak solutions.

2A. Monotone operators. Let \mathcal{H} be a real Hilbert space endowed with an inner product $(\cdot, \cdot)_{\mathcal{H}}$. We recall that a map \mathcal{A} defined on a subset $\mathcal{D} = D(\mathcal{A})$ with values in \mathcal{H} is monotone if

$$\forall (U, \widehat{U}) \in \mathcal{D} \times \mathcal{D}, \quad (\mathcal{A}U - \mathcal{A}\widehat{U}, U - \widehat{U})_{\mathcal{H}} \geq 0.$$

In addition \mathcal{A} is called maximal monotone if

$$\forall F \in \mathcal{H}, \exists U \in D(\mathcal{A}), \quad \mathcal{A}U + U = F.$$

The following result is well known; see [Brezis 1973].

Proposition 2.1. *If \mathcal{A} is maximal monotone, for each $T > 0$, each $U_0 \in D(\mathcal{A})$ and $F = F(t) \in W^{1,1}(0, T; \mathcal{H})$ there is a unique function $U \in W^{1,1}(0, T; \mathcal{H})$ with $U(t) \in D(\mathcal{A})$ for almost all $t \in (0, T)$, $U(0) = U_0$ and such that for almost all $t \in (0, T)$*

$$U'(t) + \mathcal{A}U(t) = F(t). \tag{2-1}$$

In addition if for some $\widehat{U}_0 \in D(\mathcal{A})$ and $\widehat{F} \in W^{1,1}(0, T; \mathcal{H})$ we consider the solution $\widehat{U} \in W^{1,1}(0, T; \mathcal{H})$ with $\widehat{U}(t) \in D(\mathcal{A})$ for almost all $t \in (0, T)$, $\widehat{U}(0) = \widehat{U}_0$ of

$$\widehat{U}'(t) + \mathcal{A}\widehat{U}(t) = \widehat{F}(t),$$

then the difference satisfies the inequality

$$\forall t \in [0, T], \quad |U(t) - \widehat{U}(t)| \leq |U_0 - \widehat{U}_0| + \int_0^t |F(s) - \widehat{F}(s)| ds.$$

This proposition allows one to define by density, for any $U_0 \in \overline{D(\mathcal{A})}$ and $F = F(t) \in L^1(0, T; \mathcal{H})$, a weak solution of (2-1) such that $U(0) = U_0$ [Brezis 1973].

2B. Functional setting. Throughout this article we let H and V be two Hilbert spaces with norms respectively denoted by $\|\cdot\|$ and $|\cdot|$. We assume that V is densely and continuously embedded into H . Identifying H with its dual H' , we obtain $V \hookrightarrow H = H' \hookrightarrow V'$. We denote inner products by (\cdot, \cdot) and duality products by $\langle \cdot, \cdot \rangle$; the spaces in question will be specified by subscripts. The notation $\langle f, u \rangle$ without any subscript will be used sometimes to denote $\langle f, u \rangle_{V', V}$. The duality map: $V \rightarrow V'$ will be denoted by A . We observe that A is characterized by the property

$$\forall (u, v) \in V \times V, \quad \langle Au, v \rangle_{V', V} = (u, v)_V.$$

2C. Weak solutions. We consider the dissipative evolution equation

$$\ddot{u} + Au + g(\dot{u}) = h(t), \tag{2-2}$$

where $g \in C(V, V')$ satisfies

$$\forall (v, w) \in V \times V, \quad \langle g(v) - g(w), v - w \rangle \geq 0. \tag{2-3}$$

We consider the (generally unbounded) operator \mathcal{A} defined on the Hilbert space $\mathcal{H} = V \times H$ by

$$D(\mathcal{A}) = \{(u, v) \in V \times V : Au + g(v) \in H\}$$

and

$$\forall (u, v) \in D(\mathcal{A}), \quad \mathcal{A}(u, v) = (-v, Au + g(v)).$$

Lemma 2.2. *The operator \mathcal{A} is maximal monotone.*

Proof. Let $U = (u, v)$ and $\widehat{U} = (\hat{u}, \hat{v})$ be two elements of $D(\mathcal{A})$. We have

$$\begin{aligned} (\mathcal{A}U - \mathcal{A}\widehat{U}, U - \widehat{U})_{\mathcal{H}} &= -(u - \hat{u}, v - \hat{v})_V + (Au + g(v) - A\hat{u} - g(\hat{v}), v - \hat{v})_H \\ &\quad - (u - \hat{u}, v - \hat{v})_V + \langle Au + g(v) - A\hat{u} - g(\hat{v}), v - \hat{v} \rangle_{V', V} \end{aligned}$$

since $Au + g(v) \in H$ and $A\hat{u} + g(\hat{v}) \in H$ while v, \hat{v} are in V . This reduces to

$$(\mathcal{A}U - \mathcal{A}\widehat{U}, U - \widehat{U})_{\mathcal{H}} = \langle g(v) - g(\hat{v}), v - \hat{v} \rangle_{V', V} \geq 0.$$

Hence \mathcal{A} is monotone. To prove that \mathcal{A} is maximal monotone we are left to show that for any $F = (\varphi, \psi) \in \mathcal{H}$ the system

$$u - v = \varphi, \quad Au + g(v) + v = \psi$$

has a solution $U = (u, v) \in D(\mathcal{A})$. This is equivalent to finding a solution $v \in V$ of

$$Av + g(v) + v = \psi - A\varphi \in V'.$$

But now the operator $\mathcal{C} \in C(V, V')$, defined by

$$\forall v \in V, \quad \mathcal{C}v = Av + g(v) + v,$$

is continuous and coercive since $V \rightarrow V'$ is the sum of a monotone operator and the coercive duality map. Therefore by Corollary 14, p. 126 from [Brezis 1968], \mathcal{C} is surjective. Finally \mathcal{A} is maximal monotone as claimed. \square

As a consequence of Proposition 2.1, for any $h \in L^1_{\text{loc}}(\mathbb{R}^+, H)$ and for each $(u_0, u_1) \in V \times H$ there is a unique weak solution

$$u \in C(\mathbb{R}^+, V) \cap C^1(\mathbb{R}^+, H)$$

of (2-2) such that $u(0) = u_0$ and $\dot{u}(0) = u_1$. This solution can be recovered on each compact interval $[0, T]$ by approximating the initial data by elements of the domain, approximating the forcing term h by C^1 functions and passing to the limit; the limit is independent of the approximating elements so chosen. The next result shows that in fact the approximation can even be made uniform on \mathbb{R}^+ .

2D. Density of strong solutions.

Lemma 2.3. *For any $h \in L^2_{\text{loc}}(\mathbb{R}^+, H)$, for each $(u_0, u_1) \in V \times H$ and for each $\delta > 0$ there exists $(w_0, w_1) \in D(\mathcal{A})$ and $k \in C^1(\mathbb{R}^+, H)$ for which the solution $w \in W^{1,1}_{\text{loc}}(\mathbb{R}^+, V) \cap W^{2,1}_{\text{loc}}(\mathbb{R}^+, H)$ of*

$$\ddot{w} + Aw + g(\dot{w}) = k(t), \quad w(0) = w_0, \quad \dot{w}(0) = w_1,$$

satisfies

$$\forall t \geq 0, \quad \|u(t) - w(t)\| + |\dot{u}(t) - \dot{w}(t)| \leq \delta,$$

and in addition

$$\forall t \in \mathbb{R}^+, \quad \int_t^{t+1} |k(s) - h(s)|^2 ds \leq 2\delta.$$

Proof. It suffices to use the last result of Proposition 2.1 by observing that for any $h \in L^2_{\text{loc}}(\mathbb{R}^+, H)$ we can find $k \in C^1(\mathbb{R}^+, H)$ such that

$$\forall n \in \mathbb{N}, \quad \int_n^{n+1} |k(s) - h(s)|^2 ds \leq \delta 2^{-2n-2}.$$

Choosing $(w_0, w_1) \in D(\mathcal{A})$ such that

$$\|w_0 - u_0\| + |w_1 - v_1| \leq \delta 2^{-1}$$

the result follows immediately \square

3. A general ultimate bound

We now give a quite different proof, in a slightly more general case, of a result stated in [Haraux 1985, Remark 1.2(b), p. 167]. We assume that $h \in S^2(\mathbb{R}^+, H)$ with

$$S^2(\mathbb{R}^+, H) = \left\{ f \in L_{\text{loc}}^2(\mathbb{R}^+, H) : \sup_{t \in \mathbb{R}^+} \int_t^{t+1} |f(s)|^2 ds < \infty \right\}$$

and we set

$$\|h\|_{S^2(\mathbb{R}^+, H)} = \left(\sup_{t \in \mathbb{R}^+} \int_t^{t+1} |h(s)|^2 ds \right)^{1/2}.$$

In particular if $h \in L^\infty(\mathbb{R}^+, H)$, then $h \in S^2(\mathbb{R}^+, H)$ and $\|h\|_{S^2(\mathbb{R}^+, H)} \leq \|h\|_{L^\infty(\mathbb{R}^+, H)}$.

Theorem 3.1. *Assume that $g \in C(V, V')$ satisfies the condition (2-3) and*

$$\exists \gamma > 0, \exists C_1 \geq 0, \forall v \in V, \quad \langle g(v), v \rangle \geq \gamma |v|^2 - C_1, \quad (3-1)$$

$$\exists K > 0, \exists C_2 \geq 0, \forall v \in V, \quad \|g(v)\|_{V'} \leq C_2 + K \langle g(v), v \rangle. \quad (3-2)$$

Then any solution $u \in C(\mathbb{R}^+, V) \cap C^1(\mathbb{R}^+, H)$ of (2-2) is bounded on \mathbb{R}^+ in the sense that u has bounded range in V and \dot{u} has bounded range in H . In addition we have for some constant K depending only on A and g

$$\limsup_{t \rightarrow \infty} (|\dot{u}(t)|^2 + \|u(t)\|^2) \leq K(1 + \|h\|_{S^2(\mathbb{R}^+, H)}^4).$$

Proof. The boundedness result is known for local damping operators g , see the second case of Theorem IV.2.1.1 of [Haraux 1987], and in the general case it can be proved by adapting in this case the method from [Aloui et al. 2013]. However even in the local case these results cannot provide a reasonable estimate of the ultimate bound. We start by an estimate in the case of a strong solutions; i.e., we assume

$$u \in W_{\text{loc}}^{1,1}(\mathbb{R}^+, V) \cap W_{\text{loc}}^{2,1}(\mathbb{R}^+, H).$$

The general case will follow by density. Let

$$E(t) = \frac{1}{2}(|\dot{u}|^2 + \|u\|^2).$$

Under the regularity conditions $[u_0, v_0] \in V \times V$, $g(v_0) \in H$ and $h \in W_{\text{loc}}^{1,1}(\mathbb{R}^+, H)$, the function: $t \rightarrow E(t)$ is absolutely continuous and we have, for all $t \in \mathbb{R}^+$

$$\frac{d}{dt}E(t) = (h, \dot{u}) - \langle g(\dot{u}), \dot{u} \rangle. \quad (3-3)$$

In addition $t \rightarrow (u(t), \dot{u}(t))$ is absolutely continuous and

$$\frac{d}{dt}(u(t), \dot{u}(t)) = |\dot{u}|^2 - \|u\|^2 - \langle g(\dot{u}), u \rangle + (h, u).$$

By using (3-2), we obtain

$$\frac{d}{dt}(u(t), \dot{u}(t)) \leq |\dot{u}|^2 - \|u\|^2 + \|u\|(P|h| + C_2 + K\langle g(\dot{u}), \dot{u} \rangle) \quad (3-4)$$

with

$$P = \sup\{|u| : u \in V, \|u\| = 1\}.$$

Introducing

$$\Phi(t) = 2E(t) \quad \forall t \geq 0,$$

we are reduced to estimating the upper limit bound for $\Phi(t)$. Let us introduce

$$M = \limsup_{t \rightarrow \infty} \Phi(t)$$

and let us consider a sequence of times t_n tending to infinity for which

$$\Phi(t_n) \geq M - \frac{1}{n}.$$

In addition, for n large enough we have $t_n \geq \tau$ and

$$\Phi(t_n - \tau) \leq M + \frac{1}{n},$$

where τ is any fixed positive number to be chosen later. Therefore by integrating (3-3) on $[t_n - \tau, t_n]$ we find

$$\int_{t_n - \tau}^{t_n} \langle g(\dot{u}), \dot{u} \rangle dt \leq \frac{1}{n} + \int_{t_n - \tau}^{t_n} \langle h, \dot{u} \rangle dt \leq \frac{1}{n} + \frac{\gamma}{2} \int_{t_n - \tau}^{t_n} |\dot{u}|^2 dt + \frac{1}{2\gamma} \int_{t_n - \tau}^{t_n} |h|^2 dt.$$

As a consequence of (3-1) we deduce

$$\int_{t_n - \tau}^{t_n} \langle g(\dot{u}), \dot{u} \rangle dt \leq \frac{2}{n} + \frac{1}{\gamma} \int_{t_n - \tau}^{t_n} |h|^2 dt + C_3, \quad (3-5)$$

$$\int_{t_n - \tau}^{t_n} |\dot{u}|^2 dt \leq \frac{C_4}{n} + C_5 \int_{t_n - \tau}^{t_n} |h|^2 dt + C_6, \quad (3-6)$$

which provide an average bound of the kinetic part independent of the initial data and the transient behavior. This is remarkable since we only used the properties

$\Phi(t_n) \geq M - 1/n$ and $\Phi(t_n - \tau) \leq M + 1/n$ to express the fact that t is large. By combining these two estimates we also find an estimate of the form

$$|\Phi(t) - \Phi(s)| \leq C_7 \left(1 + \int_{t_n - \tau}^{t_n} |h|^2 dt \right), \quad (3-7)$$

which is valid for all s, t in $[t_n - \tau, t_n]$. As a consequence if we had an L^1 estimate of the total energy instead of the kinetic part, the proof would be completed with exponent 2 instead of 4. The difficulty in fact comes from the potential energy. From (3-4), by integrating on $[t_n - \tau, t_n]$ we find

$$\begin{aligned} & \int_{t_n - \tau}^{t_n} \|u\|^2 dt \\ & \leq \int_{t_n - \tau}^{t_n} |\dot{u}|^2 dt + \int_{t_n - \tau}^{t_n} [\|u\| (P|h| + C_2 + K \langle g(\dot{u}), \dot{u} \rangle)] dt + |[u(t), \dot{u}(t)]_{t_n - \tau}^{t_n}| \end{aligned}$$

Recalling the notation $M = \limsup_{t \rightarrow \infty} \Phi(t)$ we find

$$\begin{aligned} & \int_{t_n - \tau}^{t_n} \|u\|^2 dt \\ & \leq \int_{t_n - \tau}^{t_n} |\dot{u}|^2 dt + M^{1/2} \int_{t_n - \tau}^{t_n} (P|h| + C_2 + K \langle g(\dot{u}), \dot{u} \rangle) dt + C_8 M, \end{aligned} \quad (3-8)$$

where C_8 does not depend on τ . Combining (3-6) and (3-8) we obtain

$$\begin{aligned} & \int_{t_n - \tau}^{t_n} \Phi(t) dt \\ & \leq C_5 \int_{t_n - \tau}^{t_n} |h|^2 dt + C_9 + M^{1/2} \int_{t_n - \tau}^{t_n} (P|h| + C_2 + K \langle g(\dot{u}), \dot{u} \rangle) dt + C_8 M \end{aligned} \quad (3-9)$$

and by (3-5) this implies

$$\int_{t_n - \tau}^{t_n} \Phi(t) dt \leq C_{10} \left(1 + \int_{t_n - \tau}^{t_n} |h|^2 dt \right) (1 + M^{1/2}) + C_8 M. \quad (3-10)$$

Finally by combining this last inequality with (3-7) we end up with

$$(\tau - C_8)M \leq C_{11} \left(1 + \int_{t_n - \tau}^{t_n} |h|^2 dt \right) (1 + M^{1/2}). \quad (3-11)$$

Fixing $\tau \geq 1 + C_8$, the result now follows easily since

$$\int_{t_n - \tau}^{t_n} |h|^2 dt \leq (1 + \tau) \|h\|_{S^2}^2.$$

The general case of weak solutions follows easily from density, relying on Lemma 2.3. \square

Remark 3.2. This ultimate bound has been obtained under the most general known assumption ensuring boundedness of trajectories. It seems not to depend on the kind of damping operator as long as the coerciveness and growth conditions are satisfied. We have absolutely no idea whether it has a chance to be optimal in some cases. A more natural quadratic estimate is valid in many cases, as we shall see in the next section.

4. The case of a power-like damping term

For the main result of this section, we need to introduce an additional Banach space Z such that

$$V \subset Z \subset H$$

with continuous embeddings. The norm in Z of a vector $z \in Z$ will be denoted by $\|z\|_Z$.

4A. Main result.

Theorem 4.1. *Assume that $g \in C(V, V')$ satisfies the condition (2-3) and for some $\alpha \geq 0$ we have*

$$\exists \gamma > 0, \exists C_1 \geq 0, \forall v \in V, \quad \langle g(v), v \rangle \geq \gamma \|v\|_Z^{\alpha+2} - C_1, \quad (4-1)$$

$$\exists K > 0, \exists C_2 \geq 0, \forall v \in V, \quad \|g(v)\|_{V'} \leq C_2 + K \|v\|_Z^{\alpha+1}. \quad (4-2)$$

Then any solution $u \in C(\mathbb{R}^+, V) \cap C^1(\mathbb{R}^+, H)$ of (2-2) is bounded on \mathbb{R}^+ in the sense that u has bounded range in V and \dot{u} has bounded range in H . In addition we have for some constant K depending only on A and g

$$\limsup_{t \rightarrow \infty} (|\dot{u}(t)|^2 + \|u(t)\|^2) \leq K(1 + \|h\|_{S^2(\mathbb{R}^+, H)}^2).$$

Proof. We start as in the proof of Theorem 3.1; by integrating (3-3) on $[t_n - \tau, t_n]$ we find

$$\begin{aligned} \int_{t_n - \tau}^{t_n} \langle g(\dot{u}), \dot{u} \rangle dt &\leq \frac{1}{n} + \int_{t_n - \tau}^{t_n} \langle h, \dot{u} \rangle dt \\ &\leq \frac{1}{n} + \frac{\gamma}{2} \int_{t_n - \tau}^{t_n} \|\dot{u}\|_Z^{\alpha+2} dt + C(\gamma) \int_{t_n - \tau}^{t_n} |h|^{(\alpha+2)/(\alpha+1)} dt. \end{aligned}$$

As a consequence of (4-1) we deduce, since $n \geq 1$,

$$\int_{t_n - \tau}^{t_n} \langle g(\dot{u}), \dot{u} \rangle dt \leq C(\gamma, \tau) \left(\int_{t_n - \tau}^{t_n} |h|^2 dt \right)^{(\alpha+2)/(2\alpha+2)} + C_3, \quad (4-3)$$

$$\int_{t_n - \tau}^{t_n} |\dot{u}|^2 dt \leq C_4 + C_5(\gamma, \tau) \int_{t_n - \tau}^{t_n} |h|^2 dt. \quad (4-4)$$

From (4-3) we also deduce, since $\frac{1}{2} \leq (\alpha + 2)/(2\alpha + 2)$ the important new estimate

$$\int_{t_n-\tau}^{t_n} \|\dot{u}\|_Z^{\alpha+1} dt \leq C_6 + C_7(\gamma, \tau) \left(\int_{t_n-\tau}^{t_n} |h|^2 dt \right)^{1/2} \quad (4-5)$$

and by (4-2) this implies

$$\int_{t_n-\tau}^{t_n} \|g(\dot{u})\|_{V'} dt \leq C_8 + C_9(\gamma, \tau) \left(\int_{t_n-\tau}^{t_n} |h|^2 dt \right)^{1/2}. \quad (4-6)$$

Recalling the notation $M = \limsup_{t \rightarrow \infty} \Phi(t)$ we now find

$$\int_{t_n-\tau}^{t_n} \|u\|^2 dt \leq \int_{t_n-\tau}^{t_n} |\dot{u}|^2 dt + C_9(\gamma, \tau) M^{1/2} \left(\int_{t_n-\tau}^{t_n} |h|^2 dt \right)^{1/2} + C_{10}M + C_{11}, \quad (4-7)$$

where C_{10} does not depend on τ . Then by using Cauchy–Schwarz

$$\int_{t_n-\tau}^{t_n} \|u\|^2 dt \leq C_{12}(\gamma, \tau) \int_{t_n-\tau}^{t_n} |h|^2 dt + (C_{10} + 1)M + C_{11}. \quad (4-8)$$

By choosing τ large enough we obtain, as a consequence of (4-8) and (4-4), the inequality

$$\int_{t_n-\tau}^{t_n} \Phi(t) dt \leq C_{12}(\gamma) \int_{t_n-\tau}^{t_n} |h|^2 dt + C_{13}. \quad (4-9)$$

We conclude by using

$$|\Phi(t) - \Phi(s)| \leq C_{14}(\gamma) \left(1 + \int_{t_n-\tau}^{t_n} |h|^2 dt \right), \quad (4-10)$$

which is valid for all s, t in $[t_n - \tau, t_n]$ and follows easily from (4-3) and (4-4). \square

Remark 4.2. This result is optimal. For instance if we consider an eigenvector φ of A corresponding to the eigenvalue $\lambda > 0$, then for each $k > 0$, $k\varphi$ is a stationary solution of the equation with source term $h(t) \equiv k\lambda\varphi$ for any dissipative operator g . This shows that the ultimate bound of the energy is at least quadratic with respect to the size of the source term.

4B. Examples. In this section, Ω denotes a bounded open domain of \mathbb{R}^N with C^2 boundary and $\alpha \geq 0$, $c > 0$. We consider four simple special cases.

Example 1 (the wave equation with local damping).

$$\begin{cases} u_{tt} + c|u_t|^\alpha u_t - \Delta u = h(t, x) & \text{in } \mathbb{R}_+ \times \Omega, \\ u = 0 & \text{on } \mathbb{R}_+ \times \partial\Omega. \end{cases} \quad (4-11)$$

Here $V = H_0^1(\Omega)$, $H = L^2(\Omega)$ and $Z = L^{\alpha+2}(\Omega)$. We assume $(N - 2)\alpha \leq 2$.

Example 2 (the wave equation with nonlinear averaged damping).

$$\begin{cases} u_{tt} + c\left(\int_{\Omega} u_t^2(t, x) dx\right)^{\alpha/2} u_t - \Delta u = h(t, x) & \text{in } \mathbb{R}_+ \times \Omega, \\ u = 0 & \text{on } \mathbb{R}_+ \times \partial\Omega. \end{cases} \quad (4-12)$$

Here $V = H_0^1(\Omega)$ and $H = L^2(\Omega) = Z$.

Example 3 (a clamped plate equation with nonlinear structural averaged damping).

$$\begin{cases} u_{tt} - c\left(\int_{\Omega} |\nabla u_t|^2 dx\right)^{\alpha/2} \Delta u_t + \Delta^2 u = h(t, x) & \text{in } \mathbb{R}_+ \times \Omega, \\ u = |\nabla u| = 0 & \text{on } \mathbb{R}_+ \times \partial\Omega. \end{cases} \quad (4-13)$$

Here $V = H_0^2(\Omega)$, $H = L^2(\Omega)$ and $Z = H_0^1(\Omega)$.

Example 4 (a simply supported plate equation with nonlinear structural averaged damping).

$$\begin{cases} u_{tt} - c\left(\int_{\Omega} |\nabla u_t|^2 dx\right)^{\alpha/2} \Delta u_t + \Delta^2 u = h(t, x) & \text{in } \mathbb{R}_+ \times \Omega, \\ u = \Delta u = 0 & \text{on } \mathbb{R}_+ \times \partial\Omega. \end{cases} \quad (4-14)$$

Here $V = H^2 \cap H_0^1(\Omega)$, $H = L^2(\Omega)$ and $Z = H_0^1(\Omega)$.

As a consequence of Theorem 4.1 we obtain immediately:

Corollary 4.3. *In all four examples, let $h \in S^2(\mathbb{R}^+, H)$. Then any solution $u \in C(\mathbb{R}^+, V) \cap C^1(\mathbb{R}^+, H)$ of (2-2) is bounded on \mathbb{R}^+ in the sense that u has bounded range in V and \dot{u} has bounded range in H . In addition we have for some constant K independent of h and the initial data*

$$\limsup_{t \rightarrow \infty} (|\dot{u}(t)|^2 + \|u(t)\|^2) \leq K(1 + \|h\|_{S^2(\mathbb{R}^+, H)}^2).$$

Remark 4.4. In [Aloui et al. 2013], for the four previous examples, the authors proved the existence of a unique almost periodic solution when h is an S^2 -almost periodic source. In this case (in particular if h is periodic), the ultimate bound coincides with the supremum of the energy of the almost periodic solution. Actually, if we try to estimate directly the periodic solution, some boundary (in time) term disappears but the main part of the estimate is not much simpler. In addition we know that the estimate is essentially optimal, only the multiplicative constants might be worked out if one wants a more precise inequality.

5. Partial results in the antiperiodic case

If g is odd and h is τ -antiperiodic, i.e., if we have

$$h(t + \tau) = -h(t),$$

the interesting solutions are the antiperiodic ones; see, e.g., [Haraux 1989] for existence results. Since such solutions have mean value 0, the solution can be

estimated through its time-derivative, and because the estimate of the derivative is generally much better, we can expect an improvement on the energy bound.

This idea is perfectly valid if H is finite-dimensional, since then u and \dot{u} belong to the same space, but otherwise we have a problem to reach the norm of u in V . At the present time we do not know what happens if $\dim H = \infty$. For the time being we can only prove the following partial results.

Proposition 5.1. *Assume that $V = H$, $h \in C(\mathbb{R}, H)$ is τ -antiperiodic, that $g \in C(H, H)$ satisfies the condition (2-3) and for some $\alpha \geq 0$ we have*

$$\exists \gamma > 0, \exists C_1 \geq 0, \forall v \in H, \quad \langle g(v), v \rangle \geq \gamma |v|^{\alpha+2} - C_1, \quad (5-1)$$

$$\exists K > 0, \exists C_2 \geq 0, \forall v \in H, \quad |g(v)| \leq C_2 + K |v|^{\alpha+1}. \quad (5-2)$$

Then any τ -antiperiodic solution u of (2-2) is such that

$$\sup_{t \in \mathbb{R}} (|\dot{u}(t)|^2 + |u(t)|^2) \leq C(1 + \|h\|_{L^\infty(\mathbb{R}, H)}^{2/(\alpha+1)}),$$

where C is independent of h .

Proof. The starting point is the same as for the proof of Theorem 4.1. From

$$\int_0^{2\tau} \langle g(\dot{u}), \dot{u} \rangle dt \leq C(\gamma, \tau) \left(\int_0^{2\tau} |h|^2 dt \right)^{(\alpha+2)/(2\alpha+2)} + C_3 \quad (5-3)$$

we deduce, taking account of property (5-1), the more precise estimate

$$\int_0^{2\tau} |\dot{u}|^2 dt \leq C_4 + C_5(\gamma, \tau) \left(\int_0^{2\tau} |h|^2 dt \right)^{1/(\alpha+1)}, \quad (5-4)$$

which implies, since u has mean value 0,

$$\sup_{t \in [0, 2\tau]} |u(t)|^2 \leq C(1 + \|h\|_{L^\infty(\mathbb{R}, H)}^{2/(\alpha+1)}). \quad (5-5)$$

To obtain the uniform bound on \dot{u} , the trick now consists in evaluating the maximum of $\Phi(t) = |\dot{u}|^2 + |A^{1/2}u|^2$. At a maximum point θ the derivative vanishes, which gives

$$\langle g(\dot{u}), \dot{u} \rangle = (h, \dot{u});$$

hence

$$|\dot{u}(\theta)|^2 \leq C'(1 + |h|^{2/(\alpha+1)}).$$

This implies

$$\max_{t \in [0, 2\tau]} \Phi(t) = \Phi(\theta) \leq C''(1 + \|h\|_{L^\infty(\mathbb{R}, H)}^{2/(\alpha+1)})$$

and the conclusion follows immediately. \square

Remark 5.2. This result is optimal. For instance if we consider an eigenvector φ of A corresponding to the eigenvalue $\lambda > 0$, then for each $k > 0$, we have $u_k(t) = k \cos(\lambda^{1/2}t)\varphi$ is a solution of the equation

$$\ddot{u} + Au + g(\dot{u}) = g(-k\lambda^{1/2} \sin(\lambda^{1/2}t)\varphi) =: h(t)$$

and the L^∞ norm of the source term is less than a constant times $k^{\alpha+1}$ for k large. Both u and h are antiperiodic.

We have a weaker result (intermediate between Theorem 4.1 and Proposition 5.1) which is also valid in the infinite-dimensional setting and can be stated as follows:

Proposition 5.3. *Assume that the conditions of Theorem 4.1 are satisfied with (4-2) reinforced into*

$$\exists K > 0, \exists C_2 \geq 0, \forall v \in V, \quad \|g(v)\|_{Z'} \leq C_2 + K\|v\|_Z^{\alpha+1}. \quad (5-6)$$

Then any τ -antiperiodic solution $u \in C^1(\mathbb{R}, V) \cap C^2(\mathbb{R}, H)$ of (2-2) is such that

$$\sup_{t \in \mathbb{R}} (|\dot{u}(t)|^2 + \|u(t)\|^2) \leq C(1 + \|h\|_{L^2([0, \tau], H)}^{(\alpha+2)/(\alpha+1)}),$$

where C is independent of h .

Proof. The starting point is the same as for the proof of Theorem 4.1. From the inequality

$$\int_0^{2\tau} \langle g(\dot{u}), \dot{u} \rangle dt \leq C_3 \left(1 + \int_0^{2\tau} |h|^2 dt \right)^{(\alpha+2)/(2\alpha+2)} \quad (5-7)$$

we deduce the estimate

$$\int_0^{2\tau} |\dot{u}|^2 dt \leq C_4 \left(1 + \left(\int_0^{2\tau} |h|^2 dt \right)^{1/(\alpha+1)} \right), \quad (5-8)$$

but also

$$\int_0^{2\tau} \|\dot{u}\|_Z^{\alpha+1} dt \leq C_5 \left(1 + \left(\int_0^{2\tau} |h|^2 dt \right)^{1/2} \right) \quad (5-9)$$

and by (5-6) this implies

$$\int_0^{2\tau} \|g(\dot{u})\|_{Z'} dt \leq C_6 \left(1 + \left(\int_0^{2\tau} |h|^2 dt \right)^{1/2} \right). \quad (5-10)$$

From (5-9), since u has mean value 0, we deduce

$$\sup_{t \in [0, 2\tau]} \|u(t)\|_Z \leq C_7 (1 + \|h\|_{L^2([0, 2\tau], H)}^{1/(\alpha+1)}). \quad (5-11)$$

The two last inequalities imply immediately

$$\left| \int_0^{2\tau} \langle g(\dot{u}), u \rangle dt \right| \leq C_8 (1 + \|h\|_{L^2([0, 2\tau], H)}^{(\alpha+2)/(\alpha+1)}).$$

Now, multiplying the equation by u and integrating on the period we find easily after combining with (5-8)

$$\int_0^{2\tau} \Phi(t) dt \leq C_9(1 + \|h\|_{L^2([0,2\tau], H)}^{(\alpha+2)/(\alpha+1)}),$$

with $\Phi(t) = |\dot{u}|^2 + \|u\|^2$. Since

$$\Phi'(t) = (h, \dot{u}) - \langle g(\dot{u}), \dot{u} \rangle$$

by 2τ -periodicity and the inequality $\Phi' \leq |h||\dot{u}| + C_1$, we find as a consequence of (5-8)

$$\Phi(t) \leq \int_{t-\tau}^t \Phi(s) ds + C_{10}(1 + \|h\|_{L^2([0,2\tau], H)}^{(\alpha+2)/(\alpha+1)})$$

and the conclusion follows easily by using τ -antiperiodicity. \square

Remark 5.4. This result is certainly not optimal but it is all we can prove for the moment even in the most basic examples. Our result requires additional regularity on u ; this is usually achieved by assuming some regularity on h . When g is monotone, usually the antiperiodic solution is unique and depends continuously on h in L^2 , so that the estimate will be easy to transfer to the general case in the examples. This is important since we cannot derive strong estimates on solutions which are not antiperiodic and therefore approximation by strong solutions has to be performed within the antiperiodic class.

References

- [Aloui et al. 2013] F. Aloui, I. Ben Hassen, and A. Haraux, “Compactness of trajectories to some nonlinear second order evolution equations and applications”, *J. Math. Pures Appl.* (9) **100**:3 (2013), 295–326. MR Zbl
- [Amerio and Prouse 1969] L. Amerio and G. Prouse, “Uniqueness and almost-periodicity theorems for a non linear wave equation”, *Atti Accad. Naz. Lincei Rend. Cl. Sci. Fis. Mat. Natur.* (8) **46** (1969), 1–8. MR Zbl
- [Biroli 1973] M. Biroli, “Bounded or almost periodic solution of the non linear vibrating membrane equation”, *Ricerche Mat.* **22** (1973), 190–202. MR Zbl
- [Biroli and Haraux 1980] M. Biroli and A. Haraux, “Asymptotic behavior for an almost periodic, strongly dissipative wave equation”, *J. Differential Equations* **38**:3 (1980), 422–440. MR Zbl
- [Brezis 1968] H. Brezis, “Équations et inéquations non linéaires dans les espaces vectoriels en dualité”, *Ann. Inst. Fourier (Grenoble)* **18**:1 (1968), 115–175. MR Zbl
- [Brezis 1973] H. Brezis, *Opérateurs maximaux monotones et semi-groupes de contractions dans les espaces de Hilbert*, Notas de Matemática **50**, North-Holland, Amsterdam, 1973. MR Zbl
- [Haraux 1981] A. Haraux, *Nonlinear evolution equations: global behavior of solutions*, Lecture Notes in Mathematics **841**, Springer, 1981. MR Zbl
- [Haraux 1982] A. Haraux, “Dissipativity in the sense of Levinson for a class of second-order nonlinear evolution equations”, *Nonlinear Anal.* **6**:11 (1982), 1207–1220. MR Zbl

- [Haraux 1985] A. Haraux, “Two remarks on hyperbolic dissipative problems”, pp. 161–179 in *Non-linear partial differential equations and their applications* (Paris, 1983–1984), edited by H. Brezis and J.-L. Lions, Res. Notes in Math. **122**, Pitman, Boston, 1985. MR Zbl
- [Haraux 1987] A. Haraux, *Semi-linear hyperbolic problems in bounded domains*, Math. Rep. (Chur) **3**, part 1, Harwood, London, 1987. Zbl
- [Haraux 1989] A. Haraux, “Anti-periodic solutions of some nonlinear evolution equations”, *Manuscripta Math.* **63**:4 (1989), 479–505. MR Zbl
- [Haraux 1991] A. Haraux, *Systèmes dynamiques dissipatifs et applications*, Recherches en Mathématiques Appliquées **17**, Masson, Paris, 1991. MR Zbl
- [Haraux and Zuazua 1988] A. Haraux and E. Zuazua, “Decay estimates for some semilinear damped hyperbolic problems”, *Arch. Rational Mech. Anal.* **100**:2 (1988), 191–206. MR Zbl
- [Prouse 1965a] G. Prouse, “Soluzioni quasi-periodiche dell’equazione non omogenea delle onde, con termine dissipativo non lineare, I”, *Atti Accad. Naz. Lincei Rend. Cl. Sci. Fis. Mat. Natur.* (8) **38** (1965), 804–807. MR Zbl
- [Prouse 1965b] G. Prouse, “Soluzioni quasi-periodiche dell’equazione non omogenea delle onde, con termine dissipativo non lineare, II”, *Atti Accad. Naz. Lincei Rend. Cl. Sci. Fis. Mat. Natur.* (8) **39** (1965), 11–18. MR Zbl
- [Prouse 1965c] G. Prouse, “Soluzioni quasi-periodiche dell’equazione non omogenea delle onde, con termine dissipativo non lineare, III”, *Atti Accad. Naz. Lincei Rend. Cl. Sci. Fis. Mat. Natur.* (8) **39** (1965), 155–160. MR Zbl
- [Prouse 1965d] G. Prouse, “Soluzioni quasi-periodiche dell’equazione non omogenea delle onde, con termine dissipativo non lineare, IV”, *Atti Accad. Naz. Lincei Rend. Cl. Sci. Fis. Mat. Natur.* (8) **39** (1965), 240–244. MR Zbl
- [Zuazua 1988] E. Zuazua, “Stability and decay for a class of nonlinear hyperbolic problems”, *Asymptotic Anal.* **1**:2 (1988), 161–185. MR Zbl

Received 25 Aug 2017. Revised 30 Oct 2017.

ALAIN HARAUX:

haraux@ann.jussieu.fr

Sorbonne Universités, UPMC (Paris VI) & CNRS, Laboratoire Jacques-Louis Lions, Paris, France

Tunisian Journal of Mathematics

msp.org/tunis

EDITORS-IN-CHIEF

- Ahmed Abbes CNRS & IHES, France
abbes@ihes.fr
- Ali Baklouti Faculté des Sciences de Sfax, Tunisia
ali.baklouti@fss.usf.tn

EDITORIAL BOARD

- Hajer Bahouri CNRS & LAMA, Université Paris-Est Créteil, France
hajer.bahouri@u-pec.fr
- Arnaud Beauville Laboratoire J. A. Dieudonné, Université Côte d'Azur, France
beauville@unice.fr
- Bassam Fayad CNRS & Institut de Mathématiques de Jussieu-Paris Rive Gauche, Paris, France
bassam.fayad@imj-prg.fr
- Benoit Fresse Université Lille 1, France
benoit.fresse@math.univ-lille1.fr
- Dennis Gaitsgory Harvard University, United States
gaitsgde@gmail.com
- Emmanuel Hebey Université de Cergy-Pontoise, France
emmanuel.hebey@math.u-cergy.fr
- Mohamed Ali Jendoubi Université de Carthage, Tunisia
ma.jendoubi@gmail.com
- Sadok Kallel Université de Lille 1, France & American University of Sharjah, UAE
sadok.kallel@math.univ-lille1.fr
- Minhyong Kim Oxford University, UK & Korea Institute for Advanced Study, Seoul, Korea
minhyong.kim@maths.ox.ac.uk
- Toshiyuki Kobayashi The University of Tokyo & Kavli IPMU, Japan
toshi@kurims.kyoto-u.ac.jp
- Yanyan Li Rutgers University, United States
yyli@math.rutgers.edu
- Nader Masmoudi Courant Institute, New York University, United States
masmoudi@cims.nyu.edu
- Haynes R. Miller Massachusetts Institute of Technology, United States
hrm@math.mit.edu
- Nordine Mir Texas A&M University at Qatar & Université de Rouen Normandie, France
nordine.mir@qatar.tamu.edu
- Detlef Müller Christian-Albrechts-Universität zu Kiel, Germany
mueller@math.uni-kiel.de
- Mohamed Sifi Université Tunis El Manar, Tunisia
mohamed.sifi@fst.utm.tn
- Daniel Tataru University of California, Berkeley, United States
tataru@math.berkeley.edu
- Sundaram Thangavelu Indian Institute of Science, Bangalore, India
veluma@math.iisc.ernet.in
- Saïd Zarati Université Tunis El Manar, Tunisia
said.zarati@fst.utm.tn

PRODUCTION

- Silvio Levy (Scientific Editor)
production@msp.org

The Tunisian Journal of Mathematics is an international publication organized by the Tunisian Mathematical Society (<http://www.tms.rnu.tn>) and published in electronic and print formats by MSP in Berkeley.

See inside back cover or msp.org/tunis for submission instructions.

The subscription price for 2019 is US \$TBA/year for the electronic version, and \$TBA/year (+\$TBA, if shipping outside the US) for print and electronic. Subscriptions, requests for back issues and changes of subscriber address should be sent to MSP.

Tunisian Journal of Mathematics (ISSN 2576-7666 electronic, 2576-7658 printed) at Mathematical Sciences Publishers, 798 Evans Hall #3840, c/o University of California, Berkeley, CA 94720-3840 is published continuously online. Periodical rate postage paid at Berkeley, CA 94704, and additional mailing offices.

TJM peer review and production are managed by EditFlow® from MSP.

PUBLISHED BY
 **mathematical sciences publishers**
nonprofit scientific publishing
<http://msp.org/>
© 2019 Mathematical Sciences Publishers

Tunisian Journal of Mathematics

2019 vol. 1 no. 1

Welcome to the Tunisian Journal of Mathematics Ahmed Abbes and Ali Baklouti	1
Partial resolution by toroidal blow-ups János Kollár	3
Construction of a stable blowup solution with a prescribed behavior for a non-scaling-invariant semilinear heat equation Giao Ky Duong, Van Tien Nguyen and Hatem Zaag	13
Troisième groupe de cohomologie non ramifiée des hypersurfaces de Fano Jean-Louis Colliot-Thélène	47
On the ultimate energy bound of solutions to some forced second-order evolution equations with a general nonlinear damping operator Alain Haraux	59
On the irreducibility of some induced representations of real reductive Lie groups Wee Teck Gan and Atsushi Ichino	73
Truncated operads and simplicial spaces Michael S. Weiss	109
From compressible to incompressible inhomogeneous flows in the case of large data Raphaël Danchin and Piotr Bogusław Mucha	127