

ANALYSIS & PDE

Volume 13

No. 4

2020

SATOSHI MASAKI, JASON MURPHY AND JUN-ICHI SEGATA

**STABILITY OF SMALL SOLITARY WAVES
FOR THE ONE-DIMENSIONAL NLS
WITH AN ATTRACTIVE DELTA POTENTIAL**

STABILITY OF SMALL SOLITARY WAVES FOR THE ONE-DIMENSIONAL NLS WITH AN ATTRACTIVE DELTA POTENTIAL

SATOSHI MASAKI, JASON MURPHY AND JUN-ICHI SEGATA

We consider the initial-value problem for the one-dimensional nonlinear Schrödinger equation in the presence of an attractive delta potential. We show that for sufficiently small initial data, the corresponding global solution decomposes into a small solitary wave plus a radiation term that decays and scatters as $t \rightarrow \infty$. In particular, we establish the asymptotic stability of the family of small solitary waves.

1. Introduction

We study the one-dimensional nonlinear Schrödinger equation (NLS) with an attractive delta potential. This equation takes the form

$$\begin{cases} i \partial_t u = H u + \mu |u|^p u, \\ u(0) = u_0. \end{cases} \quad (1-1)$$

Here we take $u : \mathbb{R}_t \times \mathbb{R}_x \rightarrow \mathbb{C}$, $\mu \in \mathbb{R} \setminus \{0\}$, and H is the Schrödinger operator

$$H = -\frac{1}{2} \partial_x^2 + q \delta(x),$$

where $q < 0$ (the attractive case) and δ is the Dirac delta distribution. Equation (1-1) provides a simple model describing the resonant nonlinear propagation of light through optical wave guides with localized defects [Goodman et al. 2004]. For reasons to be detailed below, we consider the L^2 -supercritical case, namely, $p \geq 4$. For technical simplicity we also assume p is an even integer.

In the repulsive case ($q > 0$), (1-1) is studied from the point of view of scattering. Banica and Visciglia [2016] proved global well-posedness and scattering in the energy space for the defocusing mass-supercritical case. Ikeda and Inui [2017] considered the focusing mass-supercritical regime and proved scattering below the ground-state threshold. In our previous work [Masaki et al. 2019], we considered (1-1) with a cubic nonlinearity and proved decay and (modified) scattering for small initial data in a weighted space; see also [Segata 2015].

Such results are not expected in the attractive case. Indeed, in the attractive case the operator H has a single eigenvalue $-\frac{1}{2}q^2$, with a one-dimensional eigenspace spanned by the L^2 -normalized eigenfunction

$$\phi_0(x) := |q|^{1/2} e^{q|x|}.$$

MSC2010: 35Q55.

Keywords: NLS, solitary waves, asymptotic stability, delta potential.

One can then prove that there exists a family of small nonlinear bound states Q , parametrized by small $z \in \mathbb{C}$, which satisfy

$$HQ + \mu|Q|^p Q = EQ, \quad (1-2)$$

with $Q = Q[z] = z\phi_0 + \mathcal{O}(z^2)$ and $E = E[|z|] = -\frac{1}{2}q^2 + \mathcal{O}(z)$. The functions $u(t) = e^{-iEt}Q$ are then small solitary wave solutions to (1-1). In particular, one does not expect small solutions simply to decay and scatter in general. Instead, we will show that for small initial data, the corresponding solution decouples into a small solitary wave plus radiation. The existence and properties of $Q[z]$ are discussed in Section 2C. In fact, in the special case of the delta potential, one can find explicit formulas for the nonlinear ground states.

Our main result is the following theorem. We write P_c for the projection onto the continuous spectral subspace of H . The notation D_j denotes the derivative with respect to z_j , where we identify $z \in \mathbb{C}$ with the real vector (z_1, z_2) . Finally, $\langle \cdot, \cdot \rangle$ denotes the standard L^2 inner product.

Theorem 1.1. *Let $\|u_0\|_{H^1} = \delta$, $q < 0$, and let $p \geq 4$ be an even integer. For δ sufficiently small, there exists a unique global solution u to (1-1) and $z(t) \in \mathbb{C}$ such that writing*

$$u(t) = Q[z(t)] + v(t), \quad (1-3)$$

where $Q[z(t)]$ is the solution to (1-2), we have the following:

- v satisfies the orthogonality conditions

$$\operatorname{Im}\langle v(t), D_j Q[z(t)] \rangle \equiv 0 \quad \text{for } j \in \{1, 2\}. \quad (1-4)$$

- v obeys the global space-time bounds

$$\|v\|_{L_t^\infty H_x^1 \cap L_t^4 L_x^\infty} + \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} + \|\partial_x v\|_{L_x^\infty L_t^2} \lesssim \delta,$$

and there exists unique $v_+ \in P_c H^1$ such that

$$\lim_{t \rightarrow \infty} \|v(t) - e^{-itH} v_+\|_{H^1} = 0.$$

- $\|z\|_{L_t^\infty} \lesssim \delta$ and there exists $z_+ \in \mathbb{C}$ satisfying $||z_+| - |z(0)|| \lesssim \delta^2$ and

$$\lim_{t \rightarrow \infty} z(t) \exp\left\{i \int_0^t E[z(s)] ds\right\} = z_+. \quad (1-5)$$

Theorem 1.1 shows that any small solution decomposes into a nonlinear bound state plus a radiation term. In particular, we have the asymptotic stability of the family of small solitary waves. The condition (1-4) makes $v(t)$ orthogonal to the nondecaying solutions of the linearization of (1-1) about the solitary wave at $z(t)$; this is an essential ingredient for establishing decay and scattering for v (see Section 3 for further discussion).

Theorem 1.1 fits in the context of the stability of small solitary waves for nonlinear Schrödinger equations with potential, for which there are many results available. An even more extensive literature exists concerning other notions of stability, stability of large solitary waves, and so on. We refer the interested reader to [Buslaev and Perelman 1992; Gustafson et al. 2004; Mizumachi 2008; Kirr and

Mizrak 2009; Kirr and Zarnescu 2009; Soffer and Weinstein 1990; 1992; 2004; Weder 2000; Tsai and Yau 2002a; 2002b; 2002c; 2002d; Cuccagna 2011; 2014; Cuccagna and Pelinovsky 2014] for a sample of the many relevant results that are available. See in particular [Datchev and Holmer 2009; Deift and Park 2011; Holmer et al. 2007a; 2007b; Holmer and Zworski 2007; 2009; Kaminaga and Ohta 2009; Fukuizumi et al. 2008] for related results in the setting of NLS with a delta potential. We will keep our focus on the discussion of small solitary waves.

Our result is closely related to those appearing in [Gustafson et al. 2004; Mizumachi 2008], both of which prove asymptotic stability of small solitary waves for NLS with a potential that supports a single negative eigenvalue, with data in H^1 and mass-supercritical nonlinearities. In [Gustafson et al. 2004], the authors relied crucially on the endpoint Strichartz estimate in three dimensions. Mizumachi [2008] addressed the one-dimensional case, in which case the usual endpoint Strichartz estimate is unavailable. His approach was to establish suitable linear estimates in “reversed” Strichartz spaces, in which case the L_t^2 endpoint comes back into play.

Theorem 1.1 is an analogue of the main result appearing in [Mizumachi 2008], which treats a class of potentials that does not include the attractive delta potential. The key to extending this type of result to the delta potential is to observe that by relying on exact identities related to the Schrödinger operator with a delta potential, one can recover the full range of linear estimates that played such an essential role in [Mizumachi 2008]. We carry this out in Section 2B. Once the requisite linear estimates are in place, one could then follow many of the remaining arguments in [Mizumachi 2008] rather directly, although this is not the route that we take. Instead, we set up the problem and prove the main result in a way that is inspired by the presentation in [Gustafson et al. 2004], which we found to be rather conceptually clear.

Our result is also closely tied to the work of Fukuizumi, Ohta, and Ozawa [Fukuizumi et al. 2008] and Kaminaga and Ohta [2009]; see also [Goodman et al. 2004]. These authors considered the problem of stability and instability of nonlinear bound states for NLS with an attractive delta potential, relying in particular on explicit formulas that they derived for the nonlinear bound states (see Section 2C below). For focusing nonlinearities, one finds that in the mass-subcritical and mass-critical cases, nonlinear bound states are all orbitally stable. In the mass-supercritical case, that there exists $E_1 < -\frac{1}{2}q^2$ such that ground states corresponding to $E \in (E_1, -\frac{1}{2}q^2)$ are orbitally stable, while those corresponding to $E \in (-\infty, E_1)$ are unstable. For defocusing nonlinearities, all nonlinear ground states are orbitally stable. Thus our main result, Theorem 1.1, extends the results of [Fukuizumi et al. 2008; Kaminaga and Ohta 2009] in the mass-supercritical case to asymptotic stability for E in a neighborhood of $-\frac{1}{2}q^2$.

Finally, we would also like to mention the result of [Deift and Park 2011], which establishes the asymptotic stability of solitons for the focusing cubic NLS with a delta potential and even initial data by making use of complete integrability and the method of nonlinear steepest descent. This result in particular extended the results appearing [Datchev and Holmer 2009; Holmer et al. 2007a; 2007b; Holmer and Zworski 2007].

As mentioned above, our previous work on the one-dimensional NLS with a repulsive delta potential [Masaki et al. 2019] considered the case of a cubic nonlinearity. It is an interesting question whether one also has asymptotic stability in the setting of an attractive potential and L^2 -subcritical nonlinearities; recall

that orbital stability was proven by [Fukuizumi et al. 2008; Kaminaga and Ohta 2009; Goodman et al. 2004]. Proving asymptotic stability would most likely require the introduction of stronger integrability conditions on the initial data; for example, this is the case in [Kirr and Mızrak 2009; Kirr and Zarnescu 2009], which proved stability of small solitary waves for NLS with potential for some mass-subcritical nonlinearities in dimensions $d \in \{2, 3\}$. In our case, we start only with H^1 data and are therefore restricted to $p \geq 4$; this is completely analogous to the situation of trying to prove small-data scattering for the standard power-type NLS. To see specific the technical points that lead to this restriction, see the estimates of the $|v|^p v$ term in the proofs of Lemmas 4.6, 4.7, and 4.8 (as well as the $\mathcal{O}(v^p Q)$ -term in Lemma 4.7).

Briefly, the proof of Theorem 1.1 goes as follows. One shows that as long as the u remains small in H^1 , there exists a unique decomposition (1-3) such that (1-4) holds. Using (1-1) and differentiating (1-4) leads to a coupled system of equations for $v(t)$ and $z(t)$. Relying largely on estimates for the linear propagator e^{-itH} and estimates on the bound states $Q[z]$ for small z , one can use these equations to close a bootstrap argument, proving that the smallness of u in H^1 (as well as the smallness of v and z in various norms) persists. Thus, one can extend the decomposition for all times; furthermore, the bounds proved on v and z suffice to establish the asymptotics claimed in Theorem 1.1. The particular choice of the orthogonality condition (1-4) guarantees that the ODE involving $z[t]$ is at least quadratic in v , which is essential for proving the necessary bootstrap estimates; see Remark 3.2 for further discussion of this point.

Outline of the paper. In Section 2 we introduce notation and gather some preliminary results. We introduce the linear operator H in Section 2A. In Section 2B, we prove a range of Strichartz and local smoothing estimates for $e^{-itH} P_c$. These match the form of the estimates of [Mizumachi 2008], who considered a class of potentials that did not include the delta potential. We are able to give rather direct proofs using the explicit formula for the resolvent. We also prove a technical result related to the comparison of the \dot{H}^1 inner product to the bilinear form given by HP_c . In Section 2C we discuss the existence and properties of small nonlinear bound states, and in Section 2D we record a local well-posedness result for (1-1). In Section 3 we set up the problem, describing in detail how to find the decomposition (1-3) satisfying (1-4). Finally, in Section 4 we carry out the main bootstrap argument and complete the proof of Theorem 1.1.

2. Preliminaries

We begin by recording some notation. We write

$$\langle f, g \rangle = \int \bar{f} g \, dx$$

for the usual L^2 inner product. Throughout the paper we will write $F(u) = \mu|u|^p u$ for the nonlinearity. We write $\mathcal{F}f$ or \hat{f} for the Fourier transform. We write $A \lesssim B$ to denote $A \leq CB$ for some $A, B, C > 0$.

Constants below may depend on the parameter q (the strength of the potential), but we will not make explicit reference to this dependence. We would like to point out that some of the implicit constants in the estimates for $e^{-itH} P_c$ below would blow up as $|q| \rightarrow 0$ (for example, when the proof relies on the

fact that $|q - i\mu| \gtrsim |q|$ for $\mu \in \mathbb{R}$). In particular, the small parameter δ appearing in the statement of the main result (Theorem 1.1) depends on q and would degenerate to zero as $|q| \rightarrow 0$.

2A. Linear theory. The linear Schrödinger equation with a delta potential is a classical model in quantum mechanics that is covered extensively in [Albeverio et al. 1988]. We consider in this paper the case of an attractive delta potential of the form

$$H = -\frac{1}{2}\partial_x^2 + q\delta(x), \quad q < 0.$$

More precisely, the operator H is defined by $-\frac{1}{2}\partial_x^2$ on its domain

$$D(H) = \{u \in H^1(\mathbb{R}) \cap H^2(\mathbb{R} \setminus \{0\}) : \partial_x u(0+) - \partial_x u(0-) = 2qu(0)\}$$

and extends to a self-adjoint operator on L^2 with purely absolutely continuous essential spectrum equal to $[0, \infty)$. If $q > 0$ (the repulsive case), then H has no eigenvalues. In $q < 0$ (the attractive case), H has a single negative eigenvalue $-\frac{1}{2}q^2$ with a one-dimensional eigenspace spanned by the L^2 -normalized eigenfunction

$$\phi_0(x) := |q|^{1/2} e^{q|x|}.$$

In this paper we restrict attention to the attractive case.

2B. Local smoothing and Strichartz estimates. In this section we prove several local smoothing and Strichartz estimates for e^{-itH} . We write P_c to denote the projection onto the absolutely continuous spectrum.

The starting point for the estimates we will prove is the following spectral resolution of the free propagator:

$$e^{-itH} P_c = \int_0^\infty e^{-it\lambda} E(\lambda) d\lambda, \quad E(\lambda) := \frac{1}{2\pi i} [R(\lambda + i0) - R(\lambda - i0)]. \quad (2-1)$$

Here $R(z) = (H - z)^{-1}$ is the resolvent, and $R(\lambda \pm i0)$ denotes the analytic continuation onto the real line from the upper/lower half plane. In fact, working with the resolvent is not the only way to carry out the requisite linear analysis for this problem. As an alternative, one could also proceed through the development of the distorted Fourier transform; e.g., this is how we proceeded in our previous work involving the delta potential [Masaki et al. 2019].

For the case of the delta potential, we have explicit formulas for the integral kernel of the resolvent, namely

$$\begin{aligned} R(\lambda + i0; x, y) &= \frac{i}{2\sqrt{\lambda}} \left[e^{i|x-y|\sqrt{\lambda}} - \frac{q}{q - i\sqrt{\lambda}} e^{i(|x|+|y|)\sqrt{\lambda}} \right], \\ R(\lambda - i0; x, y) &= \frac{i}{2\sqrt{\lambda}} \left[e^{-i|x-y|\sqrt{\lambda}} - \frac{q}{q + i\sqrt{\lambda}} e^{-i(|x|+|y|)\sqrt{\lambda}} \right] \end{aligned}$$

for $\lambda > 0$. We similarly write $E(\lambda; x, y)$ for the kernel of $E(\lambda)$. These identities can be found, for example, in [Albeverio et al. 1988, Chapter I.3], but they are also readily derived by hand. In particular,

one can recognize the first term as the free resolvent, while the second term (representing the contribution of the potential) simply fixes the boundary condition.

Typically we will focus on estimating $R(\lambda + i0)$, as the other term is similar. We write the kernel in two pieces, namely

$$R(\lambda + i0; x, y) = R_1(\lambda; x, y) + R_2(\lambda; x, y),$$

where

$$R_1(\lambda; x, y) = \frac{i}{2\sqrt{\lambda}} [e^{i|x-y|\sqrt{\lambda}} - e^{i(|x|+|y|)\sqrt{\lambda}}], \quad (2-2)$$

$$R_2(\lambda; x, y) = \frac{1}{2(q - i\sqrt{\lambda})} e^{i(|x|+|y|)\sqrt{\lambda}}. \quad (2-3)$$

We note that

$$R_1(\lambda; x, y) = \frac{i}{2\sqrt{\lambda}} \begin{cases} [e^{-ix\sqrt{\lambda}} - e^{ix\sqrt{\lambda}}]e^{iy\sqrt{\lambda}}, & y \geq x \geq 0, \\ 0, & y \geq 0 \geq x, \\ e^{-ix\sqrt{\lambda}}[e^{iy\sqrt{\lambda}} - e^{-iy\sqrt{\lambda}}], & 0 \geq y \geq x. \end{cases} \quad (2-4)$$

There are analogous formulas in the cases $x \geq y \geq 0$, $x \geq 0 \geq y$, and $0 \geq x \geq y$. We will focus on treating the three cases appearing in (2-4).

To simplify the presentation below, we will use $\tilde{\mathcal{F}}f$ to denote quantities that are similar (but not identical) to the Fourier transform of f ; in particular, we use this notation for quantities that obey the bounds

$$\|\tilde{\mathcal{F}}f\|_{L^2} \lesssim \|f\|_{L^2} \quad \text{and} \quad \|\tilde{\mathcal{F}}f\|_{L^\infty} \lesssim \|f\|_{L^1}. \quad (2-5)$$

As a typical example, we could apply this notation to a term like

$$\int_x^\infty e^{i\xi y} f(y) dy = (2\pi)^{1/2} \mathcal{F}[\chi_{(x,\infty)} f](-\xi).$$

Indeed, for this example the two bounds in (2-5) can be easily checked (and are uniform in x).

We begin with the standard one-dimensional Strichartz estimates.

Proposition 2.1 (Strichartz estimates). *The following estimates hold on any space-time slab $I \times \mathbb{R}$ with $0 \in I$:*

$$\|e^{-itH} P_c f\|_{(L_t^4 L_x^\infty \cap L_t^\infty L_x^2)(I \times \mathbb{R})} \lesssim \|f\|_{L^2},$$

$$\left\| \int_0^t e^{-i(t-s)H} P_c F(s) ds \right\|_{(L_t^4 L_x^\infty \cap L_t^\infty L_x^2)(I \times \mathbb{R})} \lesssim \|F\|_{L_t^\alpha L_x^\beta(I \times \mathbb{R})}$$

for any $(\alpha, \beta) \in [1, \frac{4}{3}] \times [1, 2]$ satisfying $\frac{2}{\alpha} + \frac{1}{\beta} = \frac{5}{2}$.

As is well known, the proof boils down to the following dispersive estimates.

Lemma 2.2 (dispersive estimates). *The following estimates hold:*

$$\|e^{-itH} P_c f\|_{L^2} \lesssim \|f\|_{L^2} \quad \text{and} \quad \|e^{-itH} P_c f\|_{L^\infty} \lesssim |t|^{-1/2} \|f\|_{L^1}.$$

Proof of Lemma 2.2. It is clear that $e^{-itH}P_c$ maps L^2 to L^2 boundedly. For the $L^1 \rightarrow L^\infty$ estimate, we start from (2-1). The desired estimate is well known for the case of the free Schrödinger equation, and hence we consider only the contribution of the potential. After a change of variables, we are left to prove

$$\sup_x \left| \int f(y) \int e^{-it\lambda^2/2 - i(|x|+|y|)\lambda} \frac{q}{q-i\lambda} d\lambda dy \right| \leq |t|^{-1/2} \|f\|_{L^1}.$$

We apply Plancherel in the $d\lambda$ integral and observe (by explicit computation) that

$$\sup_{\theta \in \mathbb{R}} \|\mathcal{F}(e^{-it\lambda^2/2 + i\theta\lambda})\|_{L^\infty} \lesssim |t|^{-1/2}.$$

Therefore the proof boils down to showing that $\mathcal{F}((q-i\lambda)^{-1}) \in L^1$. In fact, by Cauchy–Schwarz and Plancherel,

$$\|\mathcal{F}((q-i\lambda)^{-1})\|_{L^1} \lesssim \|(1-\partial_x^2)(q-i\lambda)^{-1}\|_{L^2} \lesssim 1. \quad (2-6)$$

The result follows. \square

We turn to the following weighted estimates for the linear propagator.

Proposition 2.3 (local smoothing estimates). *The following estimates hold:*

$$\|\langle x \rangle^{-3/2} e^{-itH} P_c f\|_{L_x^\infty L_t^2} \lesssim \|f\|_{L^2}, \quad (2-7)$$

$$\|\partial_x e^{-itH} P_c f\|_{L_x^\infty L_t^2} \lesssim \|f\|_{H^1}. \quad (2-8)$$

Remark 2.4. The H^1 -norm in (2-8) may be replaced by the $H^{1/2}$ -norm by utilizing the analogous estimate for the free propagator, see, e.g., [Kenig et al. 1991, Lemma 2.1], and estimating the contribution of the potential separately. We proceed with (2-8), as it admits a simple self-contained proof and is sufficient for our applications.

Proof of Proposition 2.3. We begin by reducing each estimate to one given purely in terms of the resolvent. Let

$$m = m(x, \partial_x) \in \{\langle x \rangle^{-3/2}, \partial_x\}$$

and $X = L^2$ or H^1 . We will show

$$\|m e^{-itH} P_c\|_{X \rightarrow L_x^\infty L_t^2} \lesssim \|m E(\lambda)\|_{X \rightarrow L_x^\infty L_\lambda^2}. \quad (2-9)$$

To see this, we let $G \in L_x^1 L_t^2$ and use Plancherel to estimate

$$\begin{aligned} |\langle m e^{-itH} P_c f, G \rangle_{L_{t,x}^2}| &= \left| \int e^{-it\lambda} \overline{G(t, x)} m(x, \partial_x) E(\lambda; x, y) f(y) d\lambda dy dx dt \right| \\ &= \left| \int [\mathcal{F}_t^{-1} G](\lambda, x) [m E(\lambda) f](x) dx d\lambda \right| \\ &\lesssim \|\mathcal{F}_t^{-1} G\|_{L_x^1 L_\lambda^2} \|m E(\lambda) f\|_{L_x^\infty L_\lambda^2} \\ &\lesssim \|G\|_{L_x^1 L_t^2} \|m E(\lambda)\|_{X \rightarrow L_x^\infty L_\lambda^2} \|f\|_X. \end{aligned}$$

Thus (2-9) follows.

Using (2-9), we see that (2-7) will follow from

$$\|\langle x \rangle^{-3/2} R(\lambda \pm i0) f\|_{L_x^\infty L_\lambda^2} \lesssim \|f\|_{L^2}. \quad (2-10)$$

We focus on $R(\lambda + i0)$ and write $R = R_1 + R_2$ as in (2-2) and (2-3). The contribution of (2-3) is easily handled. In fact, by a change of variables,

$$\begin{aligned} \left\| \frac{1}{q - i\sqrt{\lambda}} \int e^{i|y|\sqrt{\lambda}} f(y) dy \right\|_{L_\lambda^2}^2 &= \left\| \frac{1}{q - i\sqrt{\lambda}} \tilde{\mathcal{F}} f(\sqrt{\lambda}) \right\|_{L_\lambda^2}^2 \\ &\lesssim \int \frac{|\mu|}{q^2 + \mu^2} |\tilde{\mathcal{F}} f(\mu)|^2 d\mu \lesssim \|f\|_{L^2}^2. \end{aligned}$$

To estimate the contribution of (2-2), we split into low and high energies. We let $\chi(\lambda)$ denote a smooth cutoff to $|\lambda| \leq 1$ and write $\chi^c = 1 - \chi$. On the support of χ^c , we can argue as we did for (2-3), changing variables and estimating the contribution via

$$\int_{|\mu| \geq 1} \frac{1}{|\mu|} |\tilde{\mathcal{F}} f(\mu)|^2 d\mu \lesssim \|f\|_{L^2}^2,$$

which is acceptable.

We turn to the low-energy contribution of (2-2). Here we use (2-4); in particular, we will consider the cases $y \geq x \geq 0$ and $0 \geq y \geq x$. In the first case, we use the bound

$$|e^{ix\sqrt{\lambda}} - e^{-ix\sqrt{\lambda}}| \lesssim |x|\sqrt{\lambda}$$

and estimate

$$\begin{aligned} \left\| \langle x \rangle^{-1} \chi(\lambda) \frac{1}{\sqrt{\lambda}} [e^{ix\sqrt{\lambda}} - e^{-ix\sqrt{\lambda}}] \int_x^\infty e^{i\sqrt{\lambda}y} f(y) dy \right\|_{L_x^\infty L_\lambda^2} &\lesssim \|\chi(\lambda) \tilde{\mathcal{F}} f(\sqrt{\lambda})\|_{L_\lambda^2} \\ &\lesssim \|\sqrt{\mu} \tilde{\mathcal{F}} f(\mu)\|_{L_\mu^2(|\mu| \leq 1)} \lesssim \|f\|_{L^2}, \end{aligned}$$

which is acceptable. In the remaining case, we use Cauchy–Schwarz to estimate

$$\begin{aligned} \left\| \langle x \rangle^{-3/2} \chi(\lambda) \frac{1}{\sqrt{\lambda}} \int_x^0 [e^{i\sqrt{\lambda}y} - e^{-i\sqrt{\lambda}y}] f(y) dy \right\|_{L_x^\infty L_\lambda^2} &\lesssim \|\langle x \rangle^{-3/2} \chi(\lambda) \int_0^x |y| |f(y)| dy\|_{L_x^\infty L_\lambda^2} \\ &\lesssim \|\chi(\lambda) |x|^{3/2} \langle x \rangle^{-3/2} \|f\|_{L^2}\|_{L_x^\infty L_\lambda^2} \lesssim \|f\|_{L^2}, \end{aligned}$$

which is acceptable. This completes the proof of (2-7).

We turn to (2-8). Using (2-9), it suffices to prove

$$\|\partial_x R(\lambda \pm i0) f\|_{L_x^\infty L_\lambda^2} \lesssim \|f\|_{H^1}. \quad (2-11)$$

Let us again focus on the contribution of $R(\lambda + i0)$. This time we go back to the original form of the resolvent and estimate the two pieces separately. Let us first establish

$$\left\| \lambda^{-1/2} \partial_x \int e^{i|x-y|\sqrt{\lambda}} f(y) dy \right\|_{L_x^\infty L_\lambda^2} \lesssim \|f\|_{H^1}.$$

To see this, we firstly observe that

$$\partial_x e^{i|x-y|\sqrt{\lambda}} = -\partial_y e^{i|x-y|\sqrt{\lambda}} = \pm i\sqrt{\lambda} e^{i|x-y|\sqrt{\lambda}} \quad (2-12)$$

almost everywhere, where the sign depends on the ordering of x and y . Using the final equality in (2-12) and writing χ for a cutoff to $\lambda \leq 1$, we can change variables and use (2-5) to estimate

$$\begin{aligned} \left\| \lambda^{-1/2} \chi \partial_x \int e^{i|x-y|\sqrt{\lambda}} f(y) dy \right\|_{L_x^\infty L_\lambda^2} &\lesssim \|\chi \tilde{\mathcal{F}} f(\sqrt{\lambda})\|_{L_\lambda^2} \\ &\lesssim \| |\xi|^{1/2} \tilde{\mathcal{F}} f \|_{L_\xi^2(\{|\xi| \leq 1\})} \lesssim \|f\|_{L^2}, \end{aligned}$$

which is acceptable. Writing $\chi^c = 1 - \chi$, we can use the second equality in (2-12), integrate by parts, change variables, and use (2-5) to estimate

$$\begin{aligned} \left\| \lambda^{-1/2} \chi^c \partial_x \int e^{i|x-y|\sqrt{\lambda}} f(y) dy \right\|_{L_x^\infty L_\lambda^2} &\lesssim \|\lambda^{-1/2} \chi^c \tilde{\mathcal{F}}[f'](\sqrt{\lambda})\|_{L_\lambda^2} \\ &\lesssim \| |\xi|^{-1/2} \tilde{\mathcal{F}}[f'] \|_{L_\xi^2(\{|\xi| > 1\})} \lesssim \|f'\|_{L^2}, \end{aligned}$$

which is also acceptable.

It remains to consider the piece coming from the potential. We need to show

$$\left\| \frac{1}{\sqrt{\lambda}[q - i\sqrt{\lambda}]} \partial_x [e^{i|x|\sqrt{\lambda}}] \int e^{i|y|\sqrt{\lambda}} f(y) dy \right\|_{L_x^\infty L_\lambda^2} \lesssim \|f\|_{H^1}.$$

In fact, by a change of variables and (2-5), we can control this term by

$$\left\| \frac{1}{q - i\sqrt{\lambda}} \tilde{\mathcal{F}} f(\sqrt{\lambda}) \right\|_{L_\lambda^2} \lesssim \left\| \frac{|\xi|^{1/2}}{\langle \xi \rangle} \tilde{\mathcal{F}} f \right\|_{L_\xi^2} \lesssim \|f\|_{L^2},$$

which is acceptable.

This completes the proof of (2-11) and hence the proof of Proposition 2.3. \square

Combining the usual Strichartz estimates (Proposition 2.1) with the weighted local smoothing estimate in Proposition 2.3 yields the following corollary:

Corollary 2.5. *The following estimate holds:*

$$\left\| \int_0^t e^{-i(t-s)H} P_c F(s) ds \right\|_{L_t^4 L_x^\infty \cap L_t^\infty L_x^2} \lesssim \|\langle x \rangle^{5/2} F\|_{L_{t,x}^2}.$$

Proof. Using the Strichartz estimate Proposition 2.1, the dual estimate to (2-7), and Cauchy–Schwarz, we have

$$\begin{aligned} \left\| \int_{\mathbb{R}} e^{-i(t-s)H} P_c F(s) ds \right\|_{L_t^4 L_x^\infty \cap L_t^\infty L_x^2} &\lesssim \left\| \int_{\mathbb{R}} e^{isH} P_c F(s) ds \right\|_{L^2} \\ &\lesssim \|\langle x \rangle^{3/2} P_c F\|_{L_x^1 L_t^2} \lesssim \|\langle x \rangle^{5/2} P_c F\|_{L_{t,x}^2}. \end{aligned}$$

The desired estimate now follows from the Christ–Kiselev lemma [2001]. \square

We will also need the following inhomogeneous local smoothing estimates.

Proposition 2.6. *For any $t \geq 0$, we have*

$$\left\| \langle x \rangle^{-1} \int_0^t e^{-i(t-s)H} P_c F(s) ds \right\|_{L_x^\infty L_t^2} \lesssim \|\langle x \rangle F\|_{L_x^1 L_t^2}, \quad (2-13)$$

$$\left\| \int_0^t \partial_x e^{-i(t-s)H} P_c F(s) ds \right\|_{L_x^\infty L_t^2} \lesssim \|F\|_{L_x^1 L_t^2}. \quad (2-14)$$

Proof of Proposition 2.6. We begin with the identity

$$\begin{aligned} 2 \int_0^t e^{-i(t-s)H} P_c F(s) ds \\ = \int_{\mathbb{R}} e^{-i(t-s)H} P_c F(s) ds + \int_0^\infty e^{-i(t-s)H} P_c F(s) ds - \int_{-\infty}^0 e^{-i(t-s)H} P_c F(s) ds. \end{aligned} \quad (2-15)$$

In fact, this is a consequence of

$$\int_{\mathbb{R}} e^{-i(t-s)H} P_c F(s) ds = \int_{-\infty}^t e^{-i(t-s)H} P_c F(s) ds - \int_t^\infty e^{-i(t-s)H} P_c F(s) ds,$$

which follows from the fact that both sides solve

$$i \partial_t u = H u, \quad \text{with } u(0) = \int_{\mathbb{R}} e^{isH} P_c F(s) ds.$$

In light of (2-15), it therefore suffices to estimate

$$\int_{\mathbb{R}} e^{-i(t-s)H} P_c \chi(s) F(s) ds,$$

where $\chi \in \{1, \chi_{(0,\infty)}, \chi_{(-\infty,0)}\}$.

Similar to the proof of Proposition 2.3, we will use (2-1) and Plancherel to reduce the desired bounds to an estimate given in terms of the resolvent. In particular, we write

$$\begin{aligned} \int e^{-i(t-s)H} P_c \chi F(s) ds &= \int e^{-it\lambda} E(\lambda) \int e^{is\lambda} \chi(s) F(s) ds d\lambda \\ &= \mathcal{F}_\lambda[E(\lambda) \mathcal{F}_s^{-1}(\chi F)](t), \end{aligned}$$

where we use subscripts to denote the variable of integration in the definition of the Fourier transform. Thus, writing $m = m(x, \partial_x) \in \{\langle x \rangle^{-1}, \partial_x\}$ and $X = \langle x \rangle^{-1} L^1$ or $X = L^1$, we use Plancherel and Minkowski's inequality to estimate

$$\begin{aligned} \left\| m \int e^{-i(t-s)H} P_c \chi F(s) ds \right\|_{L_x^\infty L_t^2} &\lesssim \|m E(\lambda) \mathcal{F}_s^{-1}(\chi F)\|_{L_x^\infty L_\lambda^2} \lesssim \|m E(\lambda) \mathcal{F}_s^{-1}(\chi F)\|_{L_\lambda^2 L_x^\infty} \\ &\lesssim \|m E(\lambda)\|_{X \rightarrow L_x^\infty} \|\mathcal{F}_s^{-1}(\chi F)(\lambda)\|_X \| \lambda \|_{L_\lambda^2} \\ &\lesssim \left[\sup_\lambda \|m E(\lambda)\|_{X \rightarrow L_x^\infty} \right] \|F\|_{L_t^2 X} \\ &\lesssim \left[\sup_\lambda \|m E(\lambda)\|_{X \rightarrow L_x^\infty} \right] \|F\|_{X L_t^2}. \end{aligned}$$

The proof of (2-13) and (2-14) therefore reduces to the two estimates

$$\sup_{\lambda} \|\langle x \rangle^{-1} R(\lambda \pm i0) f\|_{L_x^\infty} \lesssim \|\langle x \rangle f\|_{L_x^1}, \quad (2-16)$$

$$\sup_{\lambda} \|\partial_x R(\lambda \pm i0) f\|_{L_x^\infty} \lesssim \|f\|_{L_x^1}. \quad (2-17)$$

We consider $R(\lambda + i0)$, the other case being similar. We decompose the kernel as $R_1 + R_2$, as in (2-2) and (2-3). The contribution of R_2 to both (2-16) and (2-17) is handled easily. In fact, since $|q - i\sqrt{\lambda}| \geq |q|$, we have

$$\|R_2(\lambda) f\|_{L_x^\infty} + \|\partial_x R_2(\lambda) f\|_{L_x^\infty} \lesssim \|\tilde{F} f\|_{L_x^\infty} \lesssim \|f\|_{L_x^1}$$

uniformly in λ .

We turn to the contribution of R_1 . The contribution to (2-17) is straightforward, as we can estimate

$$\|\partial_x R_1(\lambda) f\|_{L_x^\infty} \lesssim |\tilde{F} f(\sqrt{\lambda})| \lesssim \|f\|_{L^1},$$

uniformly in λ . For the contribution to (2-16), we recall (2-4). In particular, we need only consider the cases $y \geq x \geq 0$ and $0 \geq y \geq x$. In the first case, we estimate

$$\left| \frac{1}{\sqrt{\lambda}} (e^{-ix\sqrt{\lambda}} - e^{ix\sqrt{\lambda}}) \int_x^\infty e^{i\sqrt{\lambda}y} f(y) dy \right| \lesssim |x| \tilde{F} f(\sqrt{\lambda}),$$

and hence the desired bound holds in this regime (see (2-5)). Finally, if $0 \geq y \geq x$, we estimate

$$\left| \frac{1}{\sqrt{\lambda}} \int_x^0 [e^{i\sqrt{\lambda}y} - e^{-i\sqrt{\lambda}y}] f(y) dy \right| \lesssim \|yf(y)\|_{L^1}$$

uniformly in x and λ . Thus the desired bound holds in this regime as well. This completes the proof of Proposition 2.6. \square

Finally, let us record one additional corollary of Proposition 2.3.

Corollary 2.7. *The following estimates hold:*

$$\begin{aligned} \left\| \langle x \rangle^{-3/2} \int_0^t e^{-i(t-s)H} P_c F(s) ds \right\|_{L_x^\infty L_t^2(\mathbb{R} \times [0, T])} &\lesssim \|F\|_{L_t^1 L_x^2([0, T] \times \mathbb{R})}, \\ \left\| \partial_x \int_0^t e^{-i(t-s)H} P_c F(s) ds \right\|_{L_x^\infty L_t^2(\mathbb{R} \times [0, T])} &\lesssim \|F\|_{L_t^1 H_x^1([0, T] \times \mathbb{R})}. \end{aligned}$$

Proof. To rid ourselves of the integral over $[0, t]$ we again use the decomposition (2-15) as in Proposition 2.6 and endeavor to estimate χF , with

$$\chi \in \{1, \chi_{(0, \infty)}, \chi_{(-\infty, 0)}\}.$$

Let $m = m(x, \partial_x) \in \{\langle x \rangle^{-3/2}, \partial_x\}$ and write $X = L^2$ if $m = \langle x \rangle^{-3/2}$ and $X = H^1$ if $m = \partial_x$. Then using Proposition 2.3, the boundedness of e^{isH} on X , and Minkowski's inequality, we may estimate

$$\left\| m \int_{\mathbb{R}} e^{-i(t-s)H} P_c \chi(s) F(s) ds \right\|_{L_x^\infty L_t^2} \lesssim \left\| \int_{\mathbb{R}} e^{-isH} \chi(s) F(s) ds \right\|_X \lesssim \|F\|_{L_t^1 X}.$$

The result follows. \square

We close this section with a technical result relating the usual Sobolev spaces with those defined in terms of H . We state the result we need as follows. In the following, we let $m(\partial_x)$ denote the Fourier multiplier operator with symbol $m(\mu)$.

Lemma 2.8. *We have*

$$\langle f, HP_cg \rangle = \langle f, -\tfrac{1}{2}\partial_x^2 g \rangle + B(f, g), \quad (2-18)$$

where $B(f, g)$ is a linear combination of terms of the form

$$\langle m(\partial_x)\partial_x f, \partial_x g \rangle, \quad \text{where } m(\mu) = (q - i\mu)^{-1}.$$

Consequently, for $f = P_c f$,

$$\|\sqrt{H}f\|_{L^2} \lesssim \|f\|_{\dot{H}^1} \quad \text{and} \quad \|f\|_{\dot{H}^1} \lesssim \|\sqrt{H}f\|_{L^2} + \|f\|_{L^2}. \quad (2-19)$$

Although (2-19) has already been shown in [Duchêne et al. 2011, Section VIII D] via the $W^{1,p}$ boundedness of wave operators for H , we give a simpler proof of (2-19) by using the explicit representation of \sqrt{H} .

Proof of Lemma 2.8. By the spectral theorem and the explicit form of the resolvent, we have the identity

$$\langle f, HP_cg \rangle = \langle f, -\tfrac{1}{2}\partial_x^2 g \rangle + B(f, g),$$

where $B(f, g)$ is a linear combination of terms like

$$\iiint \frac{\lambda}{\sqrt{\lambda}(q - i\sqrt{\lambda})} e^{i(|x|+|y|)\sqrt{\lambda}} f(x)g(y) dx dy d\lambda = \int \frac{\mu^2}{2(q - i\mu)} \tilde{\mathcal{F}}f(\mu)\tilde{\mathcal{F}}g(\mu) d\mu.$$

Here we use the notation

$$\tilde{\mathcal{F}}f(\mu) = \int e^{i|x|\mu} f(x) dx.$$

This is consistent with the usage above, and in fact in this case $\tilde{\mathcal{F}}f$ can be written exactly as the sum of Fourier transforms of f and its reflection. Thus (2-18) follows from Plancherel.

We turn to (2-19). For the first estimate we simply observe that $m(\partial_x)$ maps $L^2 \rightarrow L^2$ boundedly. For the second estimate, we observe in fact that $m(\partial_x)\partial_x$ maps $L^2 \rightarrow L^2$ boundedly, and hence by Young's inequality

$$\begin{aligned} \|\partial_x f\|_{L^2}^2 &\lesssim \|\sqrt{H}f\|_{L^2}^2 + \|f\|_{L^2}\|\partial_x f\|_{L^2} \\ &\lesssim \|\sqrt{H}f\|_{L^2}^2 + \varepsilon\|\partial_x f\|_{L^2}^2 + \varepsilon^{-1}\|f\|_{L^2}^2 \end{aligned}$$

for any $\varepsilon > 0$. Choosing $\varepsilon \ll 1$ implies the desired bound. □

Remark 2.9. The multiplier $m(\partial_x)$ appearing in (2-18) actually maps $L^r \rightarrow L^r$ boundedly for any $1 \leq r \leq \infty$. Indeed, it was already proven in (2-6) that $\mathcal{F}^{-1}m \in L^1$, and hence this is a consequence of Young's inequality. In particular, we are not using any multiplier theorems and are able to access the L^1, L^∞ endpoints. In a similar way, we see that $m(\partial_x)$ is bounded on $L_x^p L_t^q$ for all $1 \leq p, q \leq \infty$. These will be useful in the proof of Lemma 4.7 below.

2C. Existence of small solitary waves. In this section we discuss the existence and properties of solutions to (1-2).

In [Fukuizumi et al. 2008], the authors considered (1-1) with a focusing nonlinearity and provided an explicit formula for the family of nonlinear bound states. Using our notation, these solutions are given by

$$Q(x) = \left(\frac{(p+2)|E|}{2|\mu|} \right)^{1/p} \cosh^{-2/p} \left(p \sqrt{\frac{|E|}{2}} |x| + \operatorname{arctanh} \left(\frac{|q|}{\sqrt{2|E|}} \right) \right),$$

where $E < -\frac{1}{2}q^2$ and $\mu < 0$. This formula is obtained by solving the relevant ODE on each side of $x = 0$ and then gluing them together at $x = 0$ to impose the jump condition $Q'(0+) - Q'(0-) = 2qQ(0)$. This approach also works in the defocusing case $\mu > 0$ (see [Kaminaga and Ohta 2009]); the resulting formula is

$$Q(x) = \left(\frac{(p+2)|E|}{2\mu} \right)^{1/p} \sinh^{-2/p} \left(p \sqrt{\frac{|E|}{2}} |x| + \operatorname{arctanh} \left(\frac{\sqrt{2|E|}}{|q|} \right) \right)$$

for $-\frac{1}{2}q^2 < E < 0$. When $E = 0$, one has the solution

$$Q(x) = \left(\frac{(p+2)|q|^2}{\mu(p|q||x|+2)^2} \right)^{1/p},$$

which belongs to L^2 provided $p < 4$.

From the explicit formulas for Q , one can observe that as E approaches $-\frac{1}{2}q^2$, the functions Q behave like a small multiple of the linear eigenfunction. It will be convenient to describe this behavior in Proposition 2.10 below. In particular, we find it convenient to follow the approach of [Gustafson et al. 2004] and parametrize the family of ground states by small $z \in \mathbb{C}$.

In the following, we write

$$D_j Q[z] = \frac{\partial}{\partial z_j} Q[z],$$

where we identify $z \in \mathbb{C}$ with the real vector (z_1, z_2) . We write $DQ[z]$ for the Jacobian $DQ[z] : \mathbb{C} \rightarrow \mathbb{C}$ with

$$DQ[z]w = D_1 Q[z] \operatorname{Re} w + D_2 Q[z] \operatorname{Im} w \quad \text{for } w \in \mathbb{C}.$$

We will prove the following.

Proposition 2.10. *There exists small enough $\delta > 0$ such that for $z \in \mathbb{C}$ with $|z| < \delta$, we have the following.*

- *There exists a unique solution $Q = Q[z]$ to (1-2) with $E = E[|z|] \in \mathbb{R}$.*
- *We may write $Q[z] = z\phi_0 + h$, where*

$$\|h\|_{H^{1,k}(\mathbb{R}) \cap H^2(\mathbb{R} \setminus \{0\})} \lesssim |z|^2, \quad \|Dh\|_{H^{1,k}} \lesssim |z|, \quad \text{and} \quad \|D^2 h\|_{L^2} \lesssim 1$$

for any $k \geq 0$.

- $E[|z|] = -\frac{1}{2}q^2 + \mathcal{O}(z)$.
- $Q[ze^{i\theta}] = Q[z]e^{i\theta}$ and $Q[|z|]$ is real-valued.

Remark 2.11. In fact, the proof will show that $h = \mathcal{O}(z^{p+1})$ and $E[|z|] = -\frac{1}{2}q^2 + \mathcal{O}(z^p)$, but we will not need this refinement in what follows. Similarly, we can control Dh and D^2h in the same norms as h , but we will not need this.

Remark 2.12. Using gauge invariance (i.e., differentiating the identity $Q[ze^{i\theta}] = Q[z]e^{i\theta}$) leads to the useful identity

$$Q[z] = -iDQ[z]iz. \quad (2-20)$$

Results similar to [Proposition 2.10](#) are proved in [[Soffer and Weinstein 1990; 1992](#)]; we will sketch a proof that follows the presentation given in the appendix of [[Gustafson et al. 2004](#)]. The key ingredient is the following estimate for the resolvent at the linear eigenvalue.

Lemma 2.13. *For any integer $k \geq 0$, $(H + \frac{1}{2}q^2)^{-1}P_c$ is bounded from L^2 to $H^2(\mathbb{R} \setminus \{0\})$ and from $H^{0,k}$ to $H^{1,k}$.*

Proof. Evaluating the resolvent at $-\frac{1}{2}q^2$, we see that the integral kernel of $(H + \frac{1}{2}q^2)^{-1}P_c$ is a linear combination of terms of the form

$$e^{|x-y|q} \quad \text{and} \quad e^{q(|x|+|y|)}.$$

Terms of the second type are straightforward to handle; one needs only observe that

$$\left| \int e^{q|y|} f(y) dy \right| \lesssim \|f\|_{L^2}$$

and that $e^{q|x|} \in H^2(\mathbb{R} \setminus \{0\}) \cap H^{1,k}$ for any k . It remains to verify that convolution with $e^{q|x|}$ sends L^2 to $H^2(\mathbb{R} \setminus \{0\})$ and $H^{0,k}$ to $H^{1,k}$ for any k . Mapping to $H^2(\mathbb{R} \setminus \{0\})$ is clear, so let us consider a weighted norm. As the derivative of $e^{q|x|}$ still decays exponentially, it is enough to work with $H^{0,k}$. The desired estimate therefore reduces to the fact that the operator with kernel $\langle x \rangle^k e^{q|x-y|} \langle y \rangle^{-k}$ maps $L^2 \rightarrow L^2$ for any k (a consequence of Schur's test, for example). \square

With [Lemma 2.13](#) in place, we turn to the proof of [Proposition 2.10](#).

Proof of Proposition 2.10. We wish to solve

$$(H - E)Q + F(Q) = 0, \quad \text{with} \quad Q = z\phi_0 + h \quad \text{and} \quad E = -\frac{1}{2}q^2 + e$$

for small enough (nonzero) z , where $h = \mathcal{O}(z^2)$ is orthogonal to ϕ_0 and $e = \mathcal{O}(z)$ is real. Expanding the equation and projecting onto and away from ϕ_0 leads to the following system for (e, h) :

$$e = z^{-1} \langle \phi_0, F(z\phi_0 + h) \rangle, \quad (2-21)$$

$$h = (H + \frac{1}{2}q^2)^{-1} \{-P_c F(z\phi_0 + h) + eh\}, \quad (2-22)$$

where z is to be small. To solve this system, let us construct (e, h) as a fixed point of the operator

$$\Phi(e, h) = (\text{RHS (2-21)}, \text{RHS (2-22)}).$$

Let us prove that Φ is a contraction on the set

$$A = \{(e, h) \in \mathbb{R} \times P_c H^1 : |e| \leq |z|, \|h\|_{H^1} \leq |z|^2\},$$

where z will be chosen sufficiently small. We will then prove the desired estimates for h and e as a priori estimates using (2-21) and (2-22).

It is straightforward to show that $\Phi : A \rightarrow A$; indeed, writing $(e_1, h_1) = \Phi(e_0, h_0)$ for some $(e_0, h_0) \in A$, we can use Lemma 2.13 to estimate

$$|e_1| \lesssim |z|^{-1} \|z^{p+1}\phi_0^{p+1} + h_0^{p+1}\|_{L^2} \lesssim |z|^p \ll |z|$$

and

$$\|h_1\|_{H^1} \lesssim \|z^{p+1}\phi_0^{p+1} + h_0^{p+1}\|_{L^2} + \|e_0 h_0\|_{L^2} \lesssim |z|^{p+1} + |z|^3 \ll |z|^2.$$

Similarly, writing $(e_1, h_1) = \Phi(e_0, h_0)$ and $(\tilde{e}_1, \tilde{h}_1) = \Phi(\tilde{e}_0, \tilde{h}_0)$, we can estimate

$$\begin{aligned} |e_1 - \tilde{e}_1| &\lesssim |z|^{-1} \|(h_0 - \tilde{h}_0)(z^p \phi_0^p + h_0^p + \tilde{h}_0^p)\|_{L^2} \\ &\lesssim |z|^{p-1} \|h_0 - \tilde{h}_0\|_{H^1} \ll \|h_0 - \tilde{h}_0\|_{H^1} \end{aligned}$$

and

$$\begin{aligned} \|h_1 - \tilde{h}_1\|_{H^1} &\lesssim |z|^p \|h_0 - \tilde{h}_0\|_{H^1} + |z|^2 |e_0 - \tilde{e}_0| + |z| \|h_0 - \tilde{h}_0\|_{H^1} \\ &\ll \|h_0 - \tilde{h}_0\|_{H^1} + |e_0 - \tilde{e}_0|. \end{aligned}$$

Thus Φ defines a contraction on A (for z small enough) and hence has a unique fixed point.

Using uniqueness and gauge invariance of the nonlinearity, we can deduce that $Q[ze^{i\theta}] = e^{i\theta} Q[z]$ and $E = E[|z|]$. Similarly, by uniqueness we can guarantee that $Q[|z|]$ is real-valued.

Next, let us estimate h in $H^2(\mathbb{R} \setminus \{0\})$ and $H^{1,k}$. Using (2-22), Lemma 2.13, and Sobolev embedding, we first estimate

$$\begin{aligned} \|h\|_{H^2(\mathbb{R} \setminus \{0\})} &\lesssim \|F(z\phi_0 + h) + eh\|_{L^2} \\ &\lesssim \|z^{p+1}\phi_0^{p+1} + h^{p+1}\|_{L^2} + |z| \|h\|_{L^2} \\ &\lesssim |z|^{p+1} + \{|z|^{2p} + |z|\} \|h\|_{L^2}, \end{aligned}$$

which (for small z) implies $\|h\|_{H^2(\mathbb{R} \setminus \{0\})} \lesssim |z|^2$. Similarly,

$$\|h\|_{H^{1,k}} \lesssim |z|^{p+1} + \{|z|^{2p} + |z|\} \|h\|_{H^{0,k}},$$

which again implies $\|h\|_{H^{1,k}} \lesssim |z|^2$.

To prove bounds for Dh , we differentiate (2-21) and (2-22). This leads to

$$De = -z^{-2} D\langle \phi_0, F(z\phi_0 + h) \rangle + z^{-1} \langle \phi_0, D[F(z\phi_0 + h)] \rangle, \quad (2-23)$$

$$Dh = \left(H + \frac{1}{2}q^2\right)^{-1} \{-P_c D[F(z\phi_0 + h)] + [De]h + e[Dh]\}. \quad (2-24)$$

Using (2-23), we can readily deduce that $|De| \lesssim 1$. Feeding this into (2-24) and estimating as above using Lemma 2.13, we find

$$\|Dh\|_{H^{1,k}} \lesssim |z|^p + |z| \|h\|_{H^{0,k}} + |z| \|Dh\|_{H^{0,k}} \lesssim |z|^p + |z|^3 + |z| \|Dh\|_{H^{0,k}},$$

which implies

$$\|Dh\|_{H^{1,k}} \lesssim |z|,$$

as desired. Differentiating (2-23) and (2-24) once more and arguing similarly yields the final estimate, namely,

$$\|D^2h\|_{L^2} \lesssim 1. \quad \square$$

2D. Local well-posedness. In this section we record a local well-posedness result for (1-1). Such results have appeared previously in the literature (e.g., in [Fukuizumi et al. 2008, Proposition 1]); we provide a proof here for the sake of completeness.

Proposition 2.14 (local well-posedness). *For any $u_0 \in H^1$, there exists a local-in-time solution to (1-1). The solution may be extended as long as the H^1 -norm does not blow up.*

Proof. We will look for u decomposed as

$$u(t) = v(t) + a(t)\phi_0 := P_c u(t) + \langle \phi_0, u(t) \rangle \phi_0.$$

Equation (1-1) then becomes a coupled system for $(v(t), a(t))$, namely,

$$\begin{aligned} i\partial_t v(t) &= H v + P_c F(v(s) + a(s)\phi_0), \\ i\partial_t a(t) &= -\frac{1}{2}q^2 a(t) + \langle \phi_0, F(v(t) + a(t)\phi_0) \rangle. \end{aligned} \quad (2-25)$$

Using an integrating factor in (2-25), we may rewrite these as

$$v(t) = e^{-itH} P_c u_0 - i \int_0^t e^{-i(t-s)H} P_c F(v(s) + a(s)\phi_0) ds, \quad (2-26)$$

$$a(t) = e^{iq^2 t/2} a(0) - i \int_0^t e^{iq^2(t-s)/2} \langle \phi_0, F(v(s) + a(s)\phi_0) \rangle ds. \quad (2-27)$$

Defining $\Phi(v, a) = (\text{RHS (2-26)}, \text{RHS (2-27)})$, we will prove that Φ defines a contraction on a suitable complete metric space. Writing $M = \|u_0\|_{H^1}$ and letting $T > 0$ to be chosen below, we define

$$B_T = \{(v, a) : \|v\|_{(L_t^\infty H_x^1 \cap L_t^4 L_x^\infty)([0, T] \times \mathbb{R})} \leq 2CM, \|a\|_{L_t^\infty([0, T])} \leq 2CM\},$$

where C encodes constants appearing in Strichartz estimates. In light of (2-19), we can freely exchange $\langle \sqrt{H} \rangle$ and $\langle \partial_x \rangle$ in what follows.

Writing $(w, b) = \Phi(v, a)$ for some $(v, a) \in B_T$, we first use Proposition 2.1 to estimate

$$\begin{aligned} \|w\|_{(L_t^\infty H_x^1 \cap L_t^4 L_x^\infty)([0, T] \times \mathbb{R})} &\lesssim \|u_0\|_{H^1} + \|F(v + a\phi_0)\|_{L_t^1 H_x^1([0, T] \times \mathbb{R})} \\ &\lesssim M + T \{ \|v\|_{L_{t,x}^\infty([0, T] \times \mathbb{R})}^p \|v\|_{L_t^\infty H_x^1([0, T] \times \mathbb{R})} + \|a\|_{L_t^\infty([0, T])}^{p+1} \|\phi_0\|_{L_x^\infty}^p \|\phi_0\|_{H_x^1} \} \\ &\lesssim M + TM^{p+1}, \end{aligned}$$

while

$$\|b\|_{L_t^\infty([0, T])} \leq |a(0)| + \|F(v + a\phi_0)\|_{L_t^1 L_x^2([0, T] \times \mathbb{R})} \|\phi_0\|_{L_x^2} \lesssim M + TM^{p+1}.$$

Thus, for $T = T(M)$ sufficiently small, Φ maps B_T to B_T . Similar estimates show that Φ is a contraction in the norm

$$d((v, a), (\tilde{v}, \tilde{a})) = \|v - \tilde{v}\|_{L_t^\infty L_x^2([0, T] \times \mathbb{R})} + \|a - \tilde{a}\|_{L_t^\infty([0, T])}$$

for T sufficiently small. The result follows. \square

3. Setting up the problem

Suppose $u : [0, T] \times \mathbb{R} \rightarrow \mathbb{C}$ is a (small) solution to (1-1). We will look for a decomposition of u of the form

$$u(t) = Q[z(t)] + v(t). \quad (3-1)$$

We view $z(t)$ as a small unknown to be specified, with Q a solution to (1-2) (see Proposition 2.10) and $v(t)$ defined through (3-1).

Using (1-1), (1-2), and (2-20), any such decomposition would lead to an evolution equation for v , namely,

$$\begin{aligned} i \partial_t v &= H v + \mathcal{N}, \\ \mathcal{N} &:= F(Q + v) - F(Q) - i D Q (\dot{z} + i E z), \end{aligned} \quad (3-2)$$

where we have written $Q = Q[z(t)]$, $E = E[|z(t)|]$, and \dot{z} denotes the time derivative. We wish to choose $z(t)$ in such a way that the solution to (3-2) is well-behaved (and such that $z(t)$ remains small).

To choose $z(t)$ and thereby fix the decomposition (3-1), we will impose the orthogonality conditions

$$\operatorname{Im} \langle u - Q[z], D_j Q[z] \rangle = 0 \quad \text{for } j \in \{1, 2\}, \quad (3-3)$$

for all $t \in [0, T]$. This condition makes $v = u - Q[z]$ orthogonal to the nondecaying solutions to the linearization of (1-1) around $e^{-iEt} Q[z]$ and agrees with the condition appearing in [Gustafson et al. 2004]. We discuss the motivation for this choice in Remark 3.2 below.

The following lemma tells us that as long as the solution $u(t)$ remains small, it is always possible to choose $z(t)$ such that (3-3) holds; moreover, this choice is unique.

Lemma 3.1. *There exists $\delta > 0$ small enough such that if $\|u\|_{H^1} \leq \delta$, then there exists unique $z \in \mathbb{C}$ such that (3-3) holds, with $|z| + \|u - Q[z]\|_{H^1} \lesssim \|u\|_{H^1}$.*

Proof of Lemma 3.1. The proof is the same as [Gustafson et al. 2004, Lemma 2.3]. The idea is that if we were to choose $v = u - \langle \phi_0, u \rangle \phi_0 = P_c u$, then we would not be too far off from satisfying (3-3). We can therefore use the inverse function theorem to find z exactly satisfying (3-3). This is made precise using Proposition 2.10. We sketch the details.

Set $\varepsilon = \|u\|_{H^1}$. Define $f : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ via

$$f_j(z) = \operatorname{Im} \langle u - Q[z], D_j Q[z] \rangle$$

for $j = 1, 2$, and set $z_0 = \langle \phi_0, u \rangle$. Note that $|z_0| \leq \varepsilon$. A computation using the expansion of $Q[z]$ in Proposition 2.10 yields

$$f(z_0) = \mathcal{O}(\varepsilon^2).$$

Similarly (using Proposition 2.10), the Jacobian of the map $z \mapsto f(z)$ is computed by

$$D_j f_k(z) = \operatorname{Im} \langle u - Q[z], D_j D_k Q[z] \rangle + \operatorname{Im} \langle D_j Q, D_k Q \rangle = j - k + \mathcal{O}(\varepsilon + |z|). \quad (3-4)$$

Therefore, by the inverse function theorem, for ε small enough we may find unique z such that $f(z) = 0$. The result follows. \square

Under the (bootstrap) assumption that $\sup_{t \in [0, T]} \|u(t)\|_{H^1} \leq \delta$ for δ small enough, we can therefore uniquely decompose $u(t)$ in the form (3-1) such that (3-3) holds for each $t \in [0, T]$.

The evolution equation for v is given by (3-2). To derive the evolution equation for z , we differentiate the orthogonality conditions (3-3). Recalling (3-1), (3-2), and self-adjointness of H , this firstly leads to

$$0 = \operatorname{Im}[i\langle v, HD_j Q \rangle + i\langle F(Q + v) - F(Q), D_j Q \rangle - \langle DQ(\dot{z} + iEz), D_j Q \rangle + \langle v, D_j DQ\dot{z} \rangle].$$

Differentiating (1-2) and observing that (2-20) and (3-3) imply $\operatorname{Im} i\langle v, Q \rangle = 0$, we may write

$$\begin{aligned} \operatorname{Im} i\langle v, HD_j Q \rangle &= \operatorname{Im}[i\langle v, ED_j Q \rangle - i\mu\langle v, D_j(|Q|^p Q) \rangle] \\ &= \operatorname{Im}\langle v, D_j DQ iEz \rangle - \operatorname{Im} i\mu\langle v, D_j(|Q|^p Q) \rangle, \end{aligned}$$

where we have used (2-20) again in the final line. Continuing from above, we arrive at the system

$$\begin{aligned} \operatorname{Im}\langle v, D_j DQ(\dot{z} + iEz) \rangle + \operatorname{Im}\langle D_j Q, DQ(\dot{z} + iEz) \rangle \\ = -\operatorname{Im} i[\langle F(Q + v) - F(Q), D_j Q \rangle - \langle v, D_j(|Q|^p Q) \rangle]. \end{aligned} \quad (3-5)$$

The inner product on the right-hand side of (3-5) is of the form $\langle G(v, Q), D_j Q \rangle$, where G is at least quadratic in v (see Section 4A). Identifying \mathbb{C} with \mathbb{R}^2 , we may write this system in the more compact form

$$A(\dot{z} + iEz) = b, \quad (3-6)$$

where A is the 2×2 real matrix with entries

$$A_{jk} = \operatorname{Im}\langle v, D_j D_k Q \rangle + \operatorname{Im}\langle D_j Q, D_k Q \rangle$$

and $b \in \mathbb{R}^2$ satisfies $b_j = \text{RHS (3-5)}$. Note that A coincides with the Jacobian matrix appearing in (3-4), and hence $A_{jk} = j - k + \mathcal{O}(\delta + |z|)$.

3A. Summary. We have set up the problem as follows: assuming that we have a sufficiently small solution u to (1-1) on a time interval $[0, T]$, we choose $z(t)$ uniquely such that (3-3) holds for each t (using Lemma 3.1). Defining $v(t) = u(t) - Q[z(t)]$ (where Q is the solution to (1-2) as in Proposition 2.10), we find that v and z solve the coupled system (3-2) and (3-6).

In the next section we will use these equations to prove bounds for v and z . In particular, this will show that u remains small, which implies that the decomposition for u can be continued for all time. Furthermore, the bounds we obtain will allow us to complete the proof of the main result, Theorem 1.1.

Remark 3.2. Let us discuss in some more detail the orthogonality condition (3-3). We begin by considering the linearization of (1-1) around a fixed solitary wave $e^{-iEt} Q$. Identifying v with the real vector $v = (\operatorname{Re} v, \operatorname{Im} v)^t$, we can write the linearized equation the form $v_t = Lv$ for an explicit real matrix of operators L . Recalling that Q solves (1-2) and employing the identity (2-20), we can connect the functions $D_j Q$ to this linearized equation. In particular (recalling the identification of \mathbb{C} and \mathbb{R}^2), one can compute

$$L^t i D_j Q = -\tilde{E} z_j [z_2(i D_1 Q) - z_1(i D_2 Q)],$$

where L^t denotes the transpose of L and we write $D_j E[|z|] = \tilde{E} z_j$. One therefore finds that the pair $\{i D_j Q\}$ spans the generalized null space of L^t . The orthogonality condition (3-3) is equivalent to the orthogonality

of v (identified with the real vector $(\operatorname{Re} v, \operatorname{Im} v)^t$) to $i D_j Q$ (identified with $(-\operatorname{Im} D_j Q, \operatorname{Re} D_j Q)^t$); here we use the usual inner product for vectors of \mathbb{R} -valued functions, i.e.,

$$(f_1, f_2)^t \cdot (g_1, g_2)^t = \int f_1 g_1 + \int f_2 g_2.$$

This condition projects v away from the nondecaying solutions to $\partial_t v = Lv$, as we now explain. We let $\{w_1, w_2\}$ be a basis for the generalized null space of L^t (denoted by N) satisfying $L^t w_1 = 0$ and $L^t w_2 = w_1$. It is not difficult to check that N^\perp is invariant under the flow $\partial_t v = Lv$. Similarly, for $v(0) \in N$, we can find a solution to $\partial_t v = Lv$ of the form $v(t) = q_1(t)w_1 + q_2(t)w_2$. In fact, explicit computation reveals that q_1 and q_2 are linear functions in t . Thus, (3-3) exactly projects v away from the nondecaying solutions of $\partial_t v = Lv$, and hence we expect that the component v should decay.

At a technical level, the key benefit of imposing (3-3) arises in the computation of the ODE (3-5) for $\dot{z} + iEz$. In particular, imposing (3-3) leads to an ODE for $\dot{z} + iEz$ that contains only quadratic and higher terms in v . This is crucial because to describe the asymptotics of z will require that we estimate $\dot{z} + iEz$ in L_t^1 , while we can only hope to estimate v in spaces as low as L_t^2 (through reversed Strichartz estimates).

In contrast, suppose that we were to impose the natural condition

$$\langle v(t), \phi_0 \rangle = 0, \quad (3-7)$$

so that $v = P_c v$. This type of condition appears in [Pillet and Wayne 1997; Weder 2000] and has the advantage of allowing for Strichartz estimates for $e^{-itH} P_c$ to be applied directly to v . In this case, one would find that the ODE for z contains a term that is *linear* in v , and hence we would have no hope of estimating in L_t^1 .

On the other hand, as $v \neq P_c v$ under the assumption (3-3), we cannot apply Strichartz estimates for e^{-itH} directly to v . However, if we recall the decomposition $Q[z] = z\phi_0 + \mathcal{O}(z^2)$, then we can see that the condition (3-3) implies $\langle v(t), \phi_0 \rangle = \mathcal{O}(z^2)$, which suggests that the portion of v parallel to ϕ_0 should be small compared to v . In fact, in Lemma 4.4 we will prove that we can control v by $P_c v$ in all relevant norms, and hence we will be able to utilize the estimates for $e^{-itH} P_c$ after all.

4. Proof of the main result

We suppose u is a solution to (1-1) satisfying

$$\sup_{t \in [0, T]} \|u(t)\|_{H^1} \leq \delta \quad (4-1)$$

for δ sufficiently small, so that we may take the decomposition

$$u(t) = Q[z(t)] + v(t), \quad \text{where } \operatorname{Im} \langle v(t), D_j Q[z(t)] \rangle \equiv 0 \text{ for } j \in \{1, 2\},$$

as outlined in the previous section. By Lemma 3.1, we also have

$$\sup_{t \in [0, T]} \{|z(t)| + \|v(t)\|_{H^1}\} \lesssim \sup_{t \in [0, T]} \|u(t)\|_{H^1} \lesssim \delta.$$

Our goal is to extend these bounds to $[0, \infty)$ and to describe the asymptotics of $z(t)$ and $v(t)$ as $t \rightarrow \infty$. To accomplish this, we will prove a bootstrap estimate using the following norms, which should all be

taken over $[0, T] \times \mathbb{R}$ or $[0, T]$. We first define

$$\|v\|_X := \|v\|_{L_t^\infty H_x^1 \cap L_t^4 L_x^\infty} + \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} + \|\partial_x v\|_{L_x^\infty L_t^2}, \quad (4-2)$$

$$\|z\|_Y := \|\dot{z} + iEz\|_{L_t^1 \cap L_t^2}. \quad (4-3)$$

Noting that

$$|z(t)| = \left| z(t) \exp \left\{ i \int_0^t E[z(s)] ds \right\} \right|,$$

we observe that

$$\|z\|_{L_t^\infty} \leq |z(0)| + \|z\|_Y. \quad (4-4)$$

As the equation for v involves $Q[z(t)]$, it will be convenient to introduce notation for norms of Q as well. In particular, we define

$$\|Q\|_Z := \|\langle x \rangle^{5/2} Q\|_{L_x^1 L_{t,x}^\infty \cap L_{t,x}^\infty} + \|\partial_x Q\|_{L_{t,x}^\infty \cap L_x^\infty L_t^2}, \quad (4-5)$$

$$\|DQ\|_W := \|\langle x \rangle DQ\|_{L_x^1 L_t^\infty} + \|\langle x \rangle DQ\|_{L_t^\infty L_x^2} + \|\partial_x DQ\|_{L_t^\infty L_x^2}, \quad (4-6)$$

where $Q = Q[z(t)]$. Using [Proposition 2.10](#), we can control these norms as long as $z(t)$ remains sufficiently small.

Lemma 4.1. *If $\|z\|_{L_t^\infty}$ is sufficiently small, then*

$$\|Q\|_Z \lesssim \|z\|_{L_t^\infty} \quad \text{and} \quad \|DQ\|_W \lesssim 1.$$

Proof. We begin with the estimate

$$\|\langle x \rangle^\ell G\|_{L_x^r L_t^\infty} \lesssim \|G\|_{L_t^\infty H_x^{1,k}} \quad \text{for any } 1 \leq r \leq \infty \text{ and } k > \ell + \frac{1}{r}, \quad (4-7)$$

which follows from Hölder's inequality and the Sobolev embedding $H^1(\mathbb{R}) \hookrightarrow L^\infty(\mathbb{R})$. In particular,

$$\|Q\|_Z \lesssim \|Q[z(t)]\|_{L_t^\infty H_x^{1,k} \cap L_t^\infty H_x^2(\mathbb{R} \setminus \{0\})} \quad \text{and} \quad \|DQ\|_W \lesssim \|DQ[z(t)]\|_{L_t^\infty H_x^{1,k}}$$

for large enough k . Here we only use $H^2(\mathbb{R} \setminus \{0\})$ to control $\partial_x Q$ in L^∞ .

The result now follows from [Proposition 2.10](#); indeed, for $\sup_{t \in [0, T]} |z(t)|$ small enough, we can write

$$Q[z(t)] = z(t)\phi_0 + h(z(t)),$$

where $h(z(t)) = \mathcal{O}(|z(t)|^2)$ and $Dh(z(t)) = \mathcal{O}(|z(t)|)$ in the norms detailed in [Proposition 2.10](#). \square

4A. Estimates for the ODE. We first consider the ODE (3-6) for z , which we recall has the form

$$A(\dot{z} + iEz) = b,$$

with $A_{jk} = j - k + \mathcal{O}(\delta + |z|)$ and

$$b_j = -\operatorname{Im} i[(F(Q + v) - F(Q), D_j Q) - \mu \langle v, D_j(|Q|^p Q) \rangle].$$

To get the error bound on A_{jk} , we use [Proposition 2.10](#) (similar to the proof of [Lemma 4.1](#)). In particular, A is invertible with uniformly bounded inverse.

Lemma 4.2. *The following estimate holds:*

$$\|z\|_Y \lesssim \|DQ\|_W \{\|v\|_X^2 \|Q\|_Z^{p-1} + \|v\|_X^{p+1}\}.$$

Proof. We examine the right-hand side of the ODE (3-6) in a more detail. First,

$$D_j(|Q|^p Q) = \frac{p}{2}|Q|^{p-2} Q^2 D_j \bar{Q} + \frac{p+2}{2}|Q|^p D_j Q,$$

while

$$F(Q+v) - F(Q) = \frac{p+2}{2} \mu v \int_0^1 |Q + \theta v|^p d\theta + \frac{p}{2} \mu \bar{v} \int_0^1 |Q + \theta v|^{p-2} (Q + \theta v)^2 d\theta. \quad (4-8)$$

Thus, we may write

$$b_j = -\operatorname{Im} i \langle G(v, Q), D_j Q \rangle,$$

where

$$G(v, Q) := \frac{p+2}{2} \mu v \int_0^1 [|Q + \theta v|^p - |Q|^p] d\theta + \frac{p}{2} \mu \bar{v} \int_0^1 [|Q + \theta v|^{p-2} (Q + \theta v)^2 - |Q|^{p-2} Q^2] d\theta.$$

In particular,

$$|G(v, Q)| = \mathcal{O}(v^2 Q^{p-1} + v^{p+1}). \quad (4-9)$$

Using the above together with Proposition 2.10 and Sobolev embedding, we may now estimate

$$\begin{aligned} \|\dot{z} + iEz\|_{L_t^2} &\lesssim \|\|v\|_{L_x^\infty}^2 \|Q\|_{L_x^\infty}^{p-1} + \|v\|_{L_x^\infty}^{p+1}\|_{L_t^2} \|DQ\|_{L_t^\infty L_x^1} \\ &\lesssim \|DQ\|_W \{\|v\|_{L_t^4 L_x^\infty}^2 \|Q\|_{L_{t,x}^\infty}^{p-1} + \|v\|_{L_t^4 L_x^\infty}^2 \|v\|_{L_{t,x}^\infty H_x^1}^{p-1}\} \\ &\lesssim \|DQ\|_W \{\|v\|_X^2 \|Q\|_Z^{p-1} + \|v\|_X^{p+1}\}, \end{aligned} \quad (4-10)$$

which is acceptable. We next estimate the L_t^1 -norm. Using (4-9), we estimate

$$\begin{aligned} \|\dot{z} + iEz\|_{L_t^1} &\lesssim \int |A^{-1} \operatorname{Im} \langle G(v, Q), DQ \rangle| dt \\ &\lesssim \|G(v, Q) DQ\|_{L_{t,x}^1} \\ &\lesssim \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2}^2 \|\langle x \rangle^{3/(p-1)} Q\|_{L_{t,x}^\infty}^{p-1} \|DQ\|_{L_x^1 L_t^\infty} + \|v\|_{L_t^4 L_x^\infty}^4 \|v\|_{L_{t,x}^\infty}^{p-3} \|DQ\|_{L_t^\infty L_x^1} \\ &\lesssim \|DQ\|_W \{\|v\|_X^2 \|Q\|_Z^{p-1} + \|v\|_X^{p+1}\}, \end{aligned}$$

which is acceptable. □

4B. Estimates for the PDE. We next consider the PDE (3-2) for v .

We will prove the following.

Proposition 4.3. *The following estimate holds:*

$$\|v\|_X \lesssim \|v(0)\|_{H^1} + \|z\|_Y \|DQ\|_W + \|v\|_X \|Q\|_Z^p + \|v\|_X^{p+1}.$$

The plan is to use Strichartz and local smoothing estimates for e^{-itH} . However, we cannot apply these estimates directly to v because the orthogonality conditions (3-3) do *not* imply that v belongs to the continuous spectral subspace of H . Nonetheless, using Proposition 2.10 and (3-3), we can prove that v can be controlled by $P_c v$.

Lemma 4.4. *There exists $\delta > 0$ small enough that the following holds: If $\|z\|_{L_t^\infty} \leq \delta$ and $v \in X$ satisfies the orthogonality condition*

$$\operatorname{Im}\langle v(t), D_j Q[z(t)] \rangle \equiv 0 \quad \text{for } j \in \{1, 2\} \quad (4-11)$$

(where $Q[z]$ is as in Proposition 2.10), then

$$\|v\|_X \lesssim \|P_c v\|_X.$$

Here X is as in (4-2) and P_c denotes the projection onto the continuous spectral subspace of H .

Proof. Writing $v = P_c v + \langle \phi_0, v \rangle \phi_0$, we see that it suffices to prove

$$\|\langle \phi_0, v \rangle \phi_0\|_X \ll \|v\|_X.$$

To this end, we use Proposition 2.10 to write $Q[z(t)] = z(t)\phi_0 + h(z(t))$, with $h(z) = \mathcal{O}(z^2)$ and $Dh(z) = \mathcal{O}(z)$ in the norms detailed in Proposition 2.10. As (4-11) yields

$$|\langle \phi_0, v(t) \rangle| \lesssim |\langle Dh, v(t) \rangle|,$$

we can estimate

$$\|\langle \phi_0, v \rangle \phi_0\|_X \lesssim \|\langle Dh, v(t) \rangle\|_{L_t^2 \cap L_t^\infty}.$$

We now claim that

$$\|\langle Dh, v(t) \rangle\|_{L_t^2 \cap L_t^\infty} \lesssim \|z\|_{L_t^\infty} \|v\|_X, \quad (4-12)$$

from which the result follows. To see this, first note that by the triangle inequality and Minkowski's inequality, we have

$$\|\langle Dh, v(t) \rangle\|_{L_t^2} \lesssim \|Dh v(t)\|_{L_t^2 L_x^1} \lesssim \|Dh v(t)\|_{L_x^1 L_t^2} \lesssim \|\langle x \rangle^{3/2} Dh\|_{L_x^1 L_t^\infty} \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2}.$$

Using (4-7), we see that this term is acceptable. Next,

$$\|\langle Dh, v(t) \rangle\|_{L_t^\infty} \lesssim \|Dh\|_{L_t^\infty L_x^2} \|v\|_{L_t^\infty L_x^2},$$

which is acceptable as well. The result follows. \square

Using Lemma 4.4, we see that it suffices to estimate the X -norm of $P_c v$. Applying P_c to (3-2), we have

$$i \partial_t P_c v = H P_c v + P_c \mathcal{N},$$

where we recall

$$\mathcal{N} = F(Q + v) - F(Q) - i D Q(\dot{z} + i E z).$$

In particular,

$$P_c v(t) = e^{-itH} P_c v(0) - i \int_0^t e^{-i(t-s)H} P_c \mathcal{N} ds. \quad (4-13)$$

We begin with the linear evolution term.

Lemma 4.5. *The following bound holds:*

$$\|e^{-itH} P_c v(0)\|_X \lesssim \|v(0)\|_{H^1}.$$

Proof. Recalling the definition of the X -norm in (4-2), we find that the lemma follows from Proposition 2.1, Proposition 2.3, and (2-19). \square

We turn to the Strichartz norms for the inhomogeneous term.

Lemma 4.6. *The following bound holds:*

$$\left\| \int_0^t e^{-i(t-s)H} P_c \mathcal{N} ds \right\|_{L_t^\infty L_x^2 \cap L_t^4 L_x^\infty} \lesssim \|z\|_Y \|DQ\|_W + \|v\|_X \|Q\|_Z^p + \|v\|_X^{p+1}.$$

Proof. Using Corollary 2.5 we first estimate

$$\begin{aligned} \left\| \int_0^t e^{-i(t-s)H} P_c [DQ(\dot{z} + iEz)] ds \right\|_{L_t^\infty L_x^2 \cap L_t^4 L_x^\infty} &\lesssim \|\langle x \rangle^{5/2} DQ(\dot{z} + iEz)\|_{L_{t,x}^2} \\ &\lesssim \|\langle x \rangle^{5/2} DQ\|_{L_t^\infty L_x^2} \|\dot{z} + iEz\|_{L_t^2} \\ &\lesssim \|DQ\|_W \|z\|_Y, \end{aligned}$$

which is acceptable.

Next we write nonlinear term in the form

$$F(Q + v) - F(Q) = F_1 + F_2 + F_3,$$

where

$$F_1 = \mathcal{O}(vQ^p), \quad F_2 = \mathcal{O}(v^2Q^{p-1} + v^pQ), \quad \text{and} \quad F_3 = \mu|v|^p v. \quad (4-14)$$

Such a decomposition is easily achieved under the assumption that $F(u) = \mu|u|^p u$ with p equal to an even integer greater than or equal to 4.

The linear term is handled as follows. Using Corollary 2.5, we have

$$\begin{aligned} \left\| \int_0^t e^{-i(t-s)H} P_c F_1 ds \right\|_{L_t^\infty L_x^2 \cap L_t^4 L_x^\infty} &\lesssim \|\langle x \rangle^{5/2} Q^p v\|_{L_{t,x}^2} \\ &\lesssim \|\langle x \rangle^{4/p} Q\|_{L_x^{2p} L_t^\infty}^p \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} \lesssim \|Q\|_Z^p \|v\|_X, \end{aligned}$$

which is acceptable.

Next, we use Proposition 2.1 to estimate

$$\left\| \int_0^t e^{-i(t-s)H} P_c F_2 ds \right\|_{L_t^\infty L_x^2 \cap L_t^4 L_x^\infty} \lesssim \|v^2 Q^{p-1}\|_{L_{t,x}^{6/5}} + \|v^p Q^2\|_{L_{t,x}^{6/5}}.$$

Using Minkowski's inequality to control $L_x^\infty L_t^4$ by $L_t^4 L_x^\infty$, we firstly estimate

$$\begin{aligned} \|v^2 Q^{p-1}\|_{L_{t,x}^{6/5}} &\lesssim \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2}^{4/3} \|v\|_{L_x^\infty L_t^2}^{2/3} \|\langle x \rangle^{2/(p-1)} Q\|_{L_x^{6(p-1)/5} L_t^\infty}^{p-1} \\ &\lesssim \|v\|_X^2 \|Q\|_Z^{p-1}, \end{aligned}$$

which (after an application of Young's inequality) is acceptable. The other term is treated similarly:

$$\begin{aligned} \|v^p Q\|_{L_{t,x}^{6/5}} &\lesssim \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2}^{4/3} \|v\|_{L_t^4 L_x^\infty}^{2/3} \|v\|_{L_{t,x}^\infty}^{p-2} \|\langle x \rangle^2 Q\|_{L_x^{6/5} L_t^\infty} \\ &\lesssim \|v\|_X^p \|Q\|_Z, \end{aligned}$$

which is again acceptable after applying Young's inequality.

Finally, the contribution of the F_3 term containing only v is estimated as follows: using [Proposition 2.1](#),

$$\begin{aligned} \left\| \int_0^t e^{-i(t-s)H} P_c(|v|^p v) ds \right\|_{L_t^\infty L_x^2 \cap L_t^4 L_x^\infty} &\lesssim \| |v|^p v \|_{L_t^{4/3} L_x^1} \\ &\lesssim \|v\|_{L_t^4 L_x^\infty}^3 \|v\|_{L_t^\infty L_x^{p-2}}^{p-2} \lesssim \|v\|_{L_t^4 L_x^\infty}^3 \|v\|_{L_t^\infty H_x^1}^{p-2} \lesssim \|v\|_X^{p+1}, \end{aligned}$$

which is acceptable. This completes the proof of [Lemma 4.6](#). \square

We next consider the $L_t^\infty \dot{H}_x^1$ norm of v . We treat this term by an energy estimate. We will make use of [Lemma 2.8](#).

Lemma 4.7. *The following estimate holds uniformly over $t \in [0, T]$:*

$$\|P_c v(t)\|_{\dot{H}^1}^2 \leq \|v(0)\|_{\dot{H}^1}^2 + \|v\|_X \|z\|_Y \|DQ\|_W + \|v\|_X^2 \|Q\|_Z^p + \|v\|_X^{p+2},$$

where norms are taken over $[0, t] \times \mathbb{R}$.

Proof. By (2-19), we have

$$\|P_c v(t)\|_{\dot{H}^1} \lesssim \|\sqrt{H} P_c v(t)\|_{L^2}^2 + \|v(t)\|_{L^2}^2.$$

As the $L_t^\infty L_x^2$ norm is controlled via [Lemma 4.6](#), it suffices to estimate $\sqrt{H} P_c v$.

To this end, we use the self-adjointness of H and (3-2) to write

$$\|\sqrt{H} P_c v(t)\|_{L_x^2}^2 = \|\sqrt{H} P_c v(0)\|_{L^2}^2 + \operatorname{Im} \int_0^t \langle \sqrt{H} P_c v(s), \sqrt{H} P_c \mathcal{N} \rangle ds,$$

where

$$\mathcal{N} = DQ(\dot{z} + iEz) + F_1 + F_2 + F_3$$

as in (4-14). In fact, we will split the term F_2 (which collects the terms of orders $v^2 Q^{p-1}$ through $v^p Q$) further by writing

$$F_2 = F_2^1 + F_2^2,$$

where F_2^1 collects terms that are linear in Q . We do this so that we can group this term with those appearing in (4-17) below (rather than (4-16)). This is necessary because when the derivative lands on Q we cannot additionally absorb weights in order to produce a $\langle x \rangle^{-3/2} v$ term in $L_x^\infty L_t^2$; indeed, we only control $\partial_x Q$ in $L_{t,x}^\infty$. Thus we must put the whole term in $L_t^1 L_x^2$; see (4-18) below.

We first observe that by (2-19), we have

$$\|\sqrt{H} P_c v(0)\|_{L^2}^2 \lesssim \|v(0)\|_{\dot{H}^1}^2,$$

which is acceptable.

We next use [Lemma 2.8](#) to write

$$\int_0^t \langle \sqrt{H} P_c v(s) \sqrt{H} P_c \mathcal{N} \rangle ds = \int_0^t \langle \partial_x v(s), \partial_x \mathcal{N} \rangle ds + \int_0^t \langle m(\partial_x) \partial_x v(s), \partial_x \mathcal{N} \rangle ds, \quad (4-15)$$

where $m(\mu) = (q - i\mu)^{-1}$ (up to the addition of similar terms). We claim that both terms in (4-15) may be controlled by

$$\|\partial_x v\|_{L_x^\infty L_t^2} \|\partial_x (F_1 + F_2^2)\|_{L_x^1 L_t^2} \quad (4-16)$$

$$+ \|\partial_x v\|_{L_x^\infty L_t^2} \|\partial_x [DQ(\dot{z} + iEz) + F_2^1 + F_3]\|_{L_t^1 L_x^2}. \quad (4-17)$$

For the first term in (4-15), this follows directly from Hölder's inequality. For the second term in (4-15), we use Hölder's inequality and the fact that $m(\partial_x)$ maps $L_x^\infty L_t^2 \rightarrow L_x^\infty L_t^2$ and $L^2 \rightarrow L^2$ boundedly (see [Remark 2.9](#)).

We turn to estimating the terms in (4-16) and (4-17).

We begin with (4-16). First, by the chain rule:

$$\begin{aligned} \|\partial_x F_1\|_{L_x^1 L_t^2} &\lesssim \|\partial_x v\|_{L_x^\infty L_t^2} \|Q\|_{L_x^p L_t^\infty}^p + \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} \|\langle x \rangle^{3/(2(p-1))} Q\|_{L_x^{p-1} L_t^\infty}^{p-1} \|\partial_x Q\|_{L_{t,x}^\infty} \\ &\lesssim \|v\|_X \|Q\|_Z^p, \end{aligned}$$

which is acceptable.

We turn to the intermediate terms in F_2^2 , which contains terms of the order $v^2 Q^{p-1}$ through $v^{p-1} Q^2$. Applying the chain and product rule and Young's inequality, we are led to estimate four types of terms in $L_x^1 L_t^2$ corresponding to these two extreme cases. When the derivative lands on a copy of v , we estimate

$$\begin{aligned} \|(\partial_x v) v Q^{p-1}\|_{L_x^1 L_t^2} &\lesssim \|\partial_x v\|_{L_x^\infty L_t^2} \|v\|_{L_{t,x}^\infty} \|Q\|_{L_x^{p-1} L_t^\infty}^{p-1}, \\ \|(\partial_x v) v^{p-2} Q^2\|_{L_x^1 L_t^2} &\lesssim \|\partial_x v\|_{L_x^\infty L_t^2} \|v\|_{L_{t,x}^\infty}^{p-1} \|Q\|_{L_x^2 L_t^\infty}^2, \end{aligned}$$

which are acceptable. When the derivative lands on a copy of Q , we instead estimate

$$\begin{aligned} \|v^2 Q^{p-2} \partial_x Q\|_{L_x^1 L_t^2} &\lesssim \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} \|v\|_{L_{t,x}^\infty} \|\langle x \rangle^{3/(2(p-2))} Q\|_{L_x^{p-2} L_t^\infty}^{p-2} \|\partial_x Q\|_{L_{t,x}^\infty}, \\ \|v^{p-1} Q \partial_x Q\|_{L_x^1 L_t^2} &\lesssim \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} \|v\|_{L_{t,x}^\infty}^{p-2} \|\langle x \rangle^{3/2} Q\|_{L_x^1 L_t^\infty} \|\partial_x Q\|_{L_{t,x}^\infty}, \end{aligned}$$

which are acceptable.

We turn to (4-17). We first have

$$\begin{aligned} \|\partial_x v\|_{L_t^\infty L_x^2} \|\partial_x [DQ(\dot{z} + iEz)]\|_{L_t^1 L_x^2} &\lesssim \|\partial_x v\|_{L_t^\infty L_x^2} \|\partial_x DQ\|_{L_t^\infty L_x^2} \|\dot{z} + iEz\|_{L_t^1} \\ &\lesssim \|v\|_X \|z\|_Y \|DQ\|_W, \end{aligned}$$

which is acceptable.

Next, we estimate the contribution of F_2^1 in (4-17), which contains terms that are linear in Q . Distributing the derivative, we are led to estimate the following terms. First,

$$\|v^p \partial_x Q\|_{L_t^1 L_x^2} \lesssim \|v\|_{L_t^4 L_x^\infty}^4 \|v\|_{L_{t,x}^\infty}^{p-4} \|\partial_x Q\|_{L_t^\infty L_x^2} \lesssim \|v\|_X^p \|Q\|_Z, \quad (4-18)$$

which is acceptable. Next,

$$\begin{aligned} \|v^{p-1}(\partial_x v)Q\|_{L_t^1 L_x^2} &\lesssim \|v\|_{L_t^4 L_x^\infty}^2 \|v^{p-3}(\partial_x v)Q\|_{L_{t,x}^2} \\ &\lesssim \|v\|_{L_t^4 L_x^\infty}^2 \|\partial_x v\|_{L_x^\infty L_t^2} \|v\|_{L_{t,x}^\infty}^{p-3} \|Q\|_{L_x^2 L_t^\infty} \lesssim \|v\|_X^p \|Q\|_Z, \end{aligned}$$

which is acceptable.

It remains to estimate the contribution of F_3 in (4-17). The purely nonlinear term $F_3 = \mu|v|^p v$ is estimated as

$$\|\partial_x v\|_{L_t^\infty L_x^2} \|\partial_x(|v|^p v)\|_{L_t^1 L_x^2} \lesssim \|v\|_{L_t^4 L_x^\infty}^4 \|\partial_x v\|_{L_t^\infty L_x^2}^2 \|v\|_{L_{t,x}^\infty}^{p-4} \lesssim \|v\|_X^{p+2},$$

which is acceptable. □

It remains to estimate the contribution of the inhomogeneous Duhamel term to the $L_x^\infty L_t^2$ components of the X -norm (see (4-2)). The key ingredients will be [Proposition 2.6](#) and [Corollary 2.7](#).

Lemma 4.8. *The following estimates hold: For $m \in \{\langle x \rangle^{-3/2}, \partial_x\}$,*

$$\left\| m \int_0^t e^{-i(t-s)H} P_c \mathcal{N} ds \right\|_{L_x^\infty L_t^2} \lesssim \|DQ\|_W \|z\|_Y + \|v\|_X \|Q\|_Z^p + \|v\|_X^{p+1}.$$

Proof. We recall that

$$\mathcal{N} = DQ(\dot{z} + iEz) + F_1 + F_2 + F_3,$$

where F_j are as in (4-14).

We first use [Proposition 2.6](#) to estimate

$$\begin{aligned} \left\| m \int_0^t e^{-i(t-s)H} P_c [DQ(\dot{z} + iEz)] ds \right\|_{L_x^\infty L_t^2} &\lesssim \|\langle x \rangle DQ(\dot{z} + iEz)\|_{L_x^1 L_t^2} \\ &\lesssim \|\langle x \rangle DQ\|_{L_x^1 L_t^\infty} \|\dot{z} + iEz\|_{L_t^2} \lesssim \|DQ\|_W \|z\|_Y, \end{aligned}$$

which is acceptable.

Next, we estimate

$$\begin{aligned} \left\| m \int_0^t e^{-i(t-s)H} P_c F_1 ds \right\|_{L_x^\infty L_t^2} &\lesssim \|\langle x \rangle Q^p v\|_{L_x^1 L_t^2} \\ &\lesssim \|\langle x \rangle^{5/(2p)} Q\|_{L_x^p L_t^\infty}^p \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} \lesssim \|Q\|_Z^p \|v\|_X, \end{aligned}$$

which is acceptable.

The contribution of F_2 is estimated by

$$\begin{aligned} \left\| m \int_0^t e^{-i(t-s)H} P_c F_2 ds \right\|_{L_x^\infty L_t^2} &\lesssim \|\langle x \rangle v^2 Q^{p-1}\|_{L_x^1 L_t^2} + \|\langle x \rangle v^p Q\|_{L_x^1 L_t^2} \\ &\lesssim \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} \|v\|_{L_{t,x}^\infty} \|\langle x \rangle^{5/(2(p-1))} Q\|_{L_x^{p-1} L_t^\infty}^{p-1} + \|\langle x \rangle^{-3/2} v\|_{L_x^\infty L_t^2} \|v\|_{L_{t,x}^\infty}^{p-1} \|\langle x \rangle^{5/2} Q\|_{L_x^1 L_t^\infty} \\ &\lesssim \|v\|_X^2 \|Q\|_Z^{p-1} + \|v\|_X^p \|Q\|_Z, \end{aligned}$$

which is acceptable (after an application of Young's inequality).

Finally, we use [Corollary 2.7](#) to estimate

$$\begin{aligned} \left\| m \int_0^t e^{-i(t-s)H} P_c F_3 ds \right\|_{L_x^\infty L_t^2} &\lesssim \| |v|^p v \|_{L_t^1 H_x^{1/2}} \\ &\lesssim \| v \|_{L_t^4 L_x^\infty}^4 \| v \|_{L_{t,x}^\infty}^{p-4} \| v \|_{L_t^\infty H_x^1} \lesssim \| v \|_X^{p+1}, \end{aligned}$$

which is acceptable. This completes the proof of [Lemma 4.8](#). \square

Finally, using [Lemmas 4.4, 4.5, 4.6, 4.7, and 4.8](#) we complete the proof of [Proposition 4.3](#).

4C. Completing the proof. In this section, we first use the estimates of the previous two sections in order to close a bootstrap estimate, which allows us to continue the decomposition of u for all time, as well as to prove the desired properties for $z(t)$ and $v(t)$ and hence complete the proof of [Theorem 1.1](#).

We let $u(t)$ be the solution to (1-1) with initial data u_0 , where $\|u_0\|_{H^1} = \delta$ for some small $\delta > 0$. By local well-posedness and [Lemma 3.1](#), we can uniquely decompose u as

$$u(t) = Q[z(t)] + v(t), \quad \text{with } \text{Im}\langle v(t), D_j Q[z(t)] \rangle \equiv 0 \text{ for } j \in \{1, 2\}, \quad (4-19)$$

at least on some time interval, with $|z(t)| + \|v(t)\|_{H^1} \lesssim \|u(t)\|_{H^1} \lesssim \delta$. On such an interval, we can now collect the estimates from the previous section. Collecting [Lemma 4.1, \(4-4\), Lemma 4.2, and Proposition 4.3](#), we have

$$\|z\|_{L_t^\infty} \lesssim \delta \implies \|Q\|_Z \lesssim \|z\|_{L_t^\infty} \text{ and } \|DQ\|_W \lesssim 1, \quad (4-20)$$

$$\|z\|_{L_t^\infty} \leq |z(0)| + \|z\|_Y,$$

$$\|z\|_Y \lesssim \|DQ\|_W \{ \|v\|_X^2 \|Q\|_Z^{p-1} + \|v\|_X^{p+1} \}, \quad (4-21)$$

$$\|v\|_X \lesssim \|v(0)\|_{H^1} + \|z\|_Y \|DQ\|_W + \|v\|_X \|Q\|_Z^p + \|v\|_X^{p+1}. \quad (4-22)$$

By a standard bootstrap argument (choosing δ small), it follows that the bounds

$$\|u(t)\|_{H^1} \lesssim \delta, \quad \|v\|_X \lesssim \delta, \quad \|z\|_{L_t^\infty} \lesssim \delta, \quad \text{and} \quad \|z\|_Y \lesssim \delta^2,$$

as well as the decomposition (4-19), persist for all time.

We turn to establishing the asymptotics $v(t)$ and $z(t)$.

First, we prove scattering in H^1 for $v(t)$. We claim that it suffices to prove scattering for $P_c v(t)$. Writing $v = P_c v + \langle \phi_0, v \rangle \phi_0$, the claim reduces to proving

$$\lim_{t \rightarrow \infty} \|\langle \phi_0, v(t) \rangle \phi_0\|_{H^1} = 0. \quad (4-23)$$

Proof of (4-23). Using the orthogonality conditions in (4-19) and using [Proposition 2.10](#) to write $Q[z(t)] = z(t)\phi_0 + h(z(t))$ (as in the proof of [Lemma 4.4](#)), we find

$$\|\langle \phi_0, v(t) \rangle \phi_0\|_{H^1} \lesssim \|Dh(z(t))\|_{L_x^{4/3}} \|v(t)\|_{L_x^4}.$$

As

$$\|Dh\|_{L_t^\infty L_x^{4/3}} \lesssim \|z\|_{L_t^\infty},$$

it suffices to prove that $\|v(t)\|_{L_x^4} \rightarrow 0$ as $t \rightarrow \infty$. To see this, we firstly observe (by interpolation of $L_t^\infty L_x^2$ and $L_t^4 L_x^\infty$) that $\|v(t)\|_{L_x^4}^4 \in L_t^2$. We will now show that $\partial_t \|v(t)\|_{L_x^4}^4$ is bounded, which implies the desired result. Using (3-2) for v and Lemma 2.8 (writing $Hv = HP_c v - q^2 \phi_0 \langle \phi_0, v \rangle$), we can firstly estimate

$$\begin{aligned} \partial_t \|v(t)\|_{L^4}^4 &\lesssim \|v\|_{L_{t,x}^\infty}^2 \|\partial_x v\|_{L_t^\infty L_x^2}^2 + \|v\|_{L_t^\infty L_x^3}^3 \|v\|_{L_t^\infty L_x^2} \|\phi_0\|_{H^1}^2 \\ &\quad + \|v\|_{L_t^\infty L_x^6}^3 \|F(Q+v) - F(Q)\|_{L_t^\infty L_x^2} + \|v\|_{L_t^\infty L_x^6}^3 \|DQ\|_{L_t^\infty L_x^2} \|\dot{z} + iEz\|_{L_t^\infty} \end{aligned}$$

uniformly in t . Using the bounds on v and $Q[z]$, we see the proof boils down to controlling $\dot{z} + iEz$ in L_t^∞ . For this, we go back to the ODE (3-6) and use the computations at the beginning of Lemma 4.2 to bound

$$\begin{aligned} \|\dot{z} + iEz\|_{L_t^\infty} &\lesssim \|(v^2 Q^{p-1} + v^{p+1})DQ\|_{L_t^\infty L_x^1} \\ &\lesssim \|DQ\|_{L_t^\infty L_x^2} \|v\|_{L_t^\infty L_x^4}^2 \{\|Q\|_{L_{t,x}^\infty}^{p-1} + \|v\|_{L_{t,x}^\infty}^{p-1}\}. \end{aligned}$$

This completes the proof of (4-23). □

It finally remains to prove scattering for $P_c v(t)$. For this we use the Duhamel formula (4-13) to show that $\{e^{itH} P_c v(t)\}$ is Cauchy in H^1 . Indeed, using the estimates from Lemmas 4.6 and 4.7, we can deduce

$$\|e^{itH} P_c v(t) - e^{isH} P_c v(s)\|_{H^1} \lesssim \|z\|_Y \|DQ\|_W + \|v\|_X \|Q\|_Z^p + \|v\|_X^{p+1},$$

where now the norms on the right-hand side are restricted to (s, t) (and not all of the components of the X -norm are L^∞ in time). Sending $s, t \rightarrow \infty$ yields the claim.

Finally, we note that $\|\dot{z} + iEz\|_{L_t^1} \lesssim \delta^2$ yields the desired bounds and asymptotics for z . This completes the proof of Theorem 1.1.

Acknowledgements

Masaki was supported by JSPS KAKENHI, grant numbers 17K14219, 17H02854, and 17H02851. Segata is partially supported by JSPS KAKENHI, grant number 17H02851. We are grateful to T. Inui for pointing out an error in the original proof of Proposition 2.3, as well as to the anonymous referee for helpful comments.

References

- [Albeverio et al. 1988] S. Albeverio, F. Gesztesy, R. Høegh-Krohn, and H. Holden, *Solvable models in quantum mechanics*, Springer, 1988. [MR](#) [Zbl](#)
- [Banica and Visciglia 2016] V. Banica and N. Visciglia, “Scattering for NLS with a delta potential”, *J. Differential Equations* **260**:5 (2016), 4410–4439. [MR](#) [Zbl](#)
- [Buslaev and Perelman 1992] V. S. Buslaev and G. S. Perelman, “Scattering for the nonlinear Schrödinger equation: states that are close to a soliton”, *Algebra i Analiz* **4**:6 (1992), 63–102. In Russian; translated in *St. Petersburg Math. J.* **4**:6 (1993), 1111–1142. [MR](#) [Zbl](#)
- [Christ and Kiselev 2001] M. Christ and A. Kiselev, “Maximal functions associated to filtrations”, *J. Funct. Anal.* **179**:2 (2001), 409–425. [MR](#) [Zbl](#)
- [Cuccagna 2011] S. Cuccagna, “The Hamiltonian structure of the nonlinear Schrödinger equation and the asymptotic stability of its ground states”, *Comm. Math. Phys.* **305**:2 (2011), 279–331. [MR](#) [Zbl](#)

- [Cuccagna 2014] S. Cuccagna, “On asymptotic stability of moving ground states of the nonlinear Schrödinger equation”, *Trans. Amer. Math. Soc.* **366**:6 (2014), 2827–2888. [MR](#) [Zbl](#)
- [Cuccagna and Pelinovsky 2014] S. Cuccagna and D. E. Pelinovsky, “The asymptotic stability of solitons in the cubic NLS equation on the line”, *Appl. Anal.* **93**:4 (2014), 791–822. [MR](#) [Zbl](#)
- [Datchev and Holmer 2009] K. Datchev and J. Holmer, “Fast soliton scattering by attractive delta impurities”, *Comm. Partial Differential Equations* **34**:7-9 (2009), 1074–1113. [MR](#) [Zbl](#)
- [Deift and Park 2011] P. Deift and J. Park, “Long-time asymptotics for solutions of the NLS equation with a delta potential and even initial data”, *Int. Math. Res. Not.* **2011**:24 (2011), 5505–5624. [MR](#) [Zbl](#)
- [Duchêne et al. 2011] V. Duchêne, J. L. Marzuola, and M. I. Weinstein, “Wave operator bounds for one-dimensional Schrödinger operators with singular potentials and applications”, *J. Math. Phys.* **52**:1 (2011), art. id. 013505. [MR](#) [Zbl](#)
- [Fukuizumi et al. 2008] R. Fukuizumi, M. Ohta, and T. Ozawa, “Nonlinear Schrödinger equation with a point defect”, *Ann. Inst. H. Poincaré Anal. Non Linéaire* **25**:5 (2008), 837–845. [MR](#) [Zbl](#)
- [Goodman et al. 2004] R. H. Goodman, P. J. Holmes, and M. I. Weinstein, “Strong NLS soliton-defect interactions”, *Phys. D* **192**:3-4 (2004), 215–248. [MR](#) [Zbl](#)
- [Gustafson et al. 2004] S. Gustafson, K. Nakanishi, and T.-P. Tsai, “Asymptotic stability and completeness in the energy space for nonlinear Schrödinger equations with small solitary waves”, *Int. Math. Res. Not.* **2004**:66 (2004), 3559–3584. [MR](#) [Zbl](#)
- [Holmer and Zworski 2007] J. Holmer and M. Zworski, “Slow soliton interaction with delta impurities”, *J. Mod. Dyn.* **1**:4 (2007), 689–718. [MR](#) [Zbl](#)
- [Holmer and Zworski 2009] J. Holmer and M. Zworski, “Breathing patterns in nonlinear relaxation”, *Nonlinearity* **22**:6 (2009), 1259–1301. [MR](#) [Zbl](#)
- [Holmer et al. 2007a] J. Holmer, J. Marzuola, and M. Zworski, “Fast soliton scattering by delta impurities”, *Comm. Math. Phys.* **274**:1 (2007), 187–216. [MR](#) [Zbl](#)
- [Holmer et al. 2007b] J. Holmer, J. Marzuola, and M. Zworski, “Soliton splitting by external delta potentials”, *J. Nonlinear Sci.* **17**:4 (2007), 349–367. [MR](#) [Zbl](#)
- [Ikeda and Inui 2017] M. Ikeda and T. Inui, “Global dynamics below the standing waves for the focusing semilinear Schrödinger equation with a repulsive Dirac delta potential”, *Anal. PDE* **10**:2 (2017), 481–512. [MR](#) [Zbl](#)
- [Kaminaga and Ohta 2009] M. Kaminaga and M. Ohta, “Stability of standing waves for nonlinear Schrödinger equation with attractive delta potential and repulsive nonlinearity”, *Saitama Math. J.* **26** (2009), 39–48. [MR](#) [Zbl](#)
- [Kenig et al. 1991] C. E. Kenig, G. Ponce, and L. Vega, “Well-posedness of the initial value problem for the Korteweg–de Vries equation”, *J. Amer. Math. Soc.* **4**:2 (1991), 323–347. [MR](#) [Zbl](#)
- [Kirr and Mızrak 2009] E. Kirr and O. Mızrak, “Asymptotic stability of ground states in 3D nonlinear Schrödinger equation including subcritical cases”, *J. Funct. Anal.* **257**:12 (2009), 3691–3747. [MR](#) [Zbl](#)
- [Kirr and Zarnescu 2009] E. Kirr and A. Zarnescu, “Asymptotic stability of ground states in 2D nonlinear Schrödinger equation including subcritical cases”, *J. Differential Equations* **247**:3 (2009), 710–735. [MR](#) [Zbl](#)
- [Masaki et al. 2019] S. Masaki, J. Murphy, and J.-I. Segata, “Modified scattering for the one-dimensional cubic NLS with a repulsive delta potential”, *Int. Math. Res. Not. IMRN* **2019**:24 (2019), 7577–7603. [MR](#) [Zbl](#)
- [Mizumachi 2008] T. Mizumachi, “Asymptotic stability of small solitary waves to 1D nonlinear Schrödinger equations with potential”, *J. Math. Kyoto Univ.* **48**:3 (2008), 471–497. [MR](#) [Zbl](#)
- [Pillet and Wayne 1997] C.-A. Pillet and C. E. Wayne, “Invariant manifolds for a class of dispersive, Hamiltonian, partial differential equations”, *J. Differential Equations* **141**:2 (1997), 310–326. [MR](#) [Zbl](#)
- [Segata 2015] J.-I. Segata, “Final state problem for the cubic nonlinear Schrödinger equation with repulsive delta potential”, *Comm. Partial Differential Equations* **40**:2 (2015), 309–328. [MR](#) [Zbl](#)
- [Soffer and Weinstein 1990] A. Soffer and M. I. Weinstein, “Multichannel nonlinear scattering for nonintegrable equations”, *Comm. Math. Phys.* **133**:1 (1990), 119–146. [MR](#) [Zbl](#)
- [Soffer and Weinstein 1992] A. Soffer and M. I. Weinstein, “Multichannel nonlinear scattering for nonintegrable equations, II: The case of anisotropic potentials and data”, *J. Differential Equations* **98**:2 (1992), 376–390. [MR](#) [Zbl](#)

- [Soffer and Weinstein 2004] A. Soffer and M. I. Weinstein, “Selection of the ground state for nonlinear Schrödinger equations”, *Rev. Math. Phys.* **16**:8 (2004), 977–1071. [MR](#) [Zbl](#)
- [Tsai and Yau 2002a] T.-P. Tsai and H.-T. Yau, “Asymptotic dynamics of nonlinear Schrödinger equations: resonance-dominated and dispersion-dominated solutions”, *Comm. Pure Appl. Math.* **55**:2 (2002), 153–216. [MR](#) [Zbl](#)
- [Tsai and Yau 2002b] T.-P. Tsai and H.-T. Yau, “Classification of asymptotic profiles for nonlinear Schrödinger equations with small initial data”, *Adv. Theor. Math. Phys.* **6**:1 (2002), 107–139. [MR](#) [Zbl](#)
- [Tsai and Yau 2002c] T.-P. Tsai and H.-T. Yau, “Relaxation of excited states in nonlinear Schrödinger equations”, *Int. Math. Res. Not.* **2002**:31 (2002), 1629–1673. [MR](#) [Zbl](#)
- [Tsai and Yau 2002d] T.-P. Tsai and H.-T. Yau, “Stable directions for excited states of nonlinear Schrödinger equations”, *Comm. Partial Differential Equations* **27**:11-12 (2002), 2363–2402. [MR](#) [Zbl](#)
- [Weder 2000] R. Weder, “Center manifold for nonintegrable nonlinear Schrödinger equations on the line”, *Comm. Math. Phys.* **215**:2 (2000), 343–356. [MR](#) [Zbl](#)

Received 31 Jul 2018. Revised 28 Jan 2019. Accepted 18 Apr 2019.

SATOSHI MASAKI: masaki@sigmath.es.osaka-u.ac.jp

Department of Systems Innovation, Graduate School of Engineering Science, Osaka University, Osaka, Japan

JASON MURPHY: jason.murphy@mst.edu

Department of Mathematics and Statistics, Missouri University of Science and Technology, Rolla, MO, United States

JUN-ICHI SEGATA: segata@m.tohoku.ac.jp

Mathematical Institute, Tohoku University, Sendai, Japan

Analysis & PDE

msp.org/apde

EDITORS

EDITOR-IN-CHIEF

Patrick Gérard

patrick.gerard@math.u-psud.fr

Université Paris Sud XI

Orsay, France

BOARD OF EDITORS

Massimiliano Berti	Scuola Intern. Sup. di Studi Avanzati, Italy berti@sissa.it	Gilles Pisier	Texas A&M University, and Paris 6 pisier@math.tamu.edu
Michael Christ	University of California, Berkeley, USA mchrist@math.berkeley.edu	Tristan Rivière	ETH, Switzerland riviere@math.ethz.ch
Charles Fefferman	Princeton University, USA cf@math.princeton.edu	Igor Rodnianski	Princeton University, USA irod@math.princeton.edu
Ursula Hamenstaedt	Universität Bonn, Germany ursula@math.uni-bonn.de	Yum-Tong Siu	Harvard University, USA siu@math.harvard.edu
Vadim Kaloshin	University of Maryland, USA vadim.kaloshin@gmail.com	Terence Tao	University of California, Los Angeles, USA tao@math.ucla.edu
Herbert Koch	Universität Bonn, Germany koch@math.uni-bonn.de	Michael E. Taylor	Univ. of North Carolina, Chapel Hill, USA met@math.unc.edu
Izabella Laba	University of British Columbia, Canada ilaba@math.ubc.ca	Gunther Uhlmann	University of Washington, USA gunther@math.washington.edu
Richard B. Melrose	Massachusetts Inst. of Tech., USA rhm@math.mit.edu	András Vasy	Stanford University, USA andras@math.stanford.edu
Frank Merle	Université de Cergy-Pontoise, France Frank.Merle@u-cergy.fr	Dan Virgil Voiculescu	University of California, Berkeley, USA dvv@math.berkeley.edu
William Minicozzi II	Johns Hopkins University, USA minicozz@math.jhu.edu	Steven Zelditch	Northwestern University, USA zelditch@math.northwestern.edu
Clément Mouhot	Cambridge University, UK c.mouhot@dpmms.cam.ac.uk	Maciej Zworski	University of California, Berkeley, USA zworski@math.berkeley.edu
Werner Müller	Universität Bonn, Germany mueller@math.uni-bonn.de		

PRODUCTION

production@msp.org

Silvio Levy, Scientific Editor

See inside back cover or msp.org/apde for submission instructions.

The subscription price for 2020 is US \$340/year for the electronic version, and \$550/year (+\$60, if shipping outside the US) for print and electronic. Subscriptions, requests for back issues from the last three years and changes of subscriber address should be sent to MSP.

Analysis & PDE (ISSN 1948-206X electronic, 2157-5045 printed) at Mathematical Sciences Publishers, 798 Evans Hall #3840, c/o University of California, Berkeley, CA 94720-3840, is published continuously online. Periodical rate postage paid at Berkeley, CA 94704, and additional mailing offices.

APDE peer review and production are managed by EditFlow® from MSP.

PUBLISHED BY

 **mathematical sciences publishers**
nonprofit scientific publishing

<http://msp.org/>

© 2020 Mathematical Sciences Publishers

ANALYSIS & PDE

Volume 13 No. 4 2020

Estimates for the Navier–Stokes equations in the half-space for nonlocalized data	945
YASUNORI MAEKAWA, HIDEYUKI MIURA and CHRISTOPHE PRANGE	
Almost-sure scattering for the radial energy-critical nonlinear wave equation in three dimensions	1011
BJOERN BRINGMANN	
On the existence of translating solutions of mean curvature flow in slab regions	1051
THEODORA BOURNI, MAT LANGFORD and GIUSEPPE TINAGLIA	
Convex projective surfaces with compatible Weyl connection are hyperbolic	1073
THOMAS METTLER and GABRIEL P. PATERNAIN	
Stability of small solitary waves for the one-dimensional NLS with an attractive delta potential	1099
SATOSHI MASAKI, JASON MURPHY and JUN-ICHI SEGATA	
Geometric regularity for elliptic equations in double-divergence form	1129
RAIMUNDO LEITÃO, EDGARD A. PIMENTEL and MAKSON S. SANTOS	
Nonexistence of global characteristics for viscosity solutions	1145
VALENTINE ROOS	
Strichartz estimates for the Schrödinger flow on compact Lie groups	1173
YUNFENG ZHANG	
Parabolic L^p Dirichlet boundary value problem and VMO-type time-varying domains	1221
MARTIN DINDOŠ, LUKE DYER and SUKJUNG HWANG	